Experimental searches for local parity violating effects in heavy ion collisions

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- Brief introduction
- The CME
- The CMW
- Summary and outlook

Parity in the strong sector

- A non-zero neutron electric dipole moment (nEDM) violates parity
- A non-zero nEDM also violates time reversal, by CPT theorem T-violation implies CP-violation
- However, experiments have consistently found nEDM consistent with zero, upper limit 2.9×10⁻²⁶ e cm from C. Baker et al, Phys. Rev. Lett. 97 (2006) 131801
- The observed absence of CP-violation in the strong sector is surprising because of natural CP-violating terms in the QCD Lagrangian



$$\mathcal{L} = -\frac{1}{4} F^a_{\mu\nu} F^{\mu\nu}_a - \frac{\theta g^2}{32\pi^2} F^a_{\mu\nu} \tilde{F}^{\mu\nu}_a + \overline{\psi} (i \not\!\!D - m e^{i\theta'\gamma_5}) \psi$$

Strong CP problem: 0 $\leq \bar{\theta} < 10^{-10}$

Topological charge and the $U(1)_A$ anomaly

The QCD vacuum is highly non-trivial!

 $U(1)_A$ anomaly:

$$\partial_{\mu}J^{\mu}_{A}=\frac{g^{2}}{32\pi^{2}}F^{a}_{\mu\nu}\tilde{F}^{\mu\nu}_{a}$$

Topological charge:

$$Q_w = rac{g^2}{32\pi^2}\int d^4x\, F^a_{\mu
u} ilde{F}^{\mu
u}_a \in \mathbb{Z}$$



•
$$Q_w = N_L - N_R$$

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• $Q_w = N_L - N_R$

- Topological charge is the change in Chern-Simons number (N_{CS})
- Instanton: tunneling through barrier (all energies/temperatures, including 0)
- Sphaleron: jumping over barrier (only sufficiently high temperatures/energies)

The goal of heavy ion physics is to create a new state of matter called the quark-gluon plasma, which we hope to create by colliding heavy nuclei at the highest possible energies

At sufficiently high temperature and/or density, the gauge coupling between quarks and gluons becomes sufficiently weak that deconfinement is achieved

Some basic information about the QGP:

- Particles produced in thermal abundances
- Hydrodynamical models describe the data very well, require fast thermalization at the parton level
- The matter is extremely hot, well in excess of the critical temperature $T_c \approx 150 \text{ MeV} (10^{12} \text{ K})$ -Stellar coronae 10^6 K
 - –Core of white dwarf 10^7 K



Centrality

- Need to characterize the overlap of the two nuclei
- The most natural choice is impact parameter
- Other things you might like to know are: how many nucleons are in the overlap region (N_{part}) how many nucleon-nucleon collisions occurred (N_{coll})



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Anisotropic flow



 $\frac{dN}{d\Delta\phi}\propto 1+2\sum_{n=1}^{\infty}v_n\cos n\Delta\phi \qquad \Delta\phi=\phi-\psi_{RP}, \ \ v_n=\langle\cos n\Delta\phi\rangle$

- Collisions that are not fully overlapping have azimuthally non-uniform shape
- Initial state spatial anisotropy creates pressure gradients that drive final state momentum anisotropy—the anisotropic flow builds up early and self quenches
- Azimuthal distribution of particles can be described as Fourier series with coefficients v_n (Voloshin and Zhang, Z. Phys. C70 (1996) 665-672) -dominant term is v_2 (called elliptic flow)
- Note that ψ_n can be different from ψ_{RP} , but for now we don't need to worry about this complication



- Hydrodynamics models describe the data well at low transverse momentum p_T
- Mass splitting from common flow velocity, plotting vs transverse kinetic energy $KE_T = \sqrt{p_T^2 + m^2} m$ makes them line up
- At higher p_T baryons and mesons group together...



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- ...scaling with the number of constituent quarks n_q groups all particles together –Seen as "smoking gun" for QGP formation

The magnetic field in heavy ion collisions



- The spectating nucleons induce a magnetic field in the overlap region
- The peak strength of the magnetic field in conventional units is roughly 10¹⁷⁻¹⁸ gauss—largest magnetic field in the known universe!
 -MRIs 10⁵ gauss
 -Magnetars 10¹⁵ gauss
- The spectators define both the magnetic field and the geometry (recall earlier slide), so $\psi_B = \psi_{RP}$



The Chiral Magnetic Effect



- Chiral imbalance induced by quantum anomaly (recall $U(1)_A$ anomaly $\partial_\mu J^\mu_A = \frac{g^2}{32\pi^2} F^a_{\mu\nu} \tilde{F}^{\mu\nu}_a \to Q_w = N_L - N_R$)
- Alignment of spins by external magnetic field induces electric current of chiral quarks

$$ec{J}_V = rac{N_c e}{2\pi^2} \mu_A ec{B}$$

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• How to measure?

$$\frac{dN}{d\Delta\phi} \propto 1 + 2\sum_{n=1}^{\infty} v_n \cos n\Delta\phi \qquad \Delta\phi = \phi - \psi_{RP}, \quad v_n = \langle \cos n\Delta\phi \rangle$$

The Fourier expansion including P-odd sine terms

$$\frac{dN}{d\Delta\phi} \propto 1 + 2\sum_{n=1}^{\infty} [v_n \cos n\Delta\phi + a_n \sin n\Delta\phi] \qquad a_n = \langle \sin n\Delta\phi \rangle$$

• Normally we ignore sine terms, but now we need them



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- Positive particles should go above the reaction plane $a_1^+ > 0$
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- However...
- Q_w fluctuates about $\langle Q_w \rangle = 0$, so the CME current changes sign event by event, and therefore $\langle a_1^{\pm} \rangle = 0$

What to do? Measure 2 particle correlation with respect to the reaction plane (Voloshin, Phys. Rev. C 70 (2004) 057901)

$$\begin{aligned} \langle \cos(\phi_a + \phi_b - 2\psi_{RP}) \rangle &= \langle \cos \Delta \phi_a \cos \Delta \phi_b \rangle - \langle \sin \Delta \phi_a \sin \Delta \phi_b \rangle \\ &= [\langle v_{1,a} v_{1,b} \rangle + B_{in}] - [\langle a_{1,a} a_{1,b} \rangle + B_{out}] \end{aligned}$$



- Backgrounds uncorrelated with RP cancel
- Same sign $\langle a_1^{\pm} a_1^{\pm} \rangle > 0$
- Opposite sign $\langle a_1^\pm a_1^\mp
 angle < 0$
- Directed flow is rapidity-odd, $\langle v_1 v_1 \rangle \approx 0$
- Optimistically, $\langle \cos(\phi_a + \phi_b - 2\psi_{RP}) \rangle = -\langle a_{1,a}a_{1,b} \rangle$

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- However...
- RP dependent backgrounds remain
- If dipole fluctuations, $\langle v_1 v_1 \rangle \neq 0$

STAR, Phys. Rev. C 81, 054908 (2010)



• Strong negative correlation for same sign, consistent with CME expectation

• Essentially no correlation of opposite sign

-Possible explanation: the large medium destroys the opposite sign correlation



STAR, Phys. Rev. C 81, 054908 (2010)

- $\bullet\,$ Strong negative correlation for same sign in both Au+Au and Cu+Cu
- Positive correlation of opposite sign for Cu+Cu despite being absent in Au+Au

 Medium in Cu+Cu is small enough that some opposite sign correlation remains?



STAR, Phys. Rev. C 81, 054908 (2010)

• No opposite sign correlation in Au+Au for any $\Delta\eta$ or any $\bar{p_T}$

• Same sign correlation gets strong for smaller $\Delta \eta$ and larger $\bar{p_T}$ -The behavior in $\Delta \eta$ matches naïve expectations, different for $\bar{p_T}$ ALICE, Phys. Rev. Lett. 110, 012301 (2013)



- ALICE results consistent with STAR results
- Naïve expectation is for weaker correlation due to shorter B-field lifetime



ALICE, Phys. Rev. Lett. 110, 012301 (2013)

• ALICE results consistent with STAR results for both $\Delta \eta$ and $\bar{p_T}$



Y. Hori, Nucl. Phys. A 904-905 (2013) 475c-479c

• Double harmonic correlator should have no CME signal, only backgrounds

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- Double harmonic correlator should have no CME signal, only backgrounds
- Difference between same sign and opposite sign consistent with zero



- $\bullet\,$ LCC: local charge conservation—charges are created in \pm pairs at a single space-time point
- Angle between pairs is collimated by the radial+anisotropic flow backround
- Simple and intuitive mechanism for generating charge-dependent angular correlations



S. Schlichting and S. Pratt, Phys. Rev. C 83 (2011) 014913

- Construct a simple model of LCC+flow using the Blastwave model
- Results show very good agreement with STAR CME correlator results (OS-SS)
- However, the absence of OS correlation and the strong SS correlation is not explained in this (simple) model
- This may indicate that the CME correlator results contain a combination of background and new physics



Y. Hori, Nucl. Phys. A 904-905 (2013) 475c-479c

• Tuning the model parameters to match the CME correlator results creates a significant mismatch between model and data two particle correlations

Before moving on to the chiral magnetic wave, we need to briefly discuss the chiral separation effect (CSE) $% \left(\left(\left({{\rm{CSE}}} \right) \right) \right) = 0$

- D.T. Son and A.R. Zhitnitsky, Phys. Rev. D 70, 074018 (2004)
- M.A. Metlitski and A.R. Zhitnitsky, Phys. Rev. D 72, 045011 (2005)
- Quantum anomalies at finite vector charge density drives the following relation

$$\vec{J}_A = rac{N_c e}{2\pi^2} \mu_V \vec{B}$$

- This effect, an axial current proportional to a vector chemical potential, is called the chiral separation effect (CSE)
- It is readily apparent that there is a strong relationship to the CME

$$ec{J}_V = rac{N_c e}{2\pi^2} \mu_A ec{B}$$

And with that, onward to the chiral magnetic wave

The Chiral Magnetic Wave



- CSE leads to separation of chiralities at opposite poles
- CME currents point in opposite directions, leading to electric quadrupole
- Gorbar, Miransky, and Shovkovy, Phys. Rev. D83, 085003 (2011)
- Kharzeev and Yee, Phys. Rev. D83, 085007 (2011)

The Chiral Magnetic Wave

Azimuthal distribution of charges

$$rac{dQ}{d\Delta\phi}=Q[1-r_e\cos(2\Delta\phi)]$$

Definition of charge asymmetry A

$$A = \frac{Q}{N^{total}} = \frac{N^+ - N^-}{N^+ + N^-}$$

Azimuthal distribution of particles

$$rac{dN^{\pm}}{d\Delta\phi} = N^{\pm} [1 + (2v_2 \mp r_e A) \cos(2\Delta\phi)]$$



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STAR, Phys. Rev. Lett. 114, 252302 (2015)



• Charge asymmetry $A_{ch} = A = (N^+ - N^-)/(N^+ + N^-)$

- Note change in x-axis scale on right plot—correction for efficiency/acceptance
- Qualitatively consistent with CMW picture



Strong, clear signal

Qualitatively consistent with STAR results



- Strong, clear signal
- Qualitatively consistent with STAR results
- Necessary to make correction

- v₂ as a function of A is very interesting, but requires efficiency correction (due to binomial sampling)
- So what else can we do? Measure the covariance: $\langle v_2 A \rangle \langle v_2 \rangle \langle A \rangle$
- A fundamental feature of statistics is that the covariance between two variables is independent of sample size, i.e. no efficiency correction for this measure is needed (we call it a "robust observable")
- It is both straightforward and very useful to generalize to arbitrary harmonic v_n
- Since v_n is a 2-point correlation, this is a 3-point correlation
- Can also generalize A to the charge of a third particle q_3 , since $\langle q_3
 angle_{\mathsf{event}} \equiv A$
- We have then $\langle v_n A \rangle \langle v_n \rangle \langle A \rangle$ and $\langle v_n q_3 \rangle \langle v_n \rangle \langle q_3 \rangle$ as observables to evaluate



What causes the increased charge separation as the collisions become more peripheral?

- Peripheral \rightarrow stronger magnetic field \rightarrow stronger CMW effect?
- \bullet Central \to more combinatoric pairs \to trivial dilution of local charge conservation (LCC) effects?
- Dependence on magnitude of v_2 or dN/dy?
- Some combination of these (and possibly other) effects?

ALICE, arXiv:1512.05739 (submitted to PRC)



- CMW quadrupole expected to affect only 2nd harmonic, LCC expected to affect all harmonics
- Small effect for 3rd harmonic, no observed effect for 4th harmonic –Note reduced y-axis scale compared to 2nd harmonic
- Higher order multipole effects for CMW or harmonic interference? LCC only?



- Trivial correlations are inversely proportional to N
- Multiplication of 2nd harmonic three-particle correlator by $dN_{ch}/d\eta$ shows correlation that still increases as the collisions become more peripheral
- This may indicate that correlation contains a non-trivial component



- Trivial correlations are inversely proportional to N
- Multiplication of 3rd and 4th harmonic three-particle correlator by $dN_{ch}/d\eta$ shows correlation that is roughly flat with centrality
- This may indicate a different nature of the correlation depending on the harmonic number

- Recall the integral correlator $\langle v_2 A
 angle \langle A
 angle \langle v_2
 angle$
- Hypothesis: $v_2^{\pm} = \bar{v_2} \mp rA/2$
- Plug and chug, rewrite to get

$$\langle v_2^{\pm} A \rangle - \langle A \rangle \langle v_2^{\pm} \rangle \approx \mp r \left(\langle A^2 \rangle - \langle A \rangle^2 \right) / 2 = \mp r \sigma_A^2 / 2.$$

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- $\bullet\,$ Note that the σ_A^2 must be the true variance, so the observed σ_A^2 must be efficiency corrected
- As with the direct measurement of v_2 vs A, the evaluation of the slope parameter in this way depends on corrections from Monte Carlo



- Reasonable agreement with STAR for mid-central collisions
- Weaker overall centrality dependence



- Reasonable agreement with STAR for mid-central collisions
- Weaker overall centrality dependence
- But is agreement with STAR really to be expected?

STAR, Phys. Rev. Lett. 114, 252302 (2015)



• Very good agreement between theory and experiment

STAR, Phys. Rev. Lett. 114, 252302 (2015)



• Very good agreement between theory and experiment

(Almost too good...)

• What kind of differential studies can we do with this correlator?

ALICE, arXiv:1512.05739 (submitted to PRC)



- Generalizing from A to q_3 as discussed, we can measure the correlator as a function of the separation between particles 1 and 3, $\Delta \eta = \eta_1 \eta_3$
- $\bullet\,$ Doing so we can directly measure the η range and dependence of the charge dependent effect
- $\bullet\,$ LCC and CMW correlations may have different η ranges, providing an additional experimental constraint
- However, we're missing something very important...

• Let us examine the three particles more carefully



- Let us examine the three particles more carefully
- 1 is the particle of interest, and we consider both ϕ_1 and q_1



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- The correlation between 1 and 2 is the harmonic coefficient



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- The correlation between 1 and 3 is the balance function



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- When removing the charge correlation between 1 and 3, all reducible correlations have been removed and the correlator is a cumulant $\langle \langle \cos(n(\phi_1 \phi_2))q_3 \rangle \rangle$
- S.A. Voloshin and R. Belmont, Nucl. Phys. A 931 (2014) 992-996



- $\langle q_3
 angle$ denotes mean charge (i.e. independent of q_1)
- $\langle q_3
 angle_1$ denotes mean charge depending on q_1
- The mean charge of the third particle is affected by the charge of the first particle due to charged pair production (the balance function)
- How does this affect the three particle correlator?



• Charge independent subtraction (charge correlation not considered)

ALICE, arXiv:1512.05739 (submitted to PRC)



• Charge dependent subtraction (charge correlation considered)

ALICE, arXiv:1512.05739 (submitted to PRC)



- Charge dependent subtraction (charge correlation considered)
- The observed effect has a large contribution from the dependence of q_3 on q_1
- Both the strength and range are significantly reduced, but a pronounced charge dependent effect remains
- How much contribution from charge conservation has been removed? Is there some way to remove all LCC effects leaving only CMW?



- Charge independent subtraction
- Moderate effect for 3rd, minimal effect for 4th



3rd harmonic

4th harmonic



- Charge dependent subtraction
- Very little effect for either

S.A. Voloshin and R. Belmont, Nucl. Phys. A 931 (2014) 992-996



- Construct a simple model of LCC+flow using the Blastwave model
- Results for the CMW correlator qualitatively and semi-quantitatively match the experimental results
- Note that the magnitude of the side "dips" very closely matches experiment, while the magnitude at $\Delta\eta\approx$ 0 is lower
- This may indicate that the CMW correlator results contain a combination of background and new physics

- Local parity violation is a fundamental feature of QCD
- In an important sense, it *must* be there, but that doesn't mean it's present in the heavy ion collisions we can measure
- In fact there are several key issues
- Does the magnetic field live long enough?
- Are the quarks formed early enough?
- Neither of those questions has been addressed yet, though work is ongoing to try to answer them
- Presence of B-field can be evinced by charge and rapidity dependent v_1
- Presence of quarks can be evinced by charge dependent v_1 in A+B collisions Can do this in Cu+Au collisions at RHIC, see e.g. T. Niida QM2015

- The biggest issue (of course) is understanding the backgrounds
- At the current time, the only viable candidate for background to the CME and CMW observables is local charge conservation on top of strong flow
- The current modeling gets some observables right but others wrong, this is very important work, and studies are ongoing
- Promising avenue of investigation: anomalous hydrodynamics, which embeds the LPV effects in a realistic hydrodynamical medium
- There's no smoking gun yet, but there's always more work to do

"The optimist regards the future as uncertain."-Eugene Wigner

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Thank you!

Additional material