

Small systems in heavy-ion collisions: creating droplets of the early universe?

Ron Belmont
University of North Carolina Greensboro

Physics Department Colloquium
North Carolina State University
Raleigh, NC
26 November 2018



THE UNIVERSITY *of* NORTH CAROLINA
GREENSBORO

Утро в сосновом лесу



Утро в сосновом лесу



Outline

- Part 0: introduction
 - History lesson, units, physical constants, some basics
- Part I: the large
 - Overview of conventional heavy ion physics with large nuclei
- Part II: the small
 - Some recent results from heavy ion physics with small nuclei
- Part III: the future
 - A brief look at the next US-based experiment in heavy ions

Key points for Part 0

- A bit of a history lesson
- Natural units
- A basic sense of scale
- Fundamentals of quantum chromodynamics (QCD)
 - quarks
 - gluons
 - hadrons and confinement

Historical Perspective

"Those who do not remember George Santayana are condemned to paraphrase him." - Unknown

- 400 BCE Democritus hypothesizes atoms
- 1687 Newton publishes *Philosophiae Naturalis Principia Mathematica*
- 1900 Planck's Law
- 1905 Einstein's 4 papers
- 1911 Rutherford scattering
- 1913 Bohr atom
- 1924 de Broglie wavelength
- 1925 Heisenberg's Matrix mechanics
- 1926 Schrödinger equation
- 1927 Dirac's relativistic quantum mechanics

Historical Perspective

"Those who do not remember George Santayana are condemned to paraphrase him." - Unknown

- 1963 Gell-Mann's Quark Model (particle zoo)
- 1965 Additional degree of freedom postulated for quarks by Han and Nambu
- 1969 Deep inelastic scattering experiments prove the existence of quarks
- 1972 Color charge and basic framework of **quantum chromodynamics**
- 1973 Asymptotic Freedom discovered by Gross, Politzer, and Wilczek
- 1975 Collins and Perry formulate a QCD plasma
- 1980 Shuryak coins term quark-gluon plasma (QGP)
- 2000 RHIC is operational
- 2010 First heavy ion collisions at LHC

What we've learned so far

Three generations of matter (fermions)			
	I	II	III
mass →	2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²
charge →	2/3	2/3	2/3
spin →	1/2	1/2	1/2
name →	u up	c charm	t top
Quarks			
mass →	4.8 MeV/c ²	104 MeV/c ²	4.2 GeV/c ²
charge →	-1/3	-1/3	-1/3
spin →	1/2	1/2	1/2
name →	d down	s strange	b bottom
Leptons			
mass →	<2.2 eV/c ²	<0.17 MeV/c ²	<15.5 MeV/c ²
charge →	0	0	0
spin →	1/2	1/2	1/2
name →	v _e electron neutrino	v _μ muon neutrino	v _τ tau neutrino
Gauge bosons			
mass →	0.511 MeV/c ²	105.7 MeV/c ²	1.777 GeV/c ²
charge →	-1	-1	-1
spin →	1/2	1/2	1/2
name →	e electron	μ muon	τ tau
mass →	80.4 GeV/c ²		
charge →	±1		
spin →	1		
name →			W [±] W boson

- **Electromagnetic force:** interactions among charged particles
- **Weak force:** flavor changes, e.g. lepton decays and nuclear beta decays
- **Strong force:** binds protons and neutrons together into nuclei (residual) and quarks together into hadrons (fundamental)
- **Gravitational force:** not part of the standard model

What we've learned so far

Which view of the world is the right one? It depends!

	slow	fast
large	Classical Physics (most of our daily life is here)	Special Relativity (effects noticeable in GPS and air travel)
small	Quantum Mechanics (solid state devices are based on small stuff)	Quantum Field Theory (only self-consistent way to combine QM and SR)

Note: GR effects **also** noticeable for GPS and air travel

Physical constants and units

- High energy physics makes physical constants very easy to remember!
- Planck's constant $\hbar = 1$
 - Shows up often quantum mechanics
- Speed of light $c = 1$
 - Shows up often in special relativity
- Boltzmann's constant $k_B = 1$
 - Shows up often in thermal physics and stat mech

Typical sizes and scales for heavy ion physics

- Mass of proton = 938.3 MeV = $1.007 \text{ amu} = 1.673 \times 10^{-27} \text{ kg}$
- Typical energy = 1 GeV = $1.602 \times 10^{-10} \text{ J}$
- Typical size = 1 fm = 10^{-15} m
- Typical time = 1 fm = $3.336 \times 10^{-24} \text{ s}$
- Typical temperature = 200 MeV = $2.321 \times 10^{12} \text{ K}$

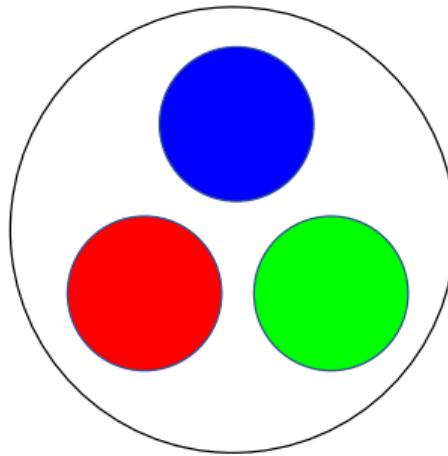
QCD as explained by approximate analogy to QED

QED	QCD	
electric charge	\leftrightarrow	color charge
electrons	\leftrightarrow	quarks
photons	\leftrightarrow	gluons
atoms	\leftrightarrow	nucleons
molecules	\leftrightarrow	nuclei

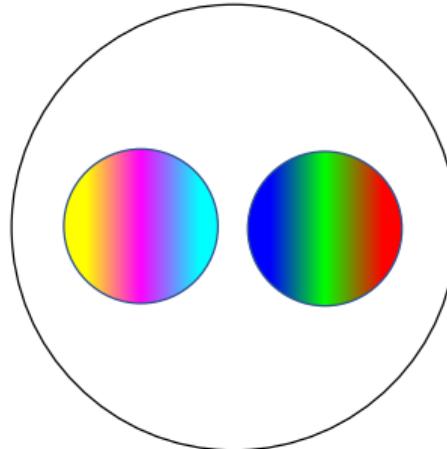
- One kind of electric charge, three kinds of color charge
- Photons do not have electric charge, gluons do have color charge
- Only one photon, eight different gluons

QCD bound states

Baryon



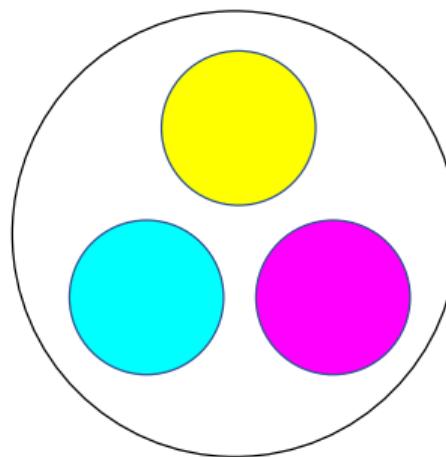
Meson



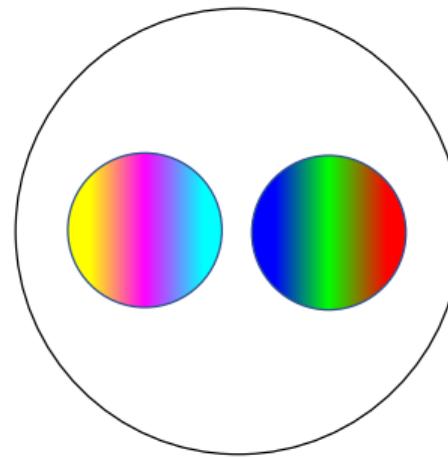
- Color-charged particles (quarks and gluons) are generically called **partons**
- QCD bound states are generically called **hadrons**, divided into **baryons** and **mesons**
- All observables must be in color singlet state—no partons can be found in isolation in nature

QCD bound states

Antibaryon



Antimeson



- Color-charged particles (quarks and gluons) are generically called **partons**
- QCD bound states are generically called **hadrons**, divided into **baryons** and **mesons**
- All observables must be in color singlet state—no partons can be found in isolation in nature

QCD Potential and confinement

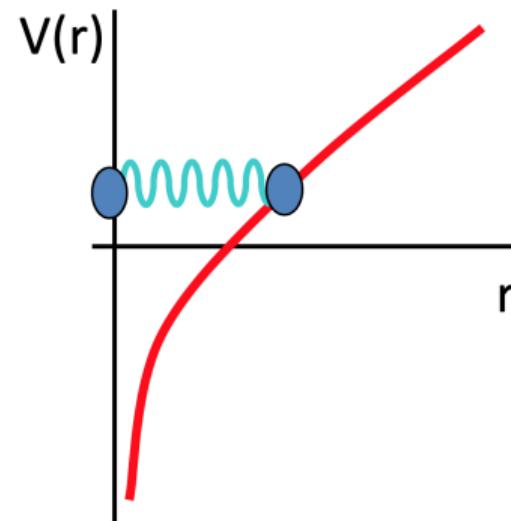
- The QED potential

$$V(r) = -\frac{\alpha_{EM}}{r}$$

- The QCD potential for $q\bar{q}$

$$V(r) = -\frac{4 \alpha_S}{3 r} + kr$$

- Coulomb part and confining part
- New pairs of quarks are created when energy exceeds mass



QCD Potential and confinement

- The QED potential

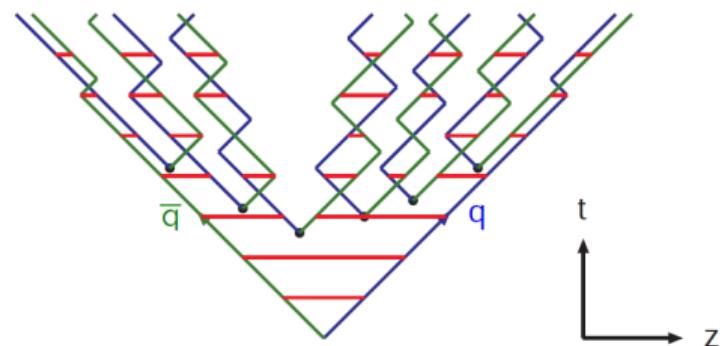
$$V(r) = -\frac{\alpha_{EM}}{r}$$

- The QCD potential for $q\bar{q}$

$$V(r) = -\frac{4 \alpha_S}{3 r} + kr$$

- Coulomb part and confining part
- New pairs of quarks are created when energy exceeds mass

Motion of quarks and antiquarks in a $q\bar{q}$ system:



gives simple but powerful picture of hadron production

The six flavors of quarks

flavor	charge	mass
down	-1/3 e	3.0–7.0 MeV
up	2/3 e	1.5–3.0 MeV
strange	-1/3 e	95 ± 25 MeV
charm	2/3 e	1.25 ± 0.09 GeV
bottom	-1/3 e	4.70 ± 0.07 GeV
top	2/3 e	174.2 ± 3.3 GeV

Summary for Part 0

- QCD is the theory of strong interactions
- Quarks and gluons (partons) are the fundamental particles
- There are three colors (and eight kinds of gluon)
- QCD exhibits confinement, meaning can only be found in bound states (hadrons)

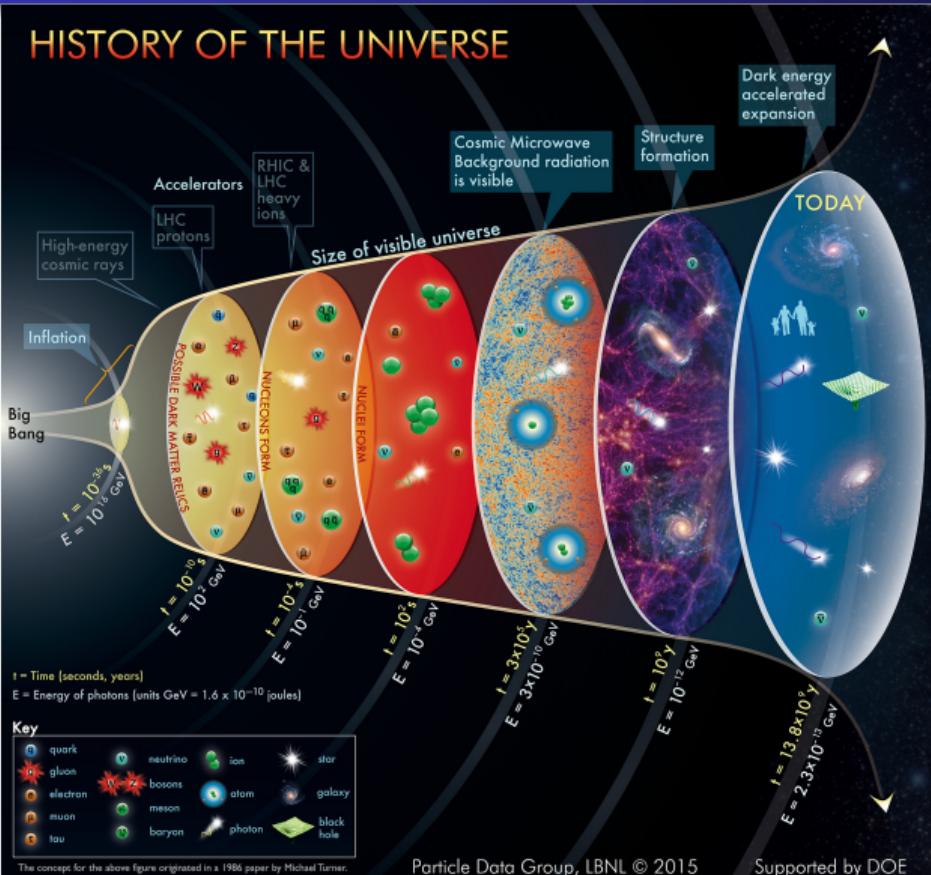
Key points for Part I

The quark gluon plasma:

- Is a phase of matter with deconfined quarks and gluons
- Existed in the very early universe
- Is created in the lab in collisions of large nuclei
 - Examples include $^{197}_{79}\text{Au} + ^{197}_{79}\text{Au}$ and $^{208}_{82}\text{Pb} + ^{208}_{82}\text{Pb}$
- Is hot and dense
- Behaves like a liquid
 - The initial-state geometry is translated into the final state

The history of the universe

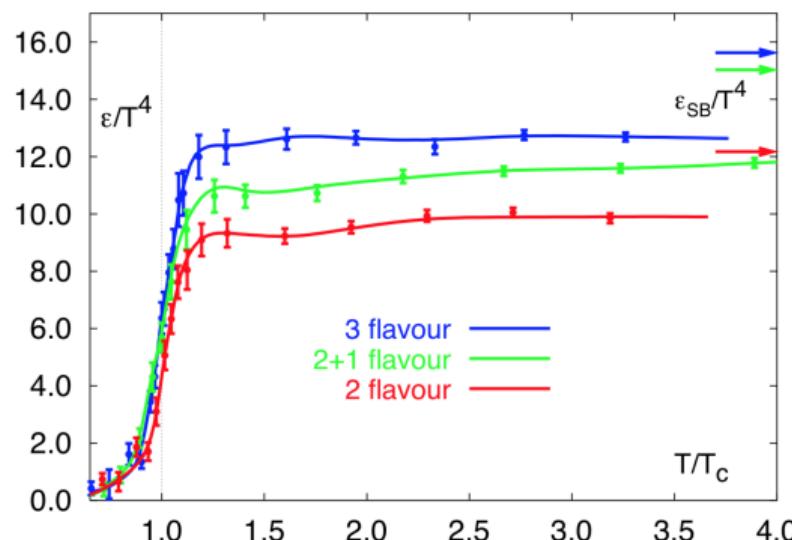
HISTORY OF THE UNIVERSE



- The early universe (few microseconds) was a quark-gluon plasma (QGP)
- The QGP is a system of deconfined quarks and gluons
- We can recreate the QGP in the lab in collisions of heavy nuclei at relativistic speeds
- Goal of heavy-ion physics: create, identify, and study the QGP

Phases of QCD matter

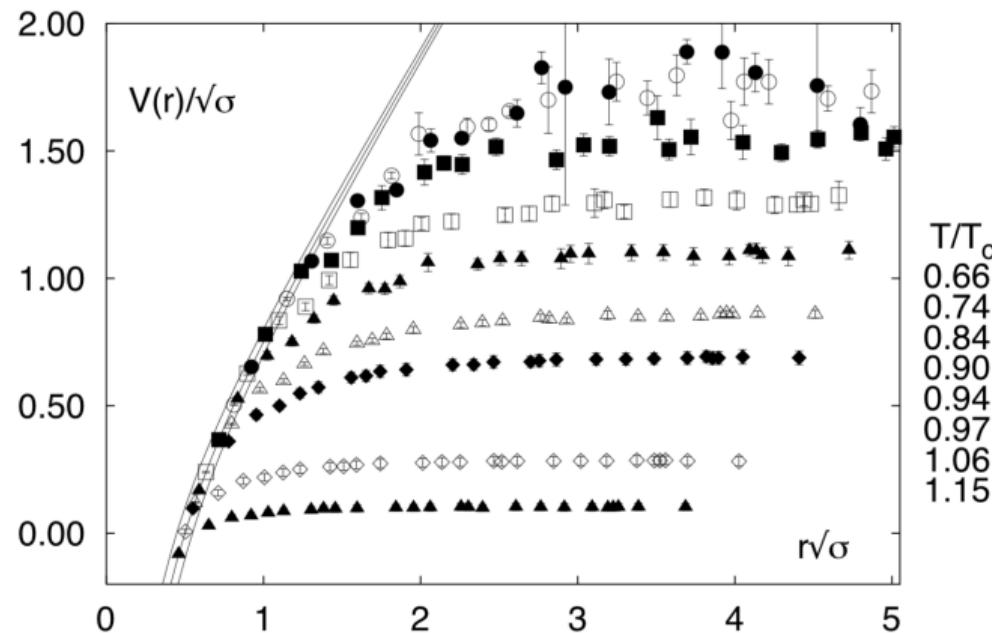
- Lattice QCD predicts a phase transition from nuclear matter to QGP
- Large increase energy density at $T_C \approx 155$ MeV due to large increase in number of degrees of freedom



$$\varepsilon_{SB} = g \frac{\pi^2}{30} T^4$$

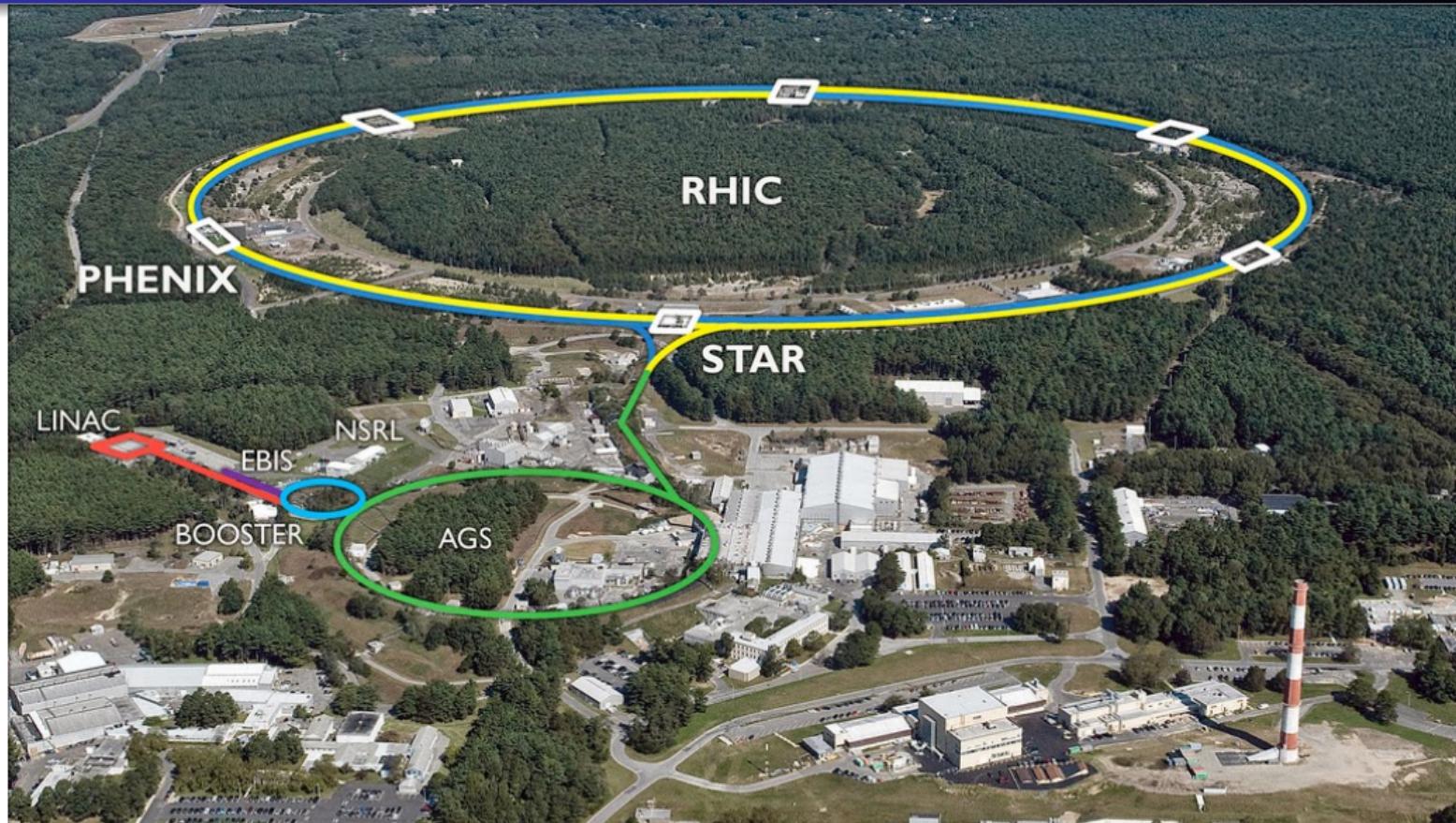
- Below T_C : $g = 3$
3 pions with spin 0
- Above T_C : $g = 37$
8 gluons with spin 1,
2 (anti)quarks with spin 1/2

Phases of QCD matter



- The confining part of gets weaker with increasing temperature
- More or less gone at the critical temperature ($T_c \approx 155$ MeV)

The Relativistic Heavy Ion Collider



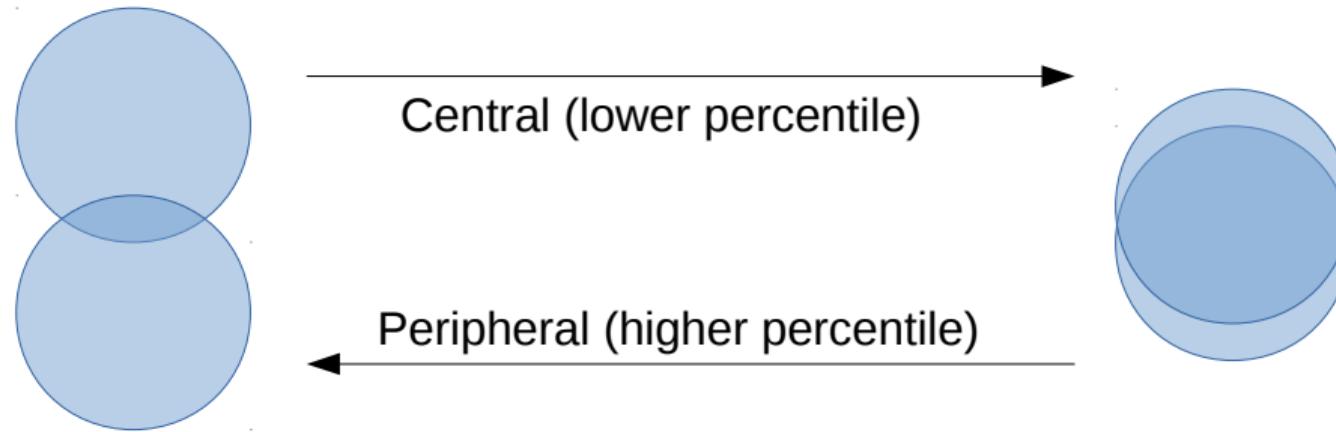
The Relativistic Heavy Ion Collider

- RHIC is the only polarized proton collider in the world
- RHIC is one of two heavy ion colliders, the other being the LHC
- RHIC is a dedicated ion collider and is designed to collide many different species of ions at many different energies—vastly more flexible than the LHC

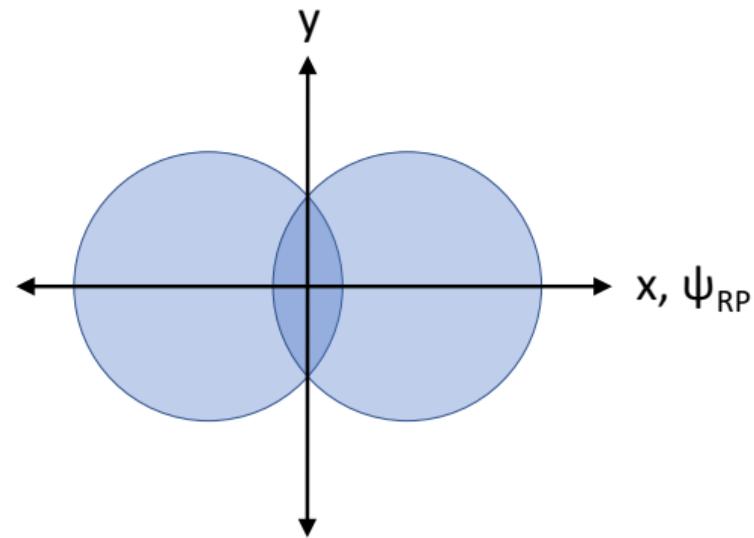
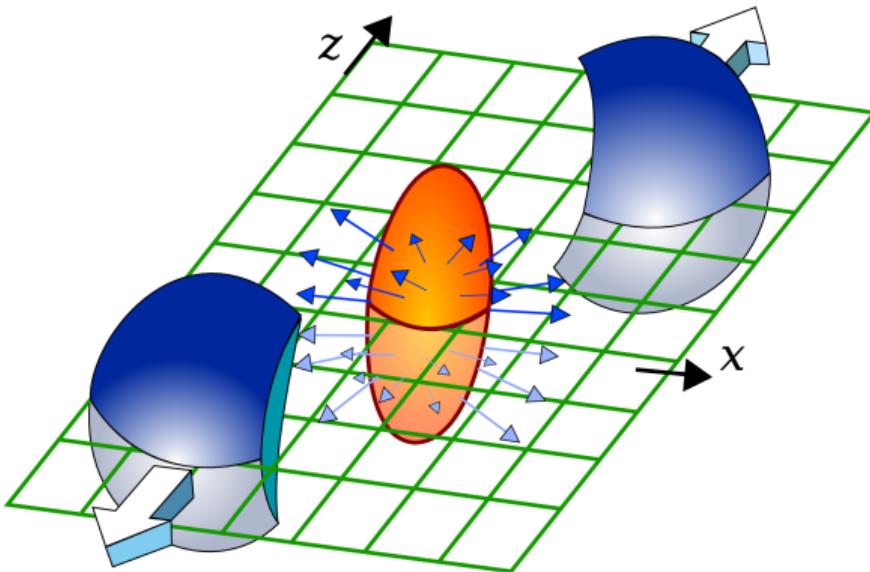
Collision Species	Collision Energies (GeV)
p↑+p↑	510, 500, 200, 62.4
p+Al	200
p+Au	200
d+Au	200, 62.4, 39, 19.6
³ He+Au	200
Cu+Cu	200, 62.4, 22.5
Cu+Au	200
Au+Au	200, 130, 62.4, 56, 39, 27, 19.6, 15, 11.5, 7.7, 5, ...
U+U	193

And lots more to come!

Centrality



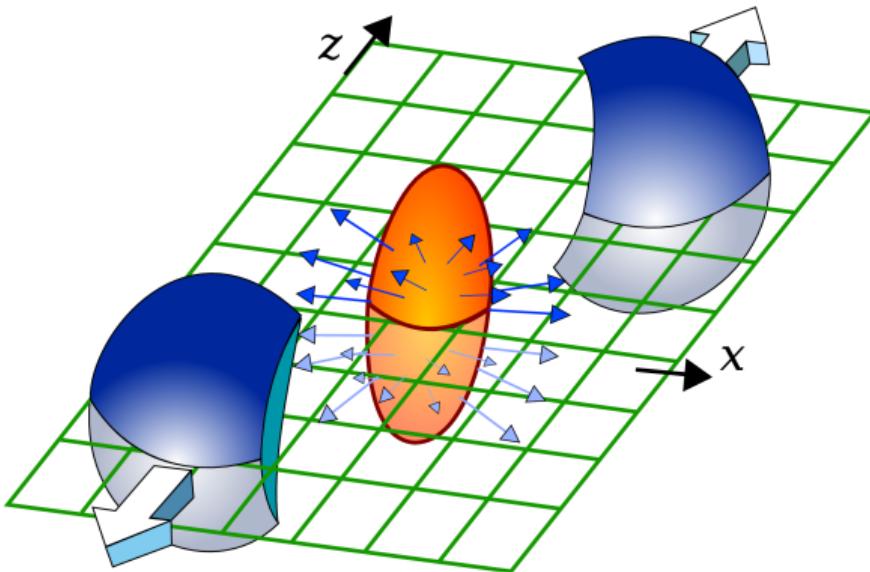
Azimuthal anisotropy measurements



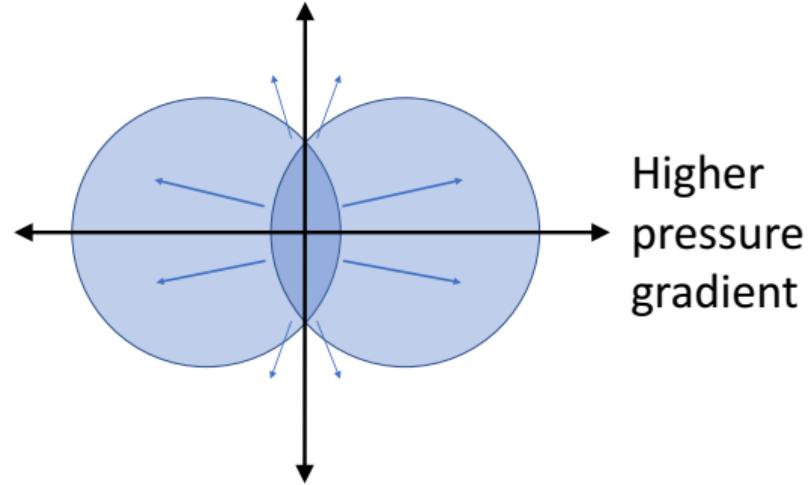
$$\frac{dN}{d\varphi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos n\varphi \quad v_n = \langle \cos n\varphi \rangle \quad \varepsilon_n = \frac{\sqrt{\langle r^2 \cos n\varphi \rangle + \langle r^2 \sin n\varphi \rangle}}{\langle r^2 \rangle}$$

- Hydrodynamics translates initial shape (ε_n) into final state distribution (v_n)
- Overlap shape approximately elliptical, expect v_2 to be the largest

Azimuthal anisotropy measurements



Lower pressure gradient



$$\frac{dN}{d\varphi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos n\varphi \quad v_n = \langle \cos n\varphi \rangle \quad \varepsilon_n = \frac{\sqrt{\langle r^2 \cos n\varphi \rangle + \langle r^2 \sin n\varphi \rangle}}{\langle r^2 \rangle}$$

- Hydrodynamics translates initial shape (ε_n) into final state distribution (v_n)
- Overlap shape approximately elliptical, expect v_2 to be the largest

Ultracold Fermi Gas

Ultracold Fermi Gases as Paradigms for Strongly Correlated Matter

J. E. Thomas, Physics Department, North Carolina State University, Raleigh NC 27695

1 Fermi Gases – A Universal Model

Ultra-cold Fermi gases with magnetically tunable interactions are a paradigm for interacting systems in nature. When tuned to a broad collisional (Feshbach) resonance, the atoms enter a unitary regime where the properties of the gas are *universal* functions of only the density n and temperature T .

2 String Theory Limit: "Perfect" Fluids

Viscosity measurement in a unitary Fermi gas is motivated by a recent conjecture derived using string theory methods. The conjecture states that over a broad class of strongly interacting fields, the ratio of shear viscosity η to the entropy density s has a universal minimum:

$$\frac{\eta}{s} \geq \frac{1}{4\pi k_B}$$

η : shear viscosity
 s : entropy density

A fluid with the minimum ratio is referred to as a **perfect fluid**. In a Fermi gas, we are able to measure both the entropy and the viscosity, connecting thermodynamics and hydrodynamics in strongly correlated systems.

3 Measurement of Entropy and Energy

Energy is measured from the cloud size by exploiting the virial theorem, which holds for a unitary gas. Entropy is measured from the cloud size after an adiabatic sweep to a weakly interacting regime.

4 Measurement of Quantum Viscosity

Viscosity η has a dimension of momentum time. In a unitary gas, the natural momentum is of order $k_B T$ and the natural area is the πk_B^2 collision cross section $4\pi k_B^2 n$. Hence, $\eta \propto k_B^2 n$. The Fermi momentum sets the length scale a_F with the interparticle spacing. Thus, the natural scale of viscosity is hL^2/n .

Measuring the aspect ratio of an elliptically expanding atom cloud as a function of time after release from the optical trap determines the shear viscosity at high temperature (red). For low temperatures, we measure the growth rate of a radial breathing mode (blue), which smoothly joins:

5 Optical Control of Feshbach Resonances

Optical fields, tuned to resonance on a singlet molecular transitions allow “dark-state” control of interactions.

- “Designer” interactions
 - control width of narrow FB resonance
 - control interactions in 3-state systems
 - Random interactions in space and time
 - Non-equilibrium Fermi gases—Fermi time
 - “Designed” dispersion relations

6 Quasi-Two-Dimensional Fermi Gases

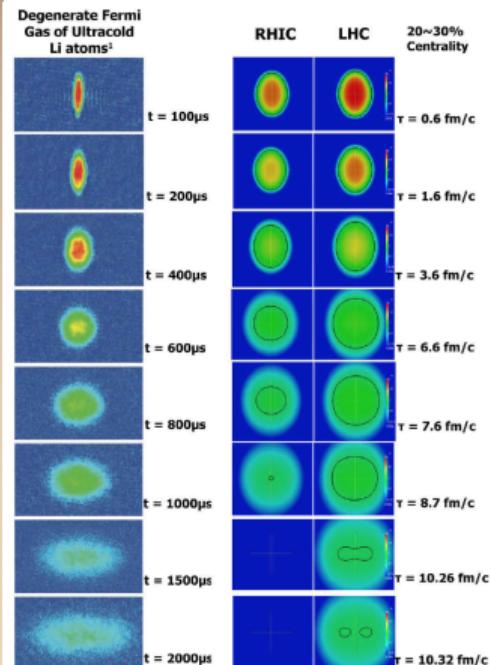
We create a 2D gas by trapping atoms in a CO_2 laser standing-wave, with a few hundred atoms per site, $5.3 \mu\text{m}$ spacing.

7 Shock Waves

Colliding Fermi gas clouds—LHC!

- Nonlinear hydrodynamics in strongly interacting quantum matter

J. Thomas et al, second floor



Ultracold Fermi Gas

J. Thomas et al, second floor

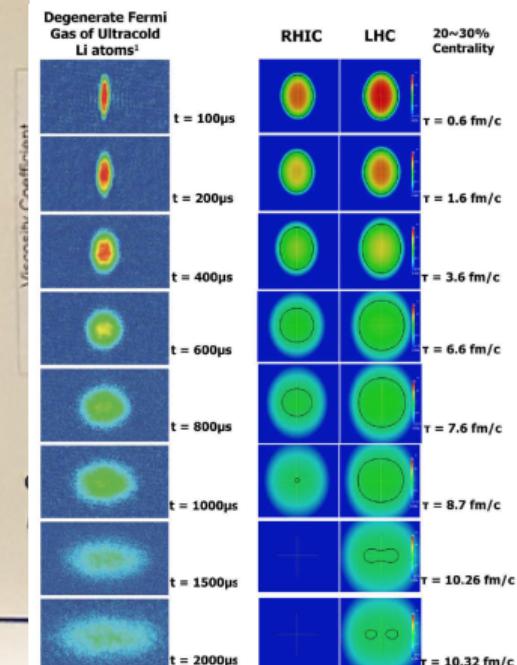
2 String Theory Limit: “Perfect” Fluids

Viscosity measurement in a unitary Fermi gas is motivated by a recent conjecture derived using string theory methods. The conjecture states that for a broad class of strongly interacting fields, the ratio of shear viscosity η to the entropy density s has a universal minimum:

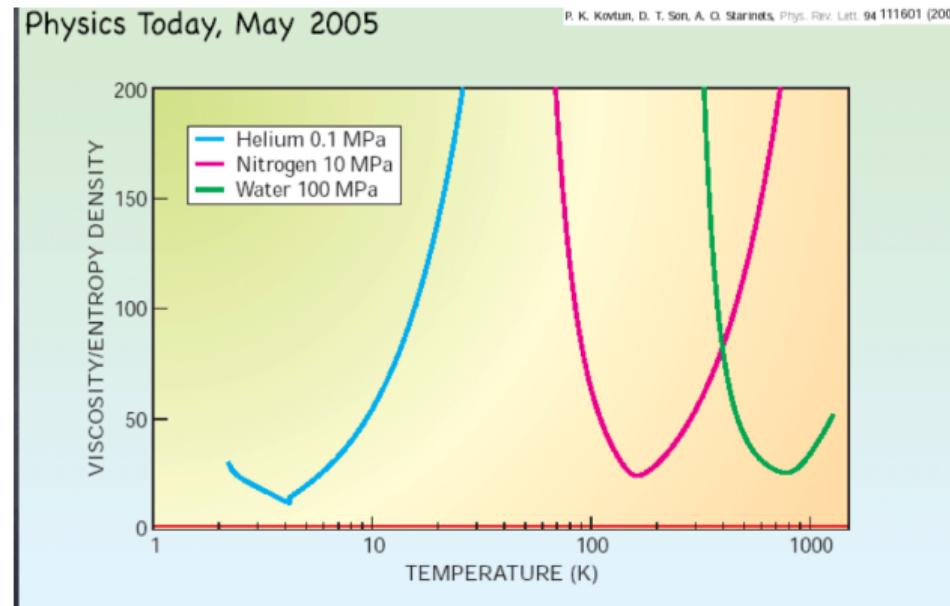
$$\frac{\eta}{s} \geq \frac{1}{4\pi} \frac{\hbar}{k_B}$$

η : Shear viscosity
 s : Entropy density

A fluid with the minimum ratio is referred to as a **perfect** fluid. In a Fermi gas, we are able to measure both the entropy and the viscosity, connecting **thermodynamics** and **hydrodynamics** in strongly correlated systems.



Viscosity (over entropy density)



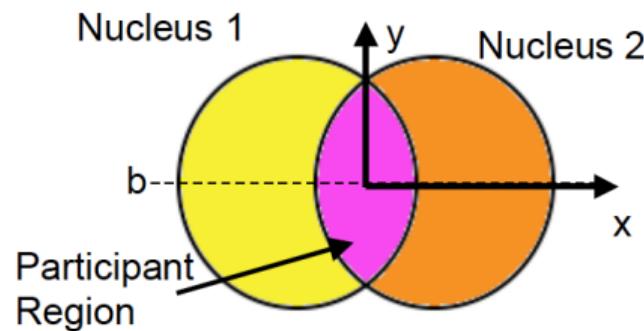
Ultracold Fermi gases: few times the lower bound

Quark-gluon plasma: very close to the lower bound

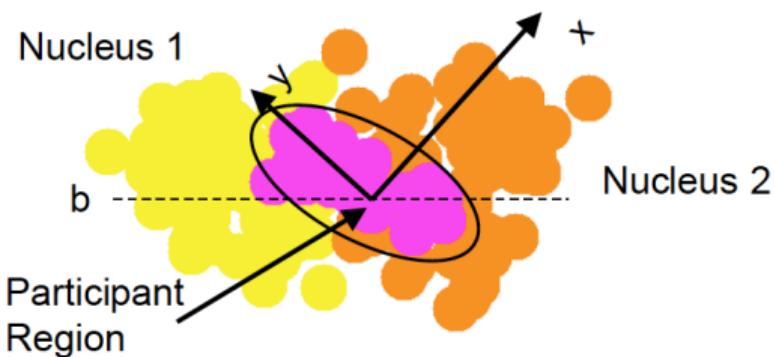
Important discovery in 2005

G. Roland, PHOBOS Plenary, Quark Matter 2005

Standard Eccentricity



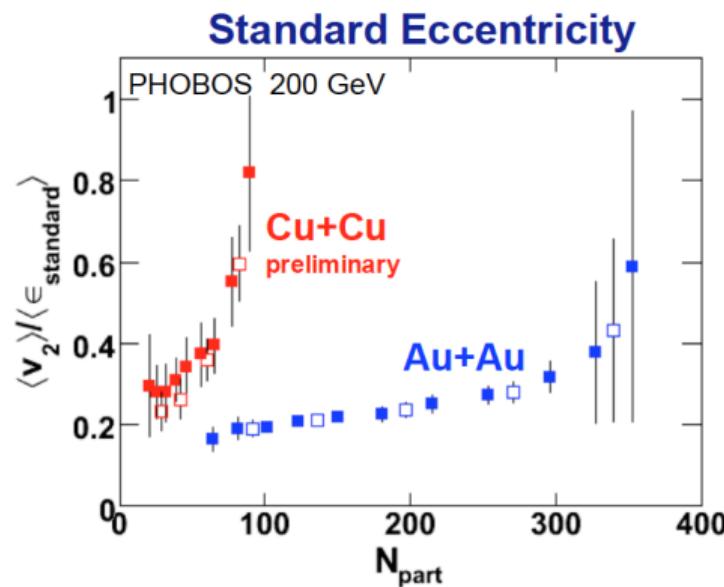
Participant Eccentricity



A nucleus isn't just a sphere

Important discovery in 2005

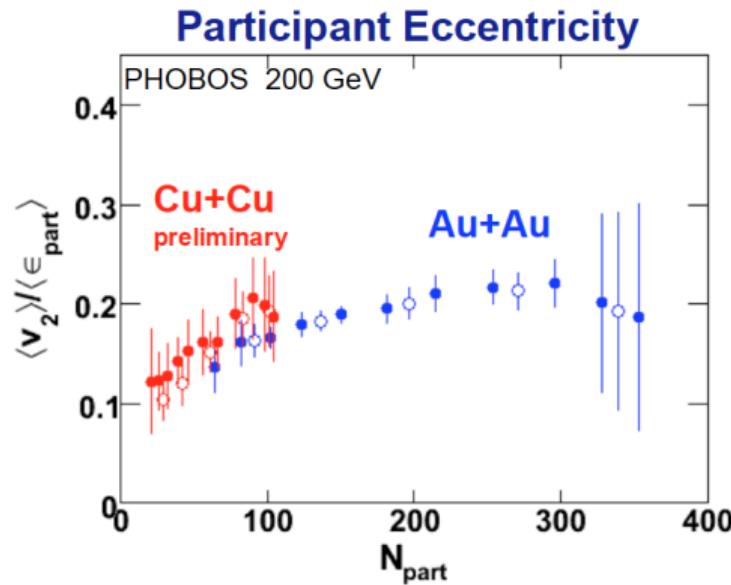
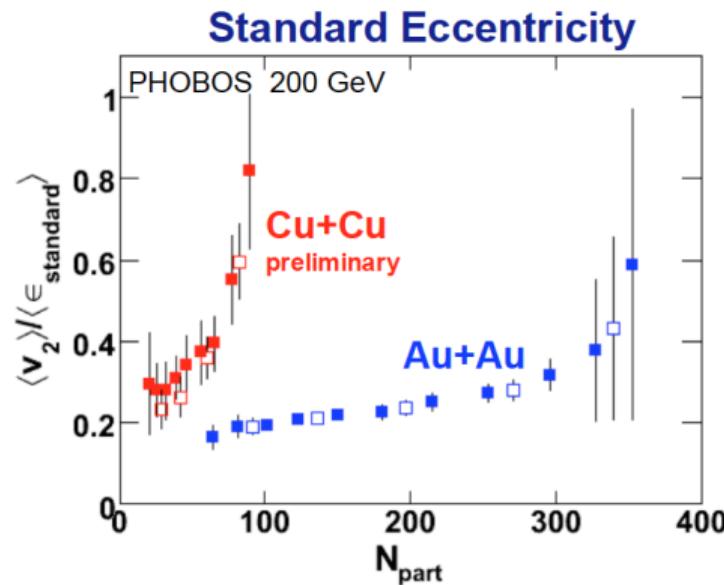
G. Roland, PHOBOS Plenary, Quark Matter 2005



A nucleus isn't just a sphere

Important discovery in 2005

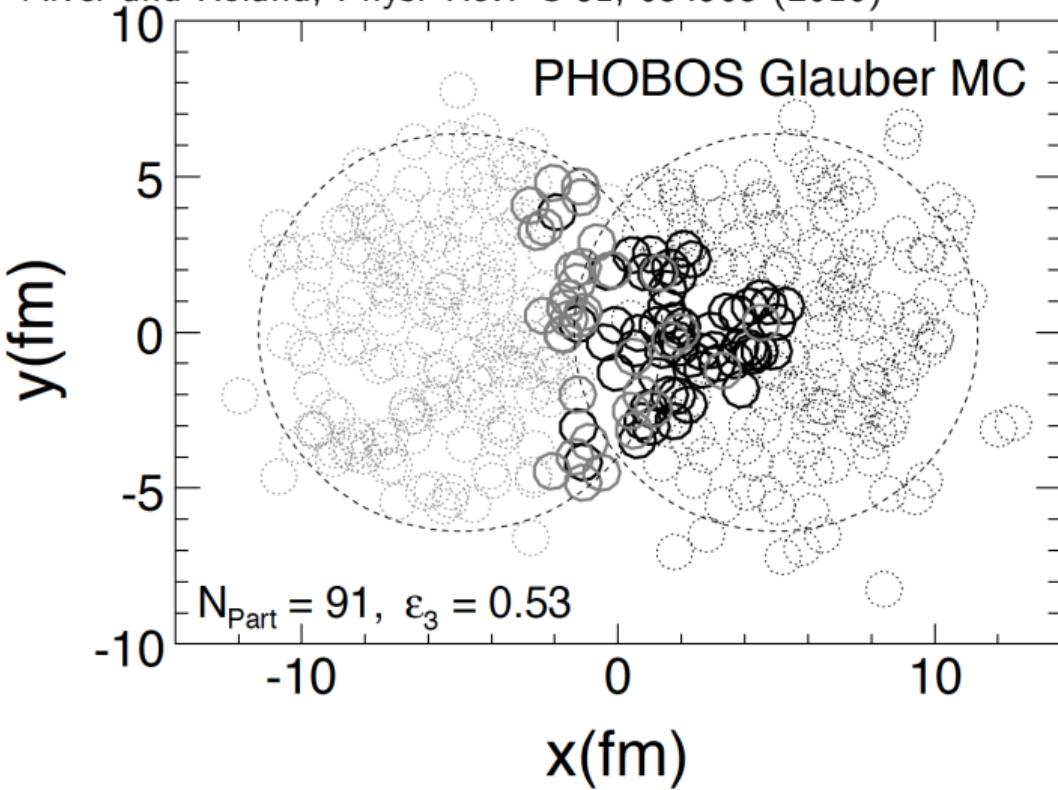
G. Roland, PHOBOS Plenary, Quark Matter 2005



A nucleus isn't just a sphere

Important discovery in 2010

Alver and Roland, Phys. Rev. C 81, 054905 (2010)



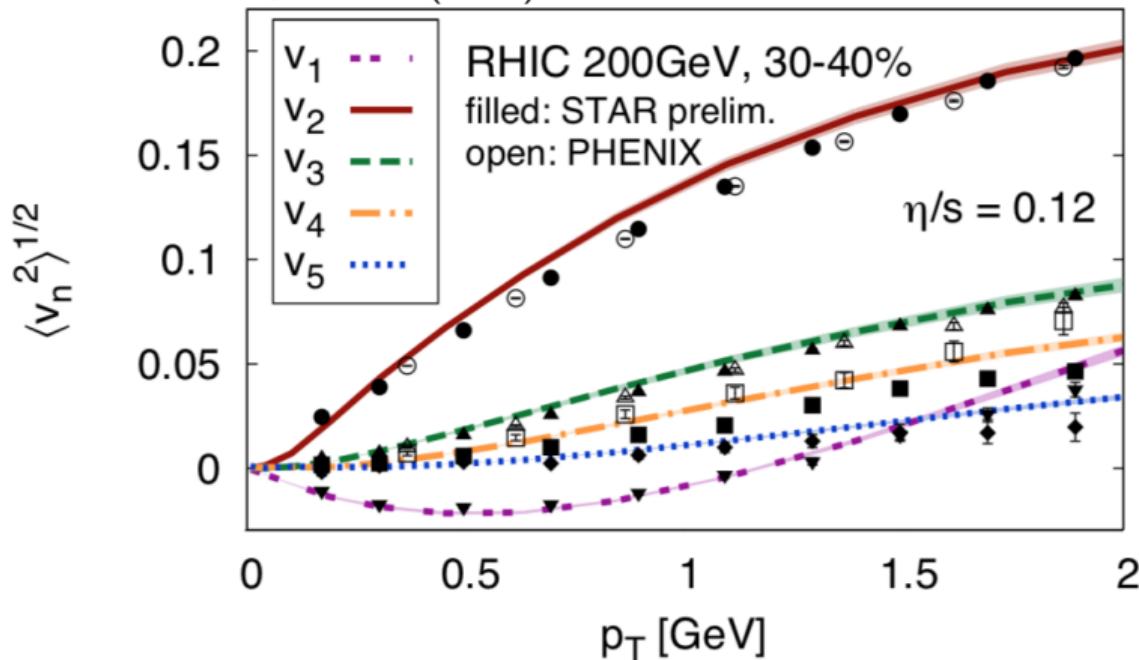
Nucleon fluctuations can produce non-zero ε_n for odd n

Symmetry planes ψ_n can be different for different harmonics

$$\varphi = \phi_{\text{lab}} - \psi_n$$

Data and theory for v_n

Gale et al, Phys. Rev. Lett. 110, 012302 (2013)

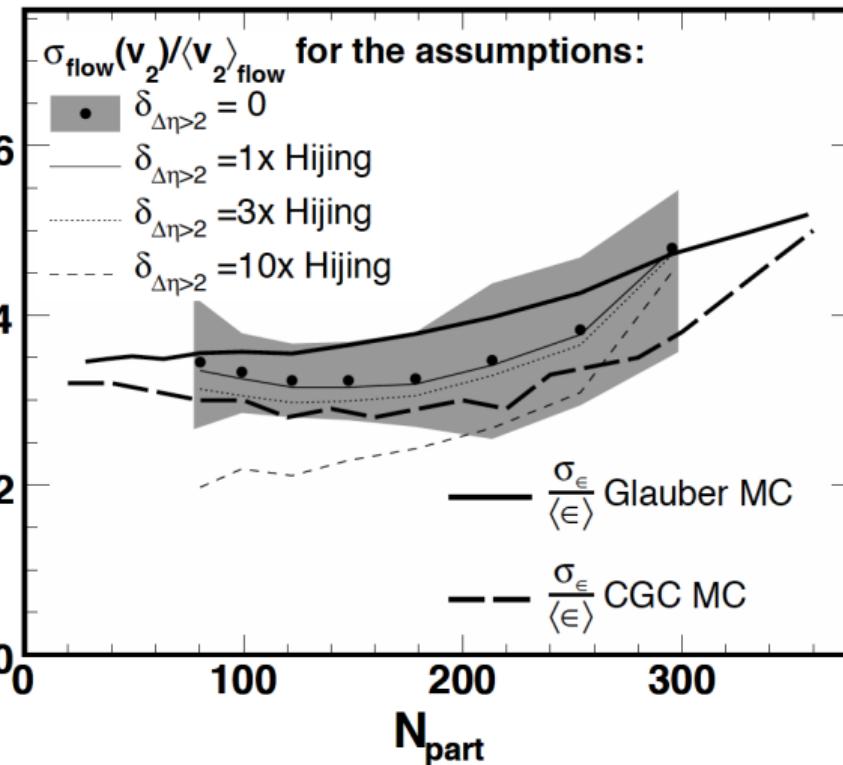


$$\frac{dN}{d\varphi} \propto 2v_1 \cos \varphi + 2v_2 \cos 2\varphi + 2v_3 \cos 3\varphi + 2v_4 \cos 4\varphi + 2v_5 \cos 5\varphi$$

Fluctuations in large systems

PHOBOS, Phys. Rev. C 81, 034915 (2010)

Relative Fluctuations



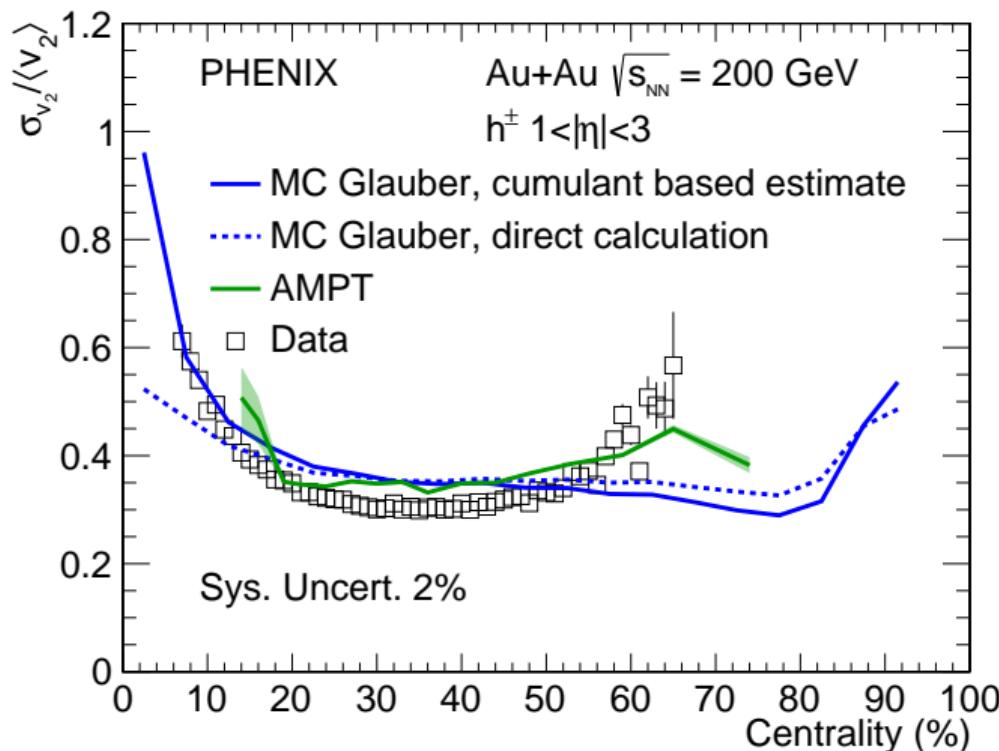
Fluctuations should also be translated, so measure $\sigma_{v2}/\langle v_2 \rangle$

$$|\eta| < 1$$

Generally good agreement with models of initial geometry

Fluctuations in large systems

PHENIX, arXiv:1804.10024 (submitted to Phys. Rev. C)



Fluctuations should also be translated, so measure $\sigma_{v_2}/\langle v_2 \rangle$

$$1 < |\eta| < 3$$

Generally good agreement with models of initial geometry

Summary for Part I

The quark-gluon plasma:

- Is a phase of matter with deconfined quarks and gluons
- Existed in the very early universe
- Is created in the lab in collisions of large nuclei
- Evolves hydrodynamically
 - The initial-state geometry is translated into the final state

Key points for Part II

- A major component of heavy ion physics nowadays is “small systems”
 - Nuclear collisions of small+large or even small+small
 - Examples include d+Au, p+Pb, and even p+p
- The matter created in small systems looks a lot like the matter in large systems
- Roughly speaking, two competing pictures
 - QGP droplets being created in small systems
 - Initial state effects from color-glass condensate (CGC)

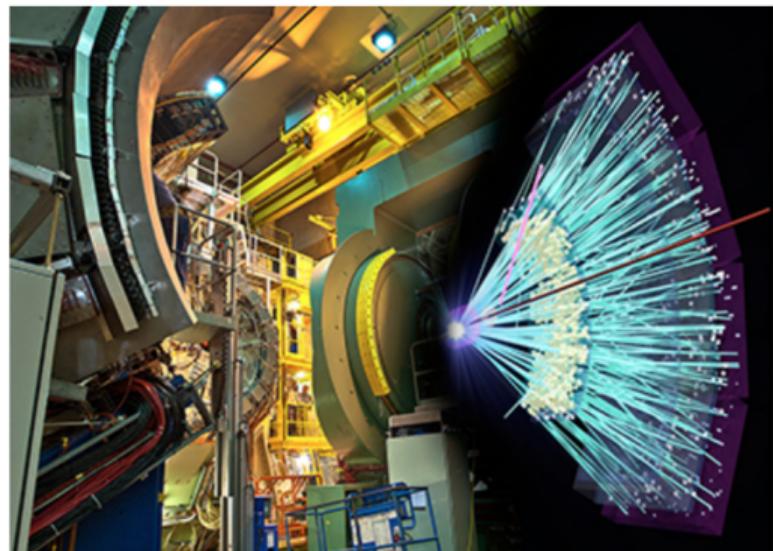
Media Attention

Physics World, September 22, 2017

Phys.org, September 18, 2017

Collider serves up drop of primordial soup

Sep 22, 2017



Tiny drop: PHENIX and reconstructed particle tracks from a QGP

PHENIX colleague **Ron Belmont** of the University of Colorado says it is still possible that the elliptical emission they have observed is due not to the formation of tiny QGPs but instead down to nuclear properties prior to collision. When accelerated close to light speed, time slows down for the heavy nuclei, which means, according to quantum chromodynamics, that they appear as a dense wall of gluons. The fact that these condensates are thicker in the centre of the nuclei might explain why particles generated in the collisions are not emitted in random directions, he says.

A very brief history of recent heavy ion physics

- 1980s and 1990s—AGS and SPS... QGP at SPS!
- Early 2000s—QGP at RHIC! No QGP at SPS. d+Au as control.
- Mid-late 2000s—Detailed, quantitative studies of strongly coupled QGP. d+Au as control.
- 2010—Ridge in high multiplicity p+p (LHC)! Probably CGC!
- Early 2010s—QGP in p+Pb!
- Early 2010s—QGP in d+Au!
- Mid 2010s and now-ish—QGP in high multiplicity p+p? QGP in mid-multiplicity p+p??
QGP in d+Au even at low energies???

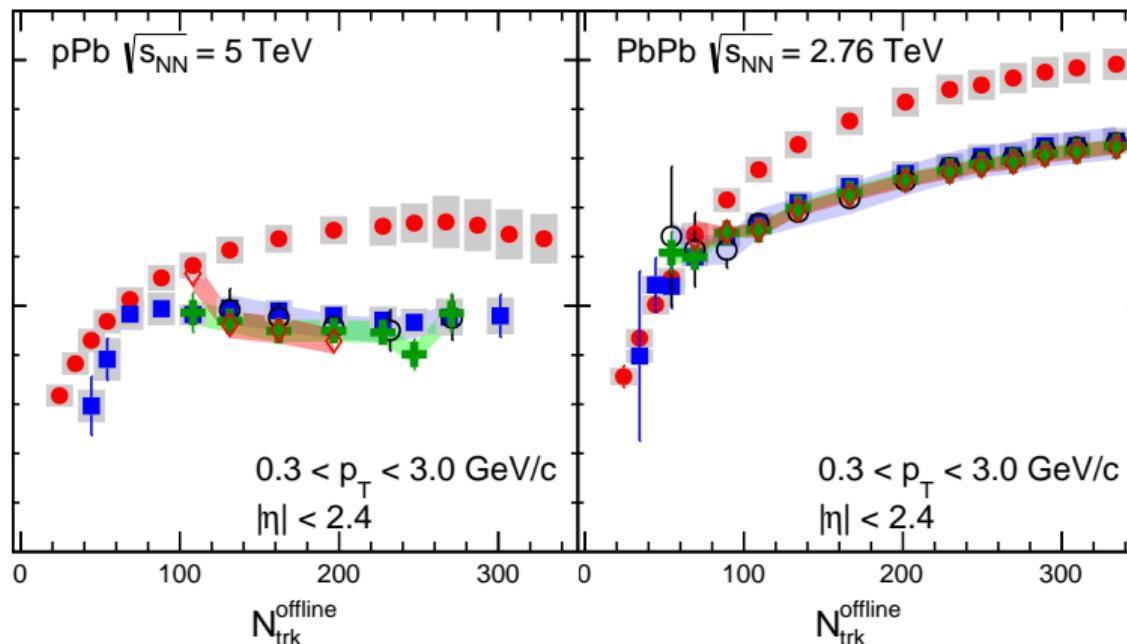
A very brief history of recent heavy ion physics

- 1980s and 1990s—AGS and SPS... QGP at SPS!
- Early 2000s—QGP at RHIC! No QGP at SPS. d+Au as control.
- Mid-late 2000s—Detailed, quantitative studies of strongly coupled QGP. d+Au as control.
- 2010—Ridge in high multiplicity p+p (LHC)! Probably CGC!
- Early 2010s—QGP in p+Pb!
- Early 2010s—QGP in d+Au!
- Mid 2010s and now-ish—QGP in high multiplicity p+p? QGP in mid-multiplicity p+p??
QGP in d+Au even at low energies???

“Twenty years ago, the challenge in heavy ion physics was to find the QGP. Now, the challenge is to not find it.” —Jürgen Schukraft, QM17

Multiparticle correlations in large and small systems

CMS, Phys. Lett. B 765 (2017) 193-220

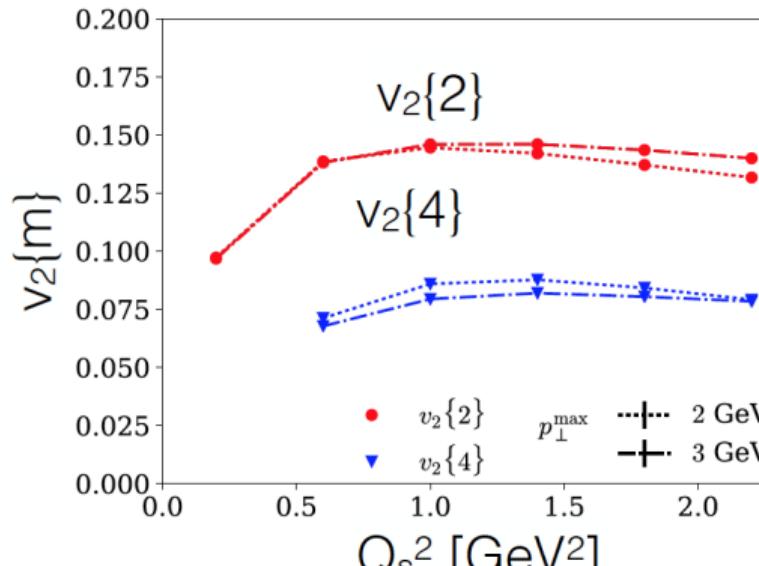


$$\begin{aligned} v_n^2\{2\} &= \langle v_n \rangle^2 + \sigma^2 \\ v_n^2\{4, 6, 8\} &\approx \langle v_n \rangle^2 - \sigma^2 \end{aligned}$$

- Multiparticle correlations reflect global correlation from geometry in Pb+Pb , Au+Au , Cu+Cu , etc
- The $p+\text{Pb}$ has a remarkably similar pattern as the Pb+Pb

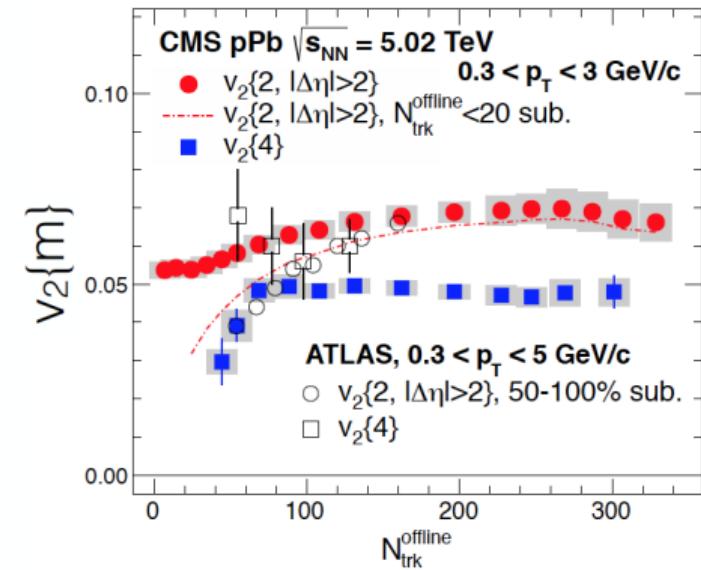
CGC results on small systems

Mark Mace, QM18



Dusling, MM, Venugopalan PRL 120 (2018)

Phys. Rev. Lett. 120, 042002 (2018)

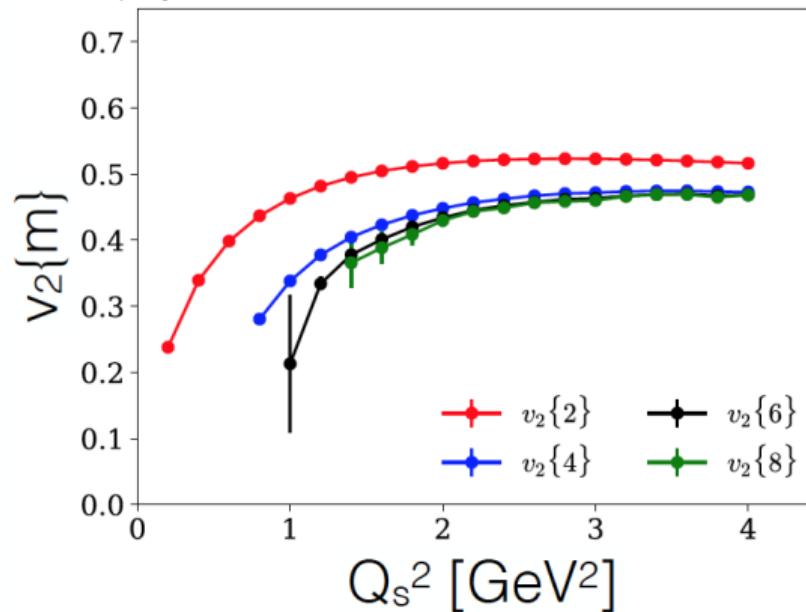


CMS PLB 724 (2013) 213

- Can reproduce $v_2\{2\}$ and $v_2\{4\}$
- Disagreement with data by a factor of 2, but qualitative features match

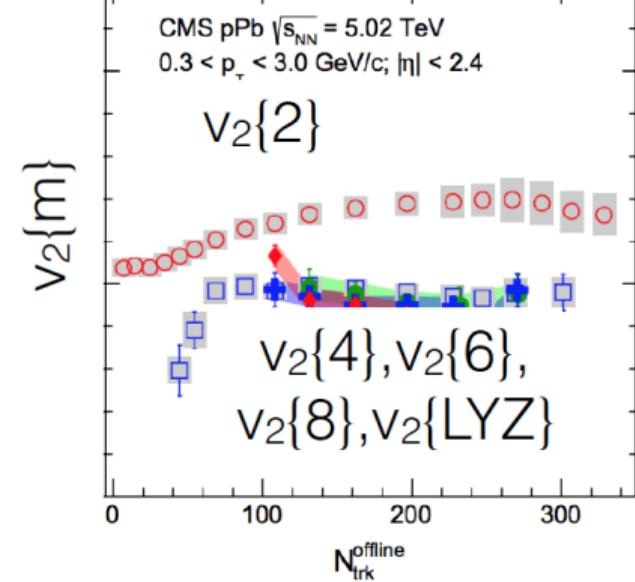
CGC results on small systems

Mark Mace, QM18



Dusling, MM, Venugopalan PRL 120 (2018)

Phys. Rev. Lett. 120, 042002 (2018)



CMS PRL 115 (2015) 012301

- Abelian calculations can produce $v_2\{2\}$, $v_2\{4\}$, $v_2\{6\}$, $v_2\{8\}$
- Disagreement with data by factor of 5, but qualitative features match

Intermission

Small systems geometry scan

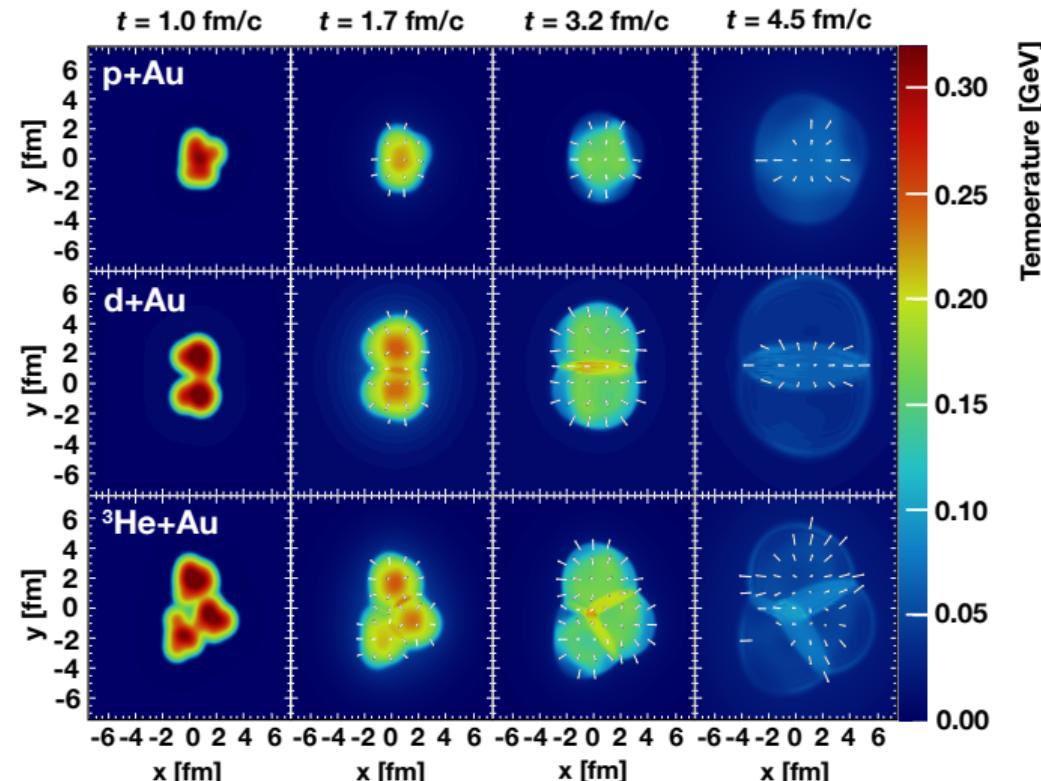
Testing hydro by controlling system geometry

arXiv:1805.02973, submitted to Nature Physics

- Hydrodynamics translates initial geometry into final state
- Test hydro hypothesis by varying initial geometry

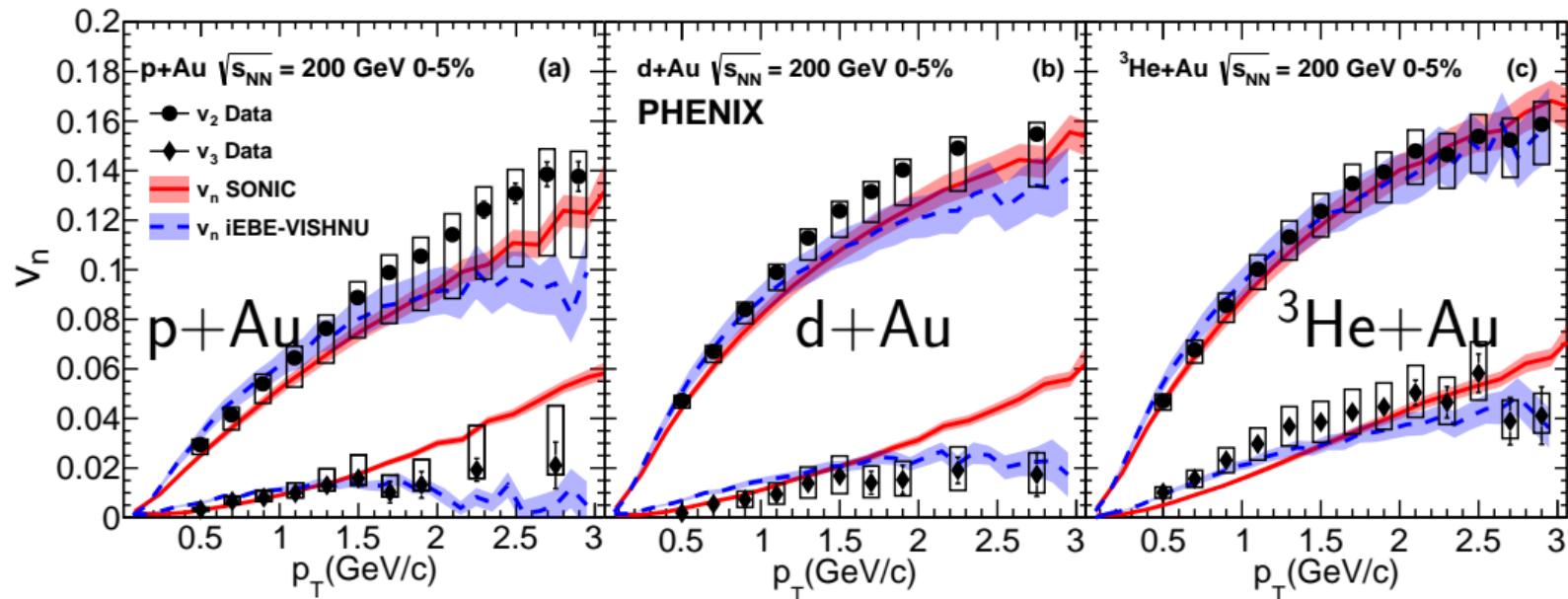
	ε_2	ε_3
p+Au	0.24	0.16
d+Au	0.57	0.17
$^3\text{He}+\text{Au}$	0.48	0.23

- $\varepsilon_2^{\text{p+Au}} < \varepsilon_2^{\text{d+Au}} \approx \varepsilon_2^{\text{He+Au}}$
- $\varepsilon_3^{\text{p+Au}} \approx \varepsilon_3^{\text{d+Au}} < \varepsilon_3^{\text{He+Au}}$



Testing hydro by controlling system geometry

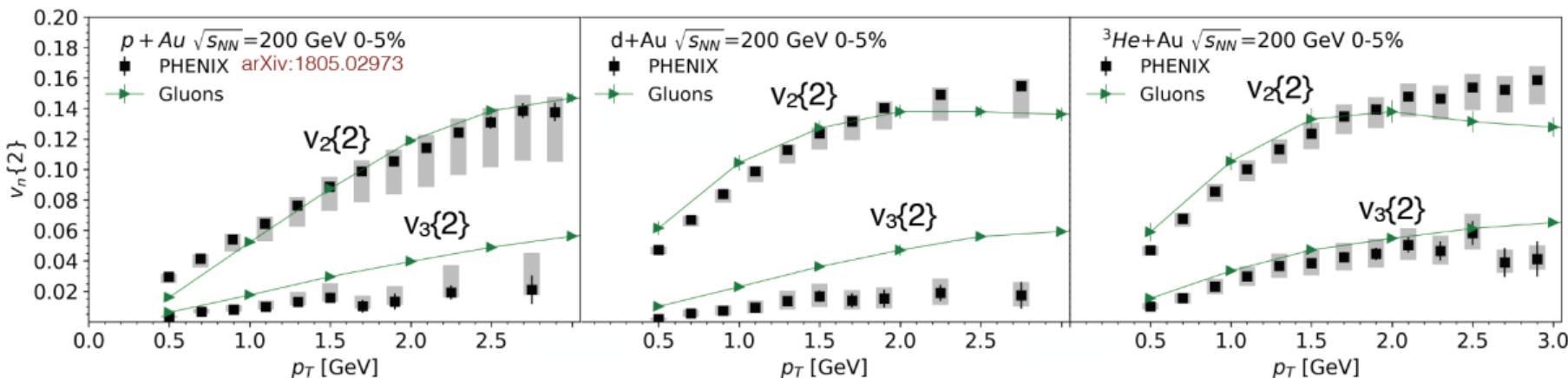
arXiv:1805.02973, submitted to Nature Physics



- v_2 and v_3 vs p_T described very well by hydro in all three systems
—Suggests QGP droplets in hydro evolution(?)

CGC results on small systems

M. Mace et al, Phys. Rev. Lett. 121, 052301 (2018)

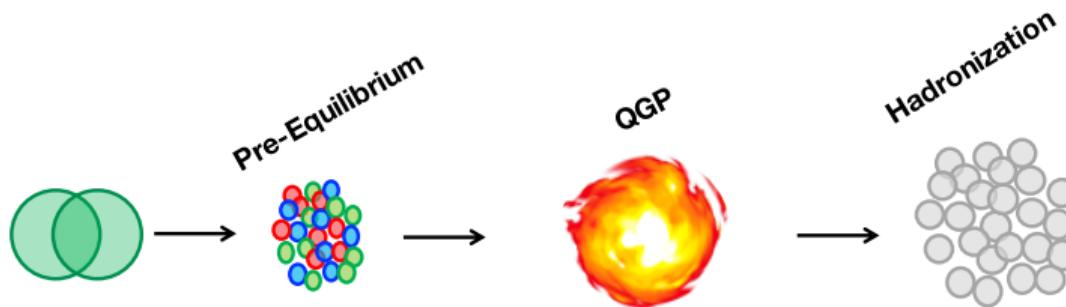


- v_2 is quite well-described
- v_3 is in the right ballpark, though hydro does better

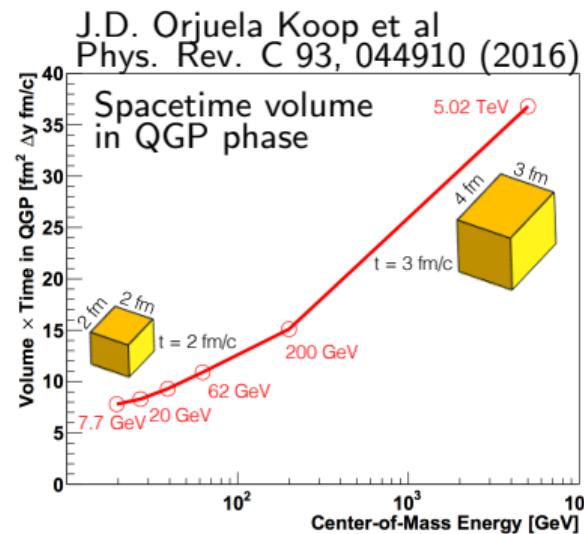
Intermission

Small systems beam energy scan

Testing hydro by controlling system size and life time



- Standard picture for A+A:
QGP in hydro evolution
- What about small systems?
And lower energies?
- Use collisions species and
energy to control system
size, test limits of hydro
applicability



d+Au beam energy scan

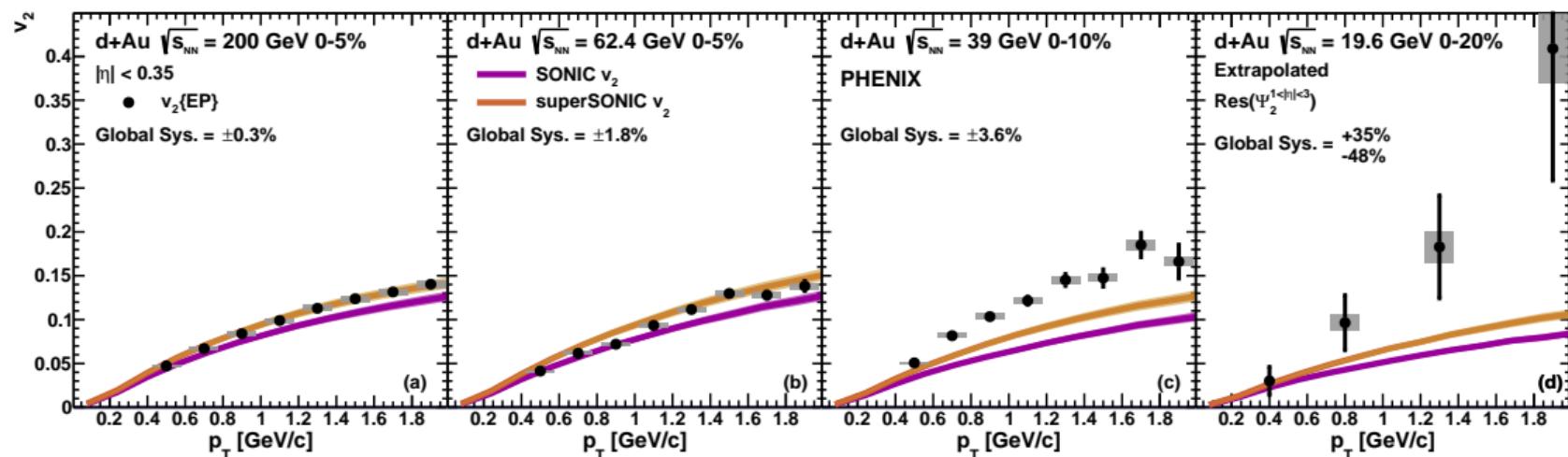
Phys. Rev. C 96, 064905 (2017)

200 GeV

62.4 GeV

39 GeV

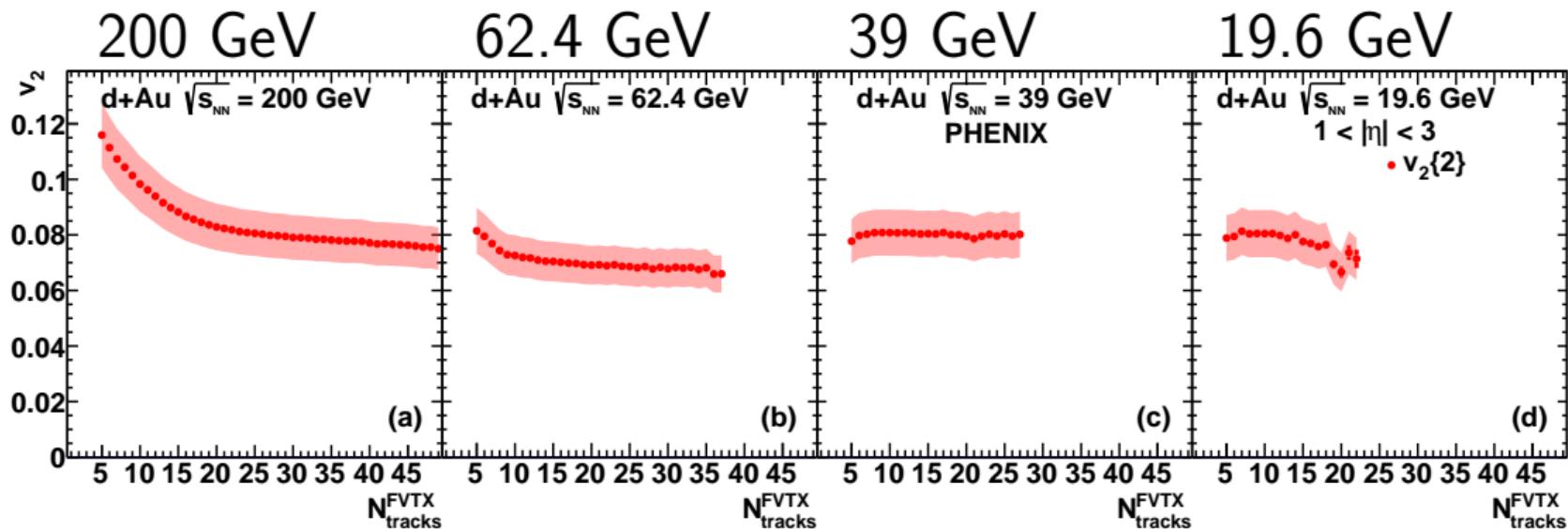
19.6 GeV



- Hydro theory agrees with higher energies very well, underpredicts lower energies
 - Breakdown of hydro?
 - Predominance of other correlations?

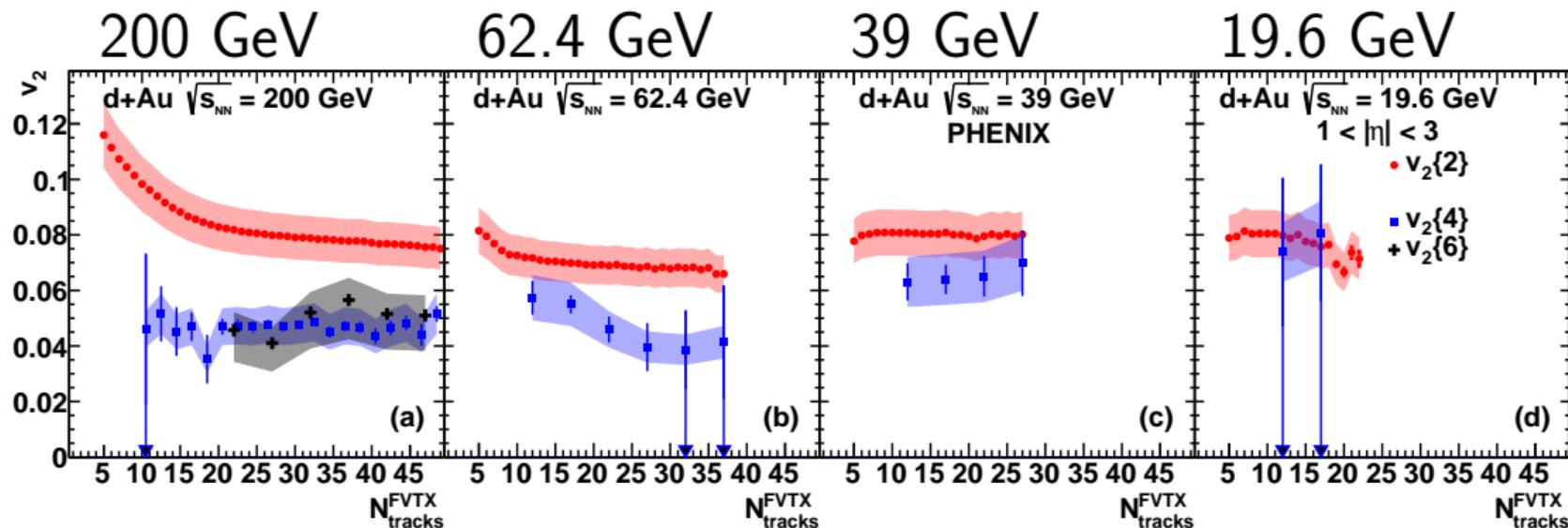
d+Au beam energy scan

Phys. Rev. Lett. 120, 062302 (2018)



d+Au beam energy scan

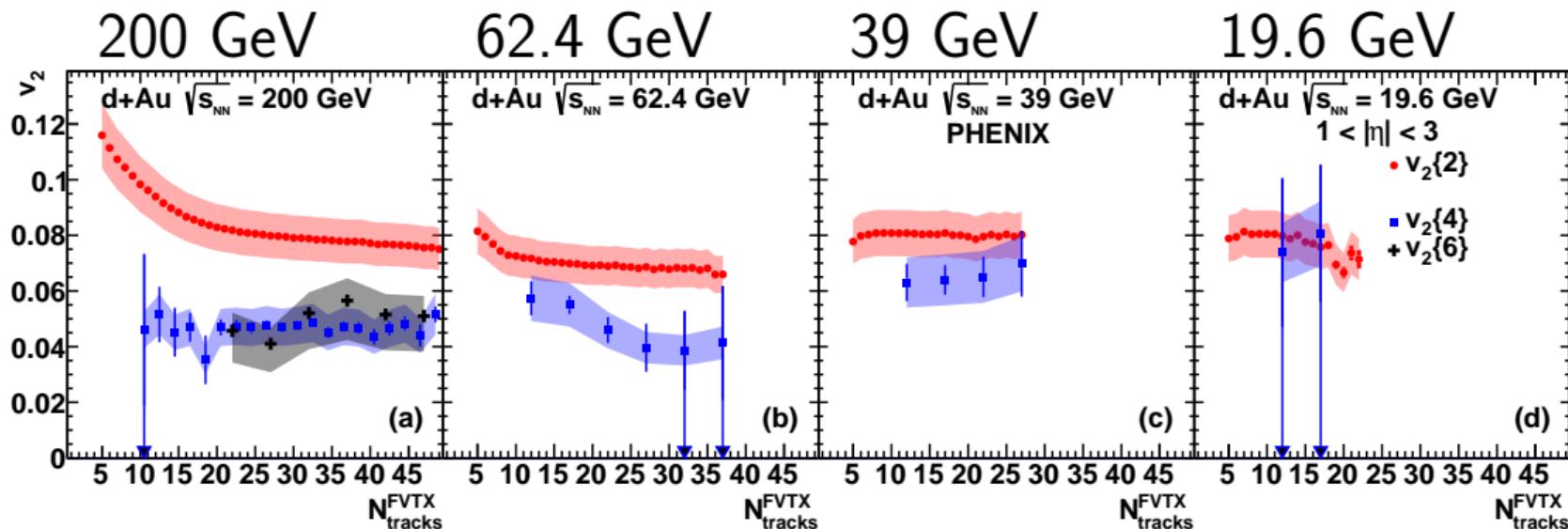
Phys. Rev. Lett. 120, 062302 (2018)



- Measurement of $v_2\{6\}$ in d+Au at 200 GeV and $v_2\{4\}$ in d+Au at all energies

d+Au beam energy scan

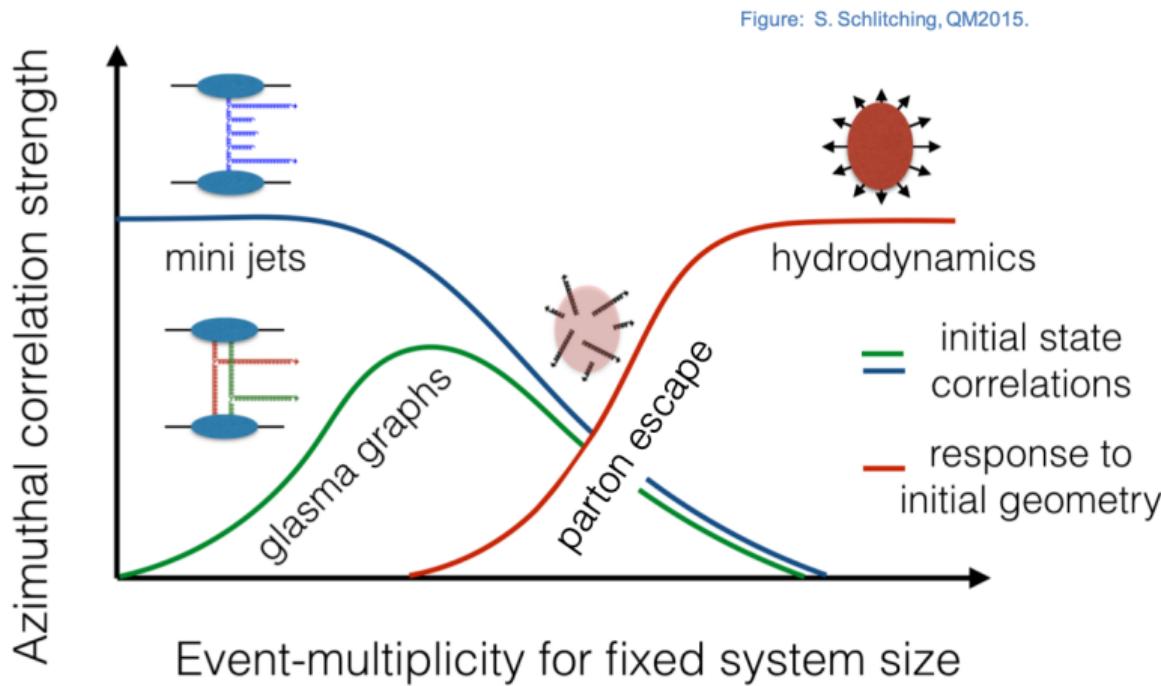
Phys. Rev. Lett. 120, 062302 (2018)



- Measurement of $v_2\{6\}$ in d+Au at 200 GeV and $v_2\{4\}$ in d+Au at all energies
- No theory comparisons available (yet...)

Competing theories

Where are we?



Summary for Part II

- Small systems is a hot topic in heavy ion physics
- We've even gotten some media attention for it
- The system created in small systems looks a like the one in large systems
- Two competing pictures: CGC and QGP
 - QGP picture ahead by points, but no knockout yet

Intermission

Part III: the future

The 2015 Long Range Plan for Nuclear Science

https://www.science.energy.gov/~/media/np/nsac/pdf/2015LRP/2015_LRPNS_091815.pdf

Recommendation I: The progress achieved under the guidance of the 2007 Long Range Plan has reinforced U.S. world leadership in nuclear science. The highest priority in this 2015 Plan is to capitalize on the investments made.

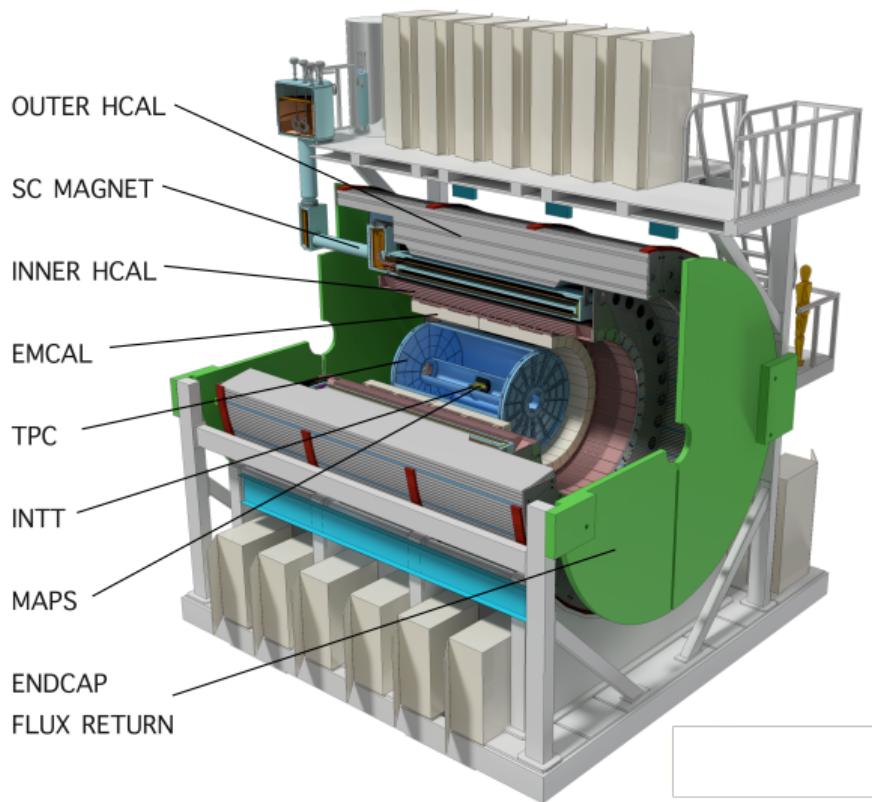
—CEBAF, FRIB, Symmetries & Neutrinos, **RHIC** (BES II & **sPHENIX**)

Recommendation II: We recommend the timely development and deployment of a U.S.-led ton-scale neutrinoless double beta decay experiment.

Recommendation III: We recommend a high-energy high-luminosity polarized EIC as the highest priority for new facility construction following the completion of FRIB.

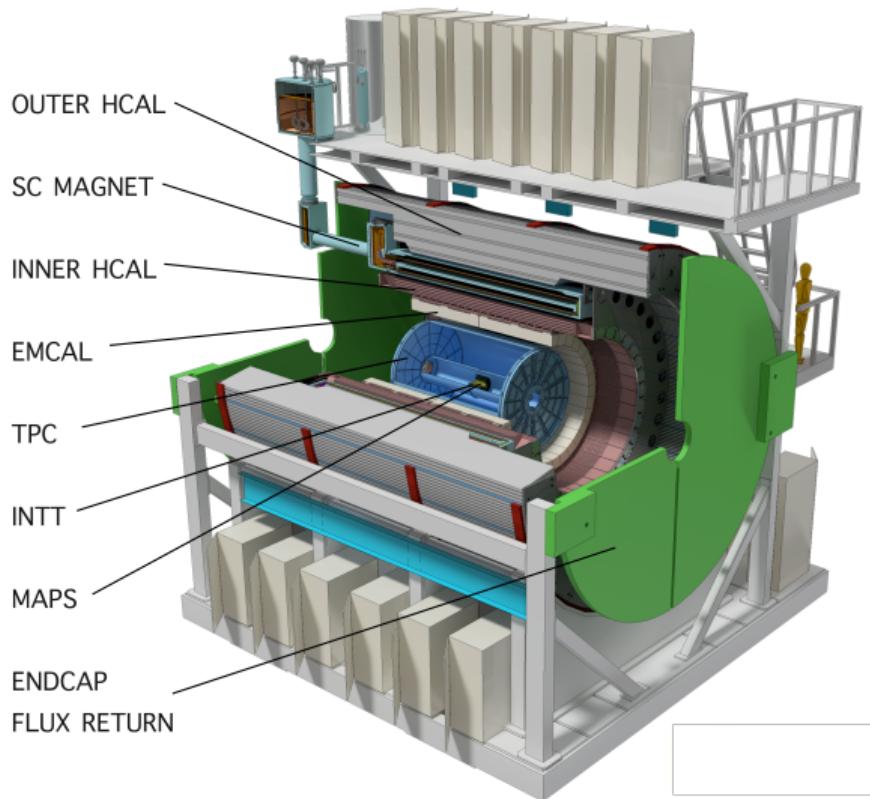
Recommendation IV: We recommend increasing investment in small-scale and mid-scale projects and initiatives that enable forefront research at universities and laboratories.

sPHENIX: QGP microscope



From the LRP: [The goal is to] probe the inner workings of QGP by resolving its properties at shorter and shorter length scales.... essential to this goal... is a state-of-the-art jet detector at RHIC, called sPHENIX.

sPHENIX: QGP microscope



Resolving power $d \propto \lambda$
de Broglie wavelength $\lambda = h/p$

p	λ
2.5 eV	500 nm
100 keV	12 pm
200 MeV	6.2 fm
1 GeV	1.2 fm
10 GeV	0.12 fm
50 GeV	0.025 fm

sPHENIX: timeline

Past and present

- Magnet purchase July 2013
- Magnet delivery April 2015
- DOE OPA CD-0 September 2016
- Order for Outer HCal steel March 2018
- DOE OPA CD-1/CD-3a August 2018

Future

- DOE OPA CD-2/CD-3b Review May 2019
- Authorization for CD-2/CD-3b July 2019
- Fabrication orders August 2019
- Installation begins April 2021
- Installation complete July 2022
- Initial commissioning complete September 2022
- First collisions January 2023

sPHENIX: magnet

STANDARD FORM 122 JUNE 1974 GENERAL SERVICES ADMINISTRATION FPMR (41 CFR) 101-32.306 FPMR (41 CFR) 101-43.315		TRANSFER ORDER EXCESS PERSONAL PROPERTY			1. ORDER NO. SLAC 2013-07-18 2. DATE July 18, 2013	
3. TO: GENERAL SERVICES ADMINISTRATION*		4. ORDERING AGENCY (Full name and address)* Brookhaven National Lab Attention: John Haggerty; haggerty@bnl.gov Upton, NY 11973-5000				
5. HOLDING AGENCY (Name and address)* SLAC National Accelerator Laboratory 2575 Sand Hill Road, MS 85A Menlo Park, CA 94025		6. SHIP TO (Consignee and destination)* Same as block 4				
7. LOCATION OF PROPERTY SLAC National Accelerator Laboratory C/O Mike Racine 2575 Sand Hill Road, MS 53 Menlo Park, CA 94025 650 926-3543 racine@slac.stanford.edu		8. SHIPPING INSTRUCTIONS BNL to arrange for shipping				
9. ORDERING AGENCY APPROVAL		10. APPROPRIATION SYMBOL AND TITLE transfer from DE-AC02-76SF00515 transfer to DE-AC02-98CH10886				
A. SIGNATURE 		B. DATE 7-19-13				
C. TITLE <i>Property Manager</i>		11. ALLOTMENT			12. GOVERNMENT B/L NO.	
13. PROPERTY ORDERED						
GSA AND HOLDING AGENCY NOS. (a)	ITEM NO. (b)	DESCRIPTION (Include noun name, FSC Group and Class, Condition Code and if available, National Stock Number) (c)	UNIT (d)	QUANTITY (e)	ACQUISITION COST	
					UNIT (f)	TOTAL (g)
	1	Administrative Transfer BaBar Solenoid and Components Date of Mfr: 1996 (See attached list)	ea	1	12,000,000.00	\$ 12,000,000.00

sPHENIX: magnet



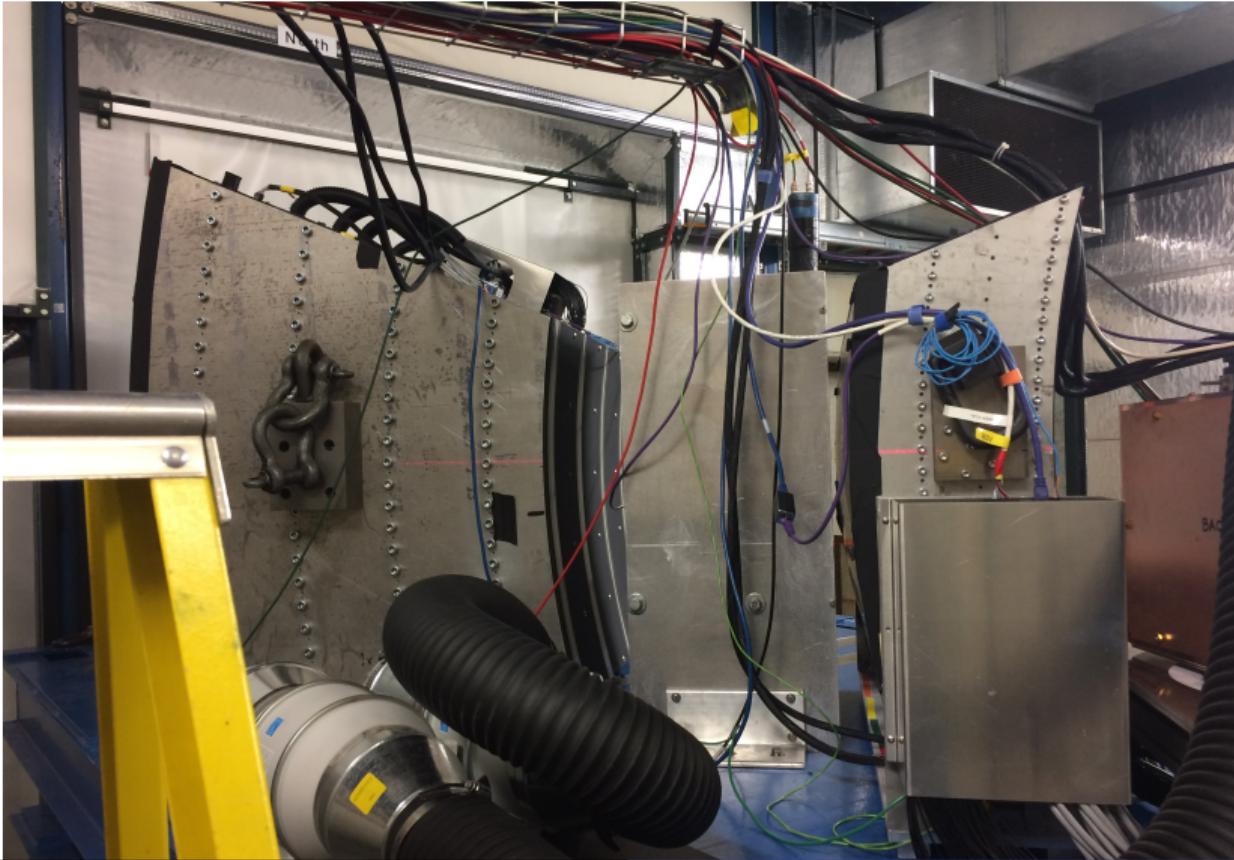
sPHENIX: magnet



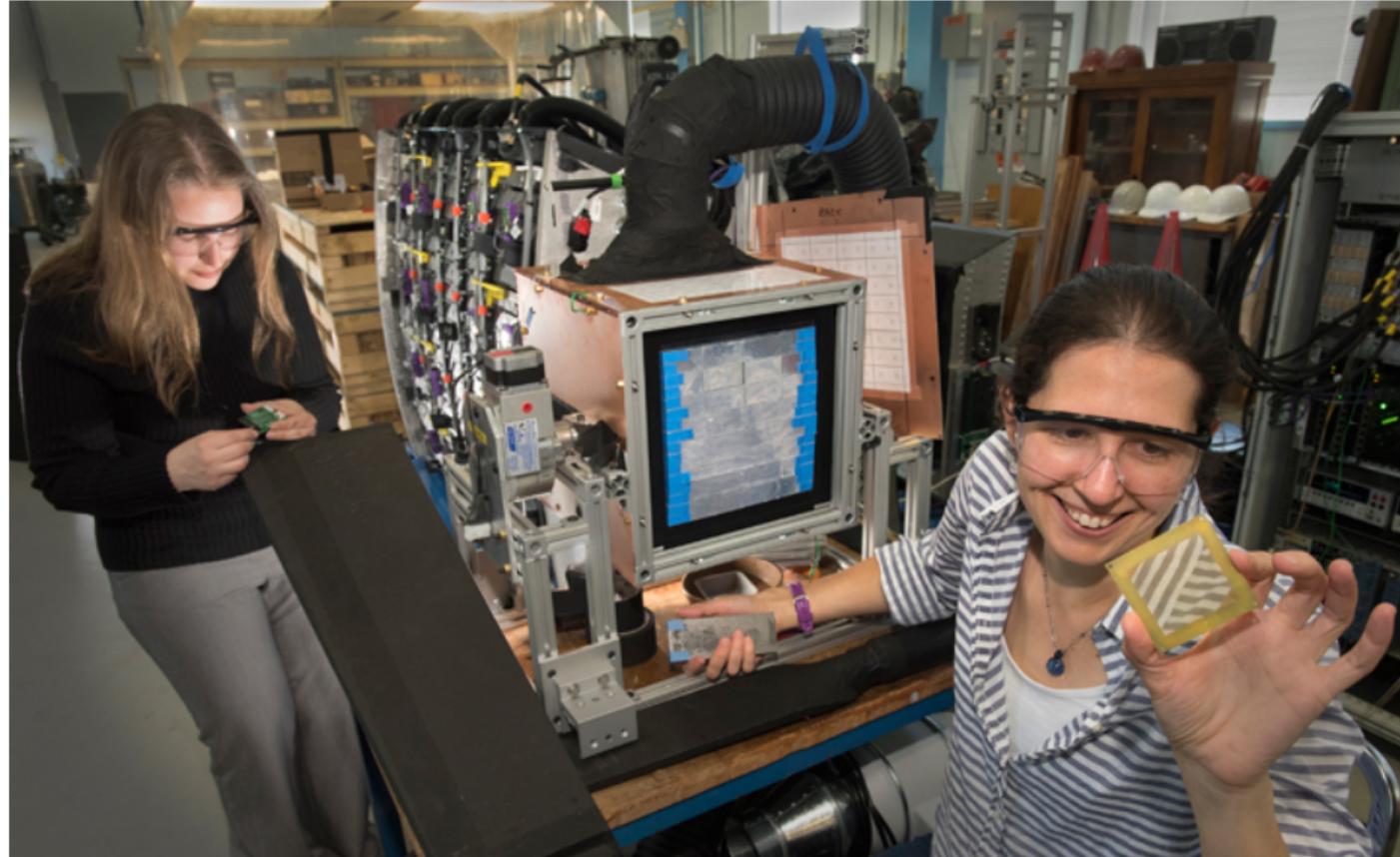
sPHENIX: beam tests



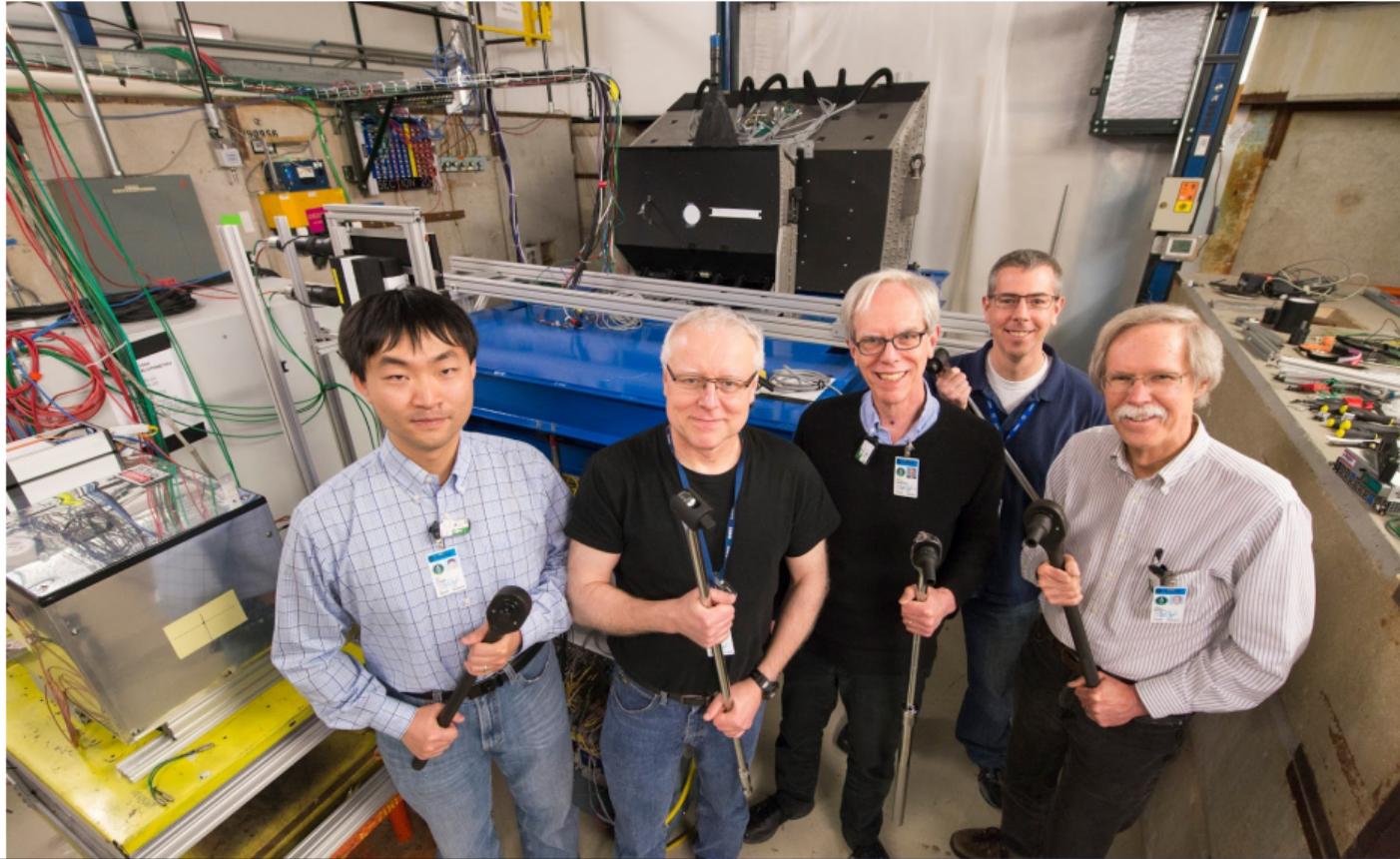
sPHENIX: beam tests



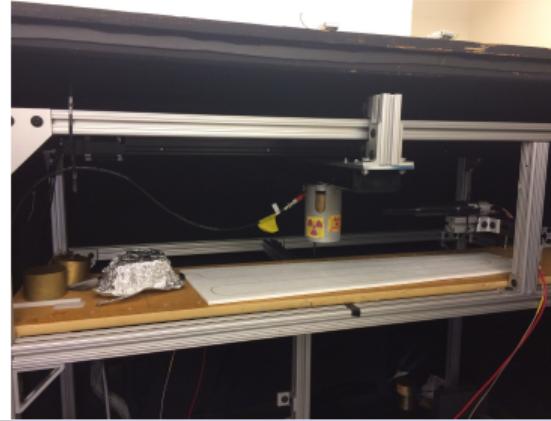
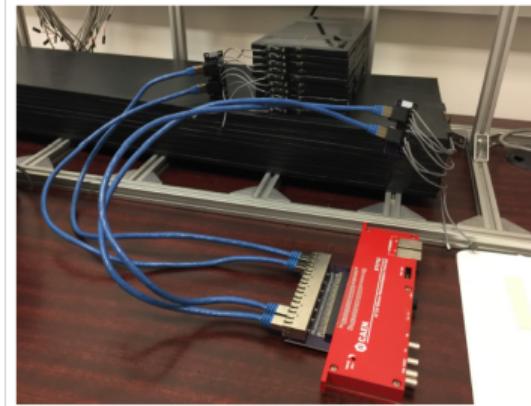
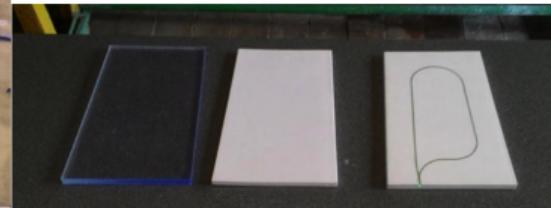
sPHENIX: beam tests



sPHENIX: beam tests



sPHENIX: HCal tiles



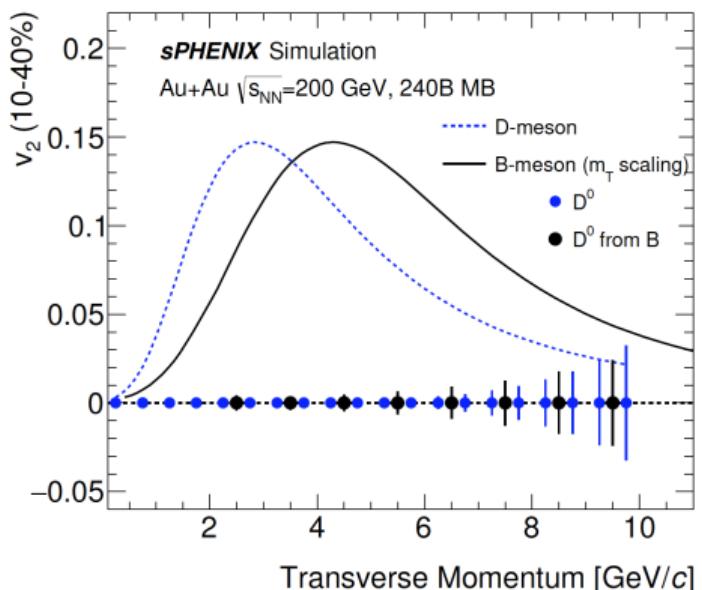
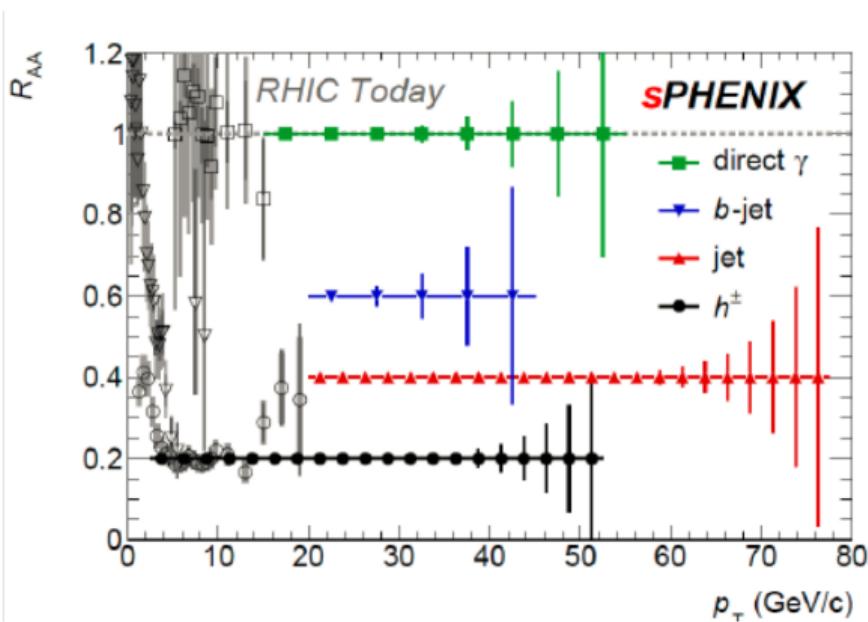
sPHENIX: jets!



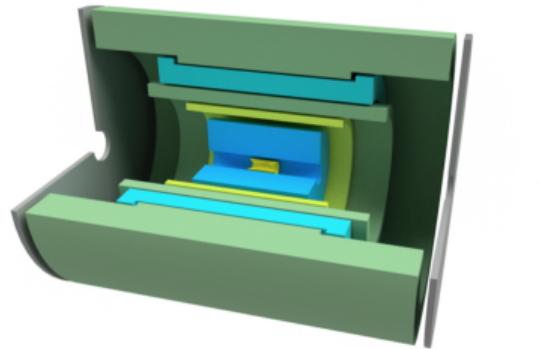
sPHENIX: heavy flavor!



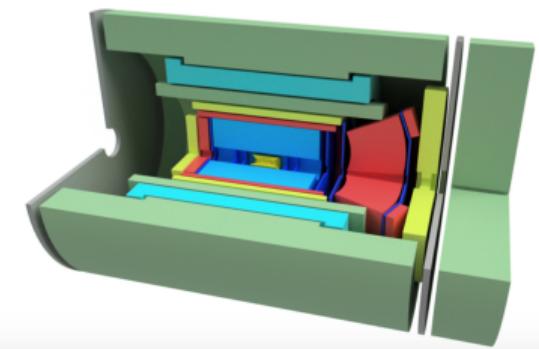
sPHENIX: projections



sPHENIX: day one EIC detector



EIC whitepaper: arXiv:1212.1701



An EIC Detector Built Around The
sPHENIX Solenoid

A Detector Design Study

Summary and outlook

- We can recreate the early universe in the lab using collisions of large nuclei
- Small systems have had a lot of surprises for us
 - Many similarities to large systems, but theoretical picture less clear
 - How to disentangle CGC from QGP?
- It's an exciting time for nuclear physics in the US

Summary and outlook

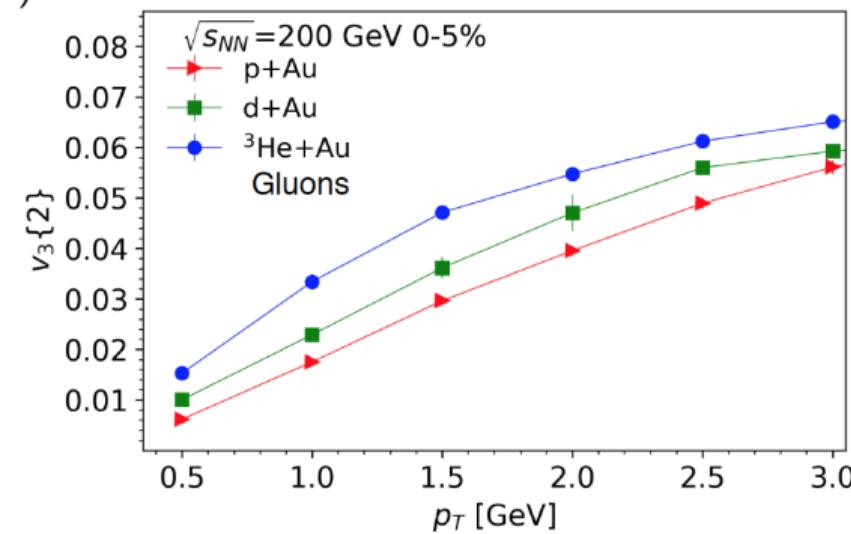
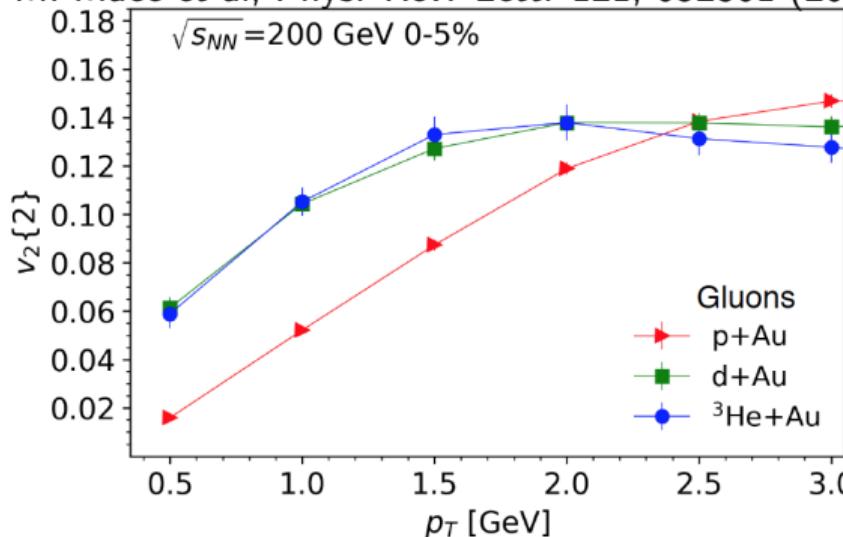
- We can recreate the early universe in the lab using collisions of large nuclei
- Small systems have had a lot of surprises for us
 - Many similarities to large systems, but theoretical picture less clear
 - How to disentangle CGC from QGP?
- It's an exciting time for nuclear physics in the US

"The optimist regards the future as uncertain." —Eugene Wigner

Extra material

CGC results on small systems

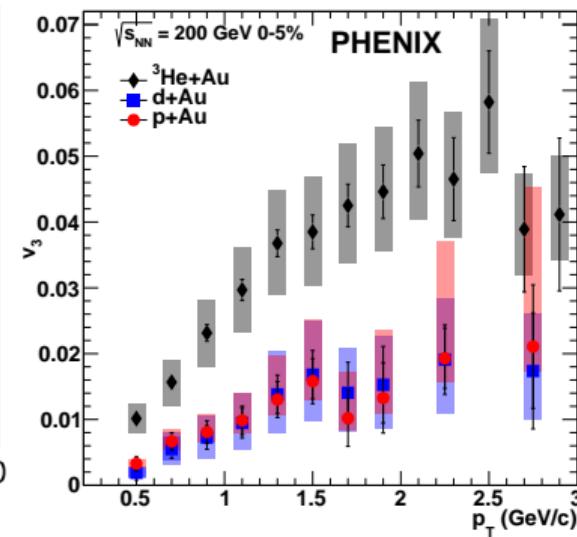
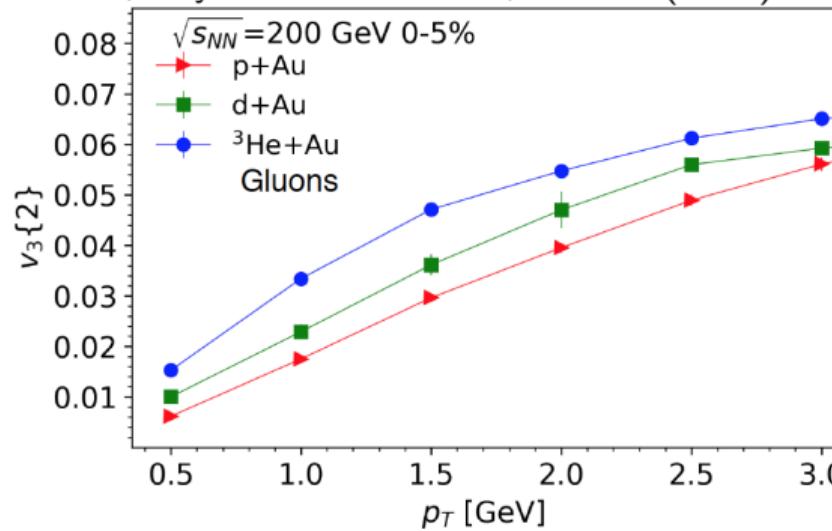
M. Mace et al, Phys. Rev. Lett. 121, 052301 (2018)



- New for QM18: v_2 and v_3 for small systems

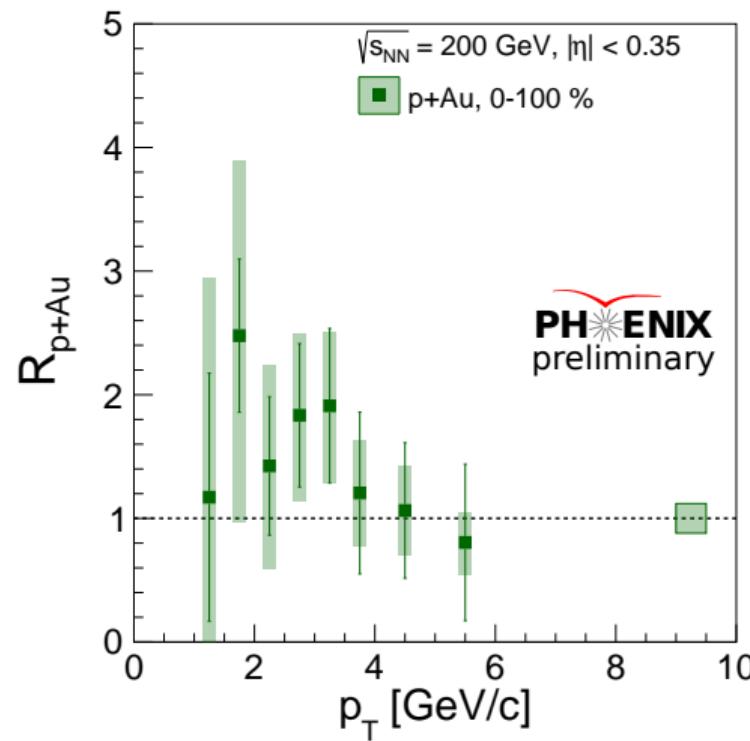
CGC results on small systems

M. Mace et al, Phys. Rev. Lett. 121, 052301 (2018)

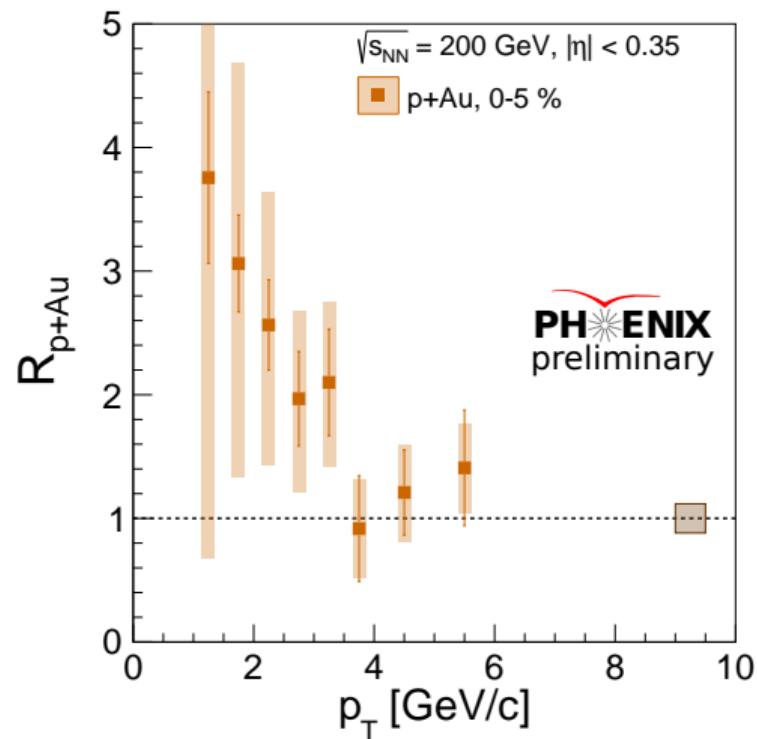
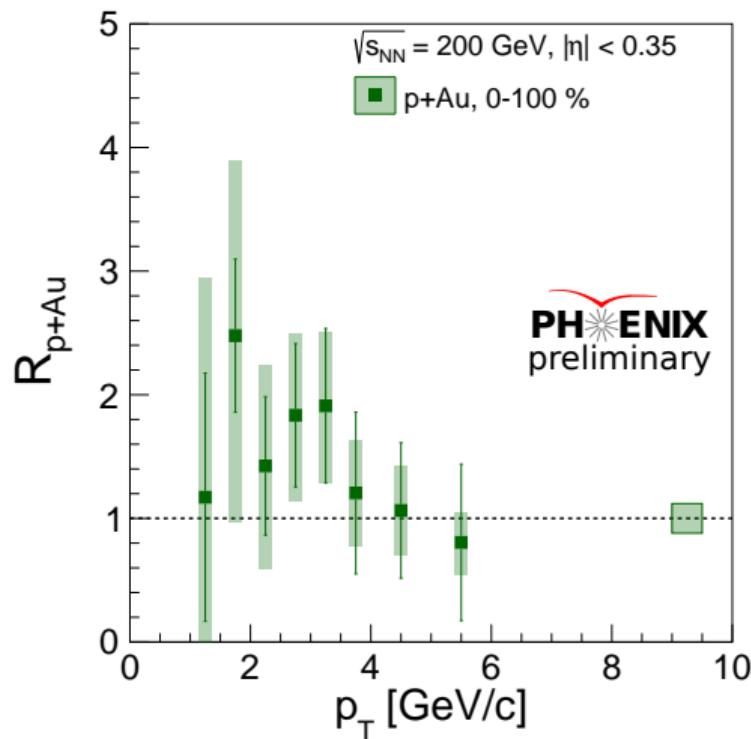


- New for QM18: v_2 and v_3 for small systems
- v_3 ordering is wrong in CGC
 - CGC: $\text{p+Au} < \text{d+Au} < {}^3\text{He+Au}$
 - Data: $\text{p+Au} \approx \text{d+Au} < {}^3\text{He+Au}$

Photons in small systems

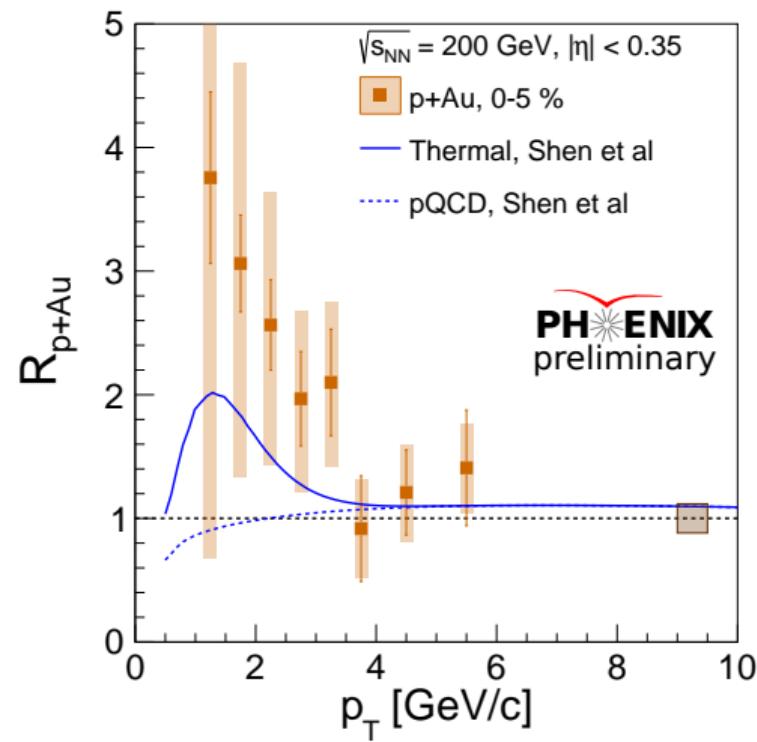
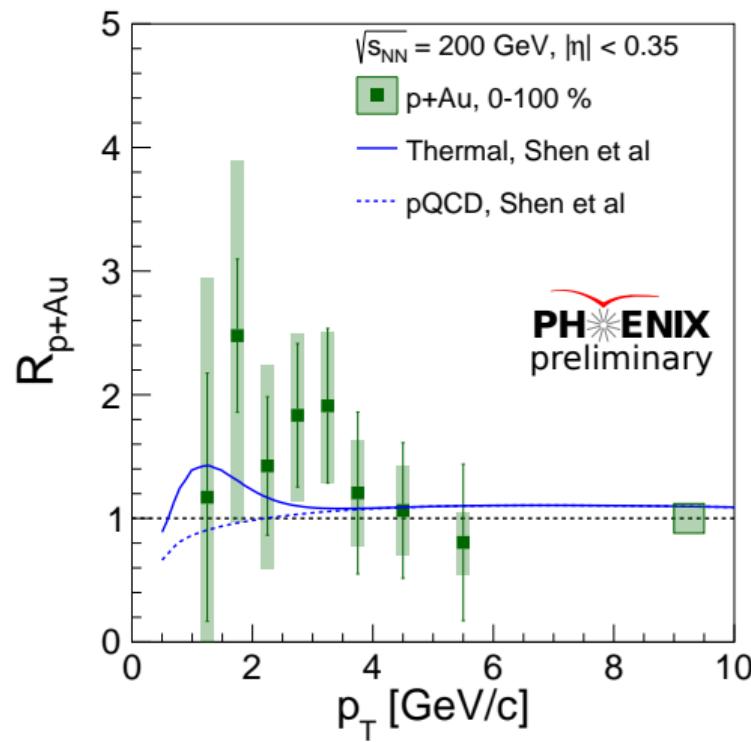


Photons in small systems



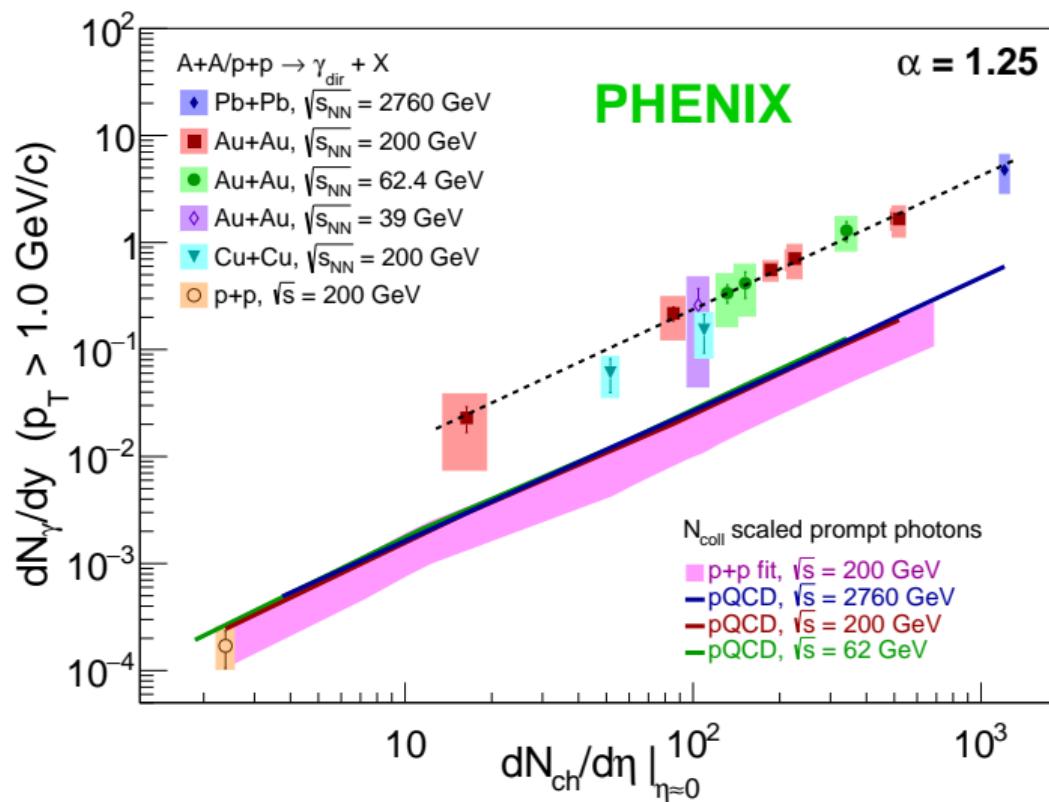
- Thermal photons in p+Au?

Photons in small systems



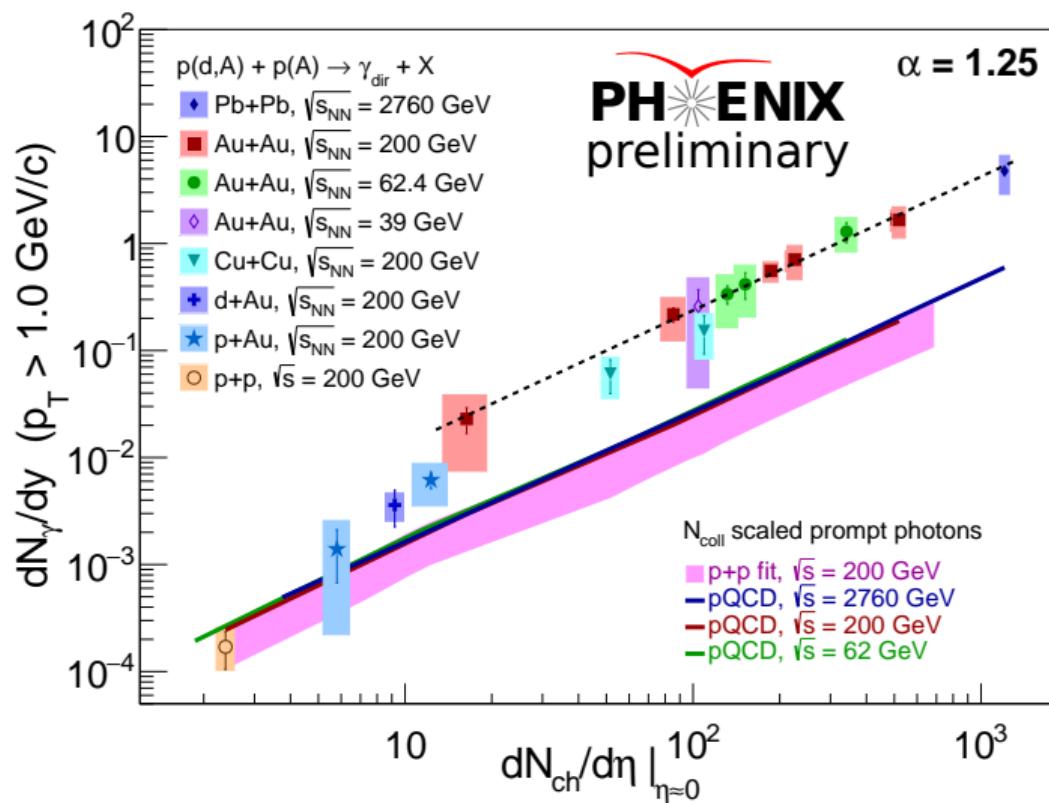
- Thermal photons in p+Au? Theory from Phys. Rev. C 95, 014906 (2017)

Photon yields



Common scaling for Au+Au and Pb+Pb at different energies; very different from N_{coll} -scaled $p+p$

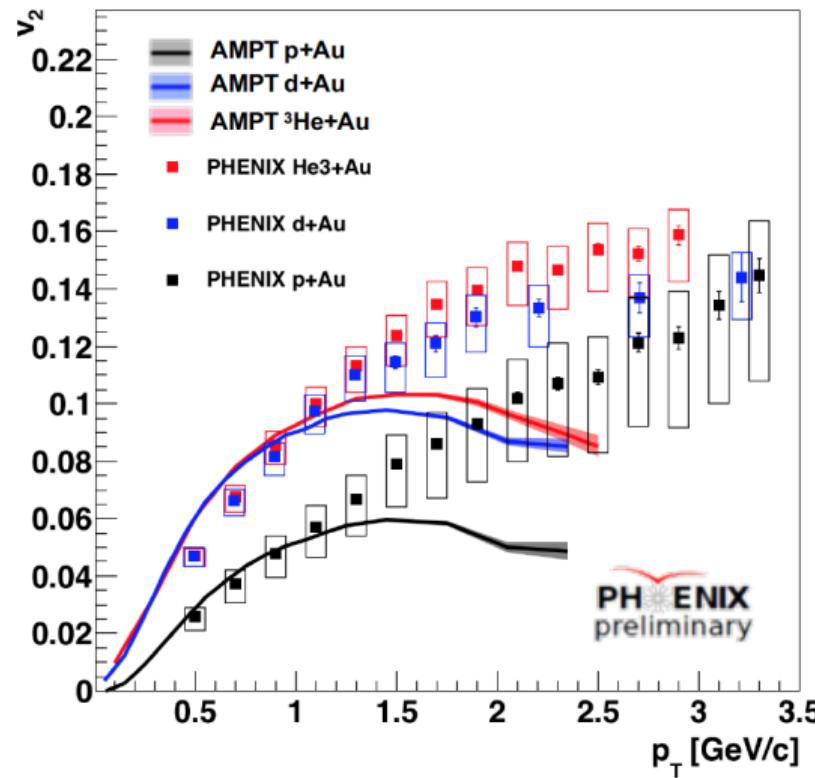
Photon yields



Common scaling for Au+Au
and Pb+Pb at different
energies; very different from
 N_{coll} -scaled $p+p$

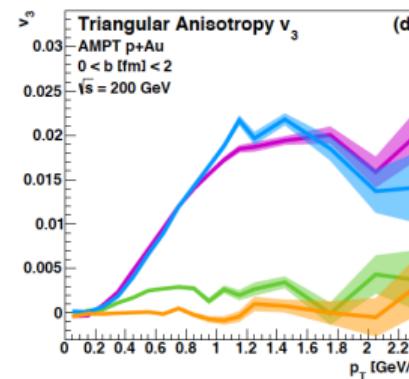
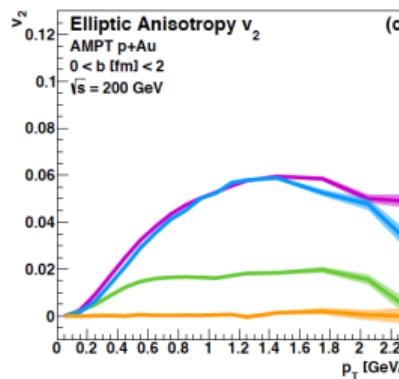
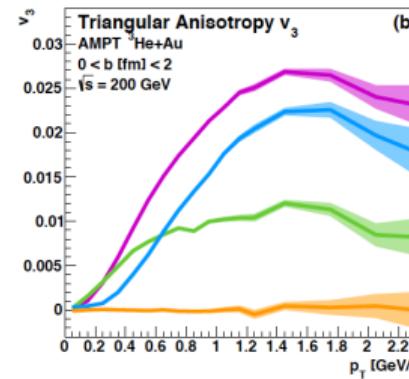
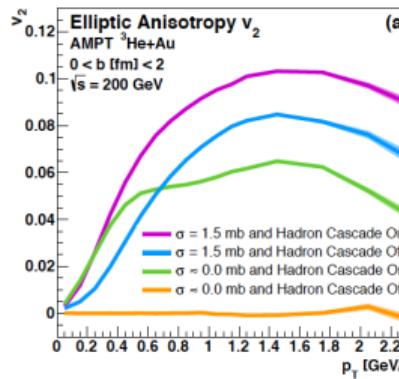
p+Au and d+Au in between

AMPT



AMPT with no scattering

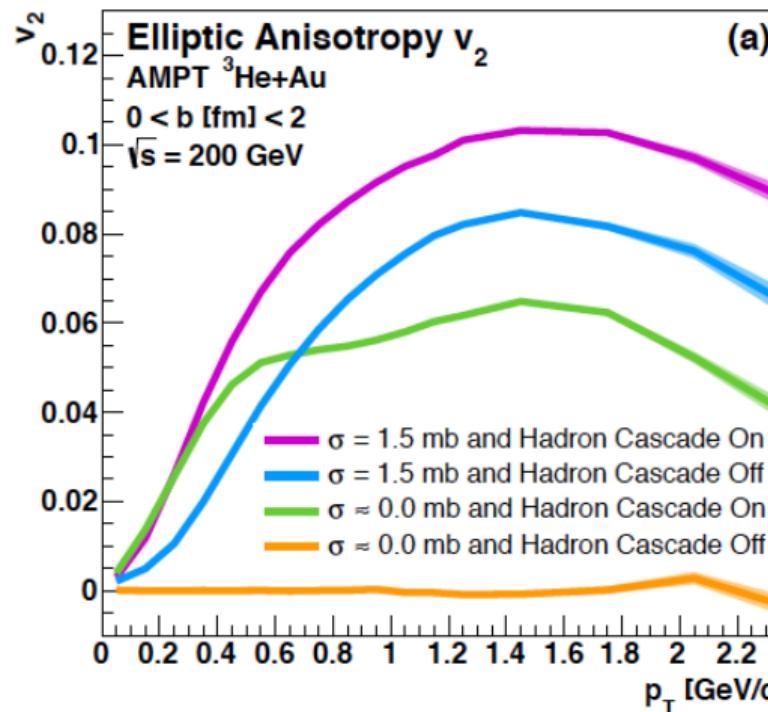
J.D. Orjuela Koop et al Phys. Rev. C 92, 054903 (2015)



- Turn off scattering in AMPT—remove all correlations with initial geometry
 $\sigma_{\text{parton}} = 0$ and $\sigma_{\text{hadron}} = 0$
- Participant plane v_2 goes to zero
- Other sources of correlation remain—non-flow

AMPT with no scattering

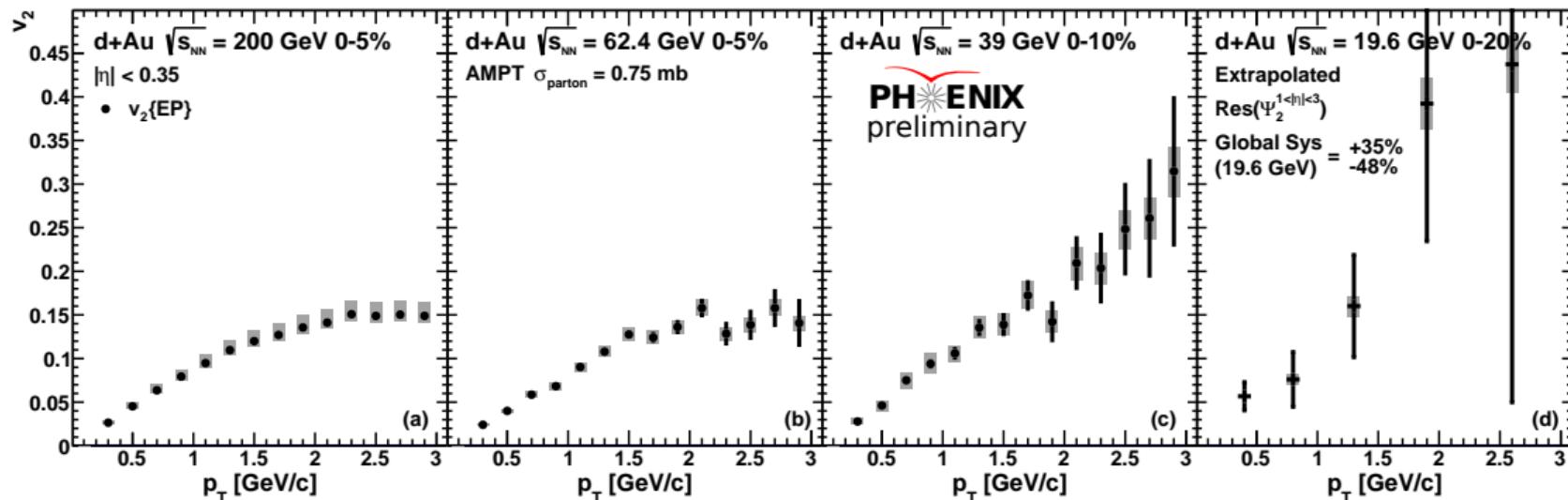
J.D. Orjuela Koop et al Phys. Rev. C 92, 054903 (2015)



- Turn off scattering in AMPT—remove all correlations with initial geometry
 $\sigma_{\text{parton}} = 0$ and $\sigma_{\text{hadron}} = 0$
- Participant plane v_2 goes to zero
- Other sources of correlation remain—non-flow

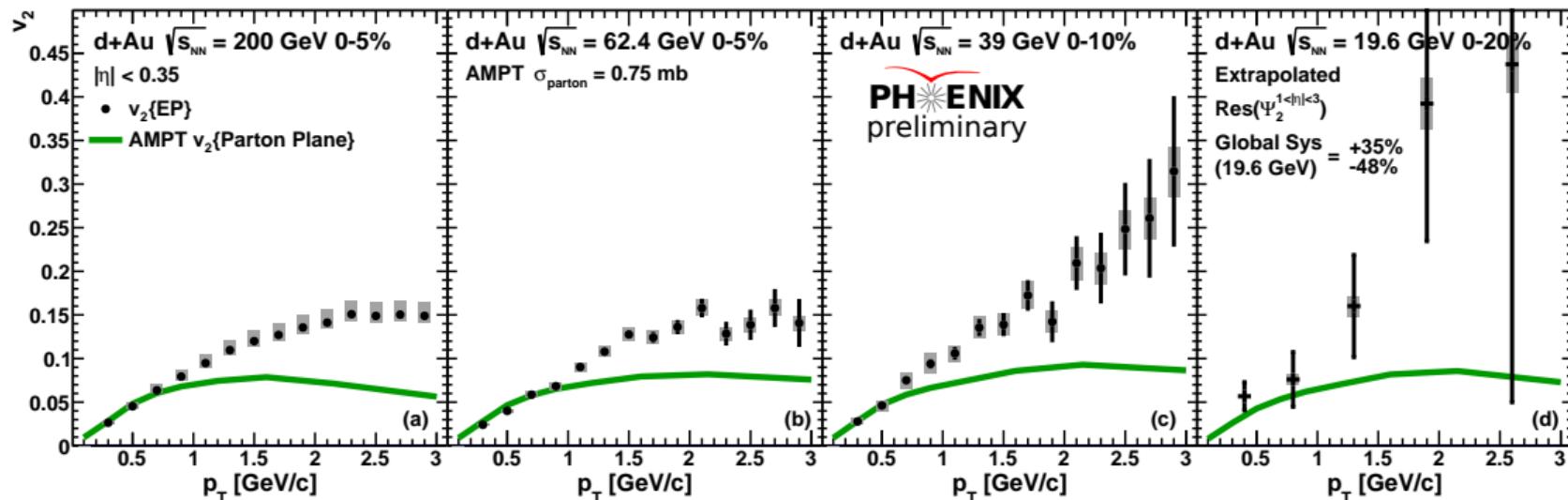
v_2 vs p_T , comparisons to AMPT

Phys. Rev. C 96, 064905 (2017)



v_2 vs p_T , comparisons to AMPT

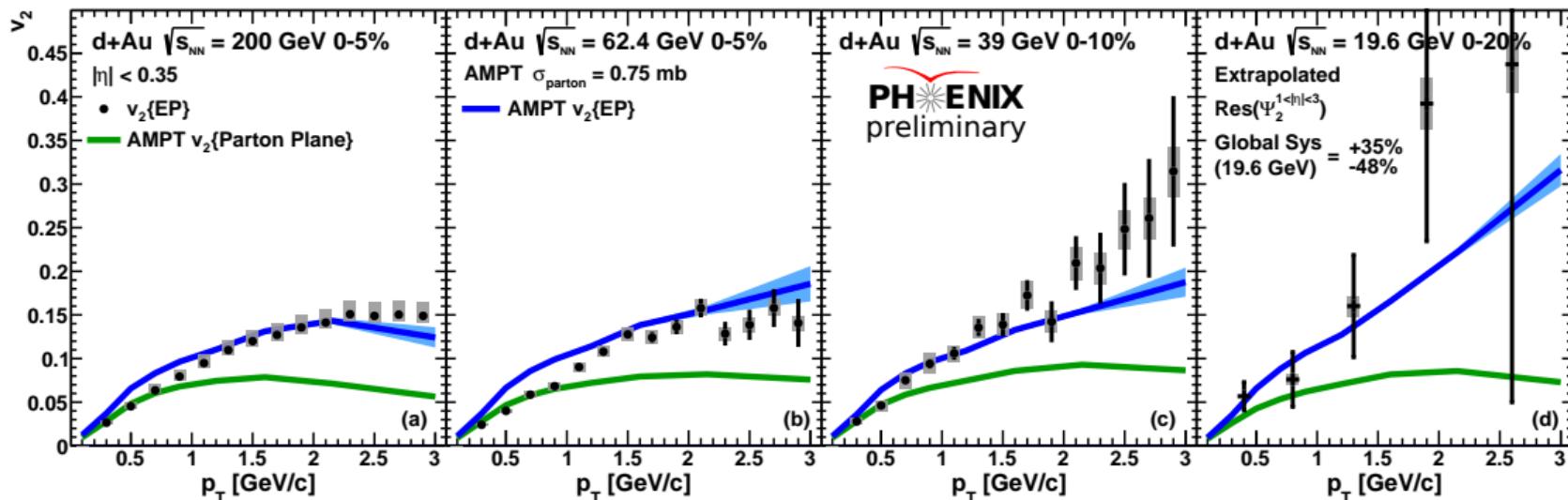
Phys. Rev. C 96, 064905 (2017)



- AMPT flow only shows good agreement at low p_T and all energies

v_2 vs p_T , comparisons to AMPT

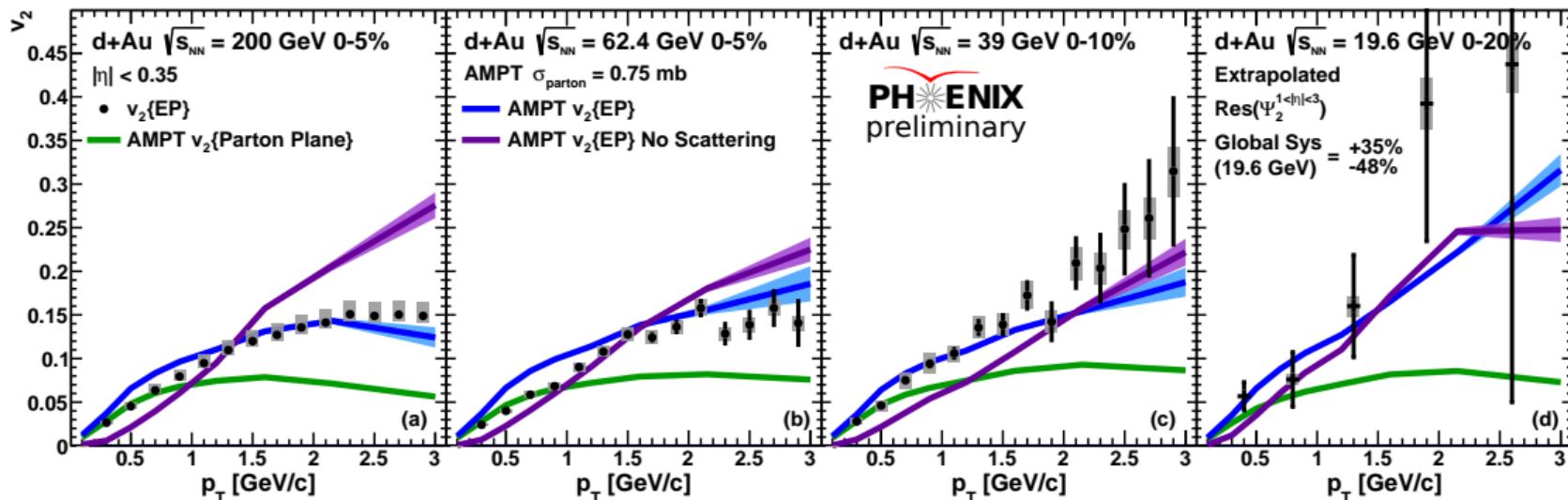
Phys. Rev. C 96, 064905 (2017)



- AMPT flow only shows good agreement at low p_T and all energies
- AMPT flow+non-flow shows reasonable agreement for all p_T and all energies

v_2 vs p_T , comparisons to AMPT

Phys. Rev. C 96, 064905 (2017)

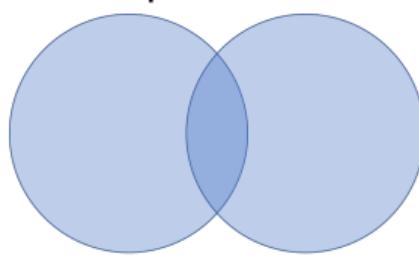


- AMPT flow only shows good agreement at low p_T and all energies
- AMPT flow+non-flow shows reasonable agreement for all p_T and all energies
- AMPT non-flow only far under-predicts for low p_T , too high for high p_T

Centrality

- Need to characterize the overlap of the two nuclei
- b (impact parameter)—separation between the centers of the two nuclei
- N_{part} —number of nucleons in the overlap region
- N_{coll} —number of nucleon-nucleon collisions

Peripheral

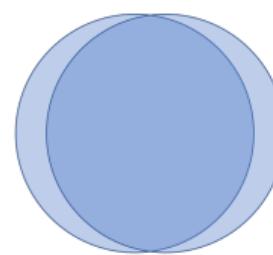


Higher b

Lower N_{part}

Lower N_{coll}

Central



Lower b

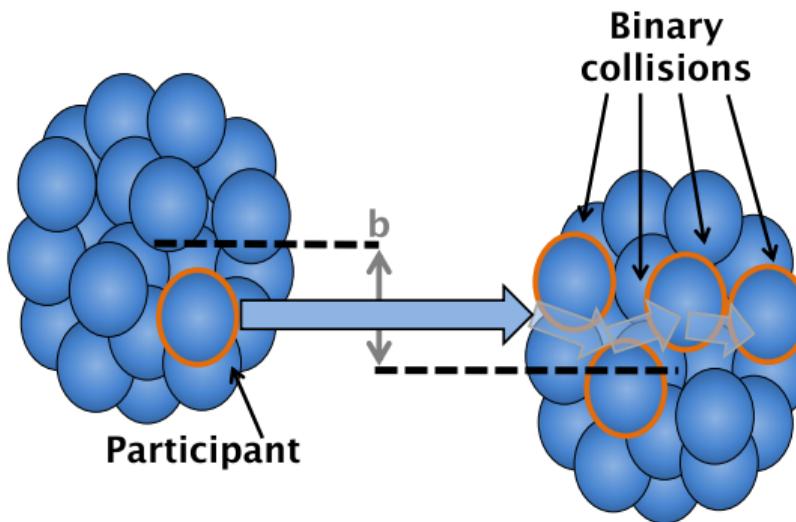
Higher N_{part}

Higher N_{coll}

Centrality	$\langle N_{part} \rangle$	$\langle N_{coll} \rangle$
Pb+Pb		
0-5%	382.7 ± 5.1	1685 ± 190
5-10%	329.7 ± 4.6	1316 ± 140
10-20%	260.5 ± 4.4	921 ± 96
20-30%	186.4 ± 3.9	556 ± 55
30-40%	128.9 ± 3.3	320 ± 32
40-50%	85.0 ± 2.6	171 ± 16
50-60%	52.8 ± 2.0	84.3 ± 7
60-70%	30.0 ± 1.3	37.9 ± 3
70-80%	15.8 ± 0.6	15.6 ± 1
p+p	$\equiv 2$	$\equiv 1$

Centrality

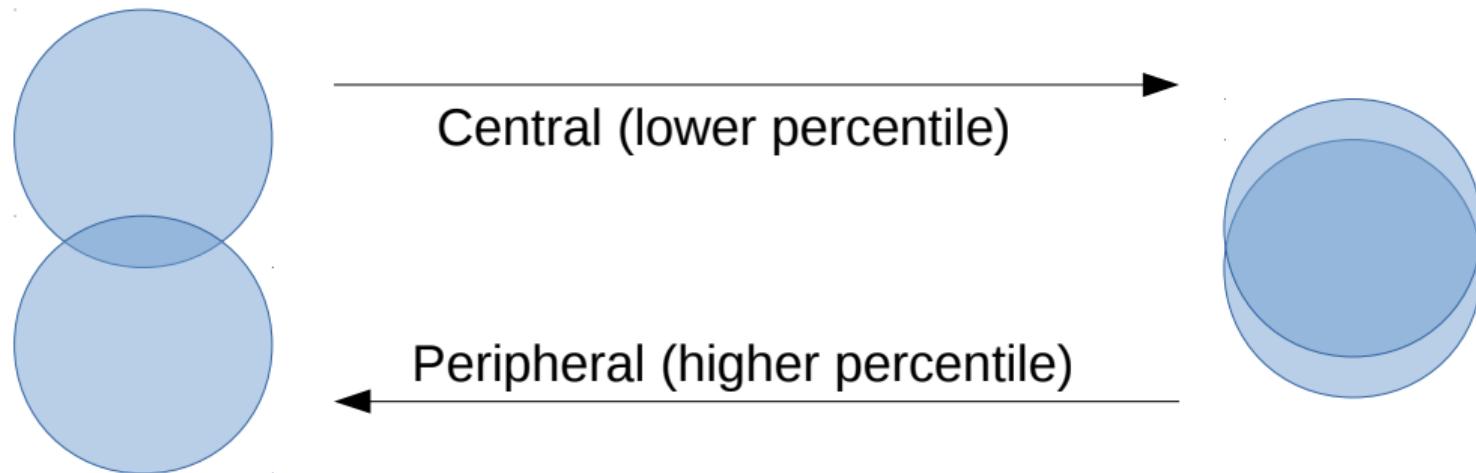
- Need to characterize the overlap of the two nuclei
- b (impact parameter)—separation between the centers of the two nuclei
- N_{part} —number of nucleons in the overlap region
- N_{coll} —number of nucleon-nucleon collisions



Centrality	$\langle N_{part} \rangle$	$\langle N_{coll} \rangle$
Pb+Pb		
0-5%	382.7 ± 5.1	1685 ± 190
5-10%	329.7 ± 4.6	1316 ± 140
10-20%	260.5 ± 4.4	921 ± 96
20-30%	186.4 ± 3.9	556 ± 55
30-40%	128.9 ± 3.3	320 ± 32
40-50%	85.0 ± 2.6	171 ± 16
50-60%	52.8 ± 2.0	84.3 ± 7
60-70%	30.0 ± 1.3	37.9 ± 3
70-80%	15.8 ± 0.6	15.6 ± 1
p+p	$\equiv 2$	$\equiv 1$

Centrality

- Need to characterize the overlap of the two nuclei
- b (impact parameter)—separation between the centers of the two nuclei
- N_{part} —number of nucleons in the overlap region
- N_{coll} —number of nucleon-nucleon collisions



The six flavors of quarks

flavor	charge	mass
down	-1/3 e	3.0–7.0 MeV
up	2/3 e	1.5–3.0 MeV
strange	-1/3 e	95 ± 25 MeV
charm	2/3 e	1.25 ± 0.09 GeV
bottom	-1/3 e	4.70 ± 0.07 GeV
top	2/3 e	174.2 ± 3.3 GeV

- No bound states with top quarks
(so heavy they decay weakly before a bound state can be formed)
- All others can form bound states
 - Any combination of quarks you can imagine is allowed
 - Though some have to be part of a linear combination
 - Sometimes a single combination can be more than one particle

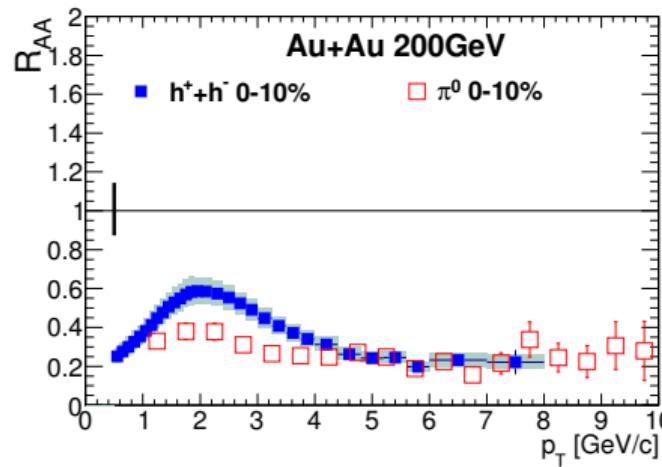
The six flavors of quarks

flavor	charge	mass
down	-1/3 e	3.0–7.0 MeV
up	2/3 e	1.5–3.0 MeV
strange	-1/3 e	95 ± 25 MeV
charm	2/3 e	1.25 ± 0.09 GeV
bottom	-1/3 e	4.70 ± 0.07 GeV
top	2/3 e	174.2 ± 3.3 GeV

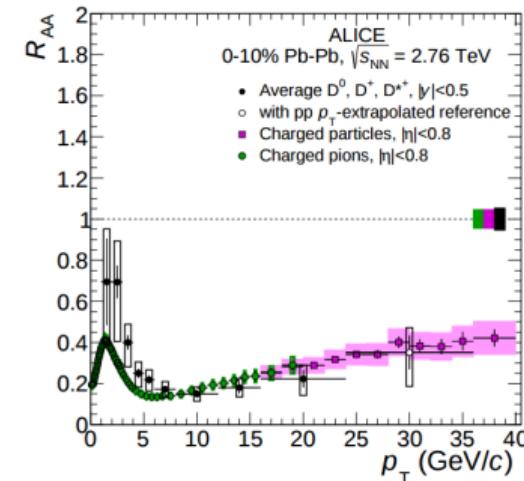
A few examples

- $p = uud$, $n = udd$
- $\pi^+ = u\bar{d}$, $\pi^- = \bar{u}d$, $\pi^0 = (u\bar{u} - d\bar{d})/\sqrt{2}$
- $\Lambda^0 = uds$, $\Lambda_c^+ = udc$, $\Lambda_b^0 = udb$
- $K^- = \bar{u}s$, $D^+ = \bar{d}c$, $B^- = \bar{u}b$
- $c\bar{c} = \eta_c, J/\psi, \chi_c, h_c$

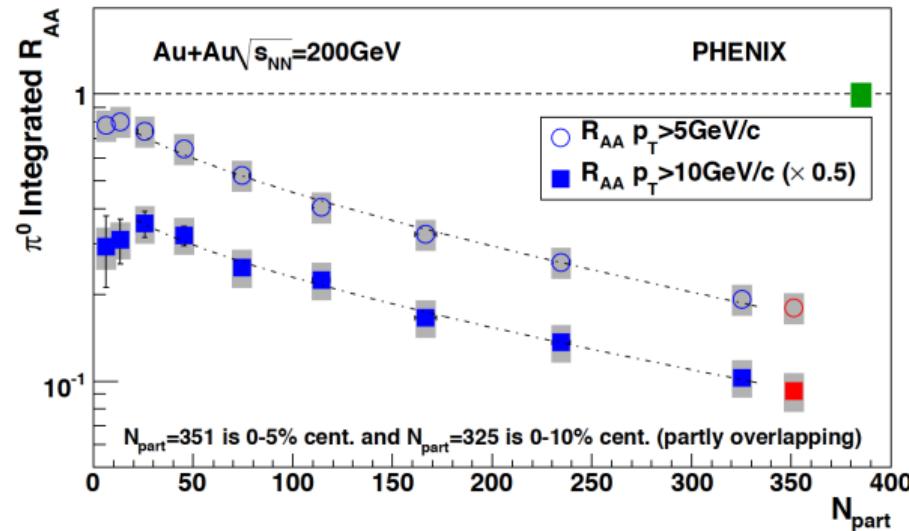
Suppression of high energy particles



- $R_{AA} = \frac{N_{\text{particles}}^{A+A}}{N_{\text{pp}}^{A+A} \times N_{\text{coll}}}$
- $R_{AA} < 1$ means particles are suppressed



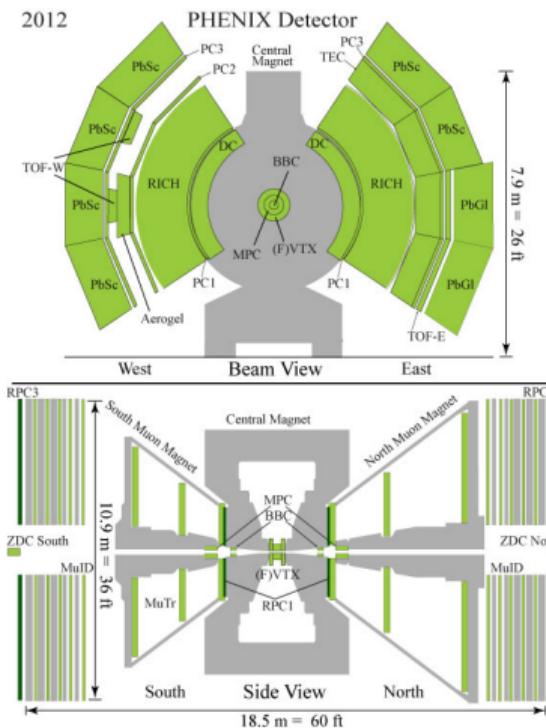
Suppression of high energy particles



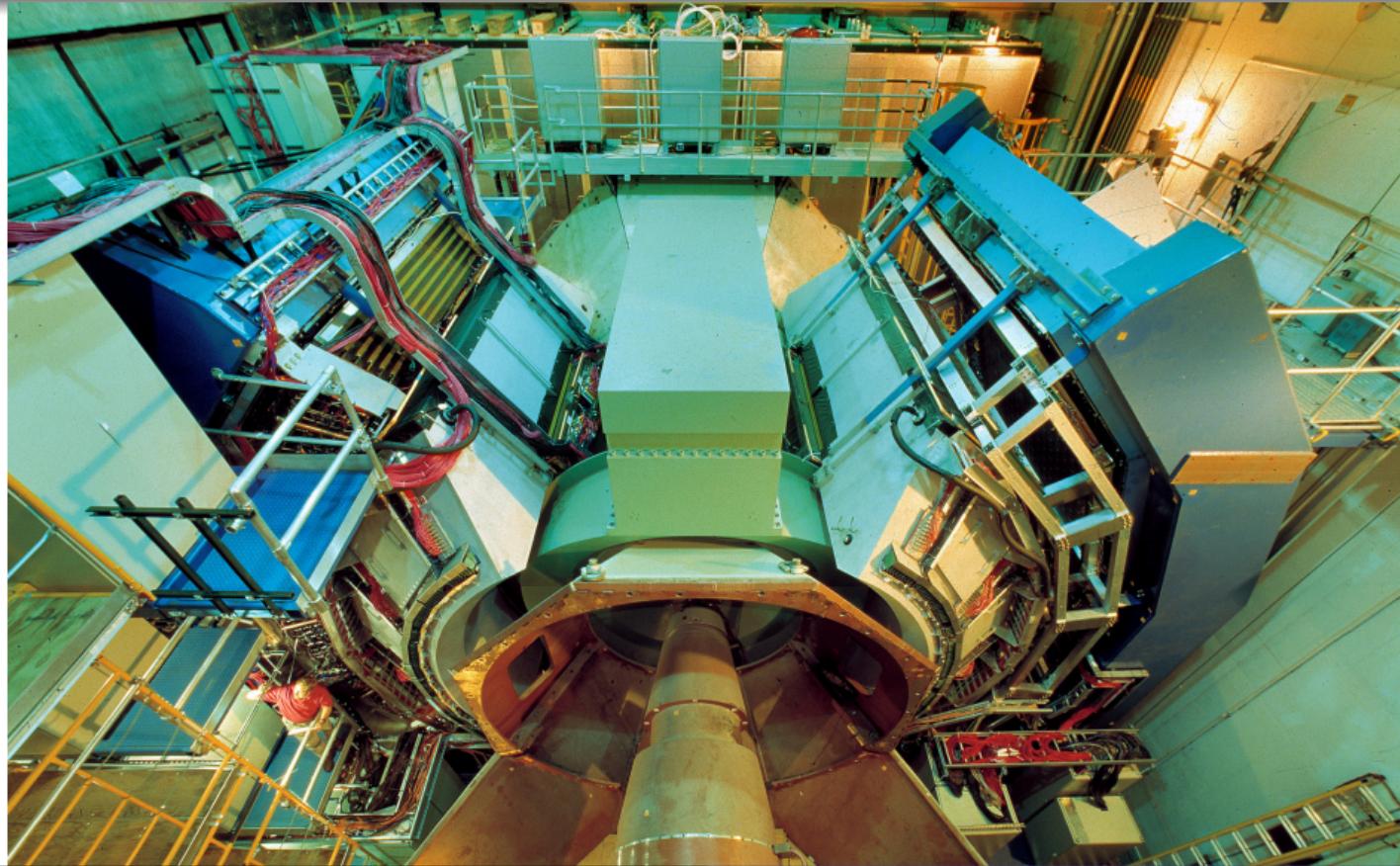
- R_{AA} decreases (more suppression) with increasing N_{part} (bigger system)
- More medium → more stuff in the way → more suppression
- System size/geometry important aspect of suppression

PHENIX

- Weighs approximately 3000 tons
- Three separate magnet systems (Central Arms and Muon North and South) weighing 1700 tons alone
- 16 detector subsystems and about 5,000,000 (silicon) plus 300,000 (other) electronics channels
- 30 feet tall, 40 feet wide, 60 feet long
- Very fast DAQ system—7 kHz, 1 GB/s
- Ideally suited for measurements of rare probes, electrons, muons, high p_T photons, etc.

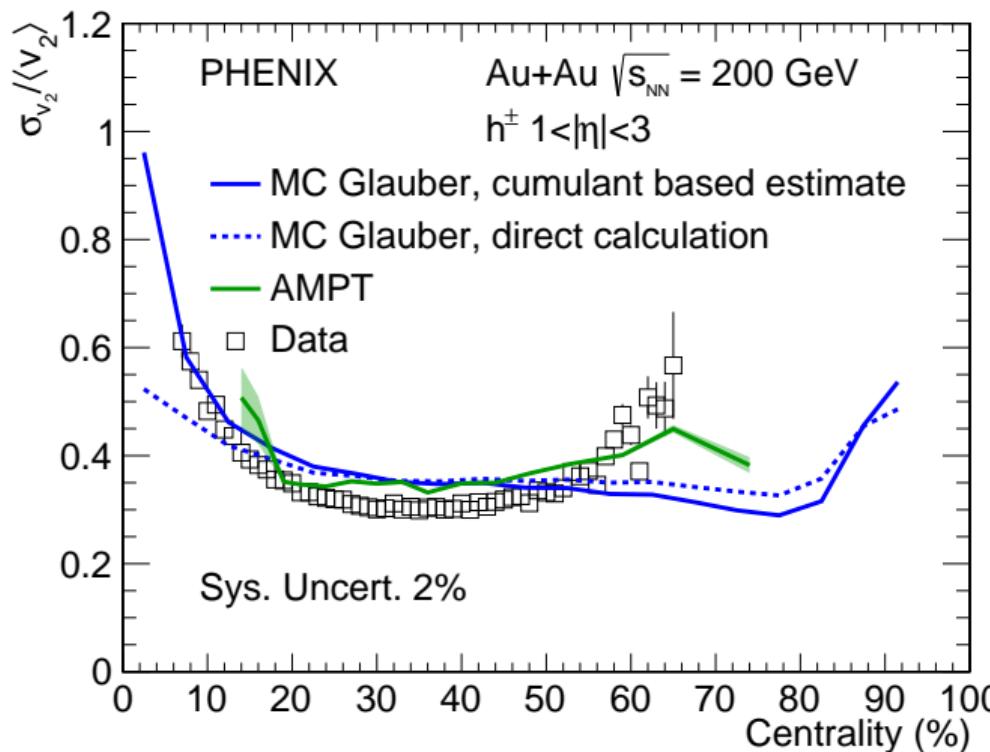


PHENIX



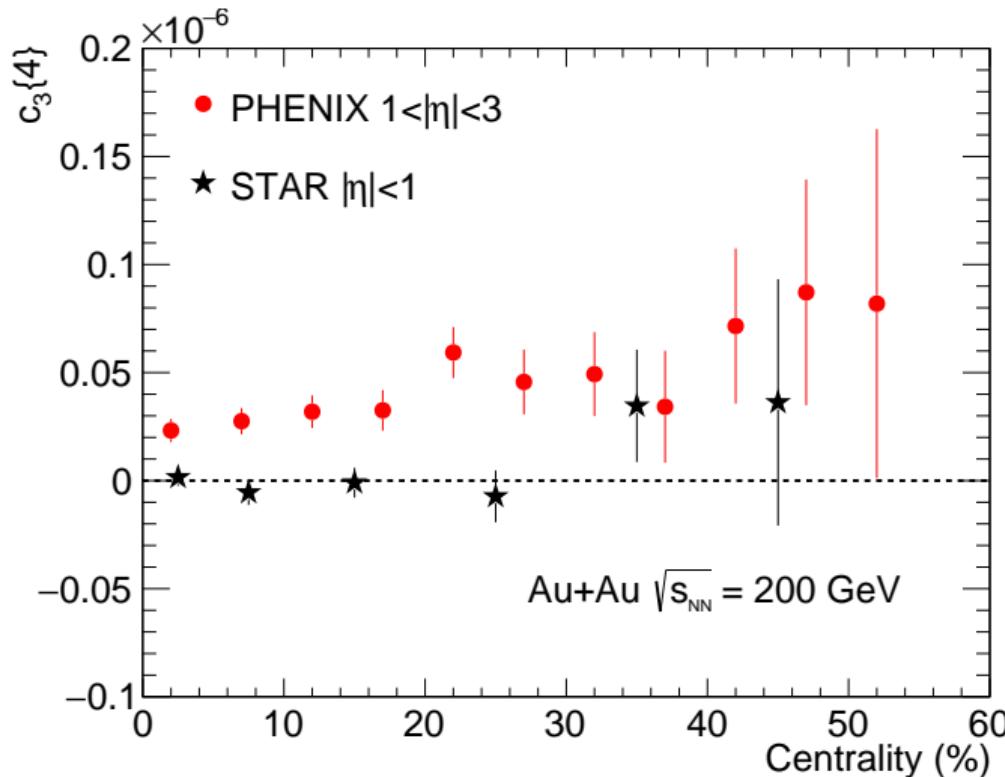
Collectivity in large systems

arXiv:1804.10024 (submitted to Phys. Rev. C)



Collectivity in large systems

arXiv:1804.10024 (submitted to Phys. Rev. C)



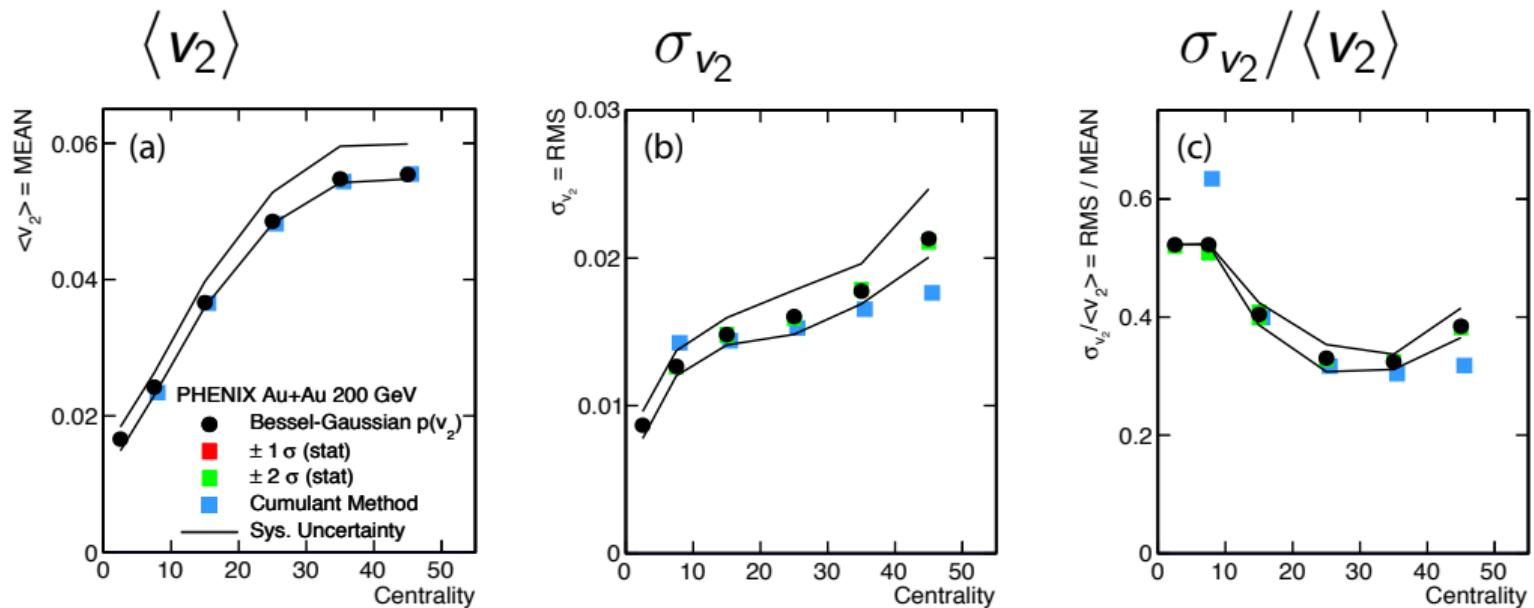
$$1 < |\eta| < 3$$

Cannot extract

$$\sigma v_3 / \langle v_3 \rangle$$

Collectivity in large systems

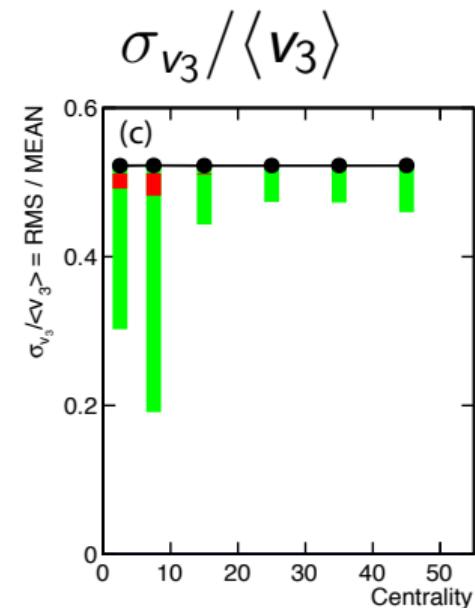
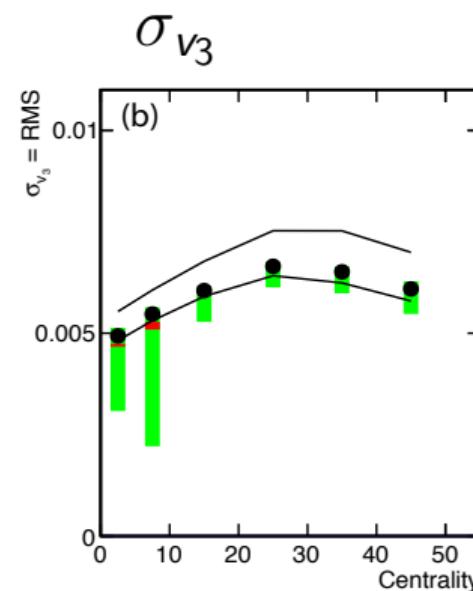
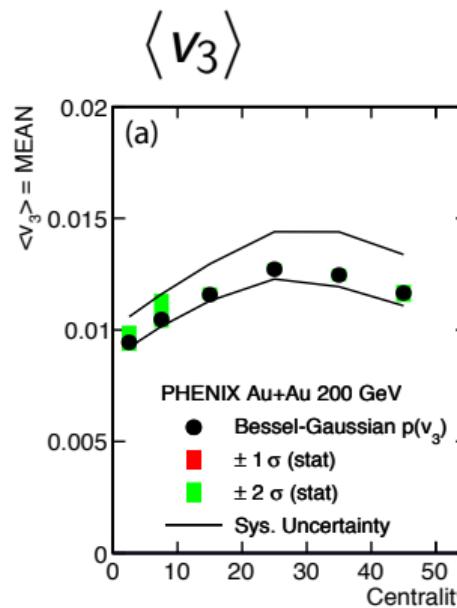
arXiv:1804.10024 (submitted to Phys Rev C)



- Can extract $\langle v_2 \rangle$ and σ_{v_2} separately using forward-fold

Collectivity in large systems

arXiv:1804.10024 (submitted to Phys Rev C)

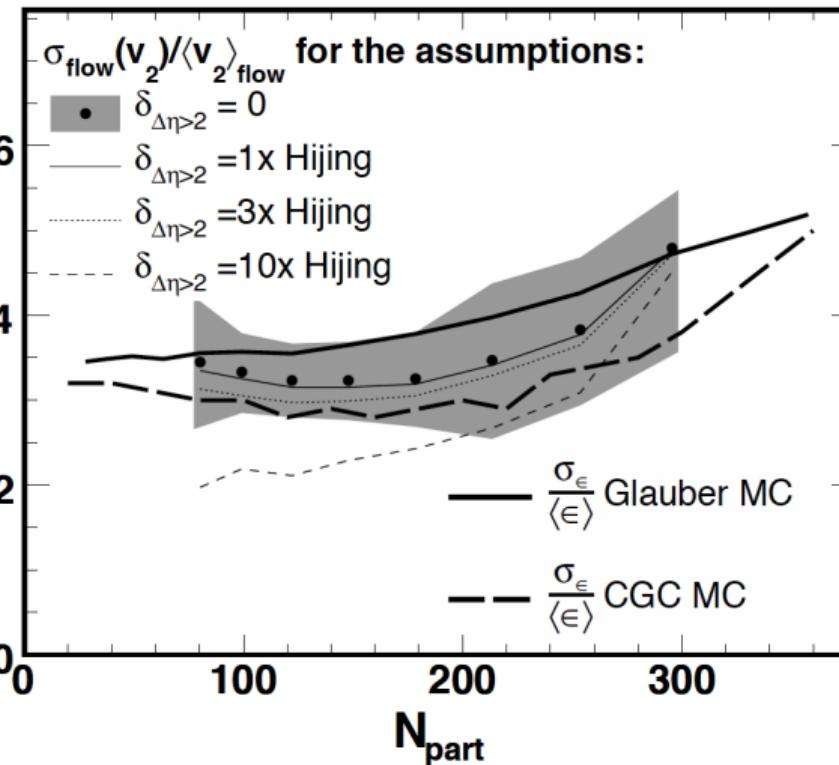


- Can extract $\langle v_3 \rangle$ and σ_{v_3} separately using forward-fold

Fluctuations in large systems

PHOBOS, Phys. Rev. C 81, 034915 (2010)

Relative Fluctuations



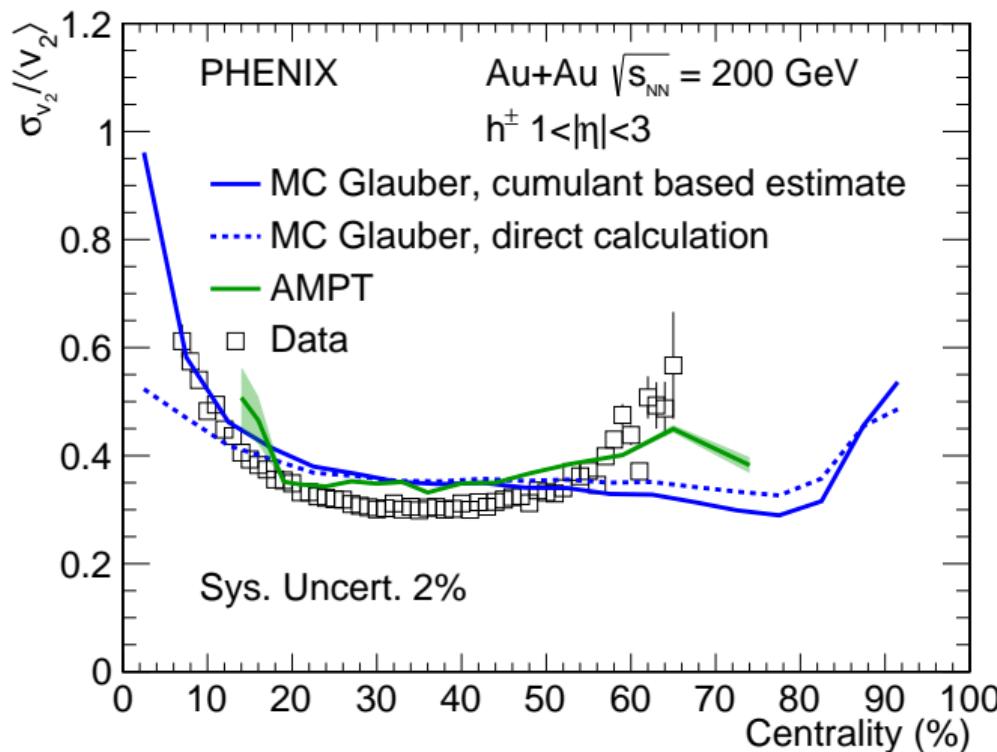
Fluctuations should also be translated, so measure $\sigma_{v_2}/\langle v_2 \rangle$

$$|\eta| < 1$$

Generally good agreement with models of initial geometry

Fluctuations in large systems

PHENIX, arXiv:1804.10024 (submitted to Phys. Rev. C)



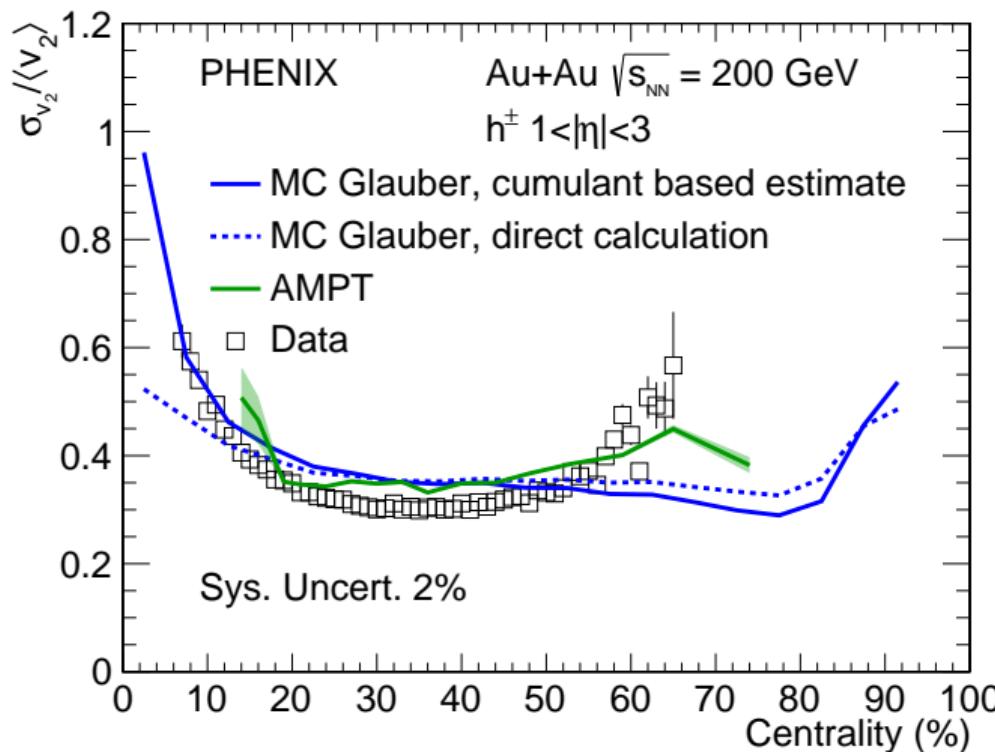
Fluctuations should also be translated, so measure $\sigma_{v_2}/\langle v_2 \rangle$

$$1 < |\eta| < 3$$

Generally good agreement with models of initial geometry

Fluctuations in large systems

PHENIX, arXiv:1804.10024 (submitted to Phys. Rev. C)



Fluctuations should also be translated, so measure $\sigma_{v_2}/\langle v_2 \rangle$

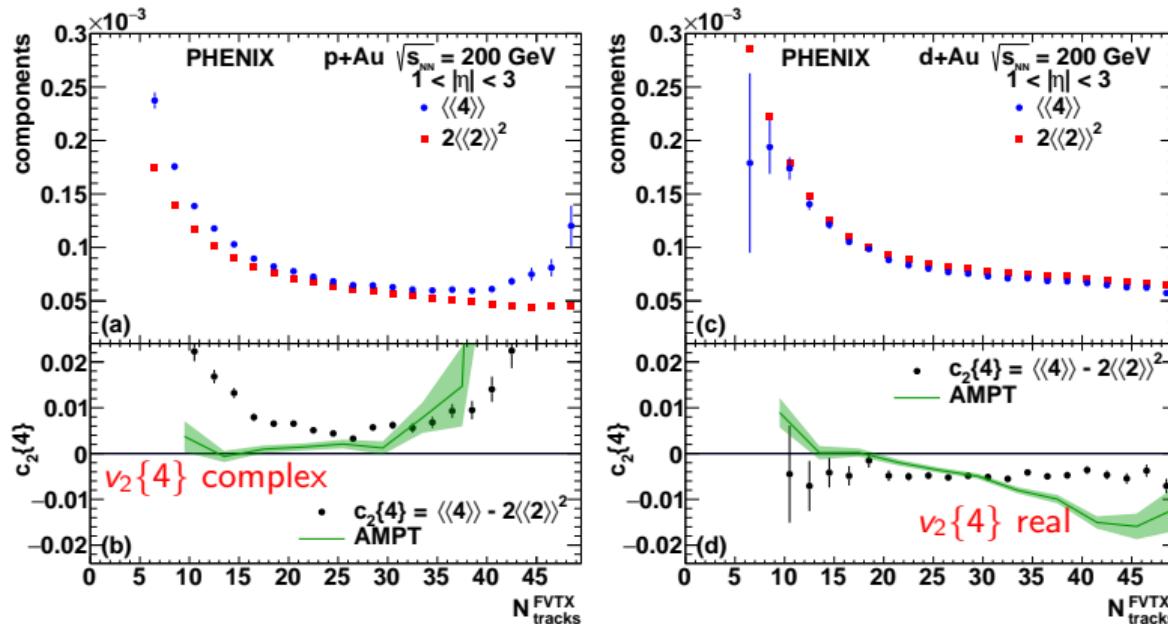
$1 < |\eta| < 3$

Central: breakdown of small-variance limit

Peripheral: non-linearity in hydro response (e.g. J. Noronha-Hostler et al Phys. Rev. C 93, 014909 (2016))

Components and cumulants in p+Au and d+Au at 200 GeV

Phys. Rev. Lett. 120, 062302 (2018)



- Is the sign of $c_2\{4\}$ a good indicator of collectivity? No.
- Fluctuations could dominate in the p+Au...

Back to basics (a brief excursion)

The (raw) moments of a probability distribution function $f(x)$:

$$\mu_n = \langle x^n \rangle \equiv \int_{-\infty}^{+\infty} x^n f(x) dx$$

The moment generating function:

$$M_x(t) \equiv \langle e^{tx} \rangle = \int_{-\infty}^{+\infty} e^{tx} f(x) dx = \int_{-\infty}^{+\infty} \sum_{n=0}^{\infty} \frac{t^n}{n!} x^n f(x) dx = \sum_{n=0}^{\infty} \mu_n \frac{t^n}{n!}$$

Moments from the generating function:

$$\mu_n = \left. \frac{d^n M_x(t)}{dt^n} \right|_{t=0}$$

Key point: the moment generating function uniquely describe $f(x)$

Back to basics (a brief excursion)

Can also uniquely describe $f(x)$ with the cumulant generating function:

$$K_x(t) \equiv \ln M_x(t) = \sum_{n=0}^{\infty} \kappa_n \frac{t^n}{n!}$$

Cumulants from the generating function:

$$\kappa_n = \left. \frac{d^n K_x(t)}{dt^n} \right|_{t=0}$$

Since $K_x(t) = \ln M_x(t)$, $M_x(t) = \exp(K_x(t))$, so

$$\mu_n = \left. \frac{d^n \exp(K_x(t))}{dt^n} \right|_{t=0}, \quad \kappa_n = \left. \frac{d^n \ln M_x(t)}{dt^n} \right|_{t=0}$$

End result: (details left as an exercise for the interested reader)

$$\mu_n = \sum_{k=1}^n B_{n,k}(\kappa_1, \dots, \kappa_{n-k+1}) = B_n(\kappa_1, \dots, \kappa_{n-k+1})$$

$$\kappa_n = \sum_{k=1}^n (-1)^{k-1} (k-1)! B_{n,k}(\mu_1, \dots, \mu_{n-k+1}) = L_n(\kappa_1, \dots, \kappa_{n-k+1})$$

Back to basics (a brief excursion)

Evaluating the Bell polynomials gives

$$\langle x \rangle = \kappa_1$$

$$\langle x^2 \rangle = \kappa_2 + \kappa_1^2$$

$$\langle x^3 \rangle = \kappa_3 + 3\kappa_1\kappa_2 + \kappa_1^3$$

$$\langle x^4 \rangle = \kappa_4 + 4\kappa_1\kappa_3 + 3\kappa_2^2 + 6\kappa_1^2\kappa_2 + \kappa_1^4$$

One can tell by inspection (or derive explicitly) that κ_1 is the mean, κ_2 is the variance, etc.

Back to basics (a brief excursion)

Subbing in $x = v_n$, $\kappa_2 = \sigma^2$, we find

$$\left(\langle v_n^4 \rangle = v_n^4 + 6v_n^2\sigma^2 + 3\sigma^4 + 4v_n\kappa_3 + \kappa_4 \right)$$

$$-\left(2\langle v_n^2 \rangle^2 = 2v_n^4 + 4v_n^2\sigma^2 + 2\sigma^4 \right)$$

→

$$\langle v_n^4 \rangle - 2\langle v_n^2 \rangle^2 = -v_n^4 + 2v_n^2\sigma^2 + \sigma^4 + 4v_n\kappa_3 + \kappa_4$$

Skewness s : $\kappa_3 = s\sigma^3$

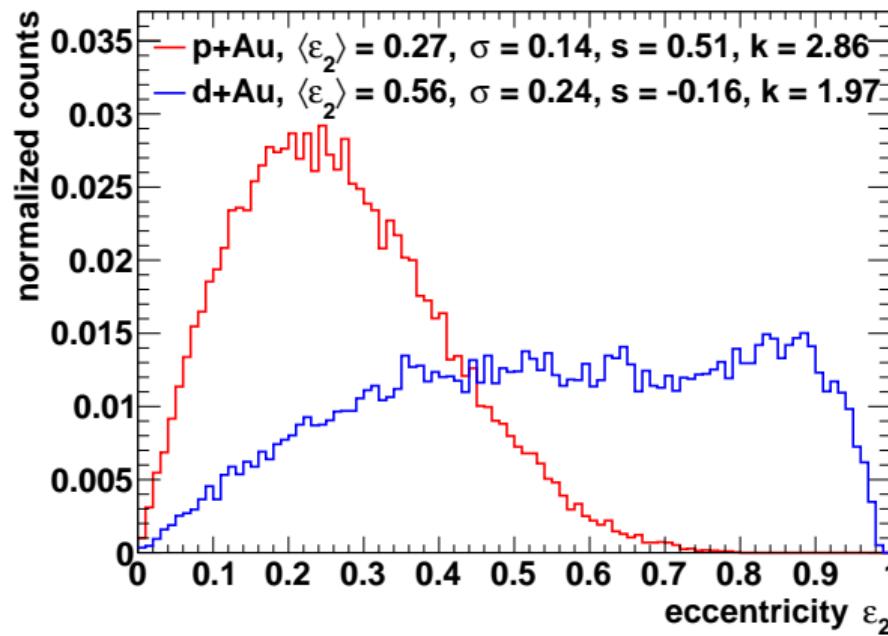
Kurtosis k : $\kappa_4 = (k - 3)\sigma^4$

$$v_n\{2\} = (v_n^2 + \sigma^2)^{1/2}$$

$$v_n\{4\} = (v_n^4 - 2v_n^2\sigma^2 - 4v_n s\sigma^3 - (k - 2)\sigma^4)^{1/4}$$

So the fully general form is a bit more complicated than we tend to think...

Eccentricity distributions and cumulants

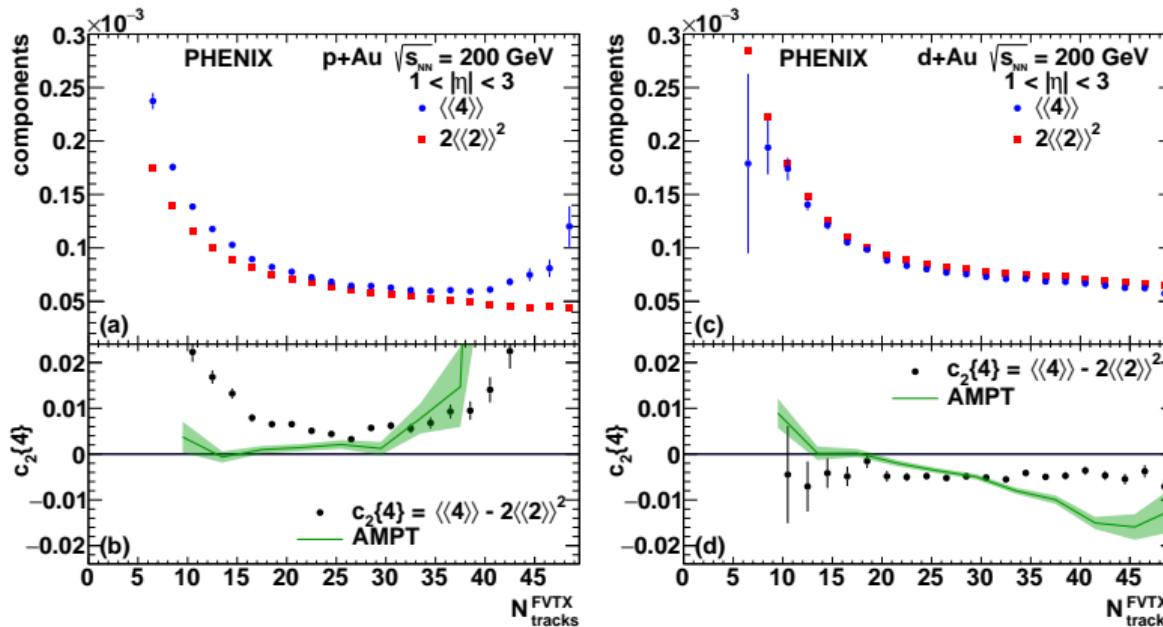


$$\varepsilon_2\{4\} = (\varepsilon_2^4 - 2\varepsilon_2^2\sigma^2 - 4\varepsilon_2 s \sigma^3 - (k - 2)\sigma^4)^{1/4}$$

	$p+Au$	$d+Au$
ε_2^4	0.00531	0.0983
$2\varepsilon_2^2\sigma^2$	0.00277	0.0370
$4\varepsilon_2 s \sigma^3$	0.00147	-0.0053
$(k - 2)\sigma^4$	0.00031	-0.0001

- the variance brings $\varepsilon_2\{4\}$ down (this term gives the usual $\sqrt{\nu_2^2 - \sigma^2}$)
- positive skew brings $\varepsilon_2\{4\}$ further down, negative skew brings it back up
- kurtosis > 2 brings $\varepsilon_2\{4\}$ further down, kurtosis < 2 brings it back up
—recall Gaussian has kurtosis = 3

Eccentricity distributions and cumulants



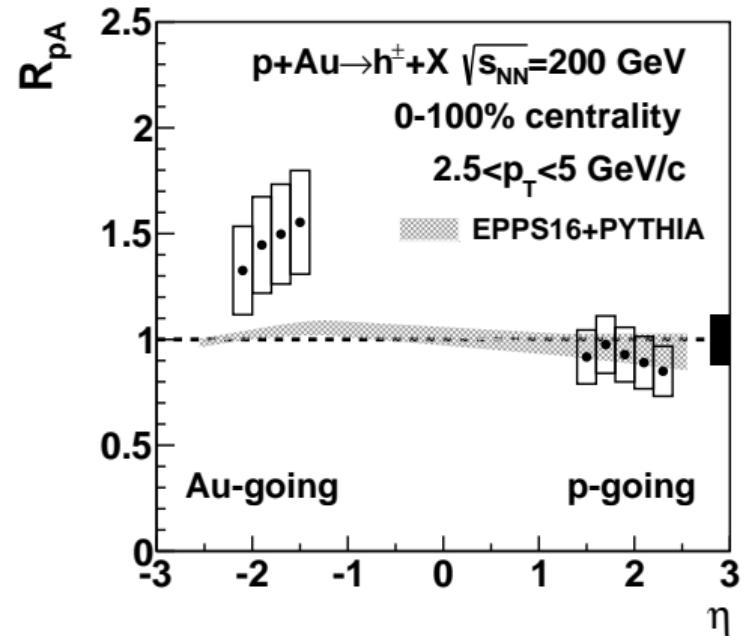
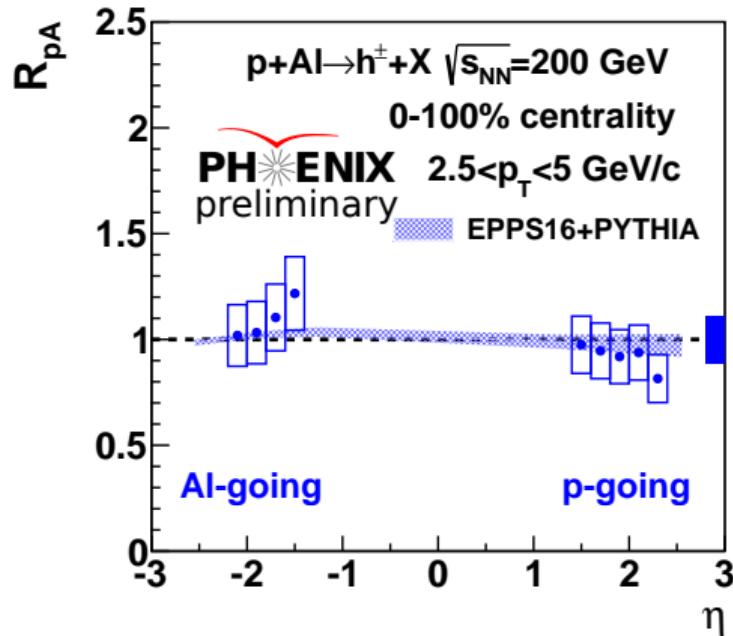
$$v_2\{4\} = (v_2^4 - 2v_2^2\sigma^2 - 4v_2 s\sigma^3 - (k-2)\sigma^4)^{1/4}$$

- Eccentricity fluctuations alone go a long way towards explaining this
- Additional fluctuations in the (imperfect) translation of ε_2 to v_2 ?

Intermission

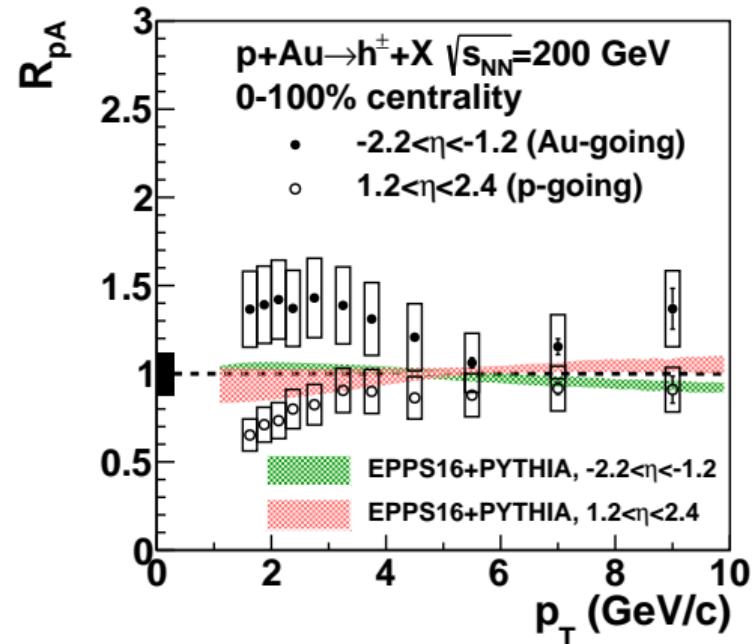
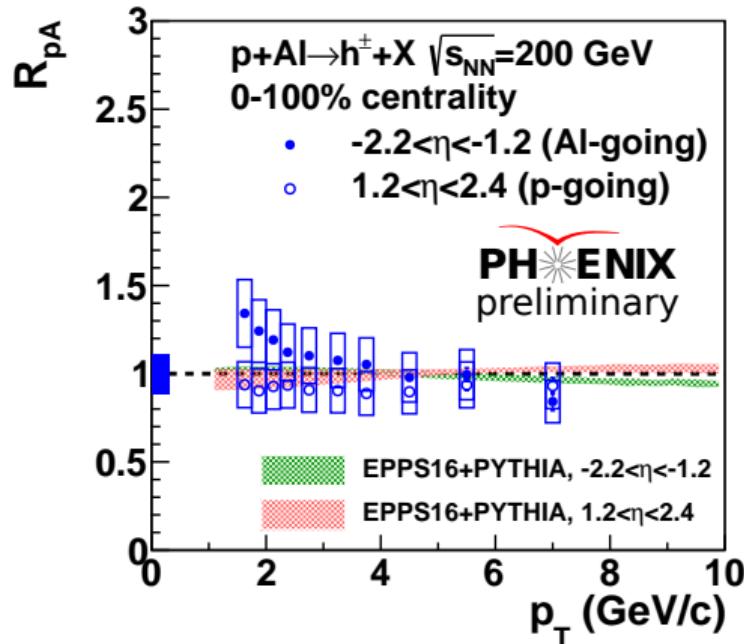
Particle production in small systems

Small systems nuclear modification



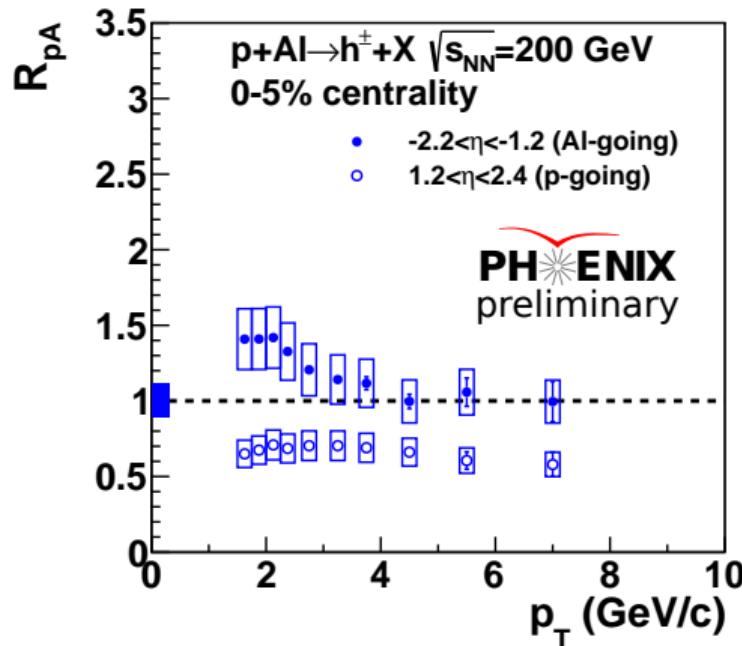
- Forward modification consistent with nPDF effects (EPPS16)

Small systems nuclear modification

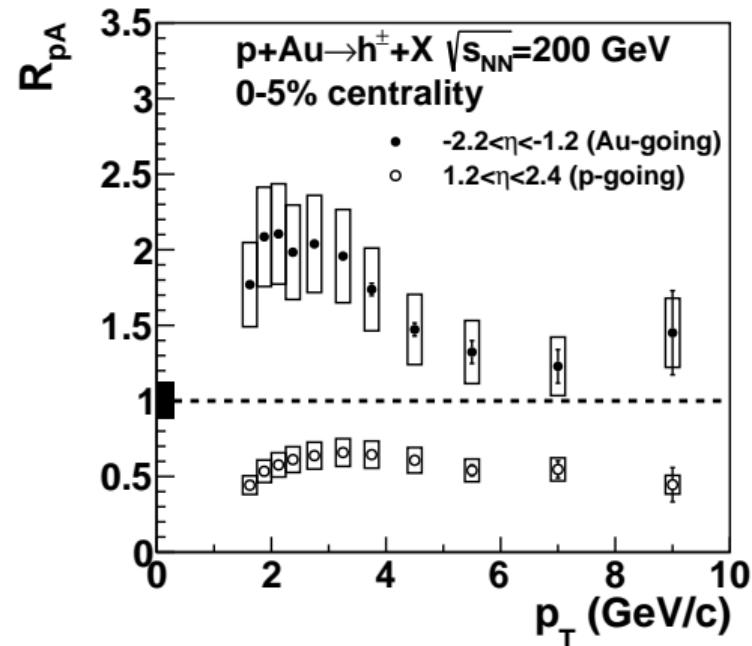


- High- p_T modification consistent with nPDF effects (EPPS16)

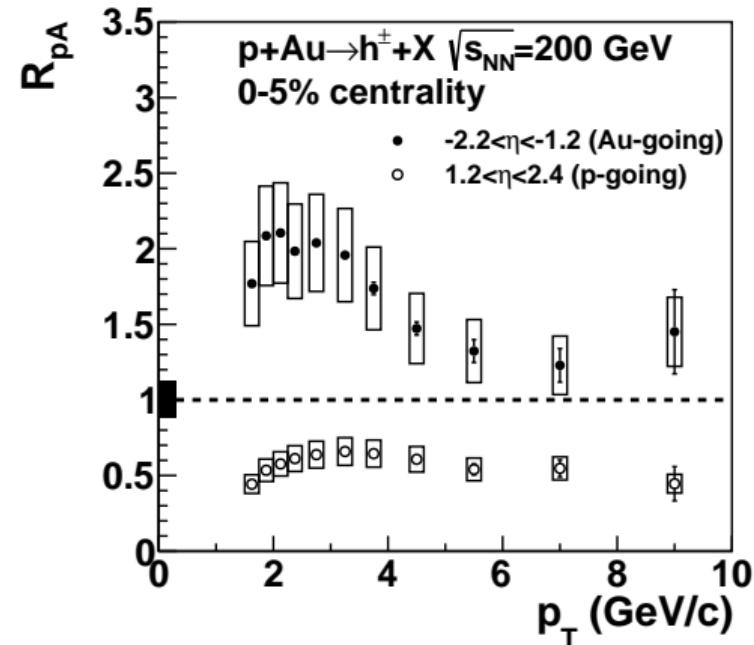
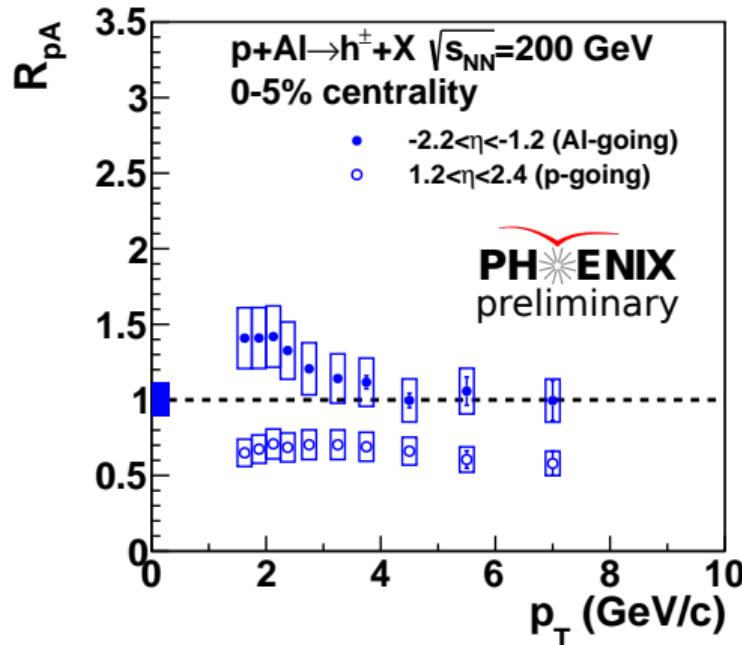
Small systems nuclear modification



- Stronger effects in central collisions

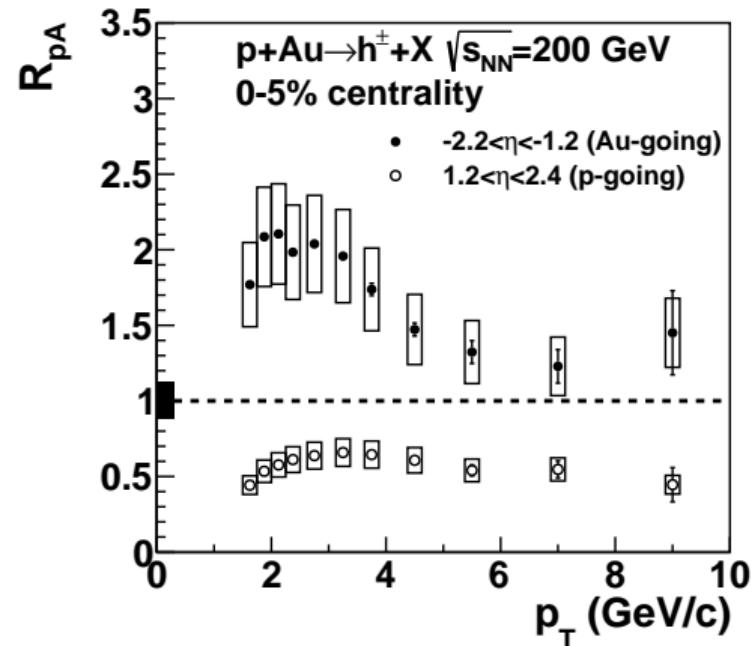
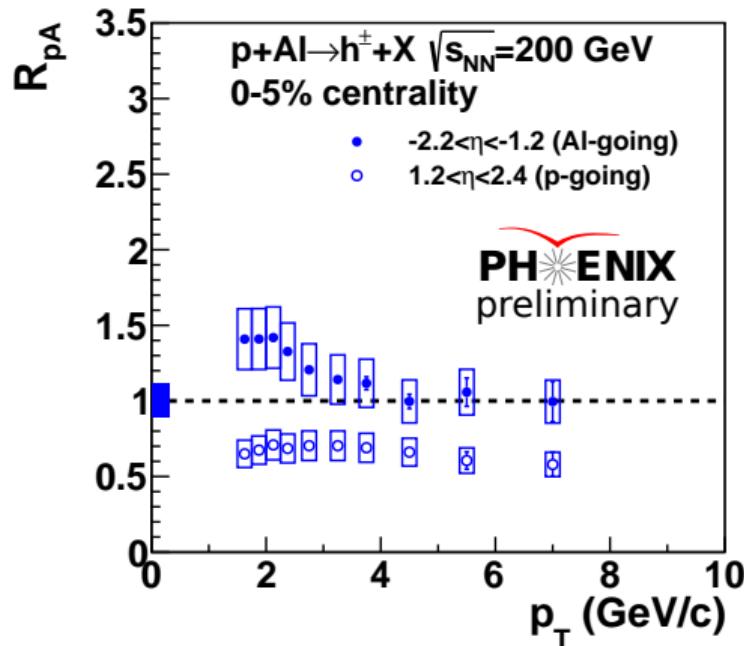


Small systems nuclear modification



- Strong enhancement for backward at intermediate p_T —why?

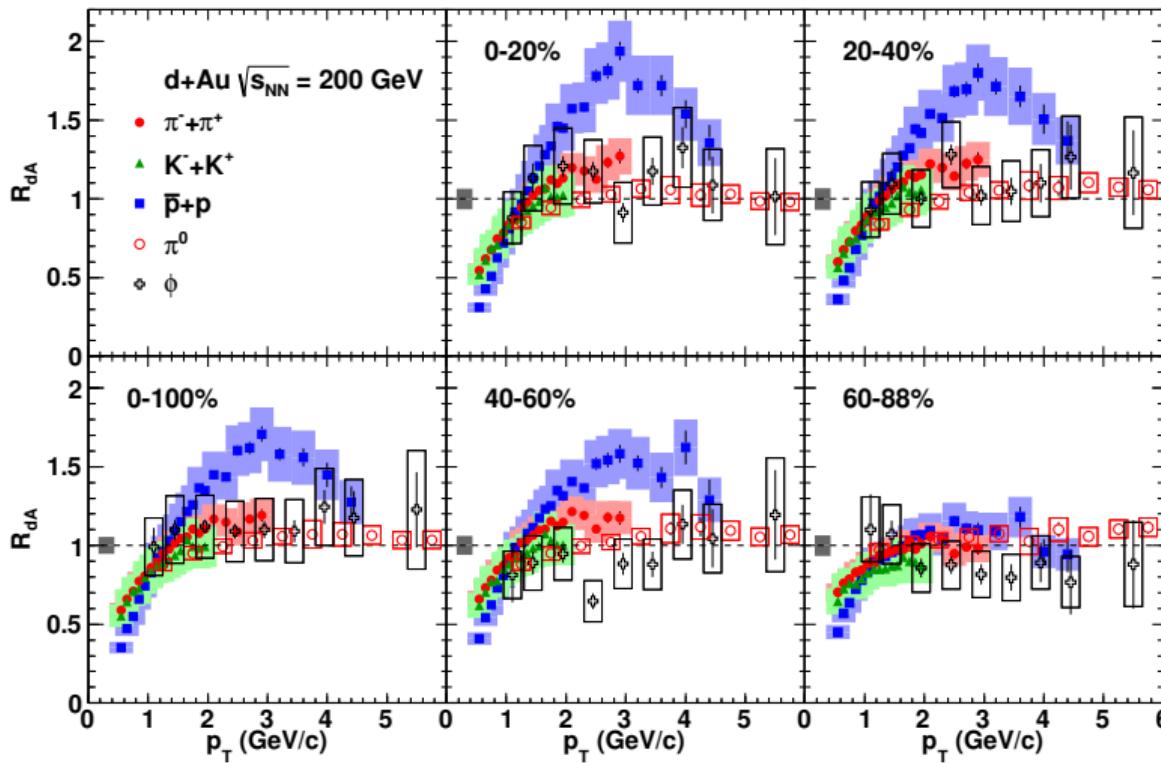
Small systems nuclear modification



- Strong enhancement for backward at intermediate p_T —why?
- Don't forget: particle species dependence of Cronin! There must be final state effect(s)...

Particle species dependence of “Cronin enhancement”

PHENIX, Phys. Rev. C 88, 024906 (2013)



$\pi^+, \pi^-, \pi^0,$
 $K^+, K^-,$
 $p, \bar{p},$
 ϕ

Protons much more strongly modified than pions

ϕ mesons confusing as always...