

**Measurement of Charged Particle Multiplicity  
at  
RHIC  
with  
PHENIX Multiplicity & Vertex Detector**

Tahsina Ferdousi  
University of California Riverside

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# Outline

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- Motivation
- Quark-gluon Plasma
- Heavy-ion Collisions at RHIC
- PHENIX Experiment at RHIC
- The Multiplicity & Vertex Detector at PHENIX
- Measurement of Charged Particle Multiplicity with MVD
- Model Comparison
- Summary

# Motivation

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- Heavy Ion Collisions at RHIC Energies are expected to produce a new Phase of matter (Quark-gluon plasma).
  - There are a number of potential signatures to be studied for the formation of the QGP.
  - The Signatures are studied simultaneously as a function of the energy density  $\varepsilon$ .

- The charged particle multiplicity  $dN/d\eta$  is related to  $\varepsilon$  by the relation

$$\varepsilon = \frac{\sqrt{m_T^2 + \langle p_T^2 \rangle}}{\pi R^2 \tau_0} \frac{3}{2} \frac{dN}{dy}$$

where  $\langle p_T \rangle$  is the mean transverse energy of the tracks produced in the collision,  $R \sim A^{1/3} \sim 6$  fm for Au and  $\tau_0$  is the formation time  $\sim 1$  fm/c

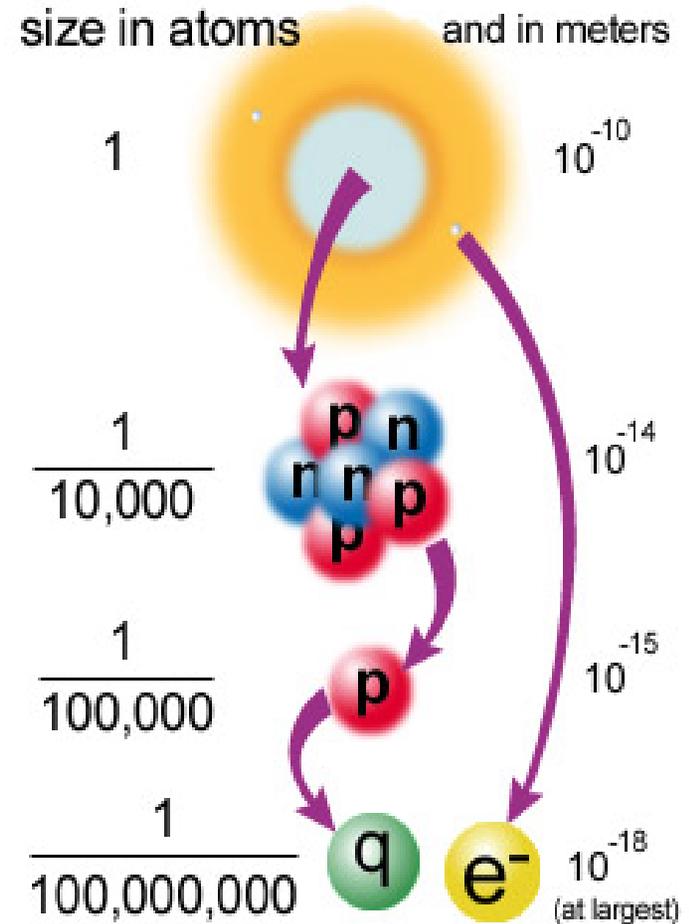
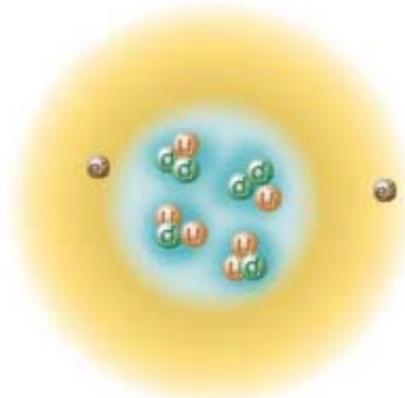
- Therefore the charged particle multiplicity can tell the initial energy density of formation.
- The potential signatures can be studied as a function of the  $dN/d\eta$  ( $dE_T/dy$ ), which are related to  $\varepsilon$ .

# Basic Building Block

An *atom* is made up of *nucleus* and *electrons*

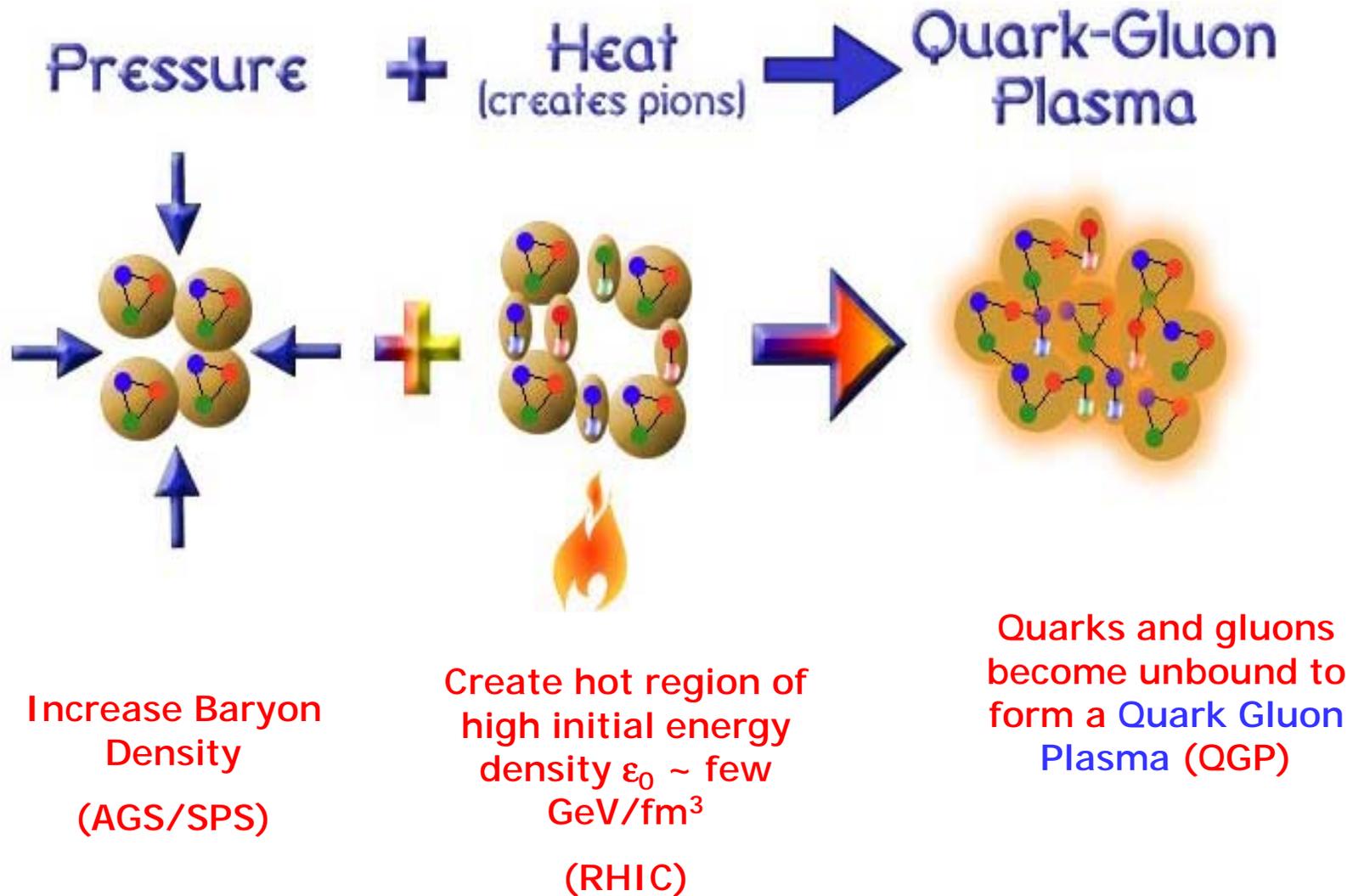
Nucleus contains *protons* and *neutrons*...

...which contain *up* and *down* *quarks*.





# How to Create QGP



# Kinematics

Transverse momentum

$$p_T = \sqrt{p_x^2 + p_y^2} = p \sin \theta$$

Transverse Mass

$$m_T = \sqrt{p_T^2 + m_0^2}$$

Rapidity

$$y = \frac{1}{2} \ln \left( \frac{E + p_z}{E - p_z} \right) = \frac{1}{2} \ln \left( \frac{1 + \beta \cos \theta}{1 - \beta \cos \theta} \right)$$

Pseudorapidity

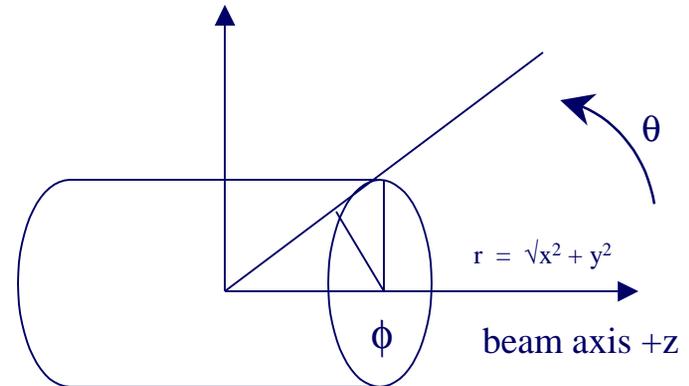
$$\eta = -\ln \left( \tan \frac{\theta}{2} \right)$$

$\eta \approx y$  when  $p \gg m_0$  and  $\theta \gg 1/\gamma$ , where  $\gamma = E/m$

Relations between momentum and (pseudo)rapidity

$$E = m_T \cosh y$$

$$p_z = m_T \sinh y = p_T \sinh \eta$$



PHENIX Coordinate System

Longitudinal : along z

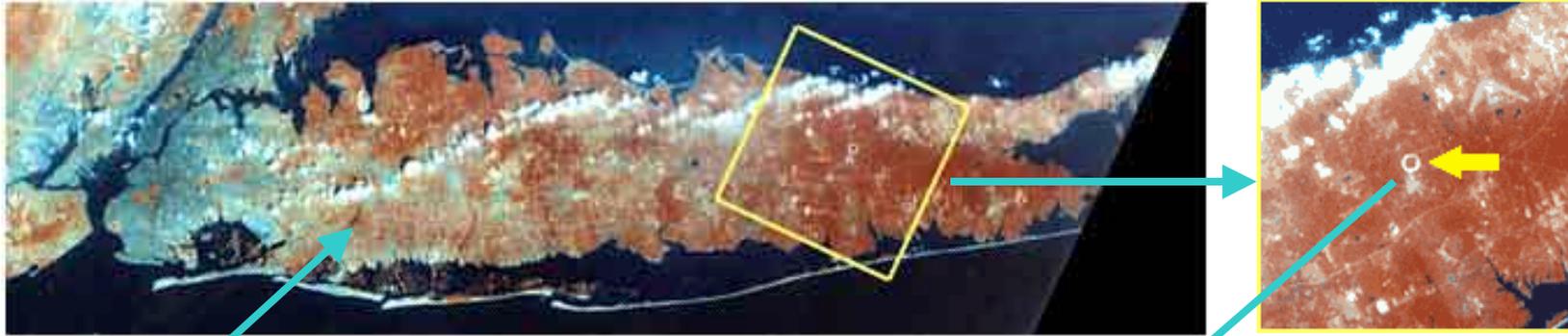
Transverse : x-y plane

Rapidity & Pseudorapidity

$$\eta = \frac{1}{2} \ln \left[ \frac{\sqrt{m_T^2 \cosh^2 y - m_0^2} + m_T \sinh y}{\sqrt{m_T^2 \cosh^2 y - m_0^2} - m_T \sinh y} \right]$$

$$\frac{dN}{d\eta dp_T} = \sqrt{1 - \frac{m_0^2}{m_T^2 \cosh^2 y}} \frac{dN}{dy dp_T}$$

# Relativistic Heavy Ion Collider (RHIC)



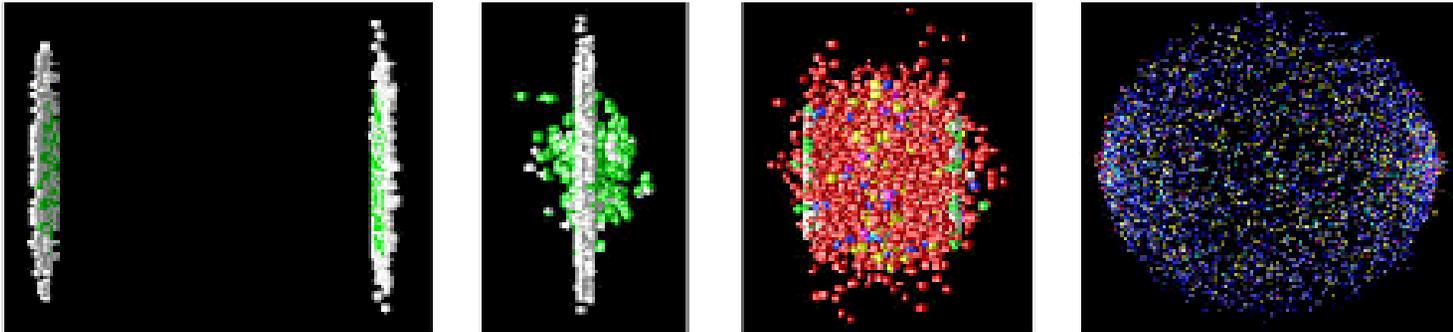
Long Island, NY

RHIC Complex at  
the Brookhaven  
National Lab



# Heavy Ion Collisions at RHIC

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Two ions approach one another. The ions are flat, instead of spherical, because they're going 99.95% of the speed of light.

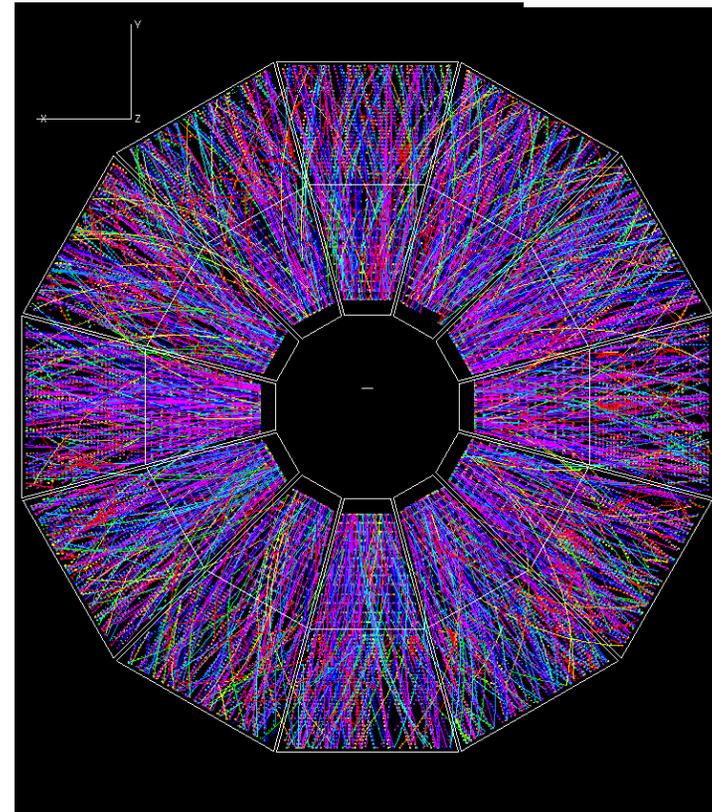
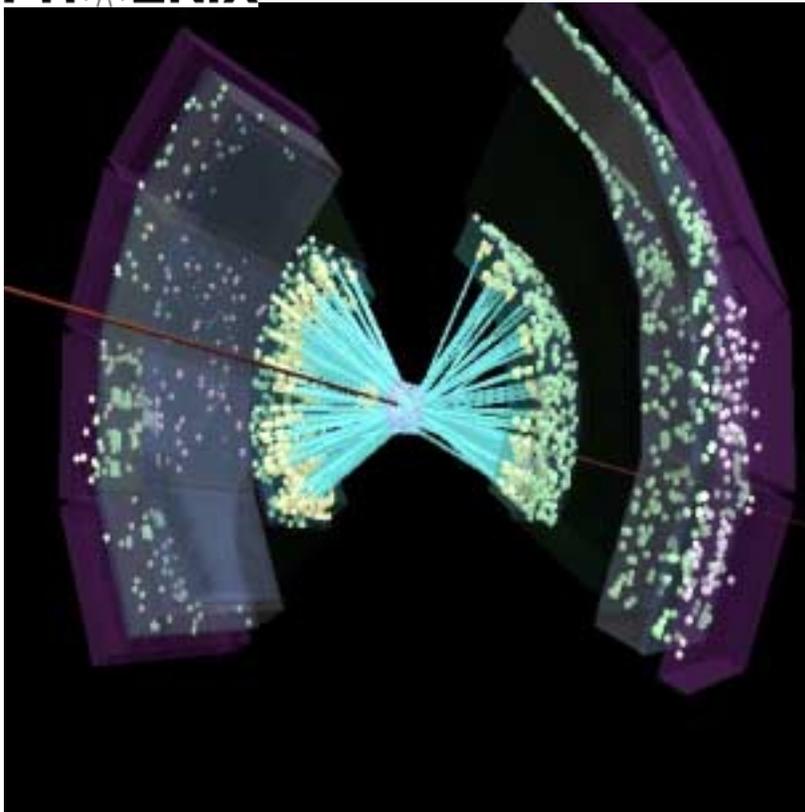
The two ions collide. If conditions are right, the collision "melts" the protons and neutrons and, for a brief instant, liberates the quarks and gluons.

Just after the collision, thousands more particles form as the area cools off. Each of these particles is a clue to what happened inside the collision.

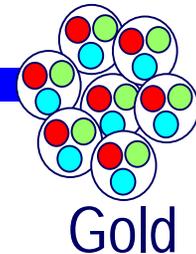
# Event Display in a Au-Au Collision

PHENIX

STAR

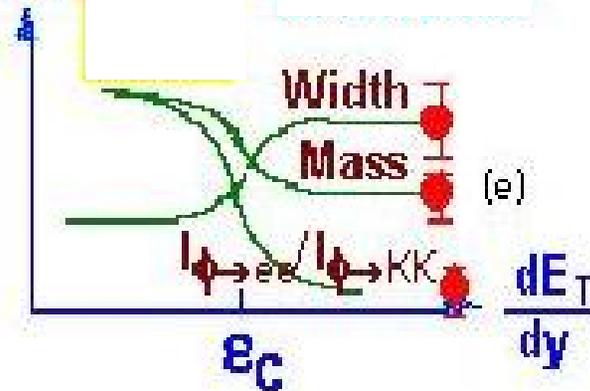


$\sqrt{s_{NN}} = 130,200 \text{ GeV}$   
(center-of-mass energy per nucleon-nucleon collision)



# Signatures of Quark-Gluon Plasma

## Phi-Meson

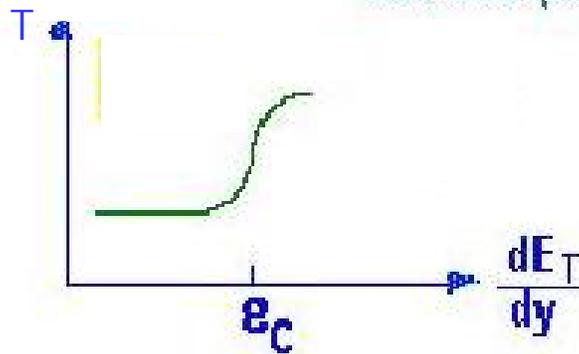


There are a number of Potential signatures for the QGP Formation.

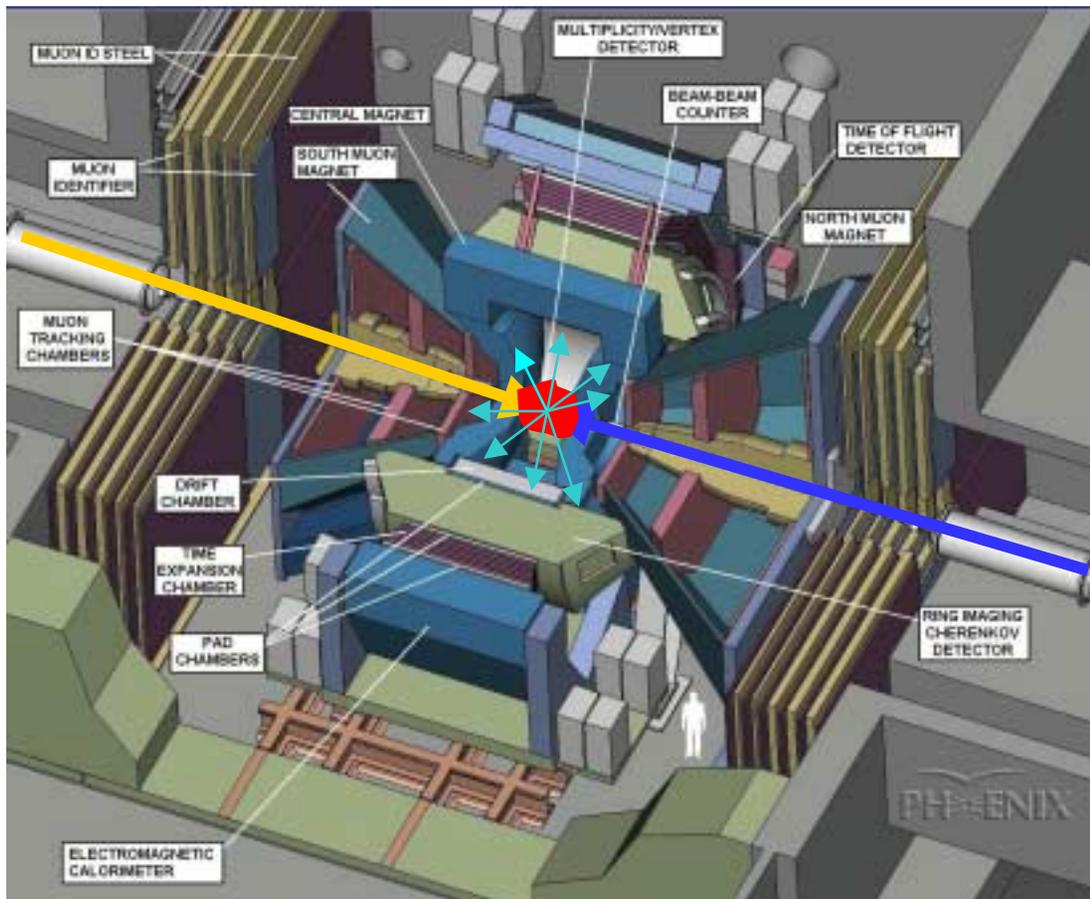
The Heavy-ion experiments study as many signatures as possible

The potential signatures are studied as a function of the energy density  $\epsilon$  or the rapidity density of the transverse energy ( $dE_T/dy$ ) or the pseudorapidity density of the charged particle multiplicity ( $dN/d\eta$ )

## Radiation from Hot Gas (thermal)



# PHENIX Experiment at RHIC



# Design Concept

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- Pair spectrometers
  - Essential for vector meson, virtual photon detection
- Good particle identification (e,  $\pi$ , K, p,  $\gamma$ ,  $\mu$  plus pairs e.g.  $\pi^0$ ,  $\omega$ ,  $\phi$ )
- Low occupancy of “pixels”
  - Essential for particle identification, systematic control of efficiency
- High data rate and selective multi-level triggering
- Two central arms
  - Electrons, photons, hadrons; Coverage:  $\Delta\phi = \pi/2$  each,  $\Delta\eta = 0.7$
- Two “muon” endcaps
  - Coverage:  $\Delta\phi = 2\pi$  each,  $\Delta\eta = 1.15-2.44$  (north), 1.15-2.25 (south) e/ $\pi$ ,  $\mu$ /h rejections better than  $10^{-3}$

# Physics Goals

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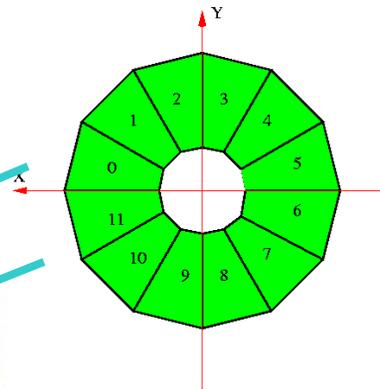
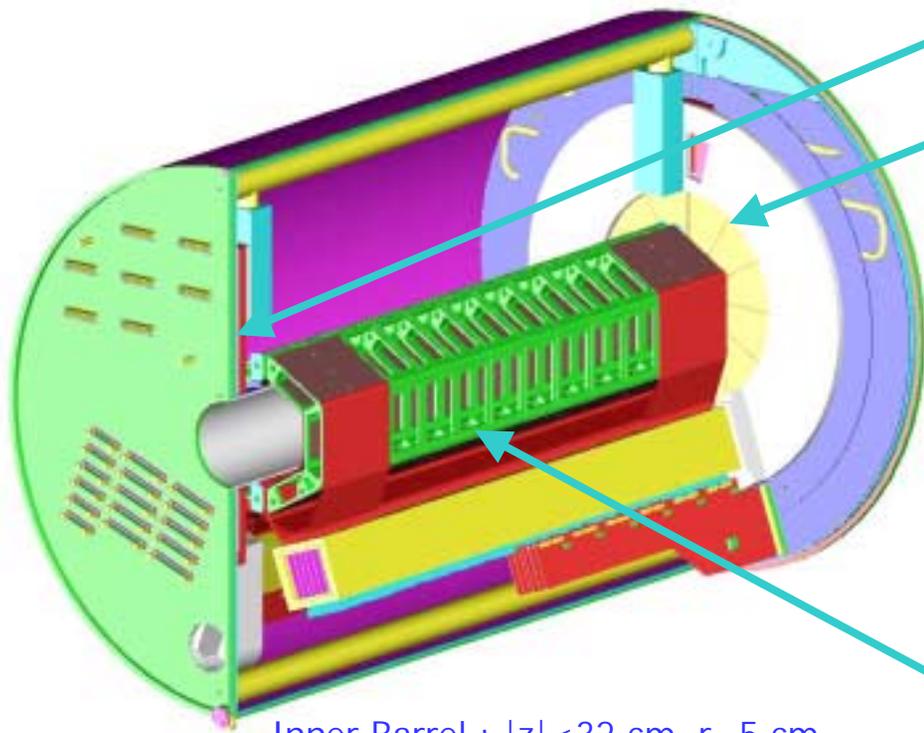
- Search for evidence of Quark-Gluon Plasma
  - Following range of probes to be covered
    - Vector mesons
    - Direct photons (real and virtual)
    - Identified hadrons
    - Collective motion and correlations
    - Global observables
  - Capability to cover  $A+A$ ,  $p+A$ ,  $p+p$  essential
- Study spin-structure of nucleon.

# Multiplicity & Vertex Detector

- Dimensions : 70 cm long, 30 cm radius
- Weight : 11 kg
- 34720 electronics channels

## Physics Goals

- Charged Particle Multiplicity
- Collision Vertex Position



End Caps :  $z = \pm 35$  cm

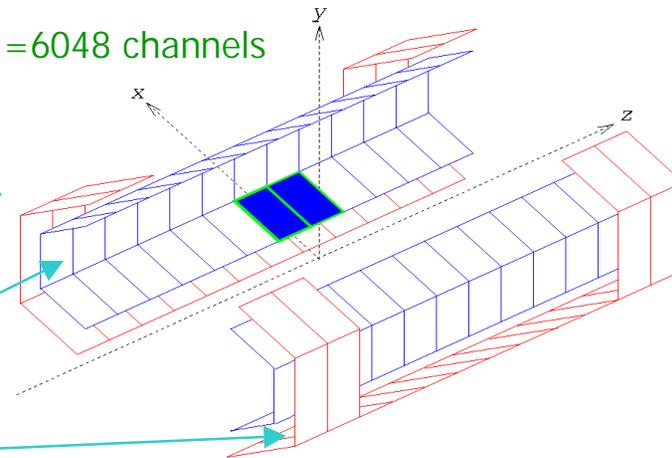
$r = 5$  to  $12$  cm

(2 end-caps \* 12 pads \* 252)

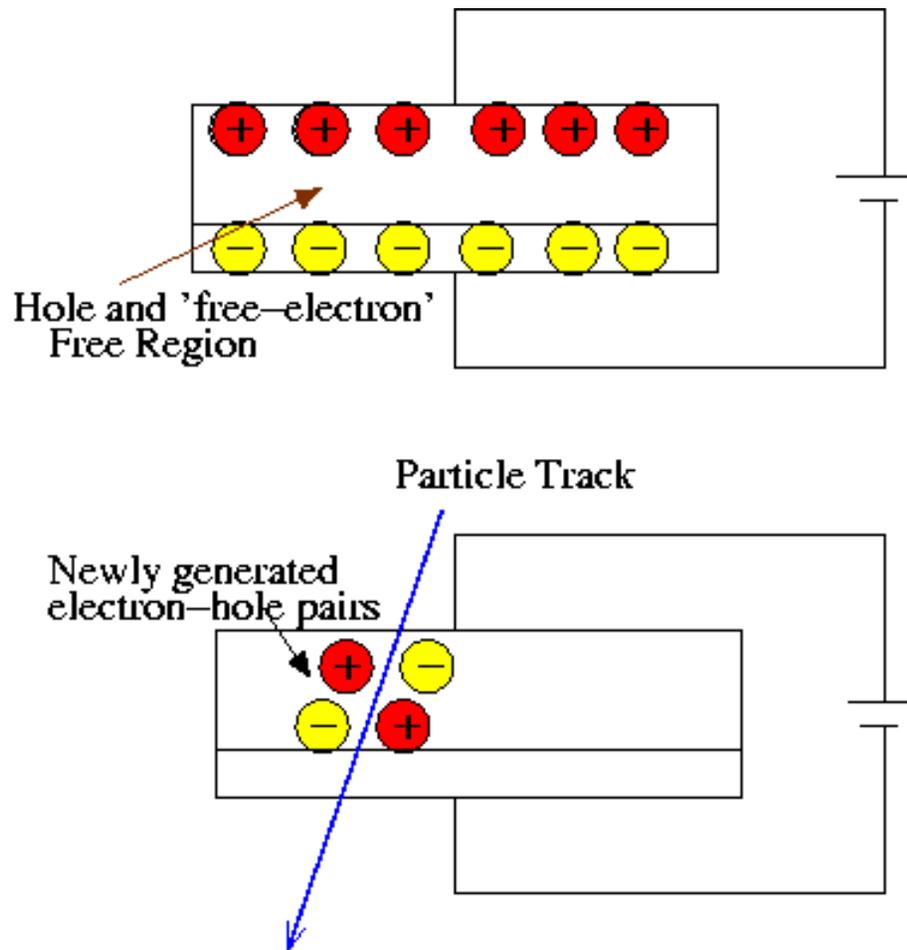
= 6048 channels

Inner Barrel :  $|z| < 32$  cm,  $r = 5$  cm  
(72 panels \* 256 strips) = 18432 channels

Outer Barrel :  $|z| < 32$  cm,  $r = 7.5$  cm  
(40 panels \* 256 strips) = 10240 channels



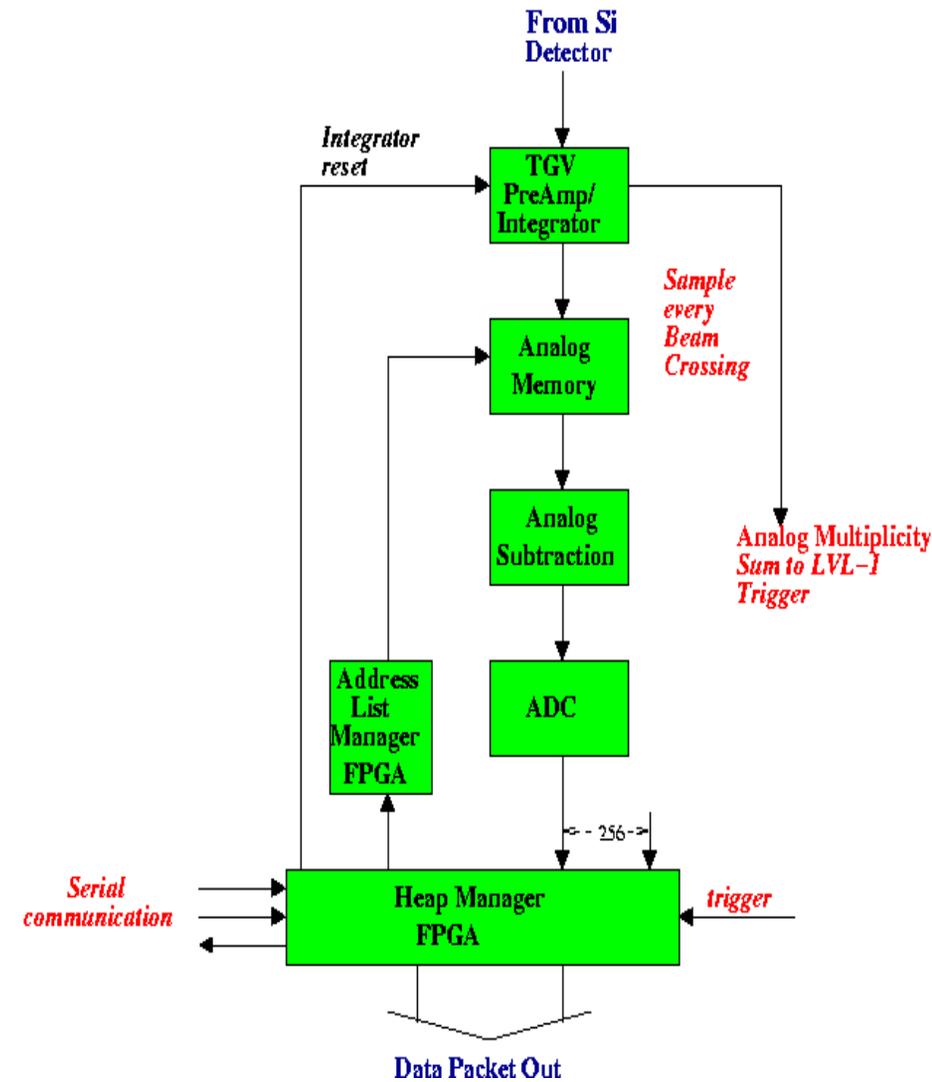
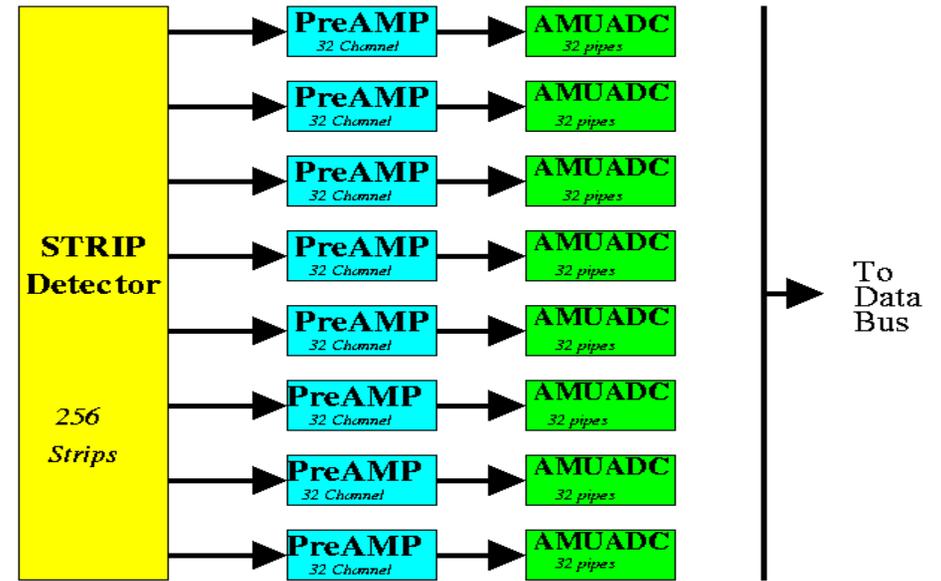
# How Does MVD (or any Silicon Detector) Work



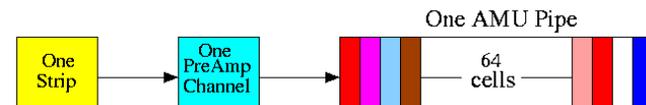
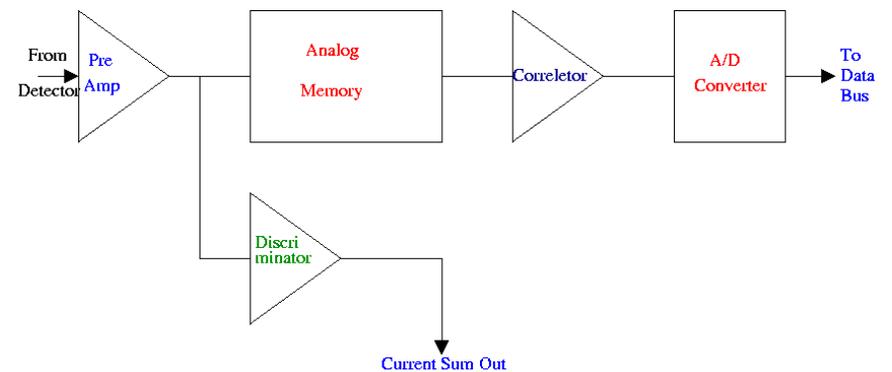
- Apply enough reverse voltage so that no free charge carriers exist
- When tracks pass through the silicon, it will deposit a fraction of its energy, which will produce free charge carriers.
- Charge carriers are attracted to the opposite terminals, and we will have signal.

# MVD Electronics Chain

Breakdown for one Strip Detector (256 strips)

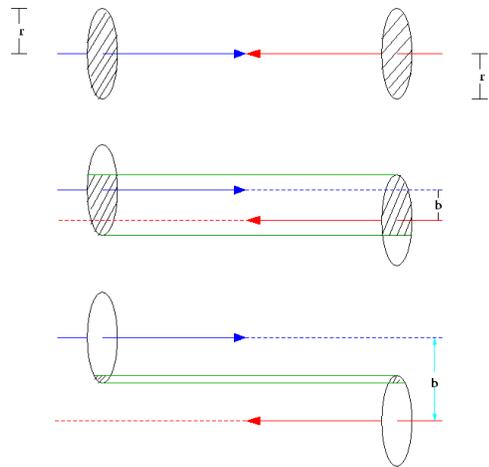


Breakdown for one strip

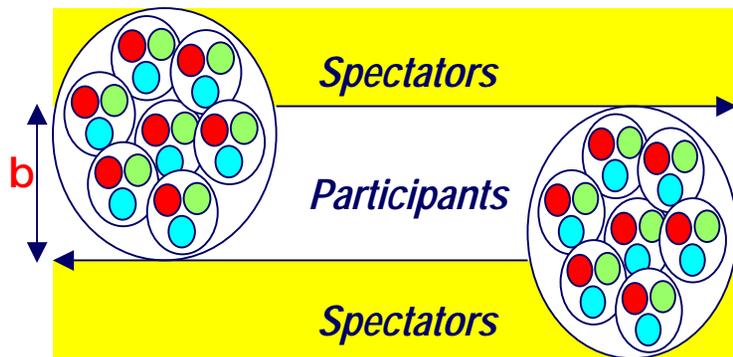
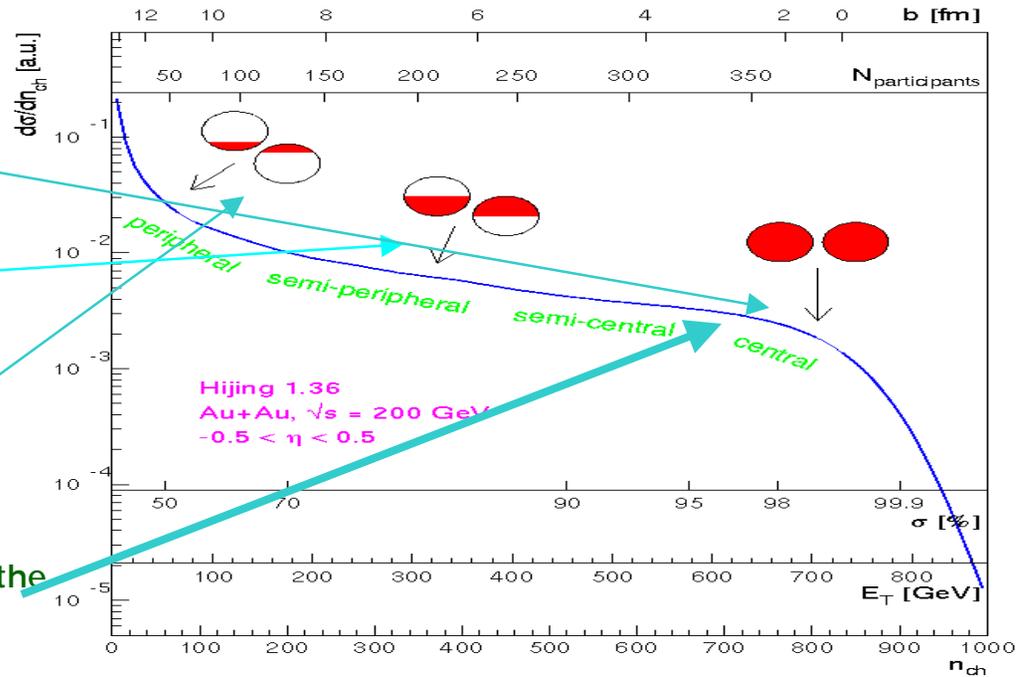




# Nuclear Collisions

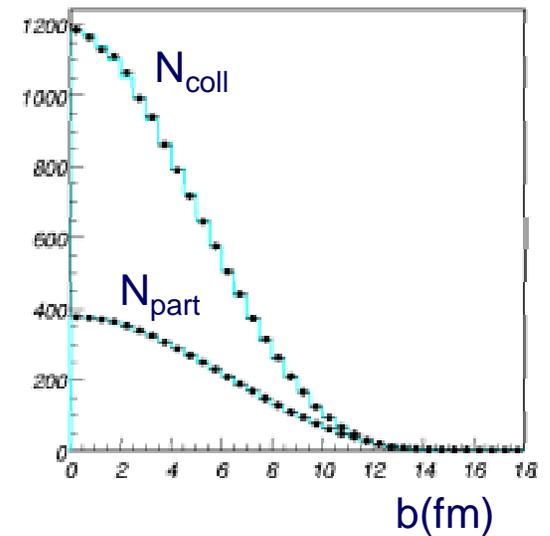


Centrality and Shape of the Multiplicity Distribution



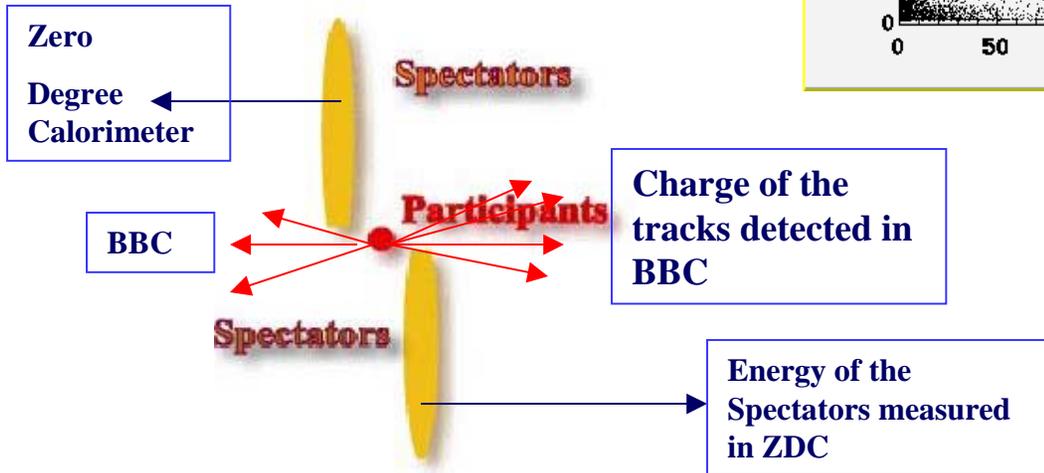
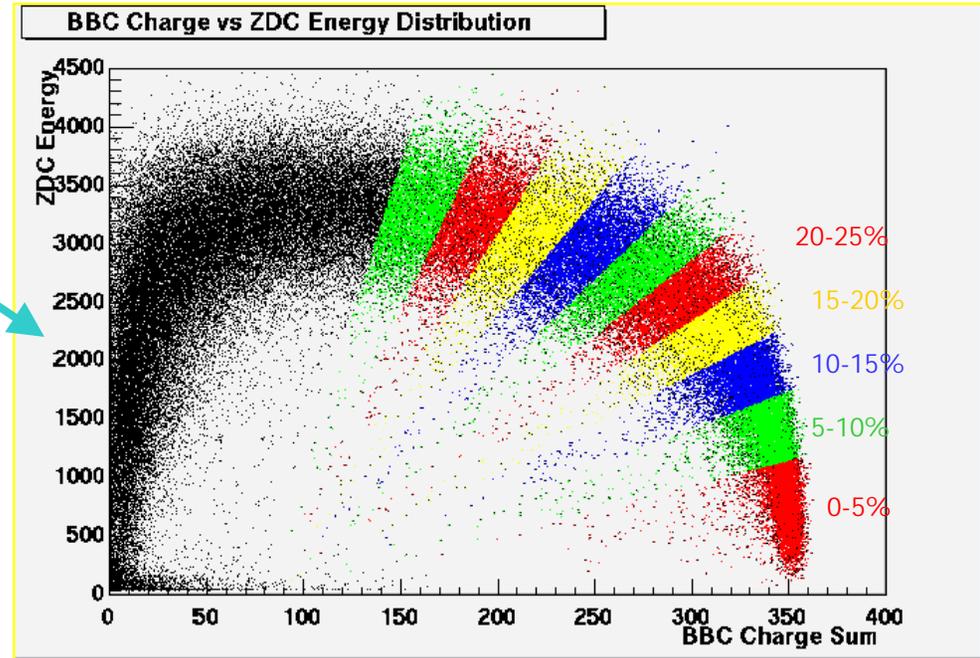
## ■ Geometrical model

- Binary collisions ( $N_{\text{coll}}$ )
- Participants ( $N_{\text{part}}$ ): Nucleons that interact
- Spectators ( $2A - N_{\text{part}}$ ): Nucleons that do not interact



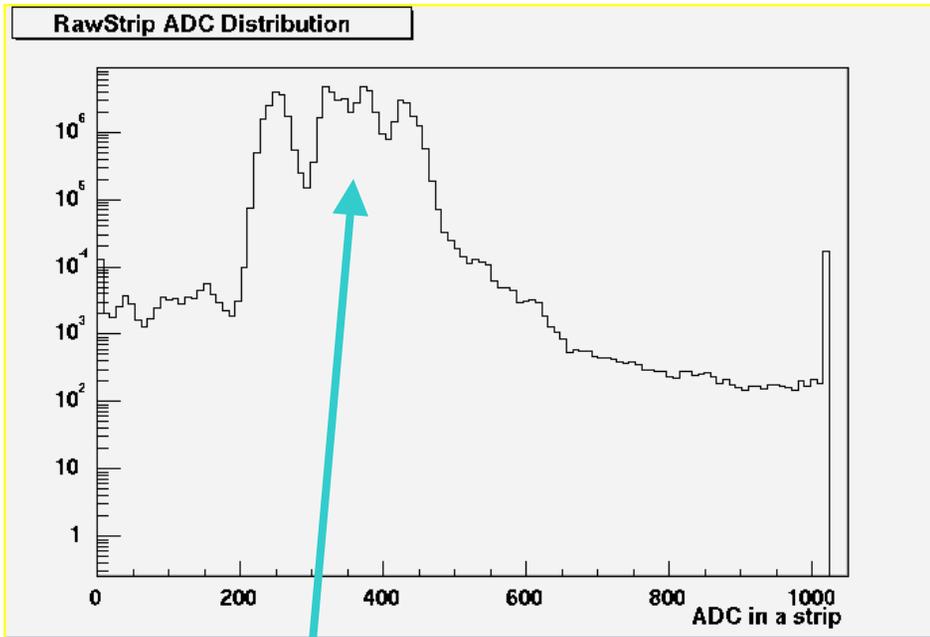
# Centrality Selection

Centrality Selection based on BBC Charge Q and ZDC Energy E



$$\phi_{\text{cent}} = \tan^{-1} \left[ \frac{(Q - Q_0) / Q_{\text{max}}}{(E - E_{\text{max}})} \right]$$

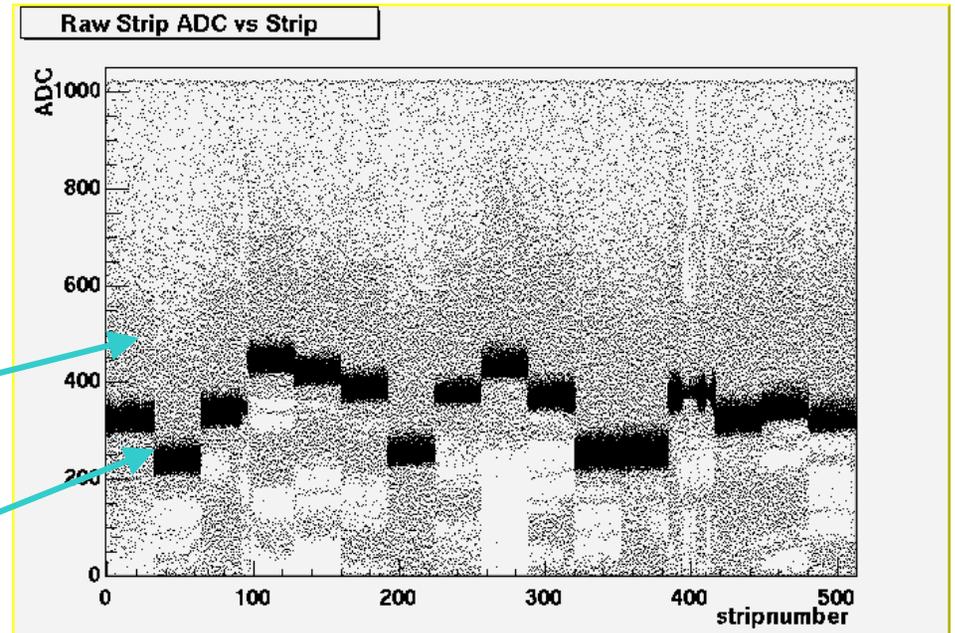
# Raw ADC Distribution for All Channels



No Calibration

## Two Major Complications:

1. Event by event jumping of pedestals
2. AMU cell dependence of pedestals

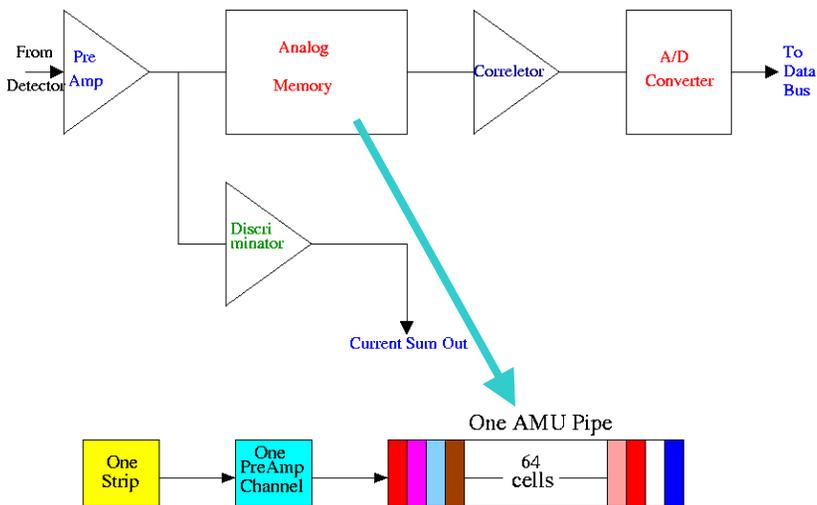


Are these Mips?

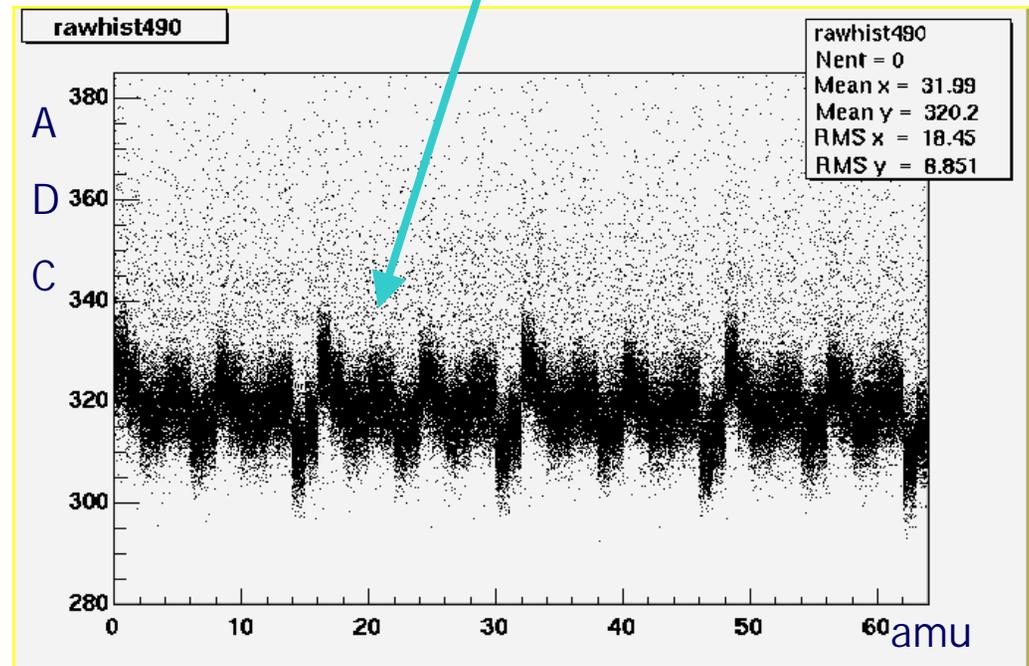
Very Wide Pedestal

# AMU Cell Dependence of Pedestals

- ▶ AMUs are analog equivalent of RAM
- ▶ Each Memory Cell is a small Capacitor
- ▶ Each Strip has one-to-one correspondence with a memory pipe
- ▶ Each pipe is 64 cell deep



Strong Dependence on AMU cell position



# Correction for the AMU Cell Dependence of Pedestals

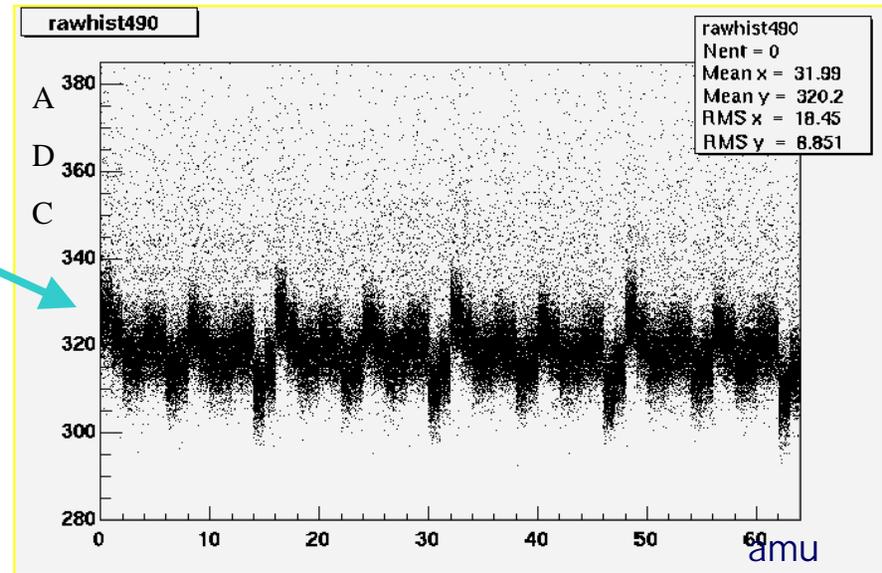
- Plot an 1-D plot for each channel for each amu.

- For each plot find the mean of the ADC

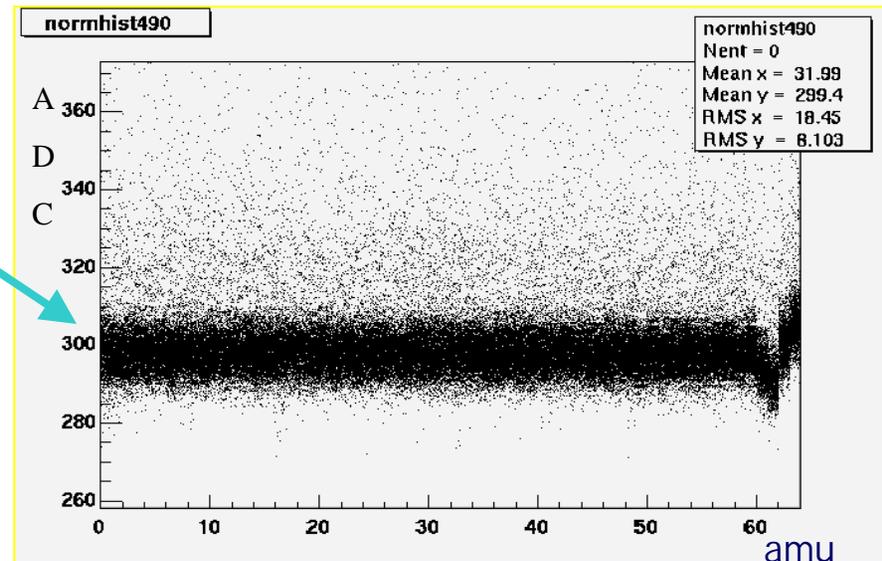
- Normalize ADC around 300.

$$\text{ADC}(\text{normalized}) = (\text{ADC}(\text{raw}) / \text{mean}) * 300$$

Before  
Correction



After  
Correction



# Event by Event Correction

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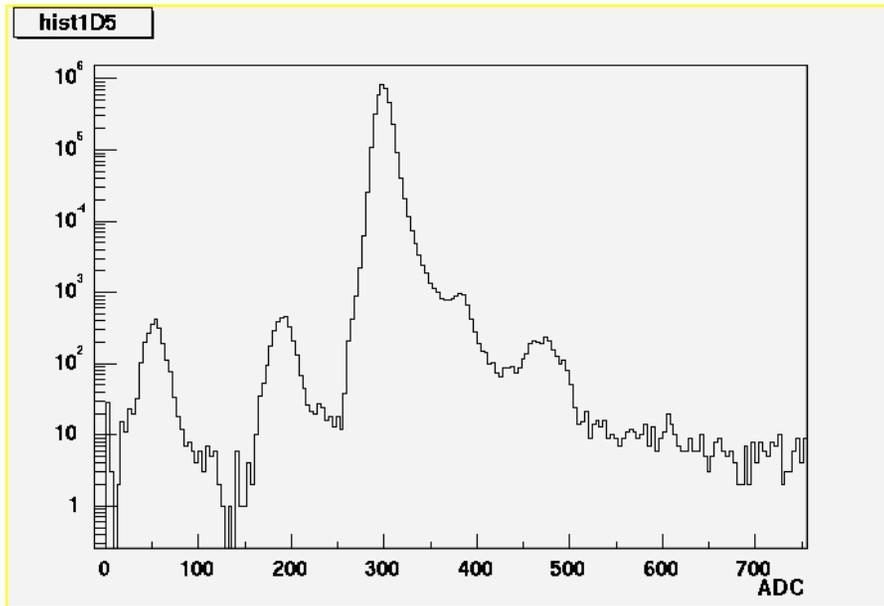
## Objective

- ✚ In each event Identify the “true” pedestals.
- ✚ Once identified, find the mean. This mean contains the amount the pedestals “jumped” in that event.
- ✚ Then subtract the mean from all the measurements in that event. This eliminates the pedestals jumping.

## Procedure

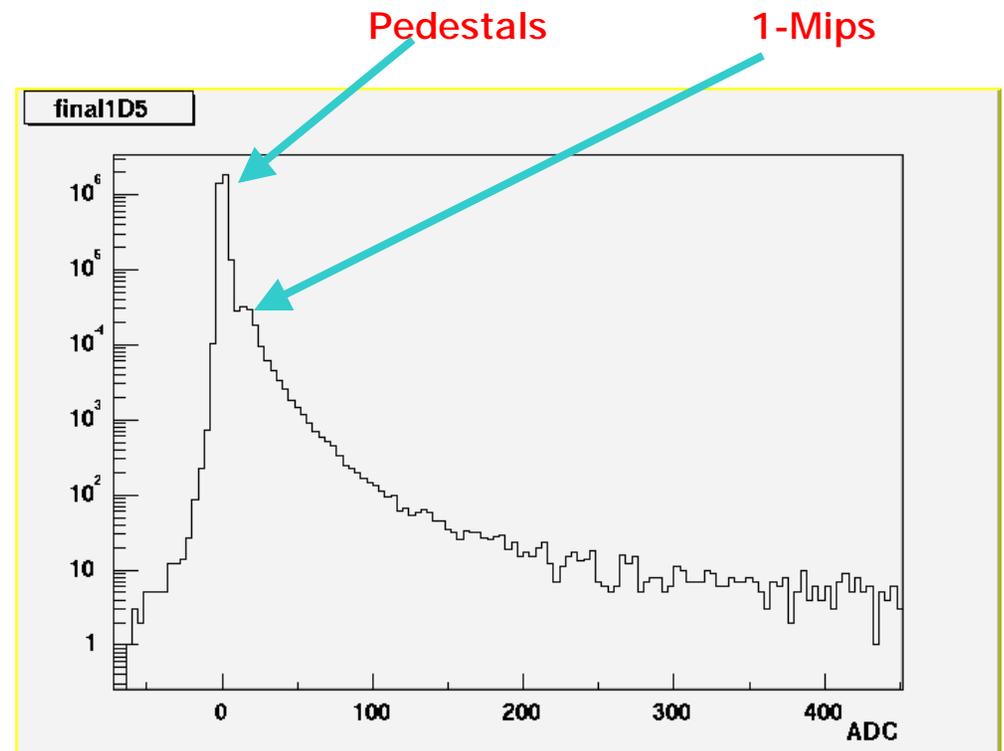
- Step1 : Find the lowest 30% entries in each event. These entries are undoubtedly pedestals. Plot them and get Mean and RMS.
- Step2 : Extend the window. We find the entries that are less than  $(\text{Mean} + 2 \times \text{RMS})$ . Plot them and get Mean and RMS.
- Step3 : Extend the window even further. Find entries that are less than  $(\text{Mean} + 3 \times \text{RMS})$  (from step2). THESE ENTRIES ARE IDENTIFIED PEDESTALS. Plot them and get the Mean.
- Step4 : Subtract this Mean from all ADC. The process is repeated for all events.

# Event by Event Correction Applied to Data



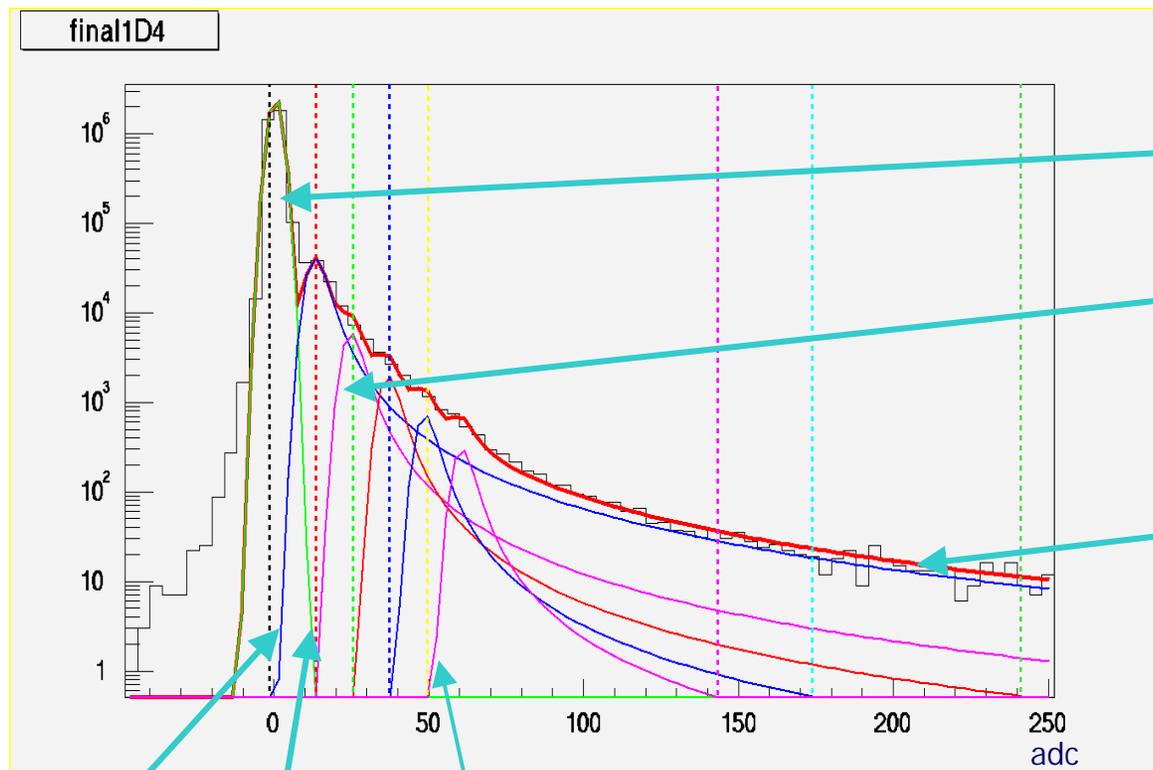
Before correction

ADC of channels 128-159  
from all Data Sample



After correction

# Fitting and Track Selection



Pedestals are Fitted With a Gaussian

Tracks are Fitted with a Landau convoluted with a gaussian

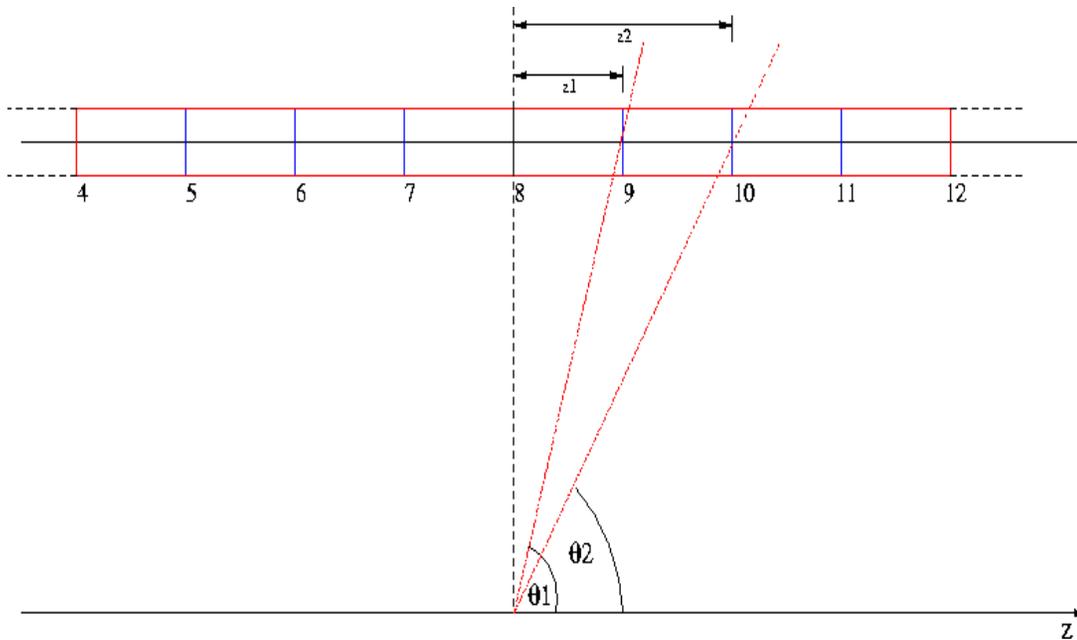
Total Distribution is Fitted with  
Total = Gaus + Gaus\*Landau

1-track  
2-tracks  
5-tracks

If  $p_0$  is the probability to be pedestal and  $p_i$  is the probability to be a  $i$ -track, then the number of track associated with an ADC is:

$$n = 0 * p_0 + 1 * p_1 + 2 * p_2 + 3 * p_3 + 4 * p_4 + 5 * p_5$$

# Strategy to Calculate $dN/d\eta$



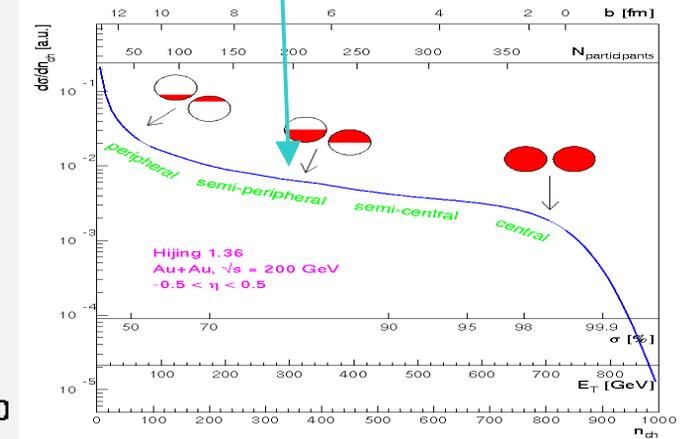
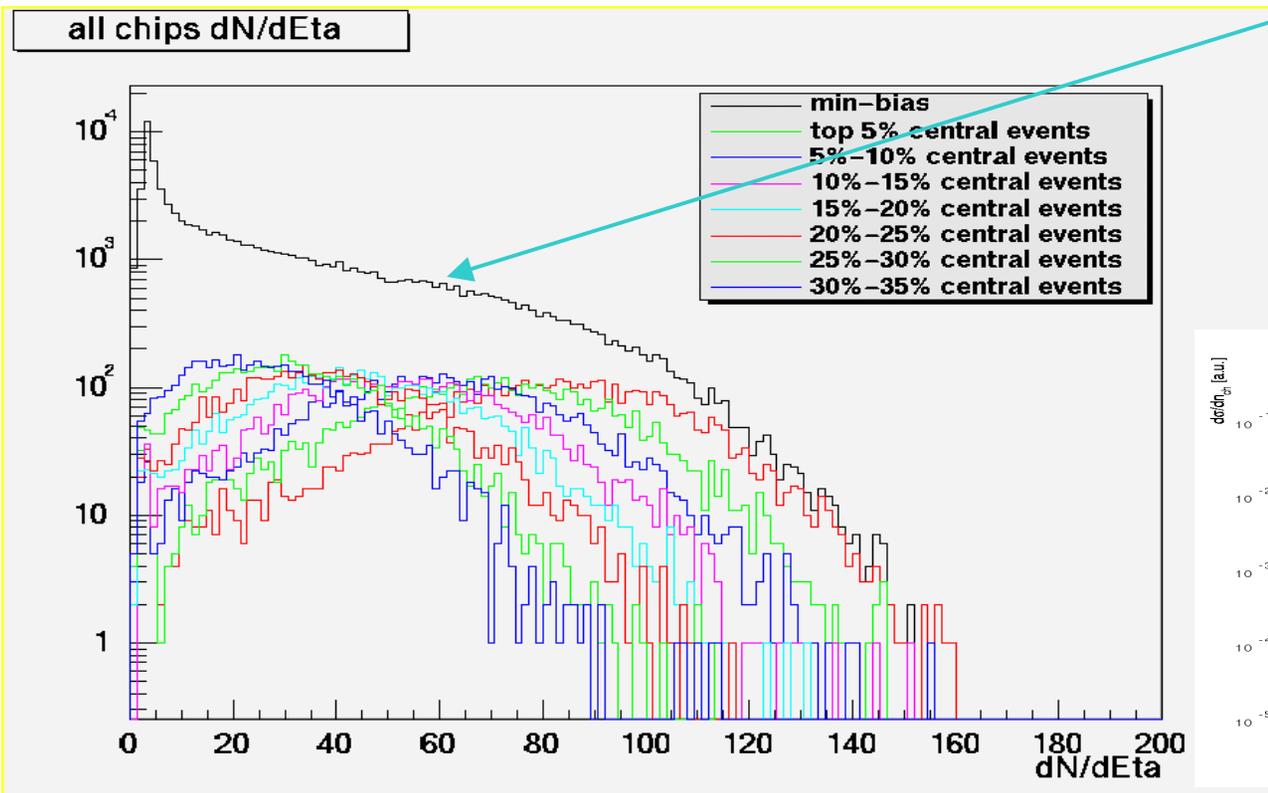
Measurement is done chip-by-chip (32 channels)

In any event

- Find the vertex position
- Find  $d\eta = \eta_{\max} - \eta_{\min}$  of the chip
- Find  $\eta$  at the center of the chip
- Find the total number of tracks in the chip in that event. This is  $dN$
- Find  $dN/d\eta$

# Charged Particle Multiplicity in MVD

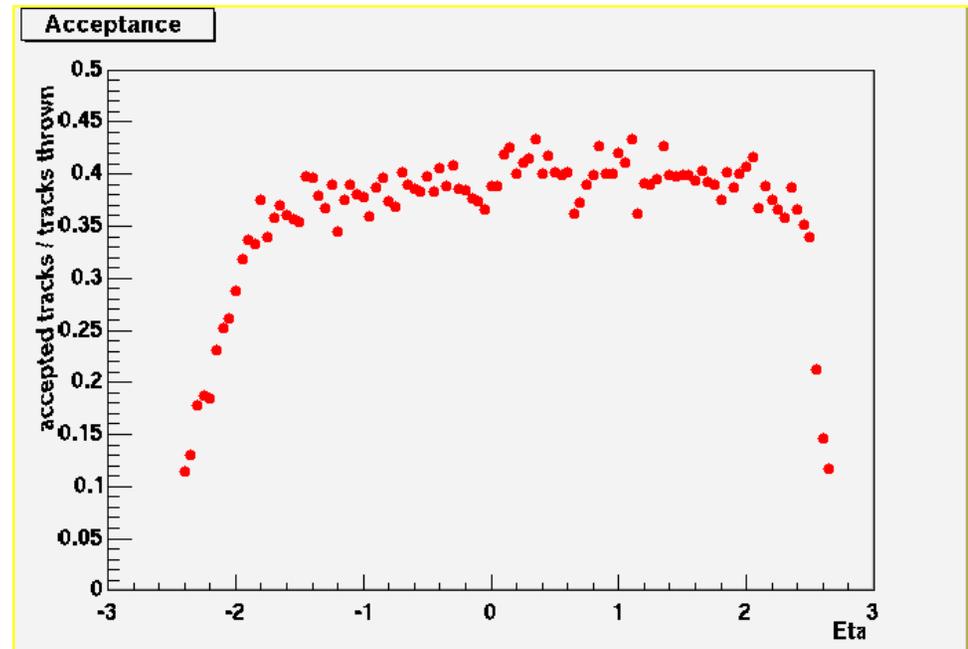
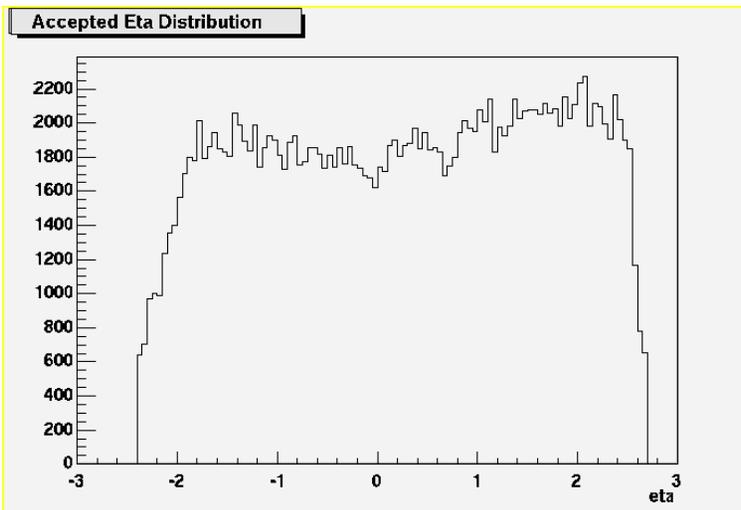
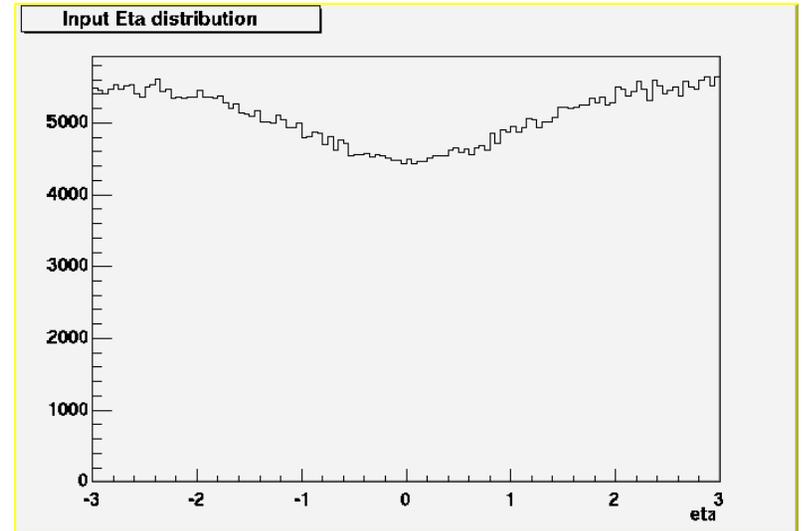
Min-bias Distribution  
Resembles the  
theoretical shape



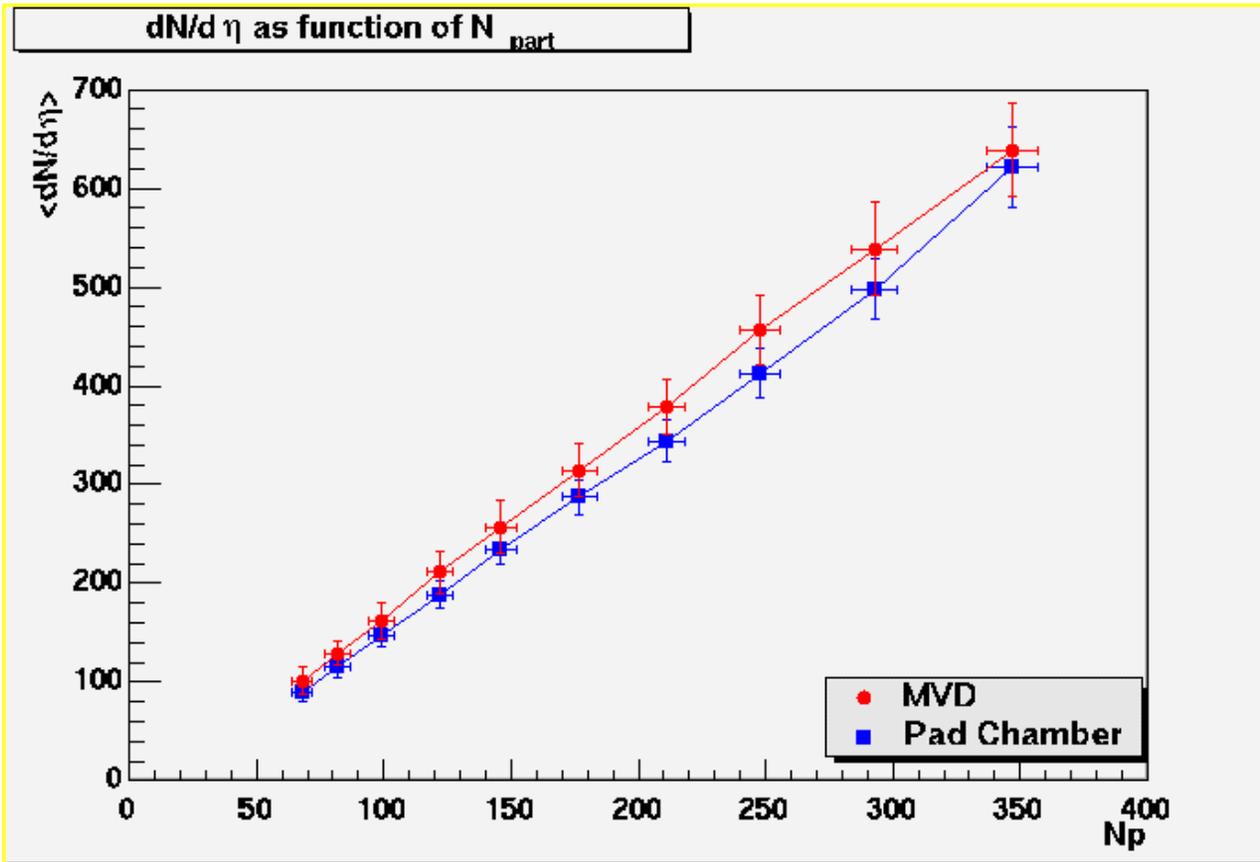
# Acceptance Correction

## Input Particles

- Flat in Rapidity  $|y| < 3$
- Random in azimuth  $0 < \phi < 2\pi$
- Exponential in  $m_t$ 
  - $(1/m_t) (dm_t/dy) = \exp(-(m_t - m_0)/T)$
  - $T = 220 \text{ MeV}$



# Charged Particle Multiplicity in MVD and Comparison with other Measurements in PHENIX

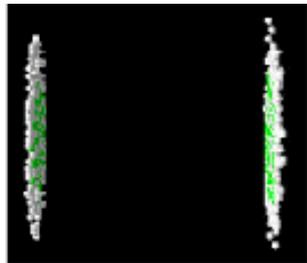


## Corrections:

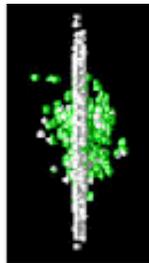
- AMU-cell dependence
- Event by event pedestal jumping
- Correction for multiple counting of tracks
- Acceptance Correction

dN/d $\eta$  as a function of number of participants in PHENIX MVD compared with the measurements from Pad Chambers

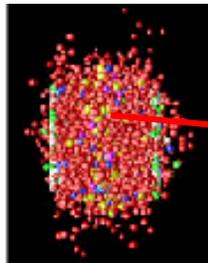
# Energy Density of the Collision



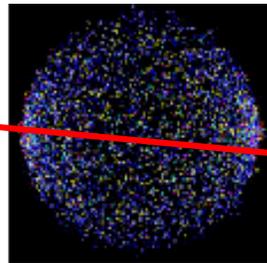
Two ions approaching



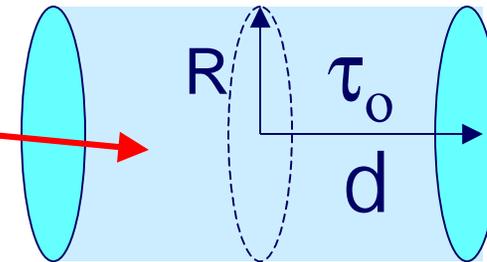
collision



Fireball



Particle Produced mainly in the transverse direction



Bjorken Estimate

Fireball expands longitudinally

After time  $\tau_0$ , fireball cools of and particles are produced

Particles are produced mainly in the transverse direction

Transverse Energy

$$E_T = \langle E \rangle N = \sqrt{m_T^2 + \langle p_T^2 \rangle} N$$

Total energy of the Fireball

$$E = \int_{-d}^d \frac{dE_T}{dz} dz = \frac{dE_T}{dy} \frac{2d}{\tau_0}$$

Volume of the Fireball

$$V = \pi R^2 (2d)$$

Energy Density

$$\varepsilon = \frac{E}{V} = \frac{1}{\pi R^2 \tau_0} \frac{dE_T}{dy} = \frac{\sqrt{m_T^2 + \langle p_T^2 \rangle}}{\pi R^2 \tau_0} \frac{dN}{dy}$$

Particles are mostly pions.

$$N \approx (N_{\pi^+} + N_{\pi^-} + N_{\pi^0}) \approx 3/2 (N_{\pi^+} + N_{\pi^-})$$

$$\varepsilon = \frac{E}{V} = \frac{1}{\pi R^2 \tau_0} \frac{dE_T}{dy} = \frac{\sqrt{m_T^2 + \langle p_T^2 \rangle}}{\pi R^2 \tau_0} \frac{3}{2} \frac{dN_{ch}}{dy}$$

$\langle p_T \rangle$  is measured at PHENIX for  $\sqrt{s} = 130$  GeV for hadrons.

$$dN/dy = dN/d\eta \cdot d\eta/dy = 1.13 dN/d\eta$$

$$R \sim A^{1/3} \sim 6 \text{ fm}$$

$\tau_0$  is the formation time  $\sim 1$  fm/c

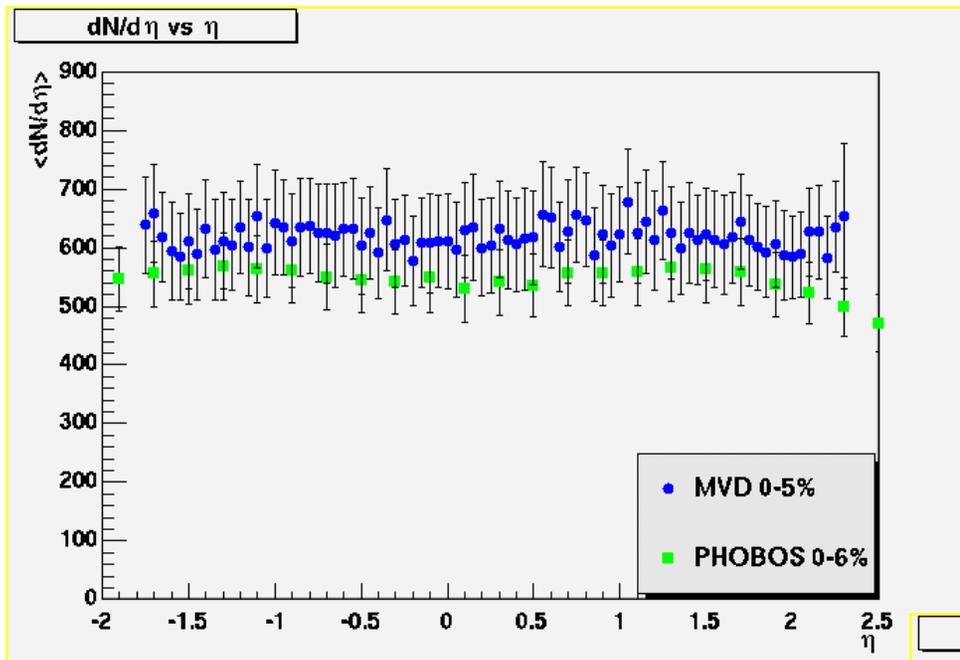
$dN/d\eta$  for the top 5% central events is 639

Which gives

$$\varepsilon_0 \tau_0 = 4.97 \text{ GeV}/(\text{fm}^2 \cdot \text{c})$$

The energy density is well above the energy density needed to induce a phase transition from lattice calculations

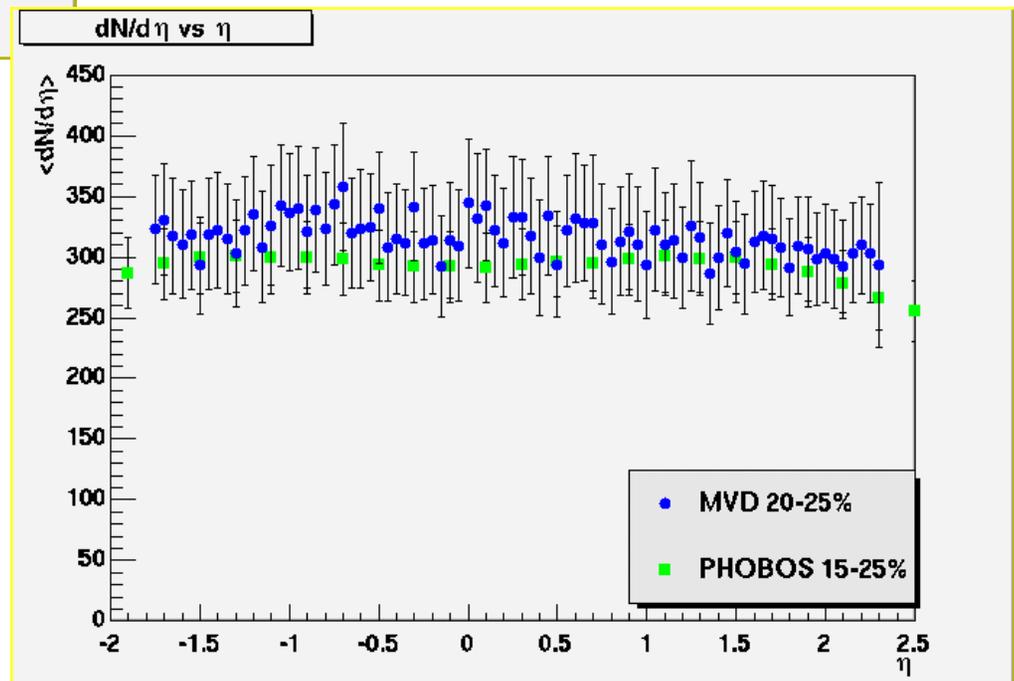
# Centrality Dependence of Multiplicity as a Function of $\eta$



$dN/d\eta$  as vs  $\eta$  plots for different centralities

As examples, plots for the most central (top 5%) events and mid-central (20%-25%) events are shown.

A comparison is made with the PHOBOS data.



# Theoretical Models

## HIJING Monte-Carlo Model (Gyulassy & Wang)

(HIJING = Heavy Ion Jet INTERaction Generator)

Particle produced in a heavy-ion collision has two components

- From soft interactions, which scales linearly with  $N_p$
- From hard process, which scales with  $N_c$ .

$$dN/d\eta = a\langle N_p \rangle + b\langle N_c \rangle$$

## Saturation Model (Kharzeev & Nardi)

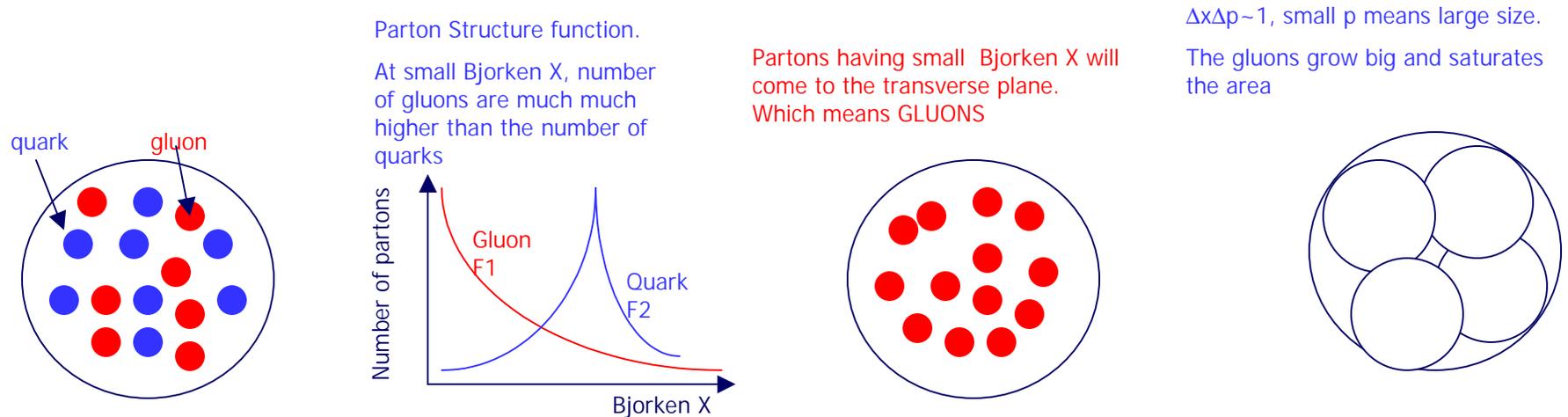
The initial parton density determines the number of produced charged particle. Naively it can be said that

$$dN/d\eta \propto \text{Number of Partons}$$

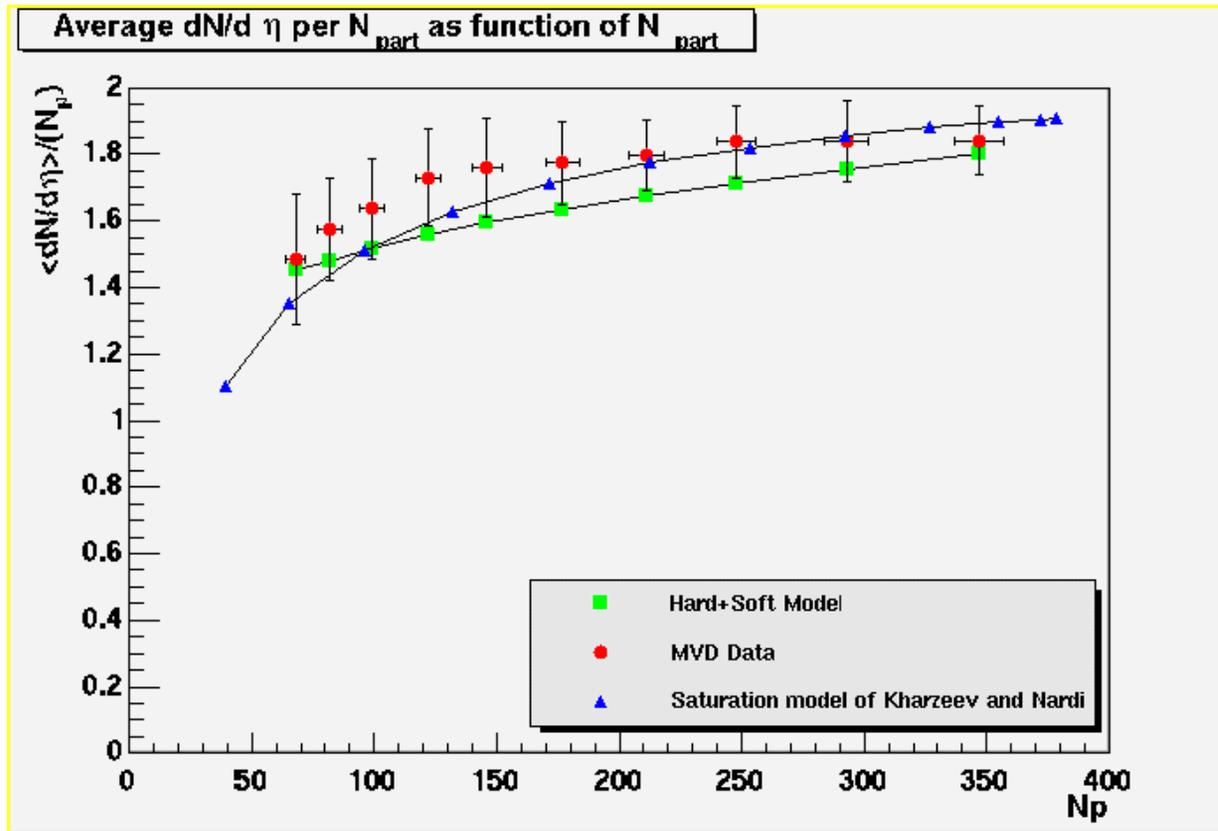
The Partons Density is Derived from the following scenario :

Nucleus is Lorentz contracted. Partons lie on a thin sheet in the transverse plane. Area of the transverse plane  $S_A \sim \pi R^2$ .

The Parton density keeps growing until all the available nuclear interaction transverse area is filled – the parton density saturates.



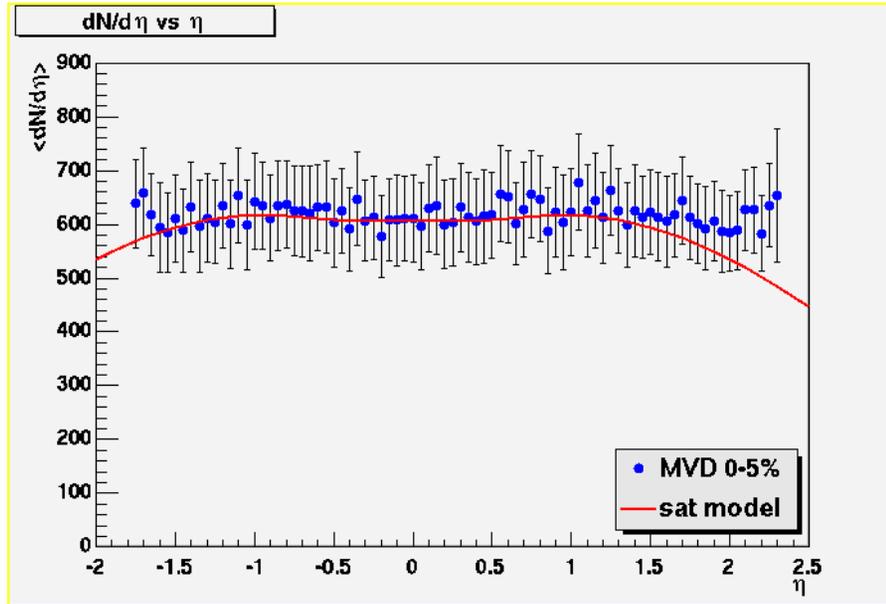
# Model Comparison of Data



Both models agrees with the data within the error bars.

The Saturation Model describes the shape of the distribution better

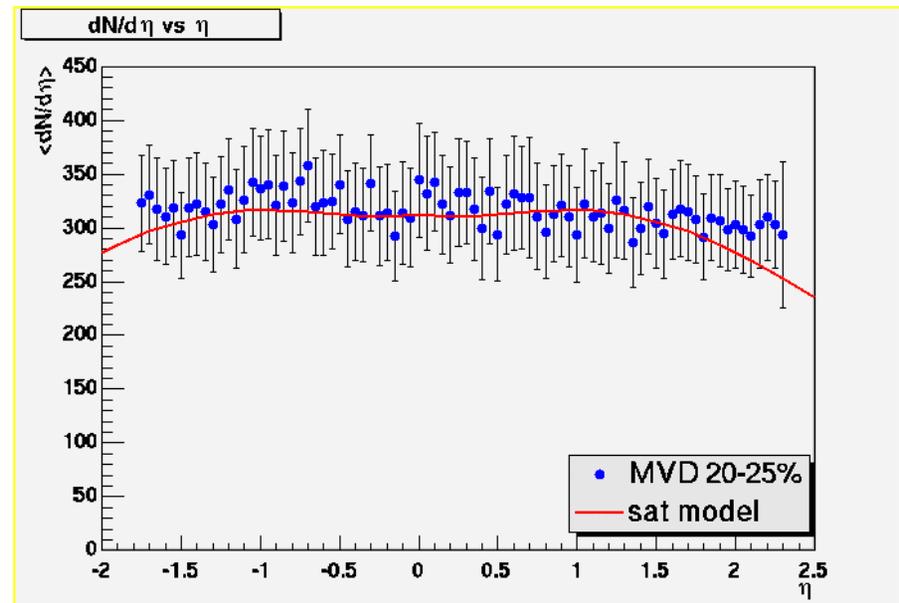
# Centrality Dependence of Multiplicity : Model Comparison



A comparison is made with the Saturation Model Calculation. The Saturation model describes the data well.

$dN/d\eta$  as vs  $\eta$  plots for different centralities

As examples, plots for the most central (top 5%) events and mid-central (20%-25%) events are shown.



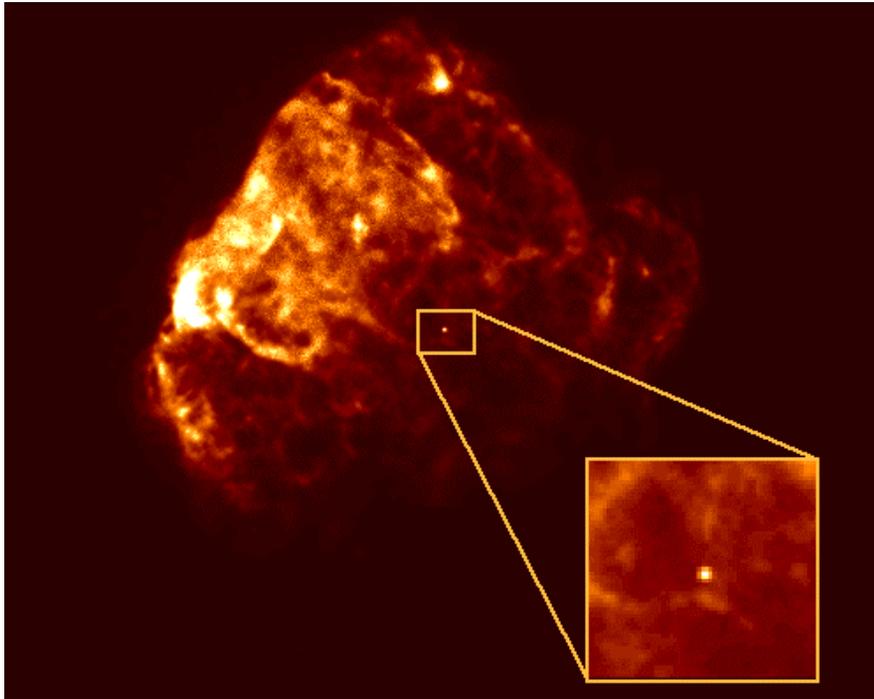
# Summary

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- The charged particle multiplicity for the Au-Au collisions at  $\sqrt{s}=130\text{GeV}$  is measured with the MVD at PHENIX at RHIC.
  - For the top 5% central events  $dN/d\eta = 639 \pm 48(\text{syst}) \pm 10(\text{stat})$
  - The energy density  $\varepsilon_0 \tau_0 = 4.97 \text{ GeV}/(\text{fm}^2\text{c})$
- The charged particle multiplicity measured with MVD agrees well with the PHENIX Pad Chamber.
  - The  $dN/d\eta$  as a function of  $\eta$  agrees well with the PHOBOS measurements.
- The centrality dependence of the charged particle multiplicity is compared with the HIJING soft+hard model and the Saturation model. Both the models fits the data within the error bar. Saturation model describes the shape of the distribution better.

# Back-up Slides

# QGP in the Universe

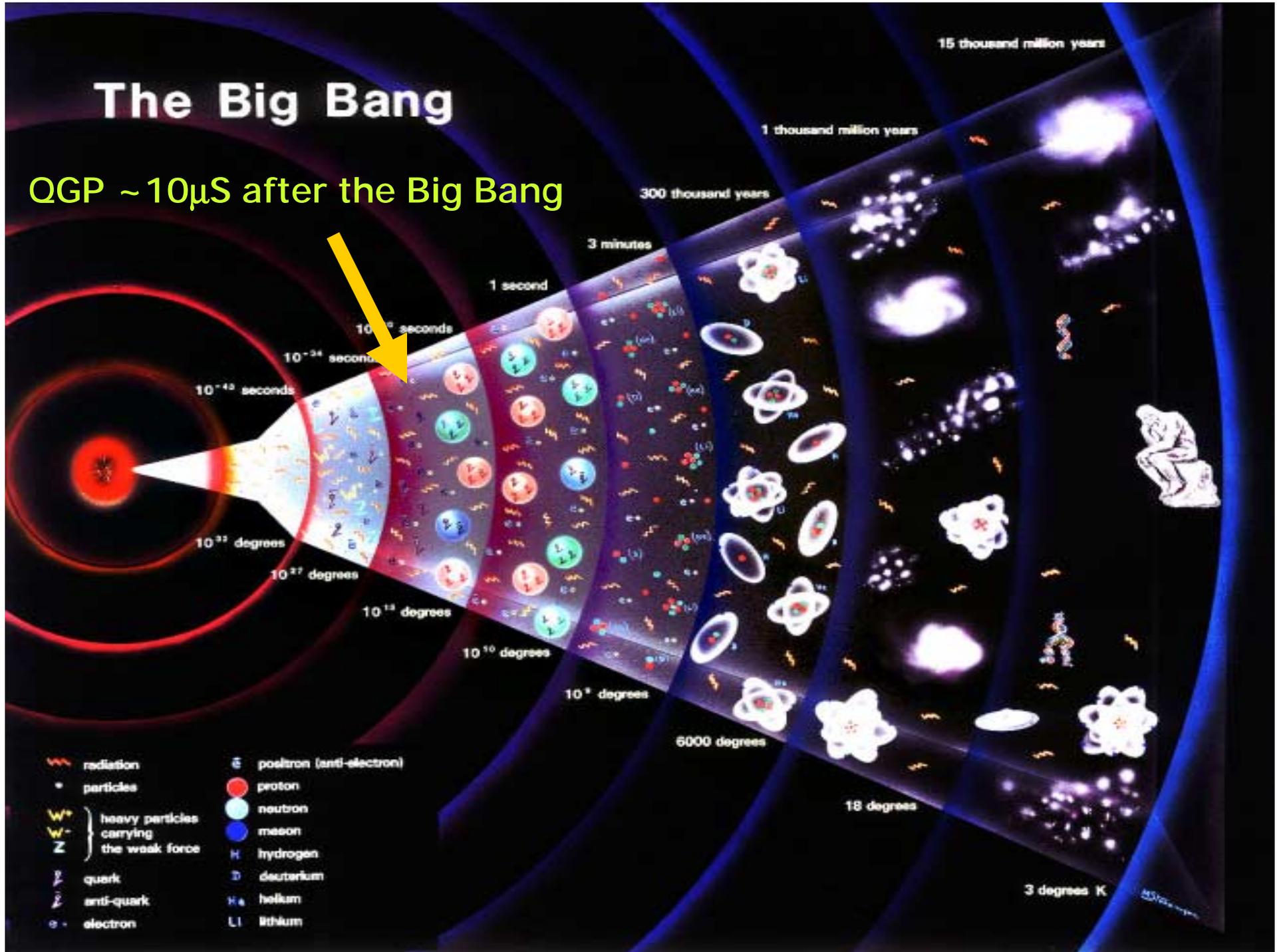


**Neutron stars are the collapsed cores of a massive star.**

**They pack the mass of the sun into the size of a city.**

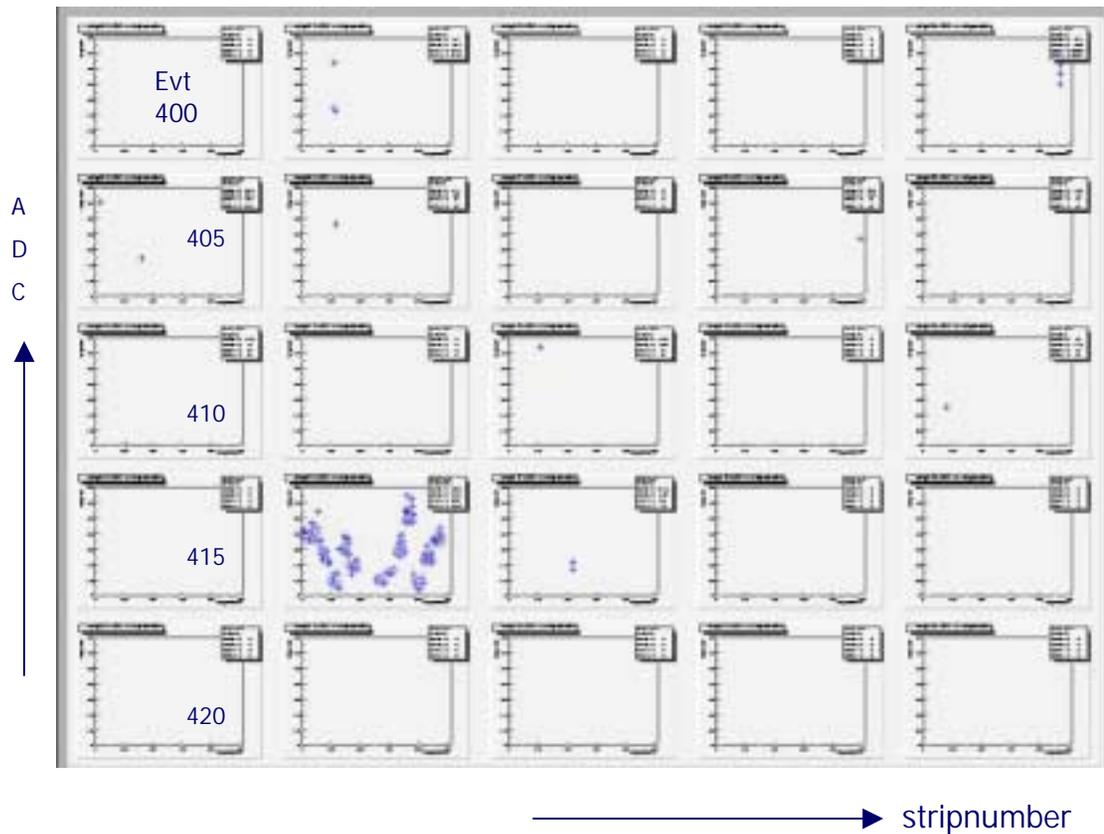
# The Big Bang

QGP  $\sim 10\mu\text{s}$  after the Big Bang



# Event by Event Pedestal Jumping

- Pedestals “jump” in event to event!!
- Jumping could be such that the pedestals can occupy the ADC range that could be that of a track (meaning they can screen the tracks)
- WE NEED TO IDENTIFY THE PEDESTALS BEFORE WE DO ANYTHING.

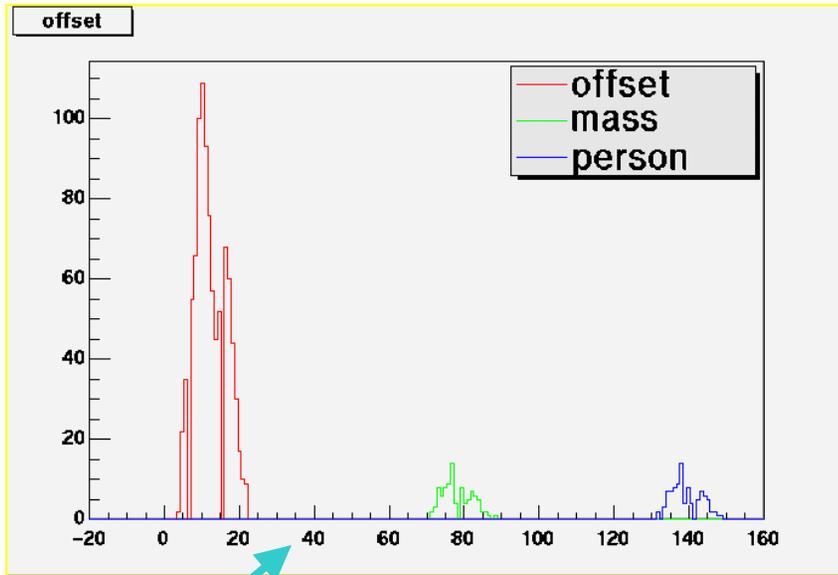


2D plot of ADC in y-axis and stripnumber in x axis.  
The plots are for events 400-424.  
Look at event 416. All the strips have high ADC.

# Event by Event Correction Procedure : Does it Work?

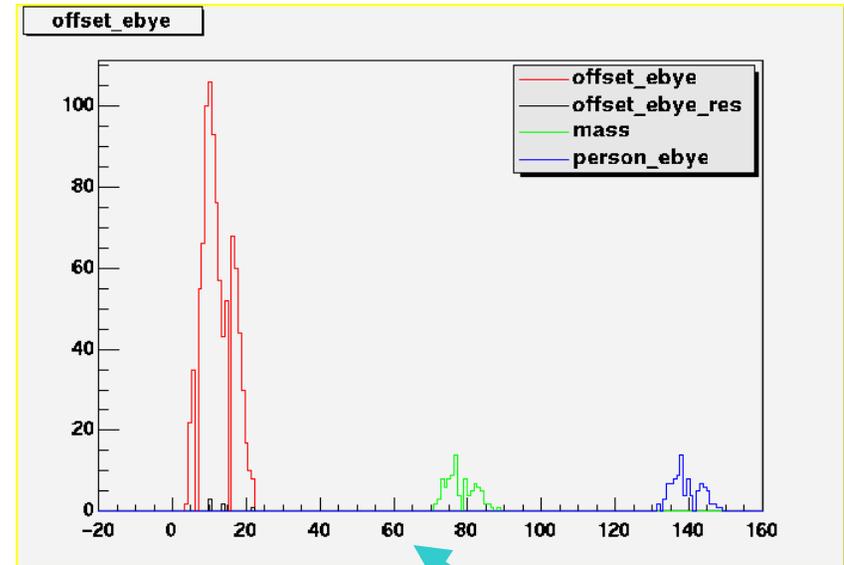
- We do a simple test
  - measure two quantities on a scale that 's offset changes in each measurement.
  - Think of an "event" where we change the offset and record offset 10 times, and measure a block and a persons weight once. So each event has 12 entries.
  - Lets do about 100 events.  
Pass the events through the event by event correction procedure.
  - If the procedure works, we will see very narrow distributions .

# Test Result



Raw

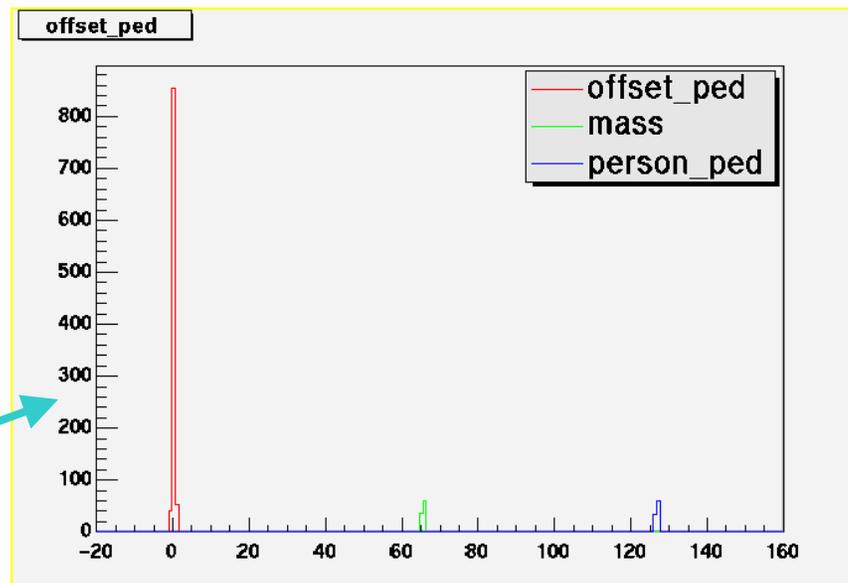
All are 1D plots of weights in x-axis



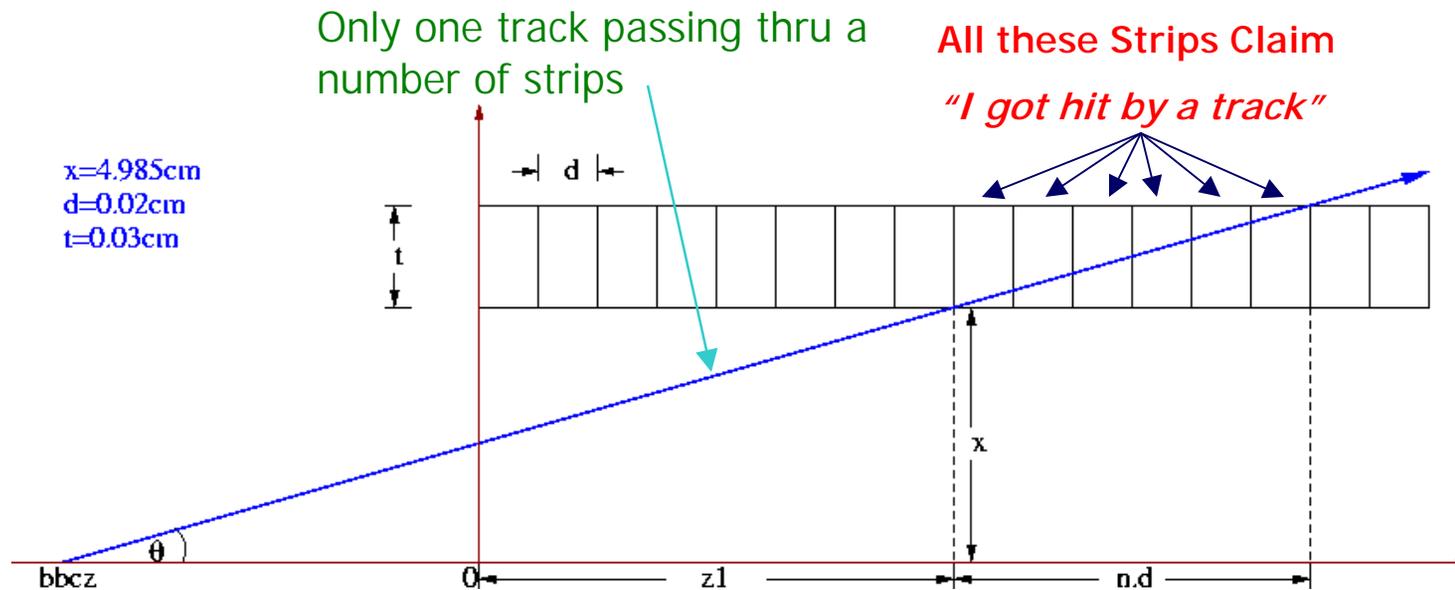
After ebye  
pedestal  
identification

**WORKS!**

After event-by-  
event pedestal  
subtraction



# Corrections for Multiple Counting of Tracks



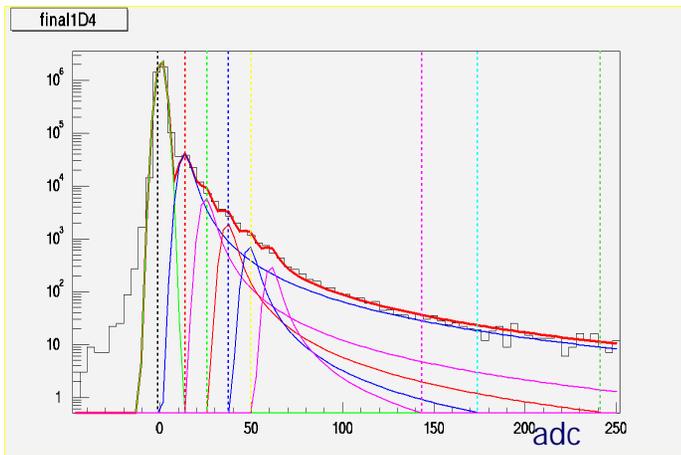
Pathlength inside MVD is  $\sqrt{t^2+n^2d^2}$

A pathlength of  $t$  corresponds to an ADC for a MIP

Pathlength  $\sqrt{t^2+n^2d^2}$  corresponds to  $\sqrt{t^2+n^2d^2}/t$  tracks

So the correction factor is  $t/\sqrt{t^2+n^2d^2} = \text{Sin}\theta$

# Systematic Error from Fitting



If  $p_0$  is the probability to be pedestal and  $p_i$  is the probability to be a  $i$ -track, then the number of track associated with an ADC is:

$$n = 1 \cdot p_1 + 2 \cdot p_2 + 3 \cdot p_3 + 4 \cdot p_4 + 5 \cdot p_5$$

Imagine that an ADC is really due to 2-tracks. Using above relation, we are undercounting the track by

$$e_2 = -1 \cdot p_1 + 0 \cdot p_2 + 1 \cdot p_3 + 2 \cdot p_4 + 3 \cdot p_5$$

We counted the number of 1-tracks, 2-tracks etc.

So, if there are  $N_2$  2-tracks in a total of  $N$  tracks, the error for such tracks are

$$\text{err}_2 = (-1 \cdot p_1 + 0 \cdot p_2 + 1 \cdot p_3 + 2 \cdot p_4 + 3 \cdot p_5) \cdot p_2 \cdot 100\%$$

Where  $p_2 = N_2/N$