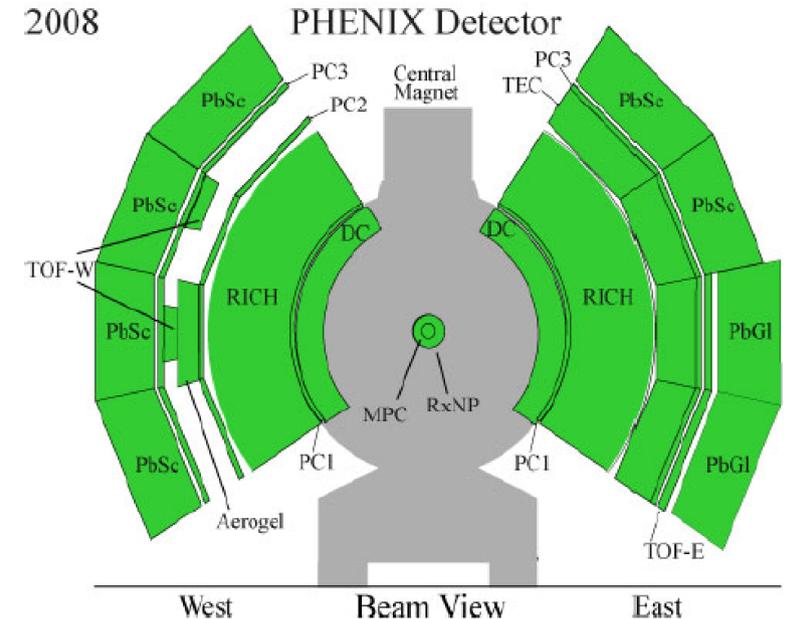
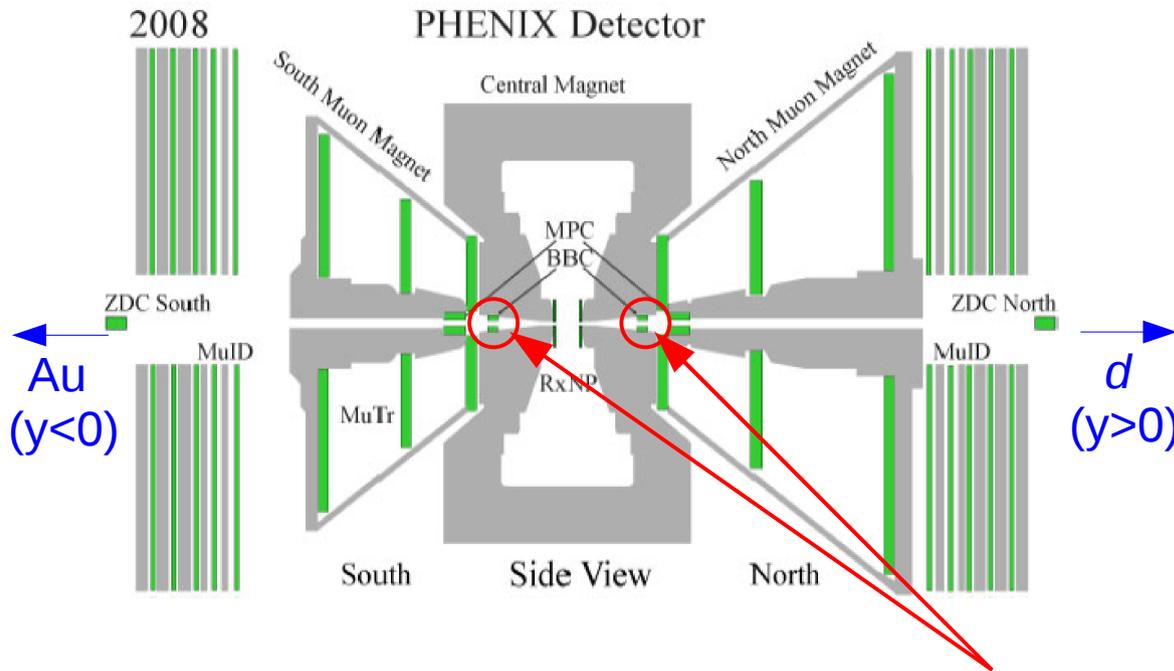


Cold Nuclear Matter Effects on J/ψ Production from PHENIX

Darren McGlinchey
Florida State University
For the PHENIX Collaboration
WWND 2012



2008 Run



Muon Arms

- Muons
- $1.2 < |\eta| < 2.4$
- $\Delta\phi = 2\pi$
- $J/\psi \rightarrow \mu^+\mu^-$

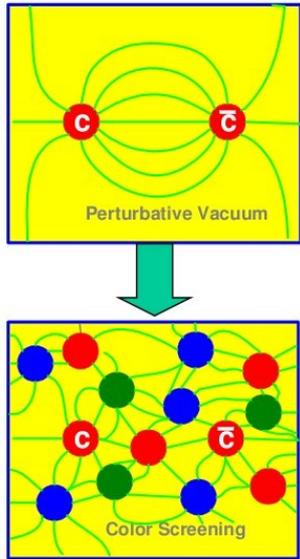
Beam-Beam Counters

- Measure Centrality (impact parameter) as a percentage of BBC charge

Central Arms

- Charged particles
- $|\eta| < 0.35$
- $\Delta\phi = \pi$
- $J/\psi \rightarrow e^+e^-$

Heavy Quarkonia



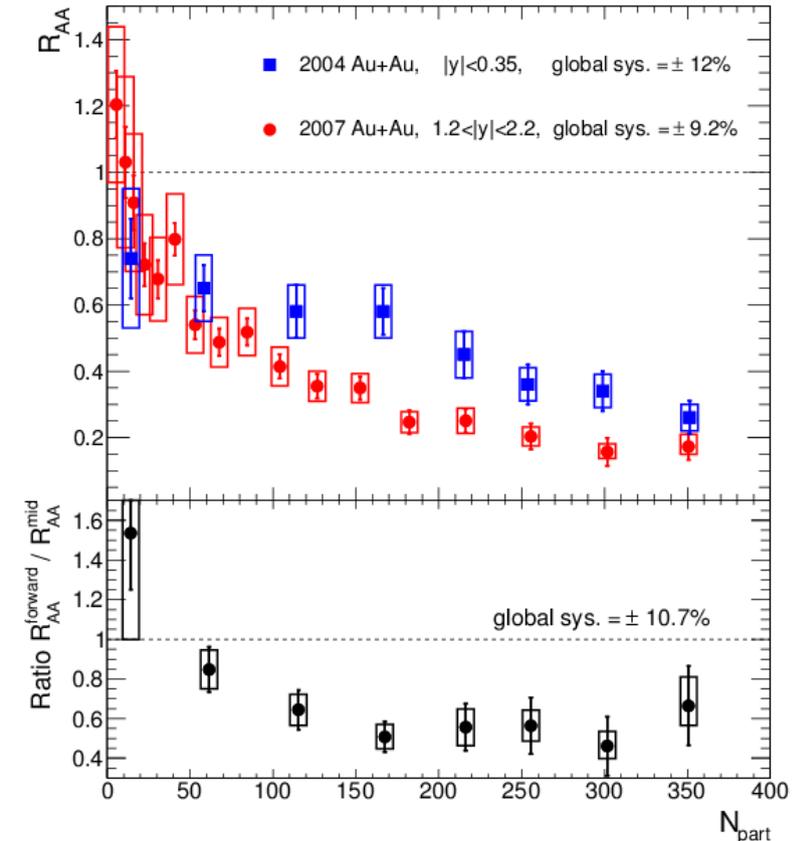
Goal: Measure the screening length in the QGP

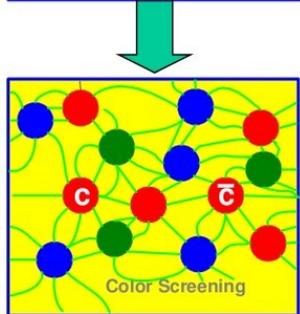
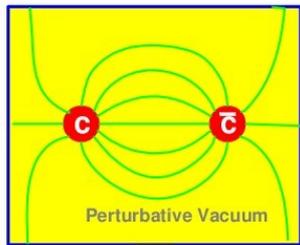
Use heavy quarkonia as a tool for measuring this directly.

Understanding A+A results alone is not trivial.

$$R_{AA} = \frac{dN^{AA}/dy}{dN^{pp}/dy N_{Coll}}$$

Phys. Rev. C 84, 054912 (2011)





Goal: Measure the screening length in the QGP

Use heavy quarkonia as a tool for measuring this directly.

Understanding A+A results alone is not trivial.

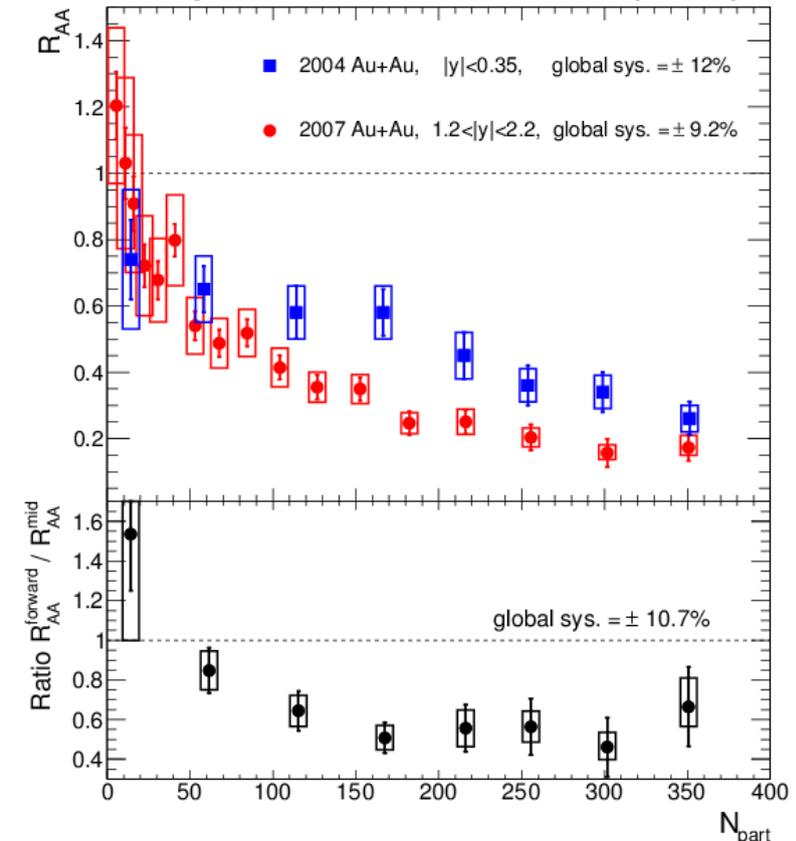
Need to understand our baseline in order to extract hot nuclear matter effects!

Need to not only measure/understand p+p production.

Need to measure/understand Cold Nuclear Matter Effects (effect of producing a J/ψ in a nuclear target)

$$R_{AA} = \frac{dN^{AA}/dy}{dN^{pp}/dy N_{Coll}}$$

Phys. Rev. C 84, 054912 (2011)



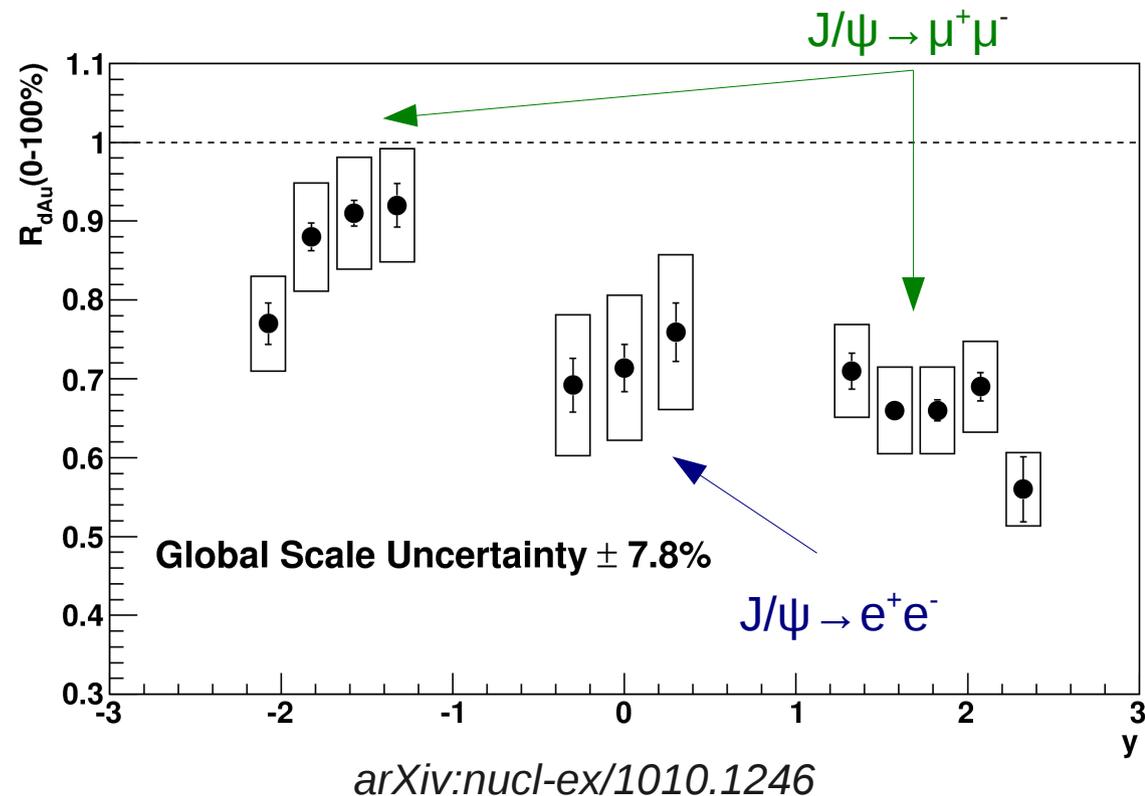
- Large statistics d+Au data taken @ 200 GeV in 2008.
- Allows direct study of Cold Nuclear Matter (CNM) effects.
- Goal: Measure J/ψ production over wide range of kinematics.
 - Rapidity dependence (*Phys. Rev. Lett.* 107, 142301 (2011)).
 - **New!** p_T dependence (*arXiv:1204.0777*).

- Minimum Bias (centrality integrated) R_{dAu} results as a function of rapidity.
- Shows increasing suppression with increasing rapidity.

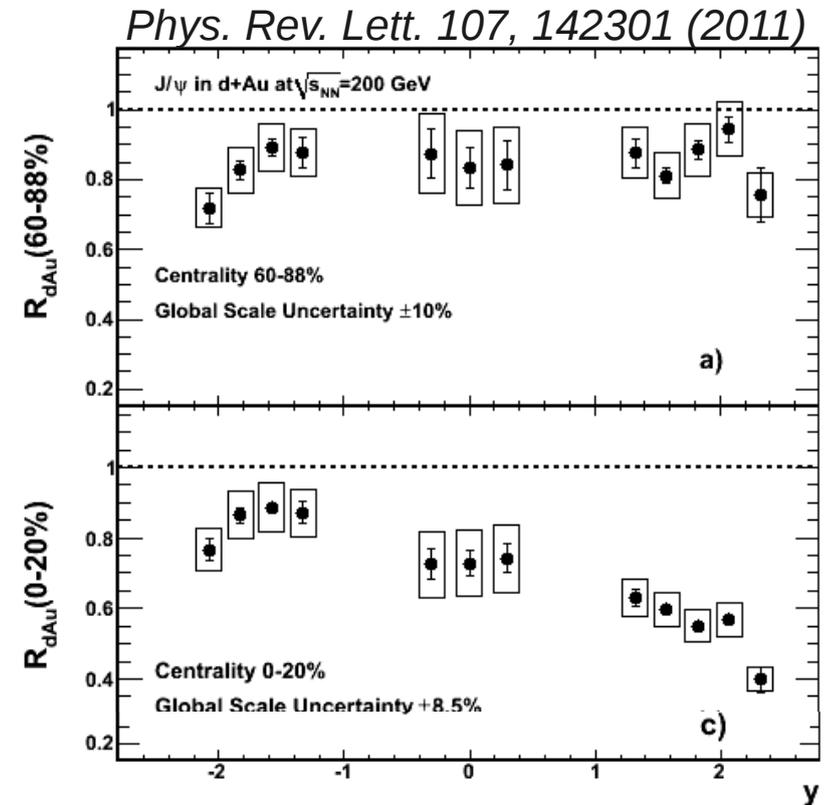
$$R_{dAu}(i) = \frac{dN_{J/\psi}^{dAu}/dy(i)}{\langle N_{coll}(i) \rangle dN_{J/\psi}^{pp}/dy}$$

Vertical Error bars – point-to-point uncorrelated errors

Boxes - point-to-point correlated errors

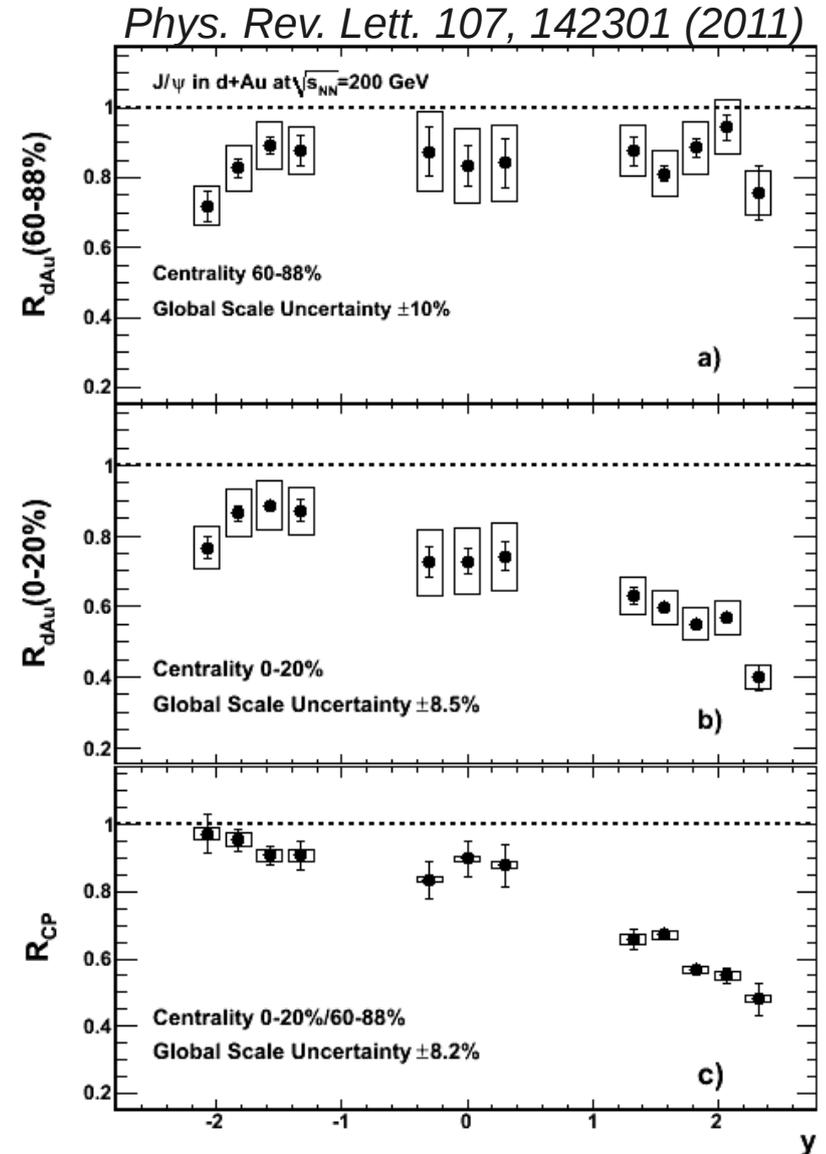


- **Top** – R_{dAu} vs y for Peripheral events
 - R_{dAu} constant w/ rapidity.
- **Bottom** – R_{dAu} vs y for Central events
 - Stronger suppression at forward rapidity.
 - Similar level of suppression at backward rapidity in central and peripheral events.



- Take the ratio of central R_{dAu} (b) to peripheral R_{dAu} (a) $\rightarrow R_{CP}$
 - Significant reduction of systematic errors.
 - Describes change in suppression over the nucleus.
 - Strong suppression at forward rapidity.

What can we learn about the geometric dependence of the modification?



- Use simple geometric model.
- Assume modification is dependent on the nuclear thickness.

$$\Lambda(r_T) = \frac{1}{\rho_0} \int dz \rho(z, r_T)$$

- Try three simple forms.
 - *Linear:* $M(r_T, a) = 1 - a\Lambda(r_T)$
 - *Quadratic:* $M(r_T, a) = 1 - a\Lambda(r_T)^2$
 - *Exponential:* $M(r_T, a) = e^{-a\Lambda(r_T)}$

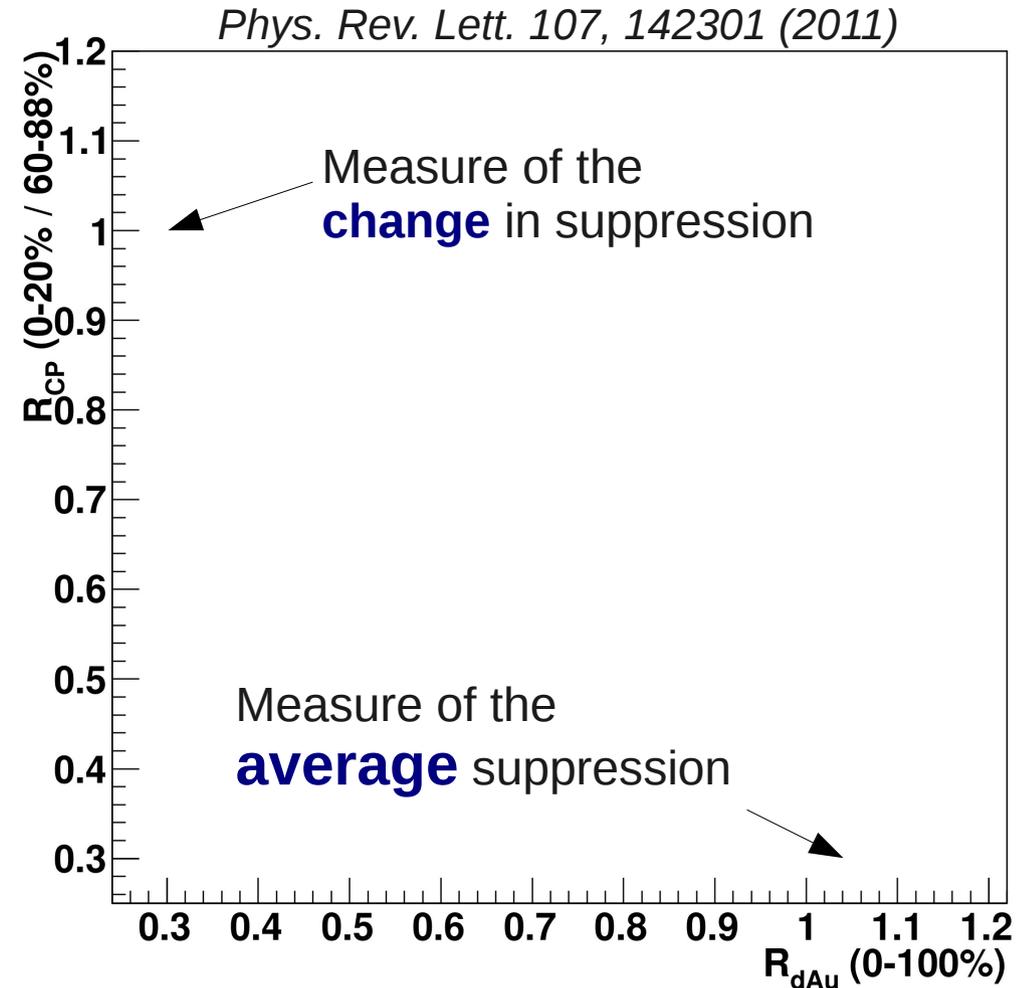
A Different Way to Look at the Centrality Dependence

- Use simple geometric model.
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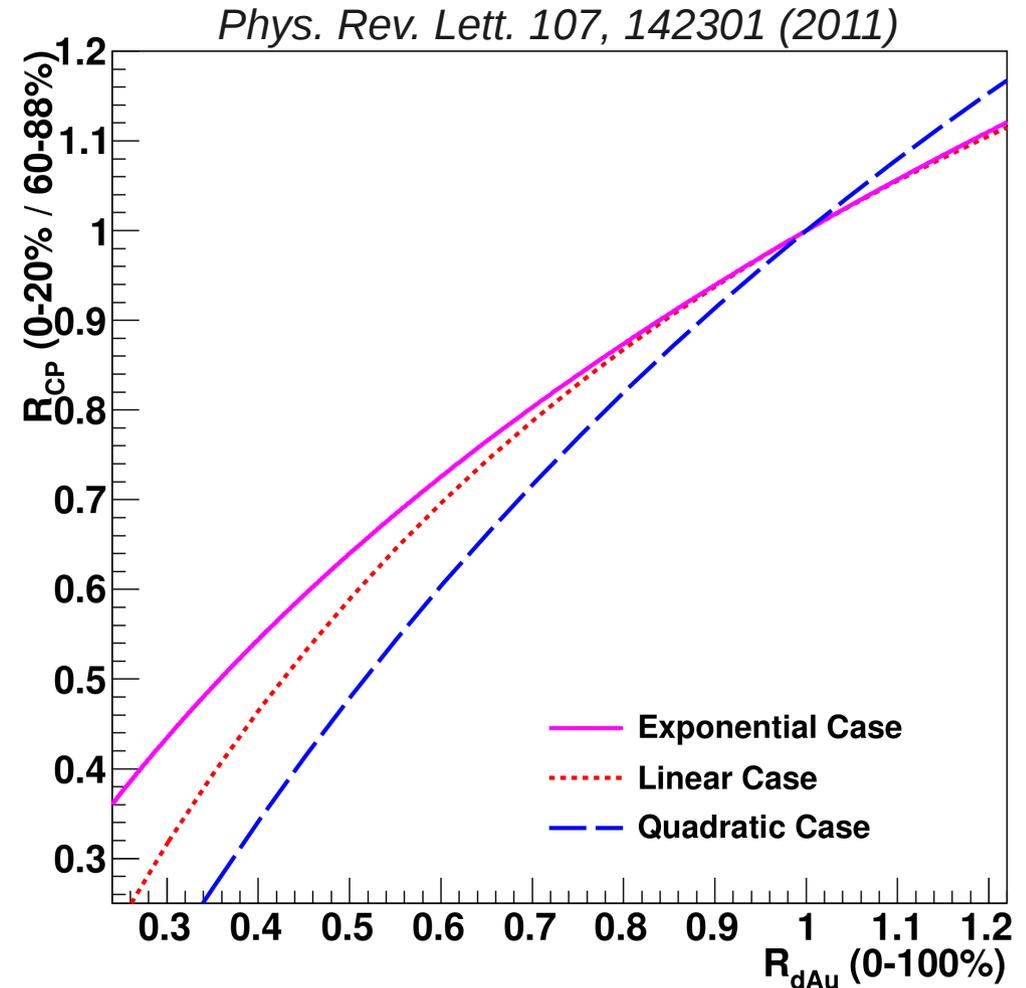


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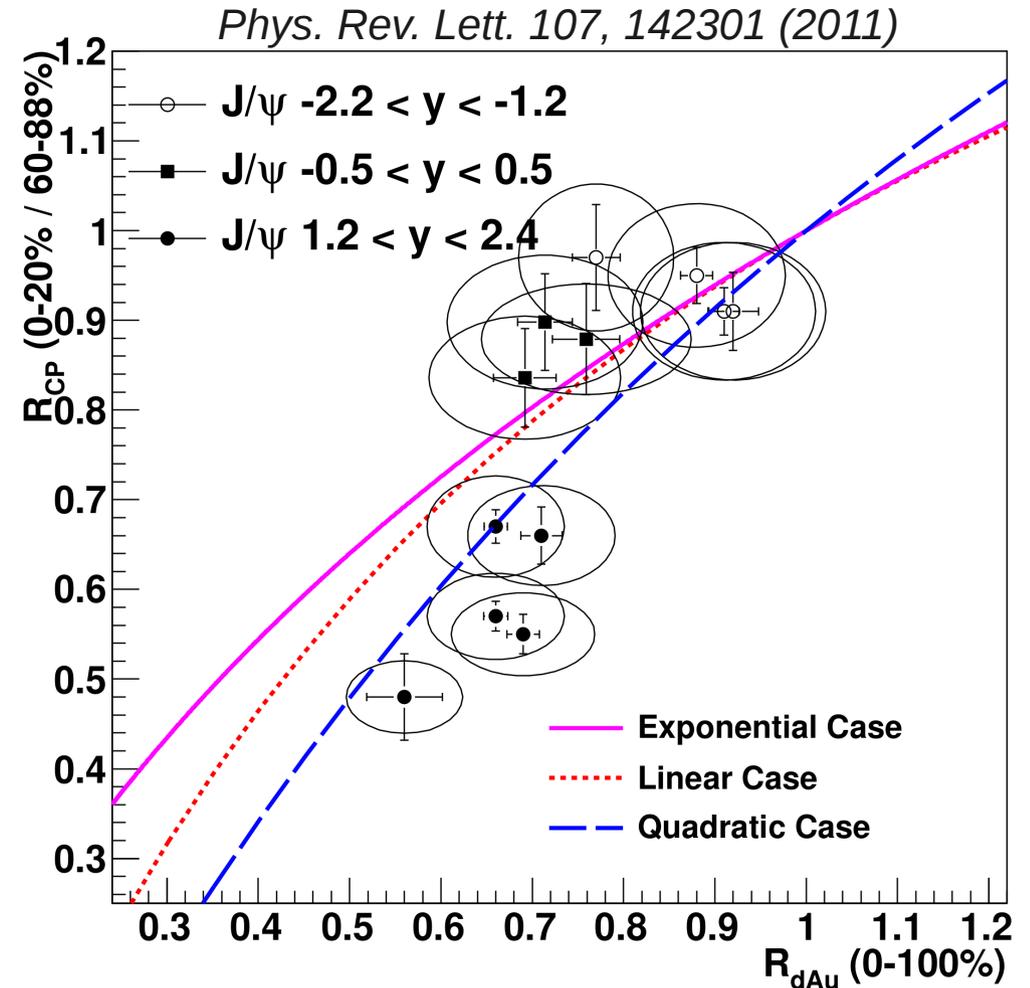
- Try three simple forms.
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- Each form maps out a unique curve on this plane.



A Different Way to Look at the Centrality Dependence

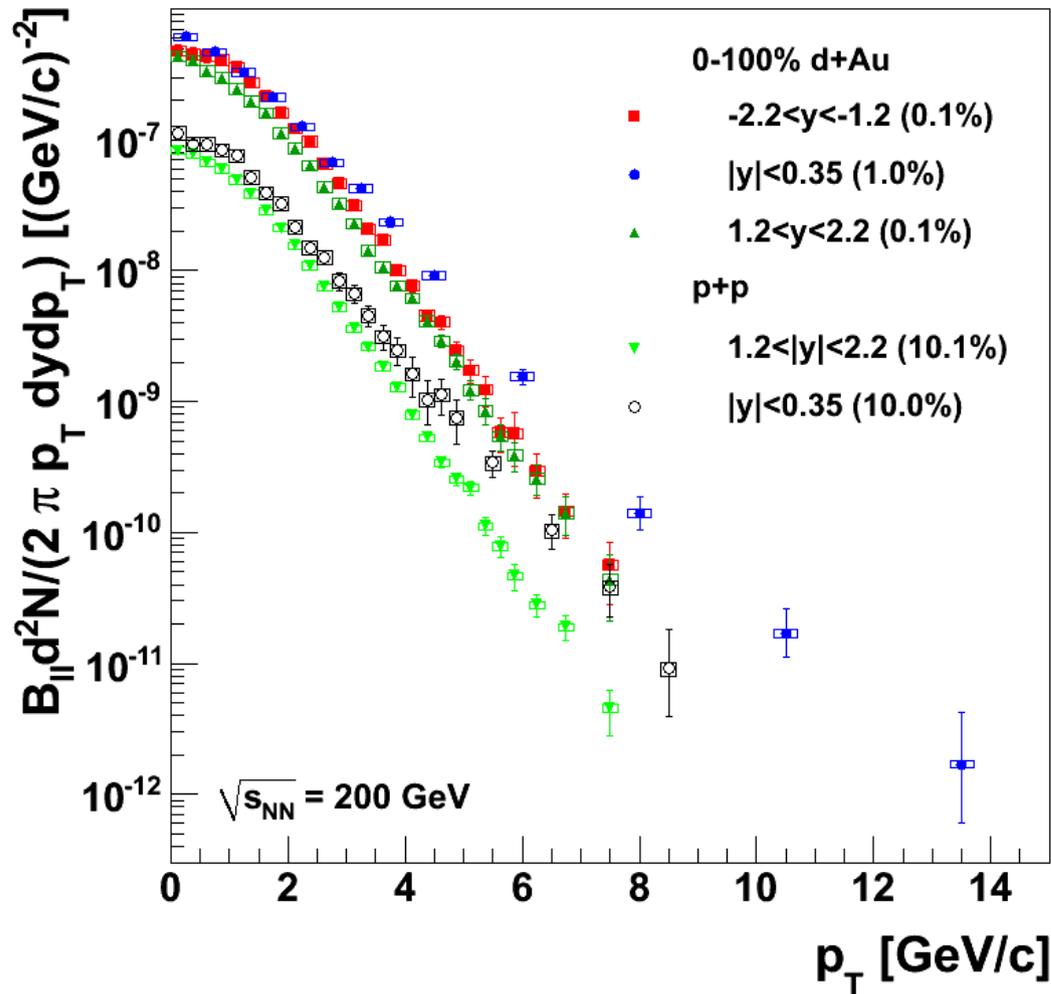
- Backward and mid rapidity data is unable to distinguish between the three cases shown here.
- Forward rapidity data requires stronger than linear or exponential dependence on the thickness.
- Use data to extract thickness dependence!

Vertical(horizontal) bars show point-to-point uncorrelated errors
Ellipses show point-to-point correlated systematic errors



J/ψ p_T Distributions

arXiv:1204.0777

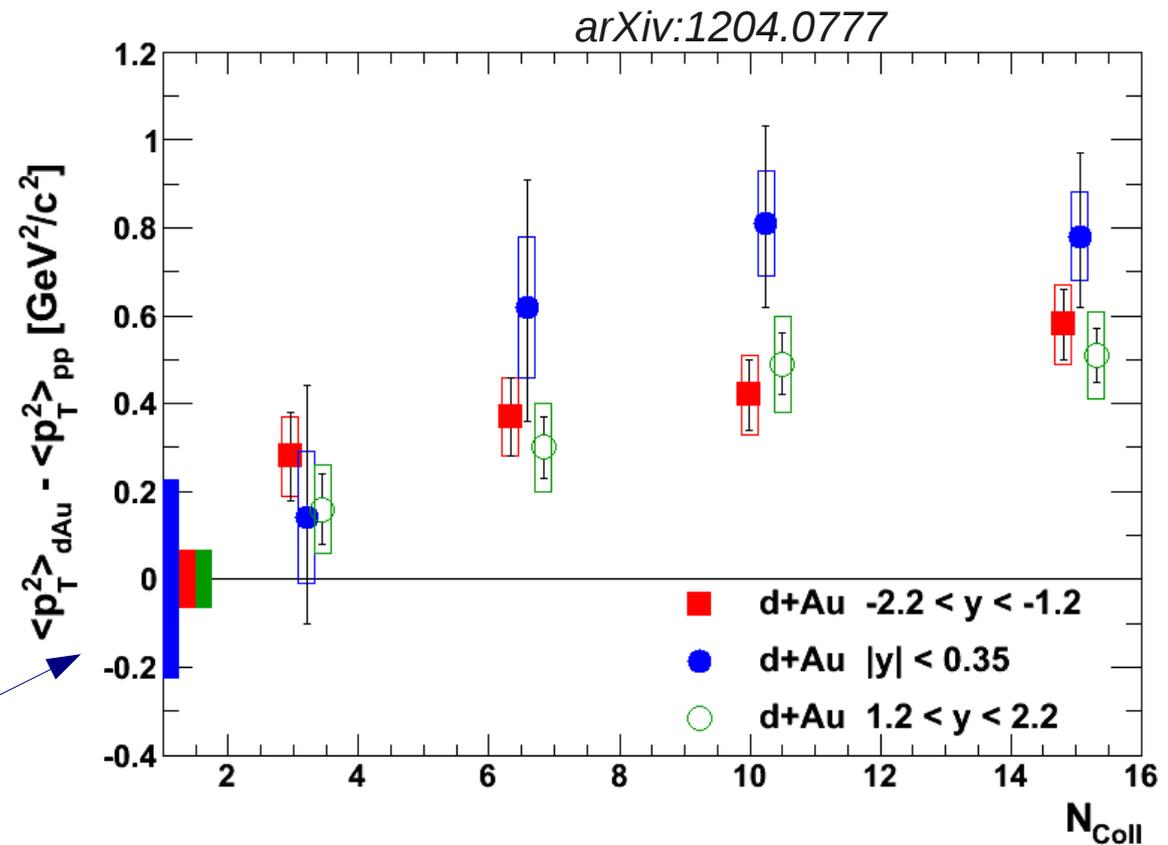


- Minimum bias J/ψ invariant yield vs p_T.
- Integrated over the rapidity of each arm.
- Precise data out to p_T = 8 GeV/c.
- Large improvement over previous d+Au PHENIX results.

$\langle p_T^2 \rangle$ in d+Au

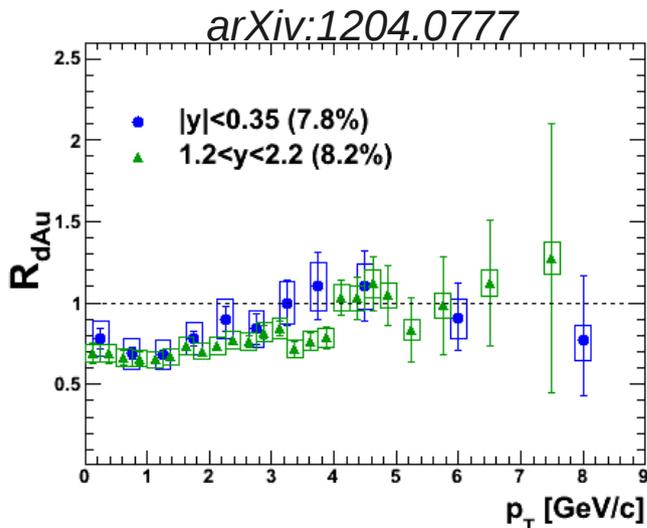
- Extract $\langle p_T^2 \rangle$ numerically from the data at each centrality.
- Subtract the p+p result from the d+Au results.
- Observe broadening in the p_T distribution which increases with increasing N_{coll} .
- Similar at all rapidities within uncertainties.

Colored boxes @ $\Delta\langle p_T^2 \rangle = 0$
 represent uncertainty on
 $\langle p_T^2 \rangle_{pp}$



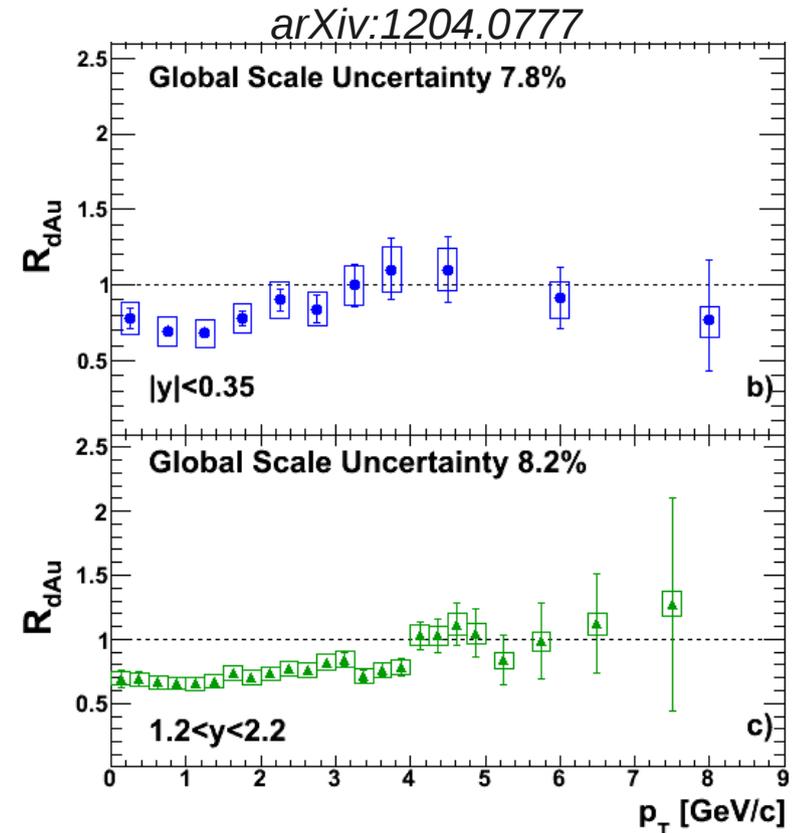
- Minimum bias R_{dAu} vs p_T .
- Similar suppression at mid & forward rapidity.
- Suppression for $p_T < 4$ GeV/c.
- $R_{dAu} \approx 1$ for $p_T > 4$ GeV/c.

$$R_{dAu}(i) = \frac{dN_{J/\psi}^{dAu} / dy(i)}{\langle N_{coll}(i) \rangle dN_{J/\psi}^{pp} / dy}$$

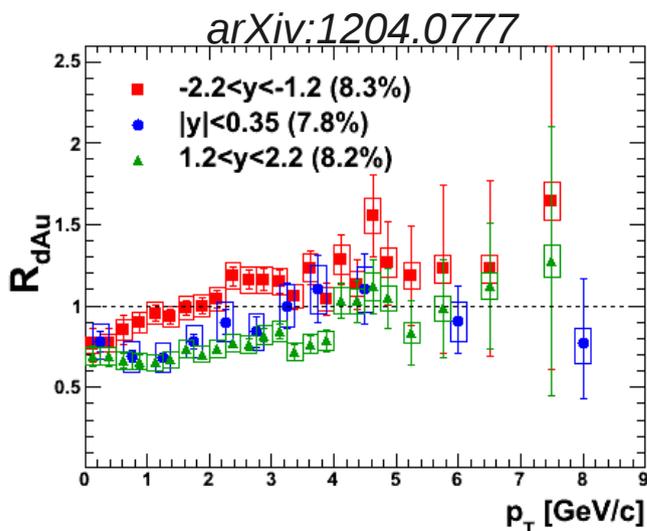


Vertical Error bars
– uncorrelated
uncertainties

Boxes – point-to-
point correlated
uncertainties.

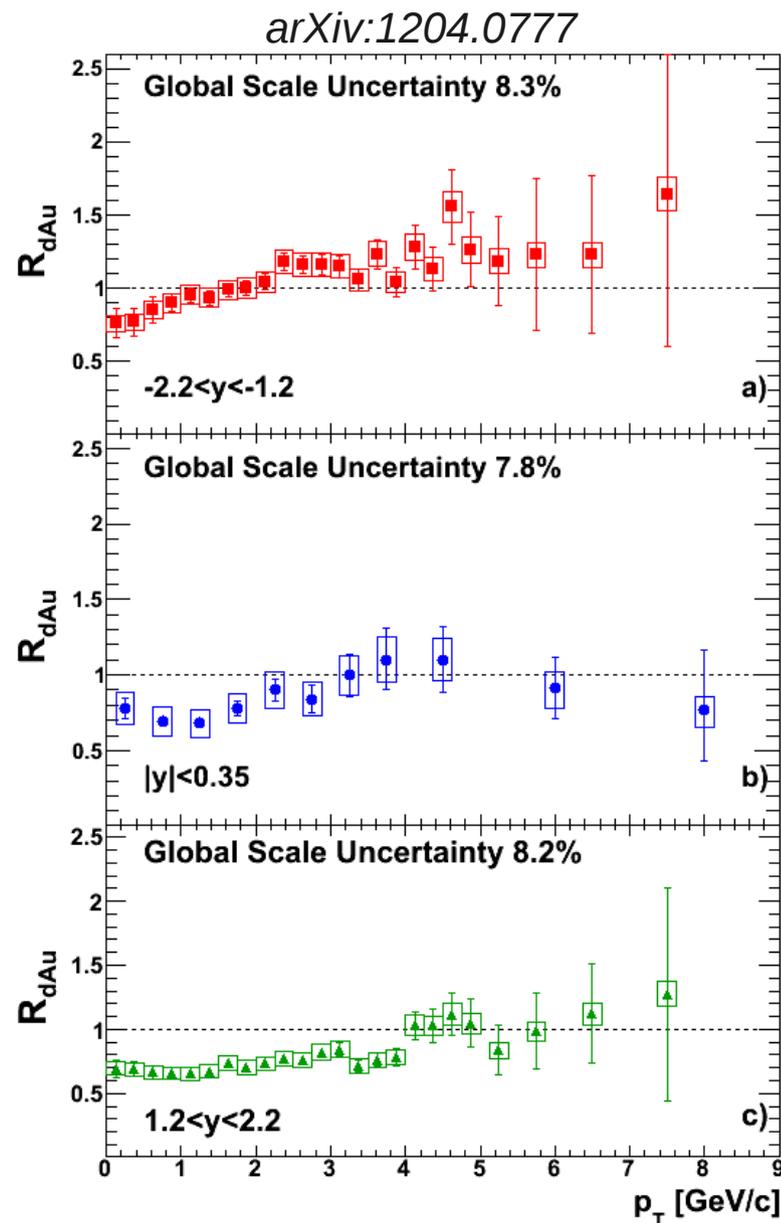


- Shape of R_{dAu} p_T distribution different at backward rapidity.
- $R_{dAu} > 1$ for $p_T > 2$ GeV/c.



Vertical Error bars
– uncorrelated
uncertainties

Boxes – point-to-
point correlated
uncertainties.

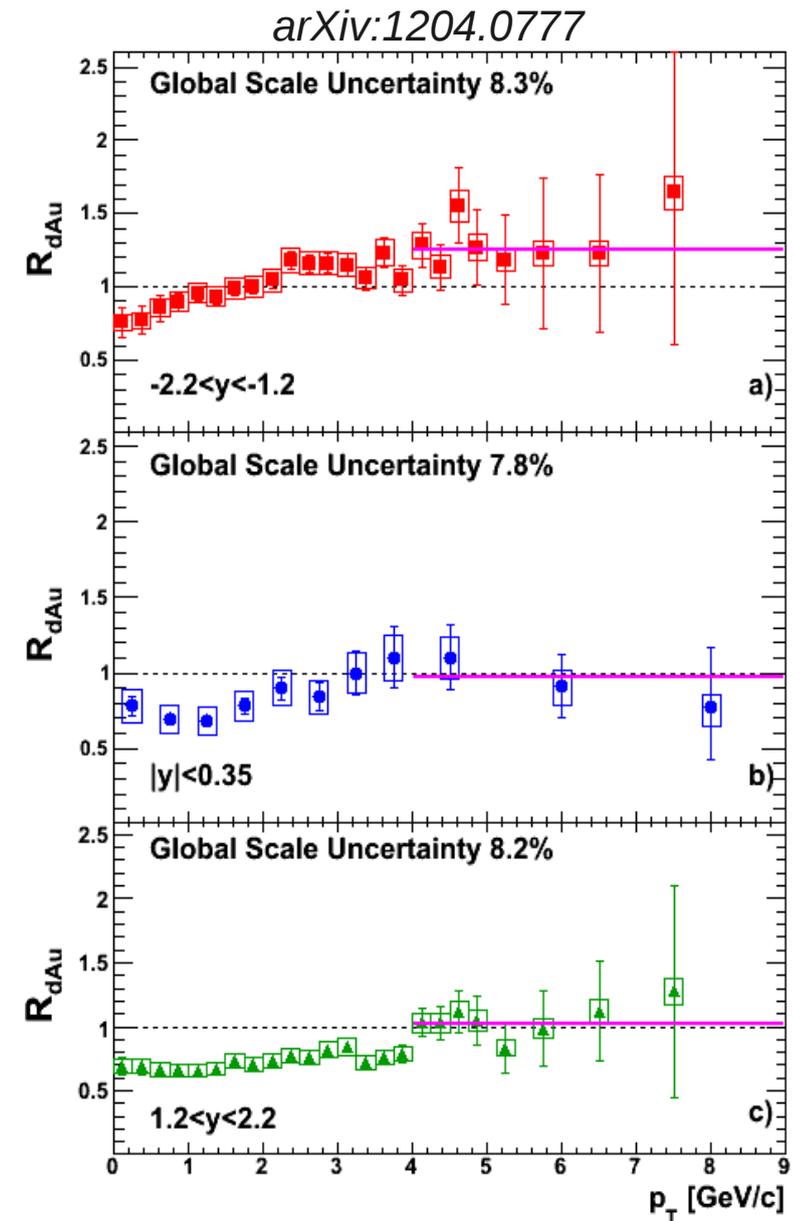


- Fit R_{dAu} w/ constant for $p_T > 4$ GeV/c to find average value.

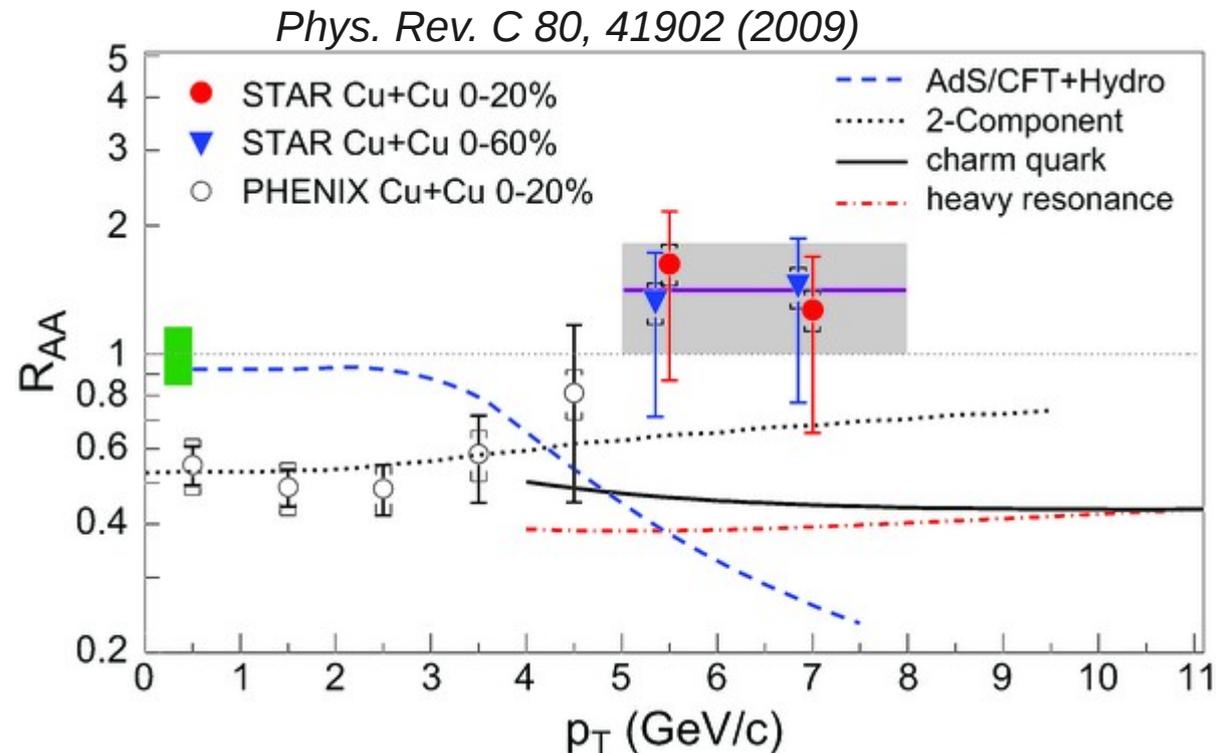
First quoted uncertainty is the fit uncertainty, second is combined Type B & C systematic from data.

| Rapidity | $R_{dAu} (p_T > 4 \text{ GeV}/c)$ |
|-------------------|-----------------------------------|
| $-2.2 < y < -1.2$ | $1.27 \pm 0.06 \pm 0.11$ |
| $ y < 0.35$ | $0.97 \pm 0.14 \pm 0.16$ |
| $1.2 < y < 2.2$ | $1.03 \pm 0.06 \pm 0.11$ |

- R_{dAu} consistent w/ 1 @ mid and forward rapidity.
- $R_{dAu} > 1$ @ backward rapidity.



- Propagating R_{dAu} to R_{AA} qualitatively.
- Expect high- p_T $R_{AA} \approx 1$ at midrapidity.
- Forward rapidity R_{AA} is a combination of forward & backward rapidity R_{dAu} .
- Expect high- p_T $R_{AA} \geq 1$ at forward rapidity.
- Need a more detailed understanding of propagation for quantitative results.

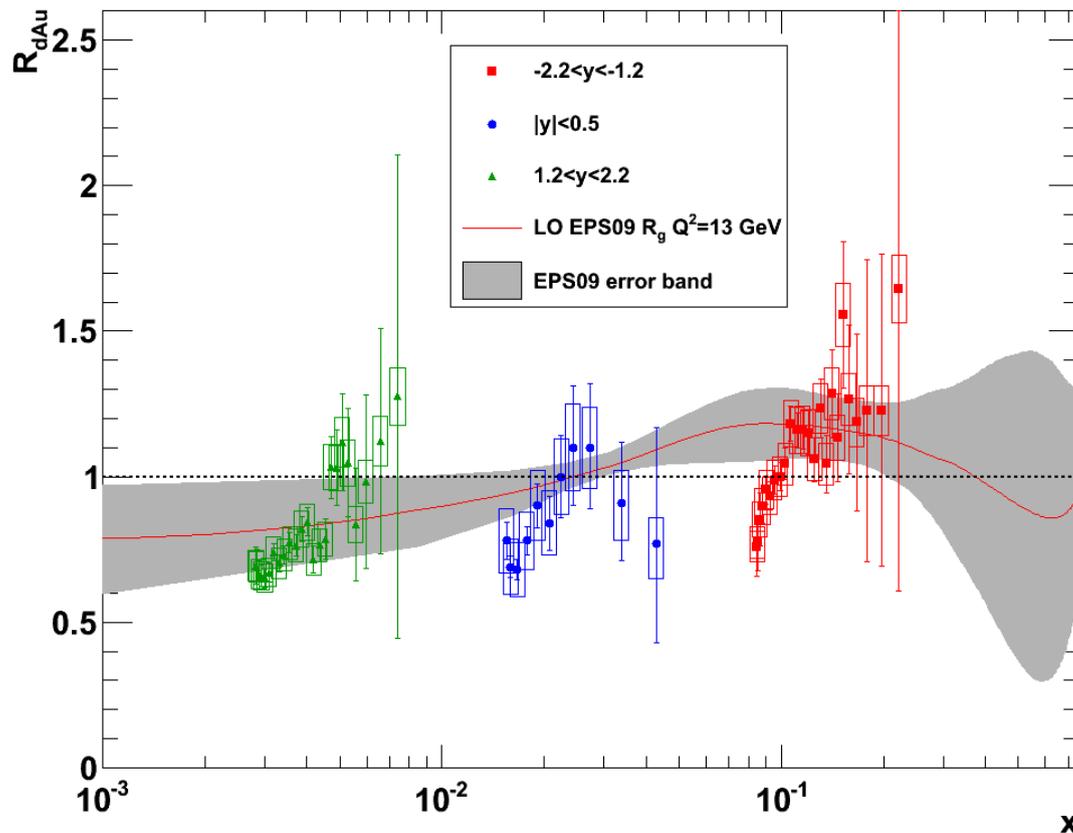


- **Nuclear Shadowing** – Modification of PDF's for nucleons bound in nuclei.
 - Parametrizations of (mostly) DIS data (ex. EKS98, nDSg, EPS09).
- **Nuclear Break-up** – Break-up of cc pair through collisions with nucleons.
 - Usually parametrized using break-up cross section.
- **Cronin Effect** – Broadening of the p_T distribution through scattering of incoming partons.
- **Initial State Energy Loss** – decrease in parton momentum due to soft scatterings while propagating through colliding nucleus.

Nuclear Shadowing (x & Q^2)

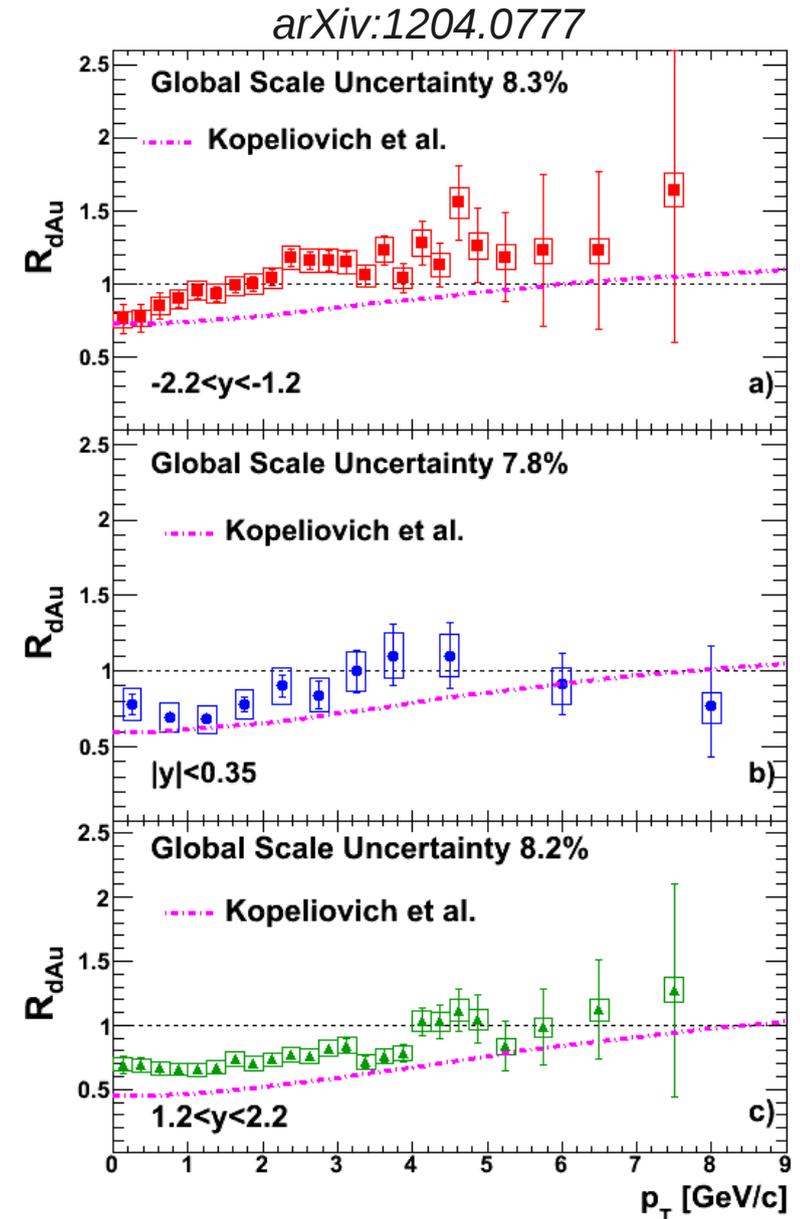
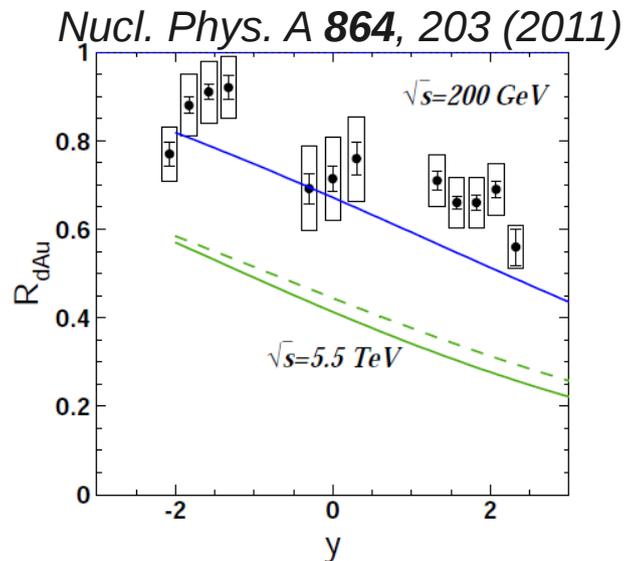
- Use $2 \rightarrow 1$ kinematics to get rough idea of the x & Q^2 regions probed by the data.
- Shown with EPS09 nPDF set.

$$x = \frac{\sqrt{M_{J/\psi}^2 + p_T^2}}{\sqrt{s}} e^{-y} \quad Q^2 = M_{J/\psi}^2 + \langle p_T \rangle^2$$

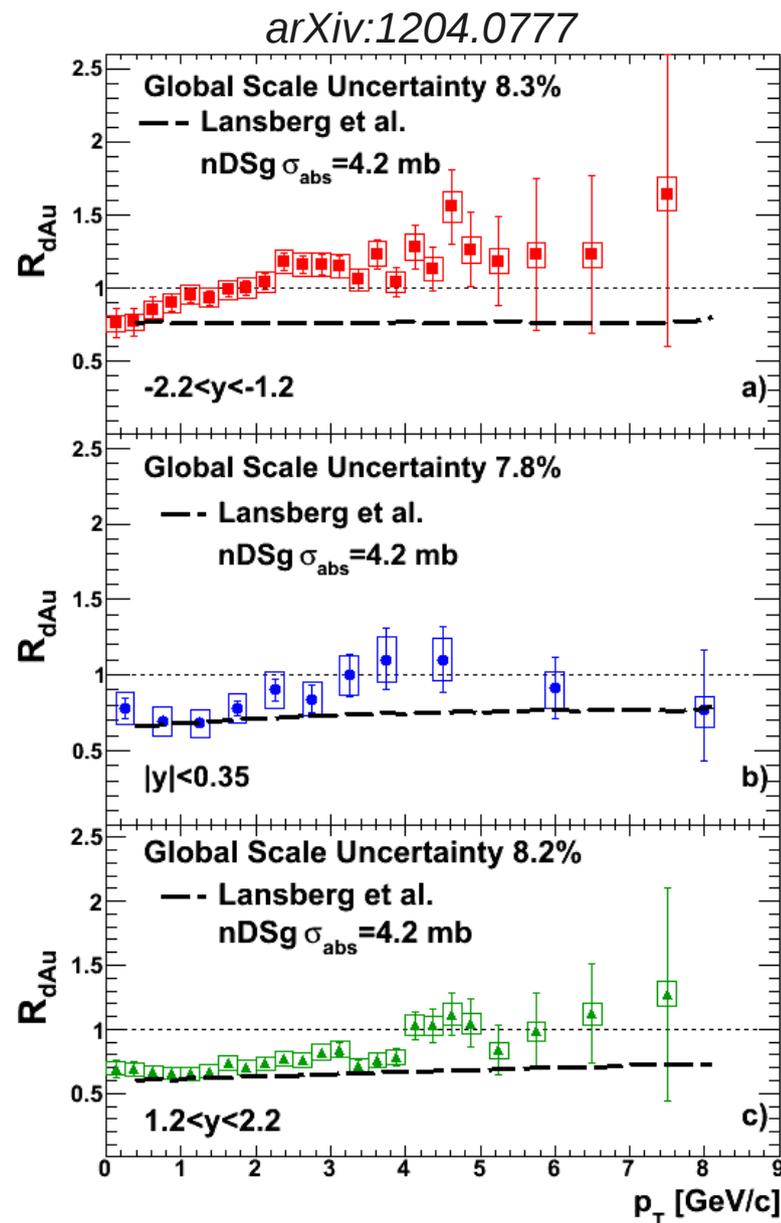
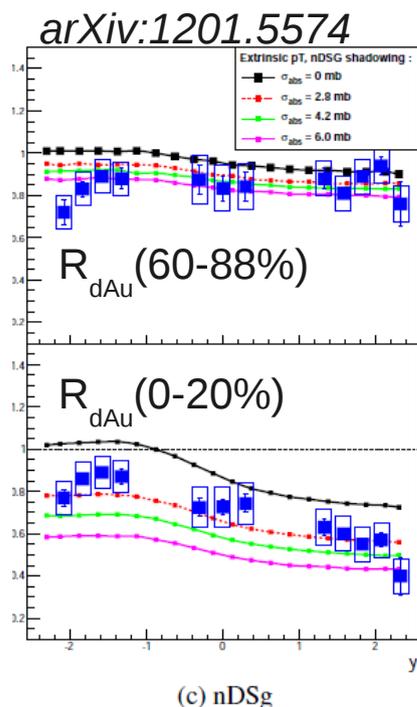


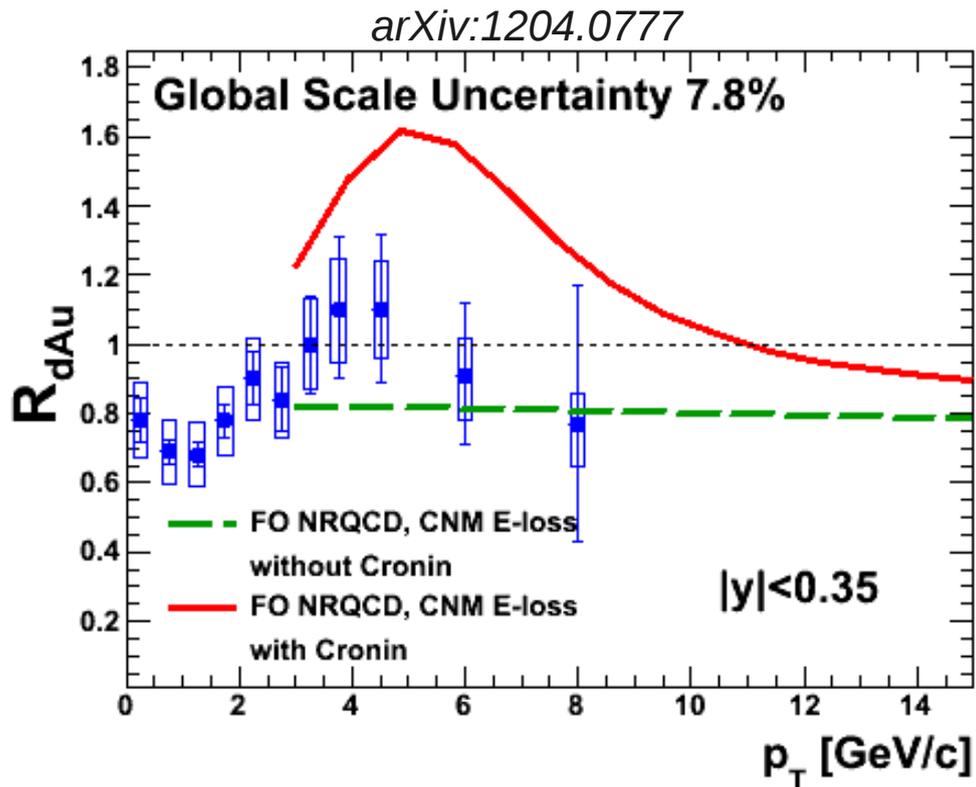
- Forward rapidity – Shadowing region, expect $R_{dAu} < 1$.
- Midrapidity – Transition region.
- Backward rapidity – Anti-shadowing region, expect $R_{dAu} > 1$

- Calculations from Kopeliovich et al. (*Nucl. Phys. A* **864**, 203 (2011))
- Nuclear Shadowing \rightarrow nDSg nPDF set.
- Nuclear Break-up \rightarrow frozen dipole approximation using parametrized break-up cross section from HERA data.
- Includes Cronin Effect.



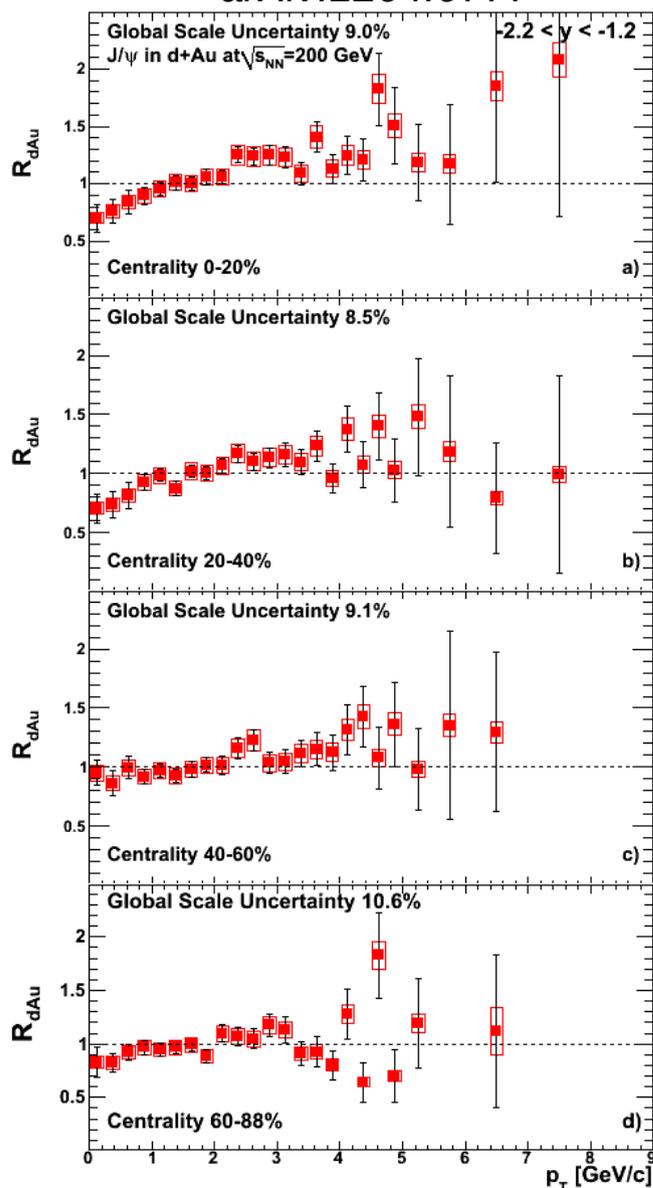
- Calculations from Lansberg et al. (*arXiv:1201.5574*)
- Nuclear Shadowing → nDSg nPDF set.
- Nuclear Break-up → effective $\sigma_{\text{abs}} = 4.2$ mb
- No Cronin Effect.





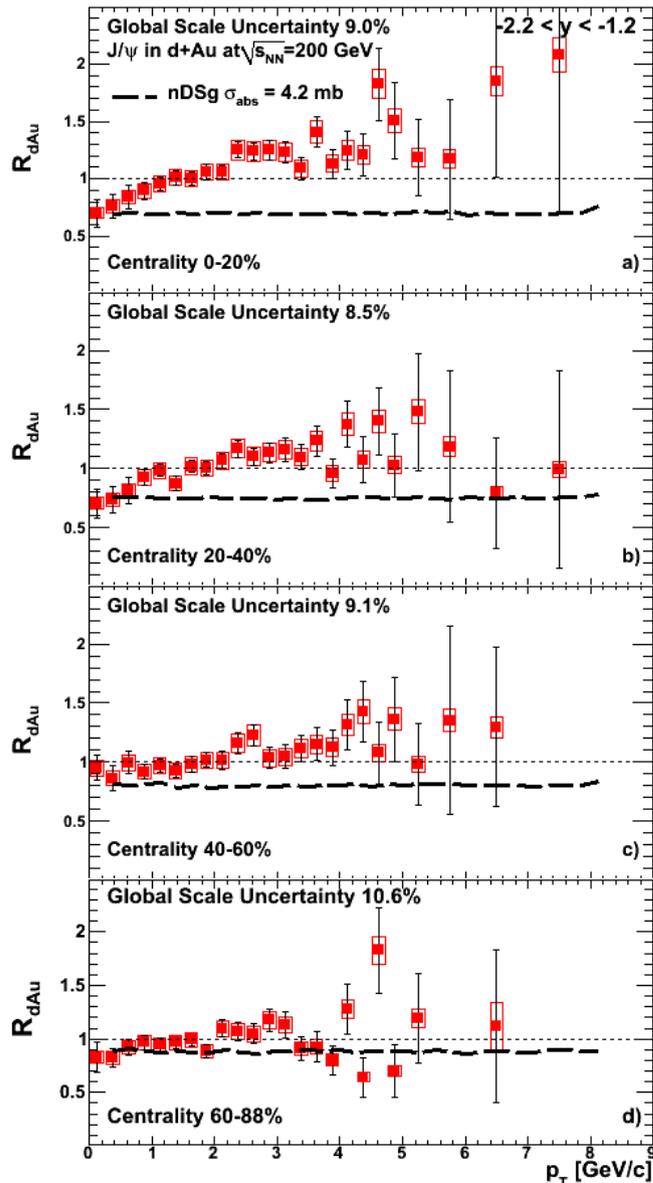
- Calculations by Sharma and Vitev (*arXiv:1203.0329*).
- Nuclear Shadowing – EKS98 w/ power suppressed coherent final state scattering
- Includes initial state energy loss.
- **Dashed Curve** – No Cronin Effect.
- **Solid Curve** – With Cronin Effect.

arXiv:1204.0777



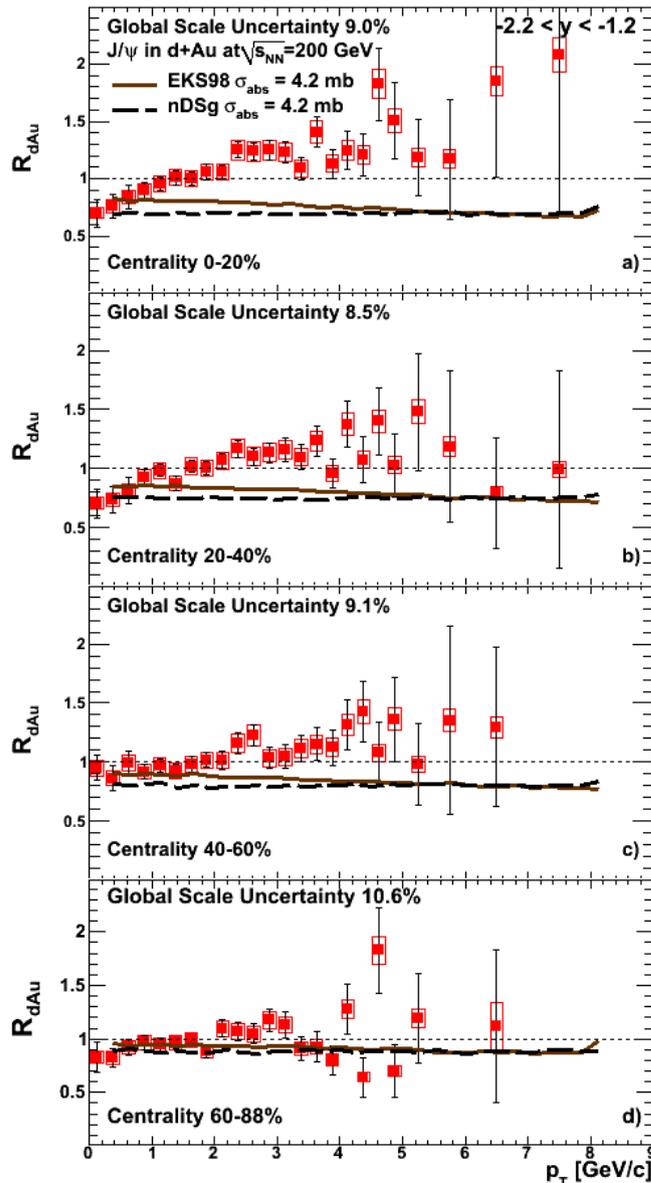
- $-2.2 < y < -1.2$
- Suppression is strongest for central collisions at low p_T .
- Shape of R_{dAu} distribution flattens when moving to peripheral events.
- $R_{dAu} \approx 1$ within uncertainties for peripheral collisions.

arXiv:1204.0777



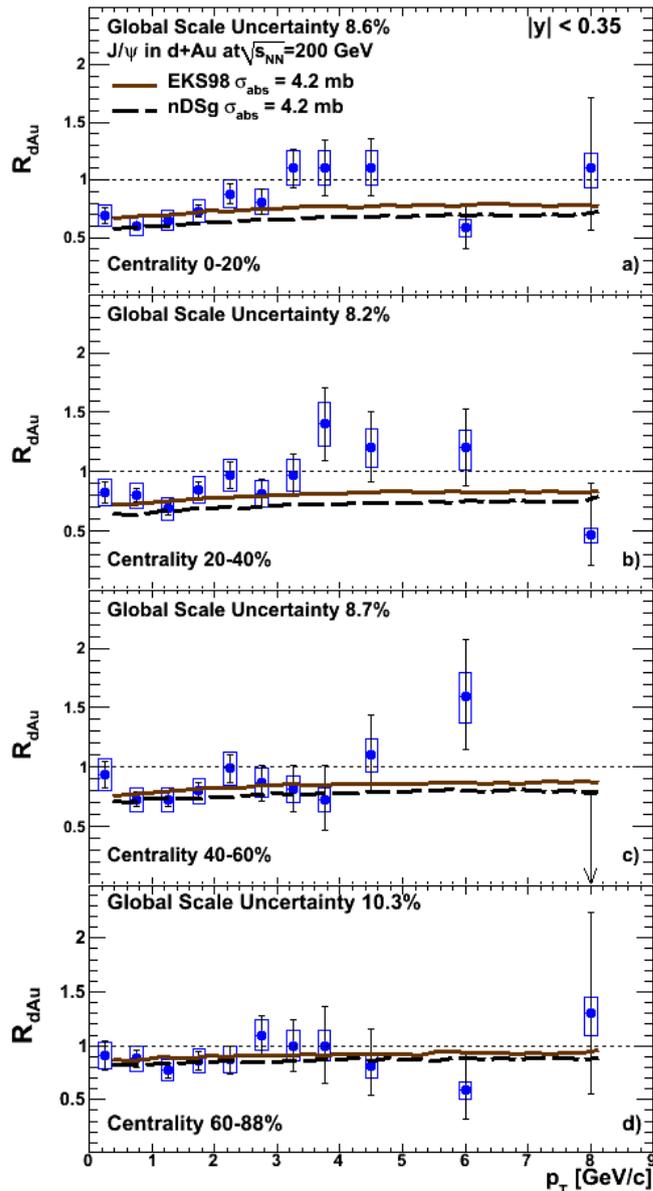
- $-2.2 < y < -1.2$
- Calculations by Lansberg et al.
 - Detailed on previous slide
 - Using nDSg nPDF set.
 - $\sigma_{abs} = 4.2$ mb.
 - Assume modification is proportional to the density.
- Flat in p_T for all centrality.
- Consistent with peripheral data.
- Disagrees strongly for more central collisions.

arXiv:1204.0777



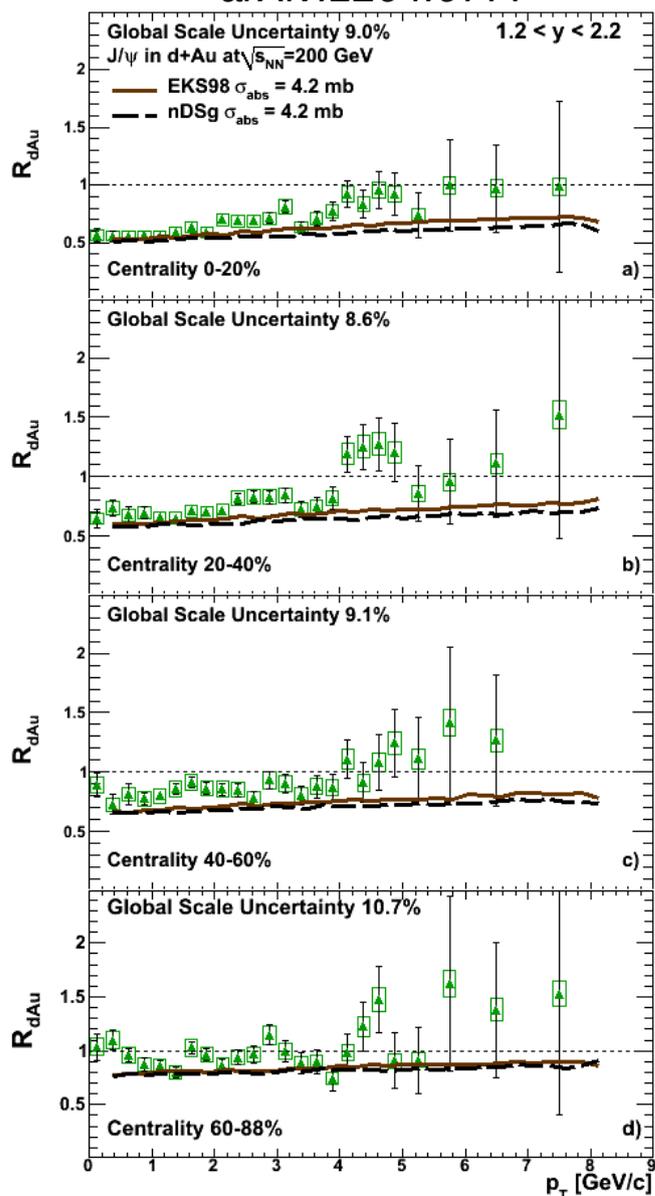
- $-2.2 < y < -1.2$
- Calculations by Lansberg et al.
- Add calculations using EKS98 nPDF set.
- Difference in the calculation due to differences in nPDF's only.
- For central collisions calculations show an increase in suppression with increasing p_T .
- Data shows the opposite trend.

arXiv:1204.0777



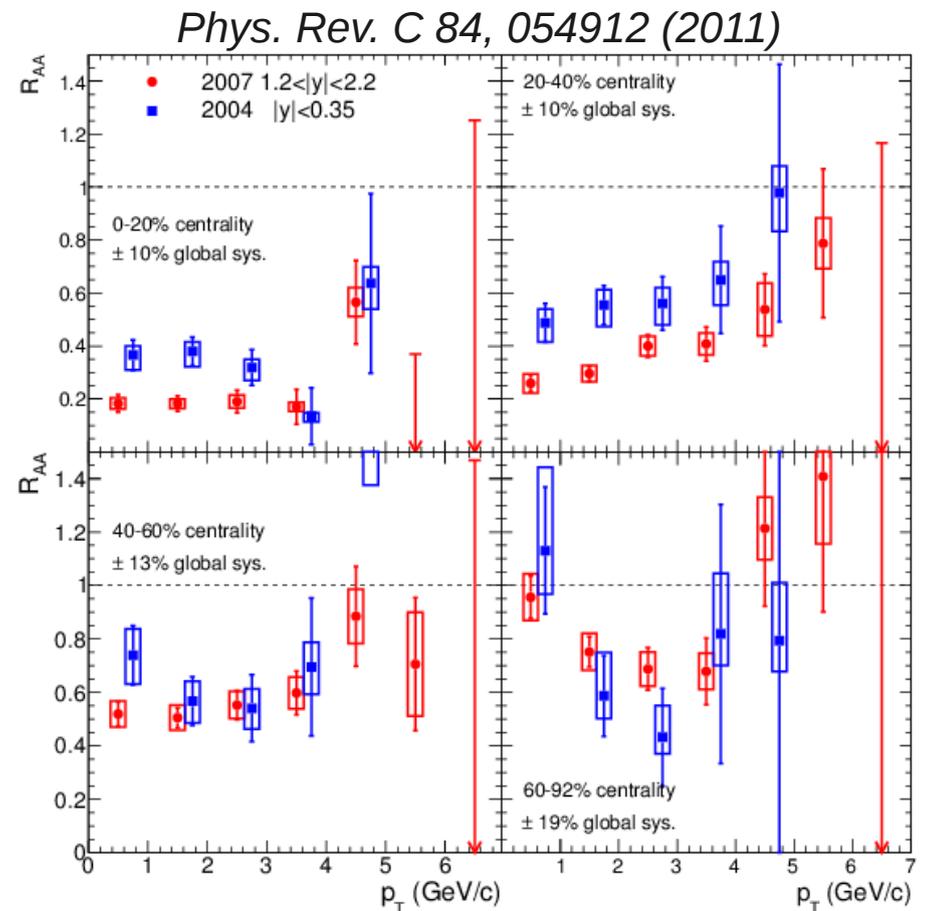
- $|y| < 0.35$
- Same set of calculations.
- Calculations w/ EKS98 and nDSg agree in shape, but not overall normalization.
- Relatively good agreement with the data at all p_T & centrality.

arXiv:1204.0777



- $1.2 < y < 2.2$
- Same set of calculations.
- Calculations w/ EKS98 and nDSg agree in at low p_T , diverge at high- p_T .
- Overall level of suppression over predicted by calculations.

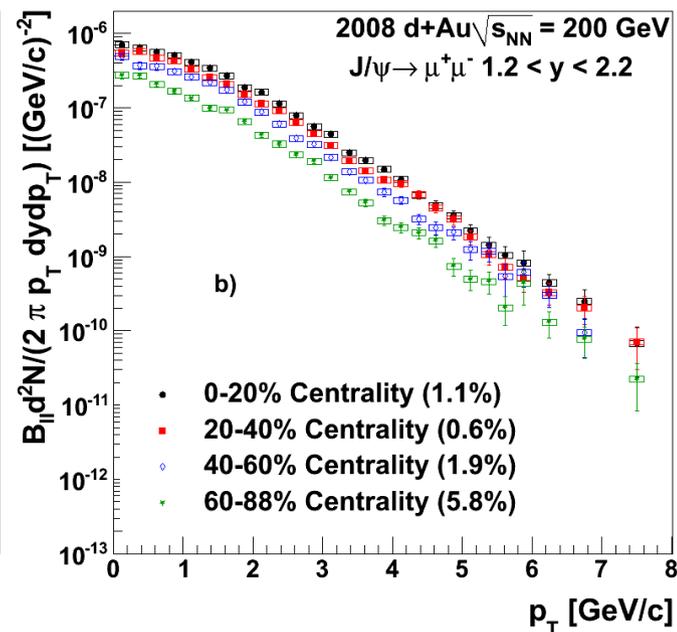
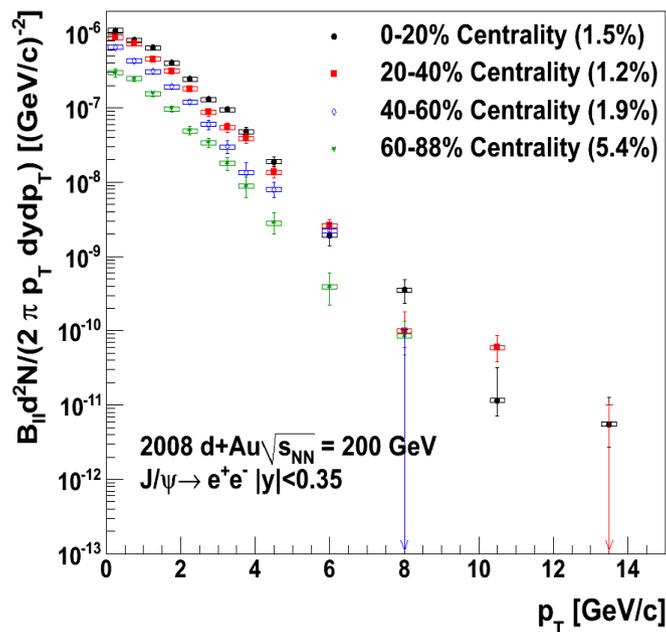
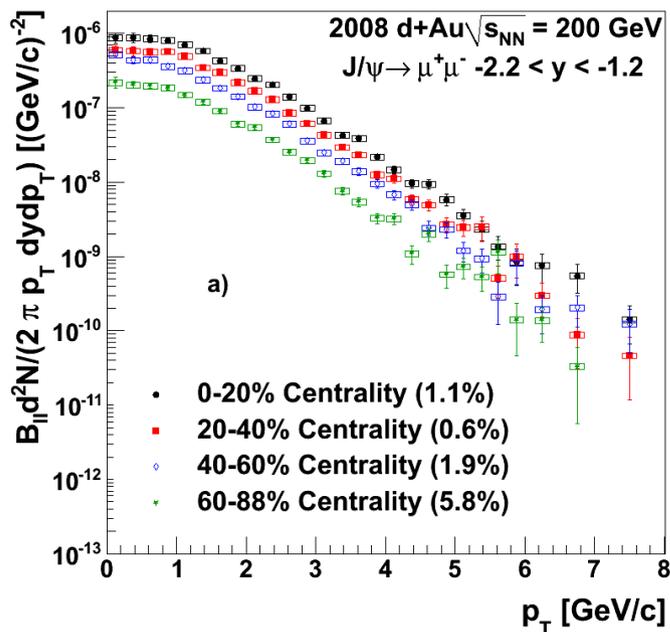
- **Midrapidity:**
 - R_{dAu} for low $p_T \approx 0.7$ for central collisions.
 - Assuming $R_{AA} = R_{dAu}^2$, $R_{AA}(\text{CNM}) \approx 0.5$
 - Observe $R_{AA} \approx 0.4$ for central collisions.
- **Forward Rapidity:**
 - R_{dAu} for low $p_T \approx 0.6$ @ forward rapidity & 0.8 @ backward rapidity for central collisions.
 - $R_{AA}(\text{CNM}) \approx 0.5$
 - Observe $R_{AA} \approx 0.2$ for central collisions.
- Crude estimate, but gives some expectation of CNM effect in A+A.

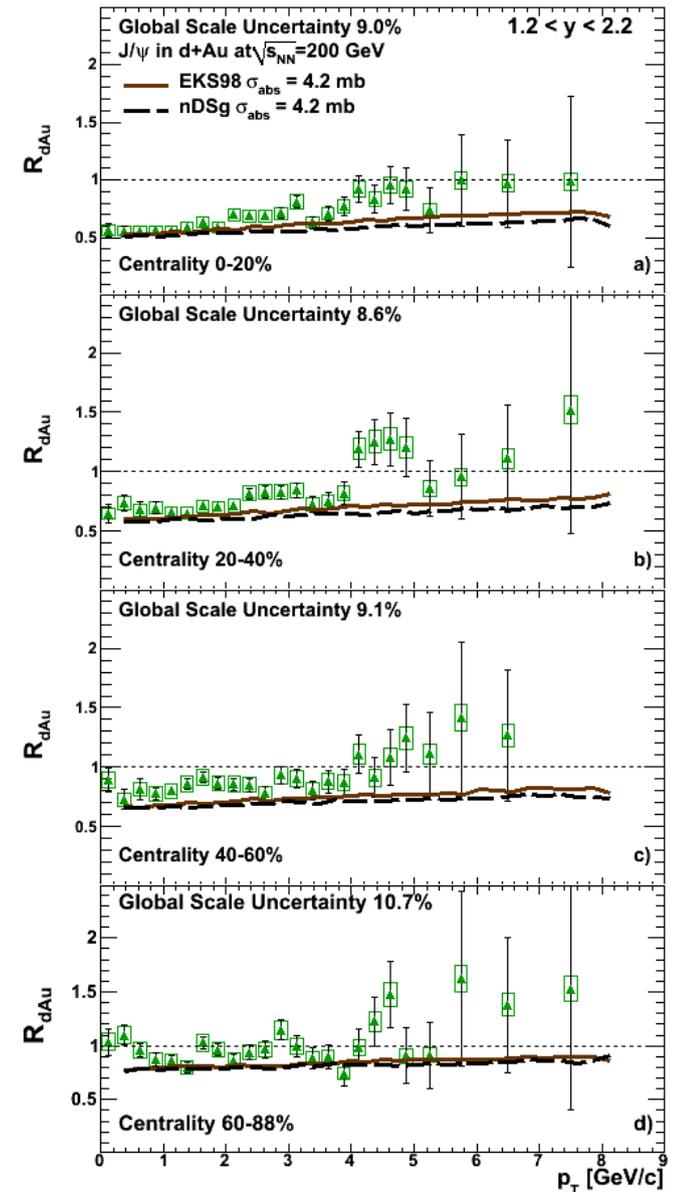
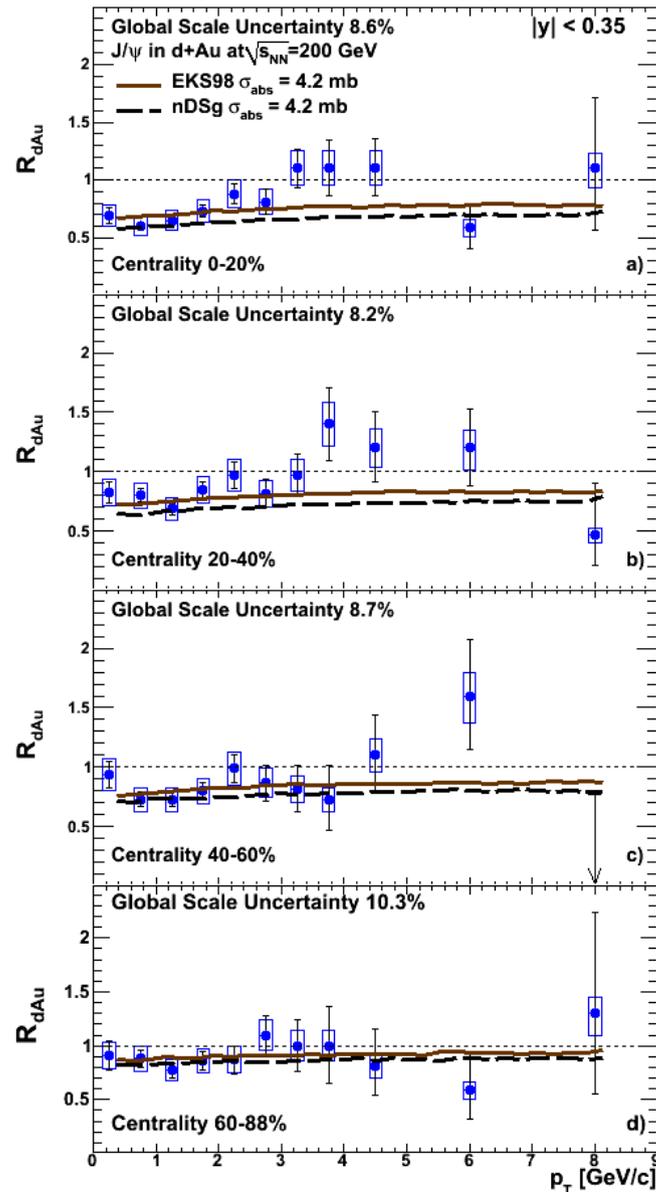
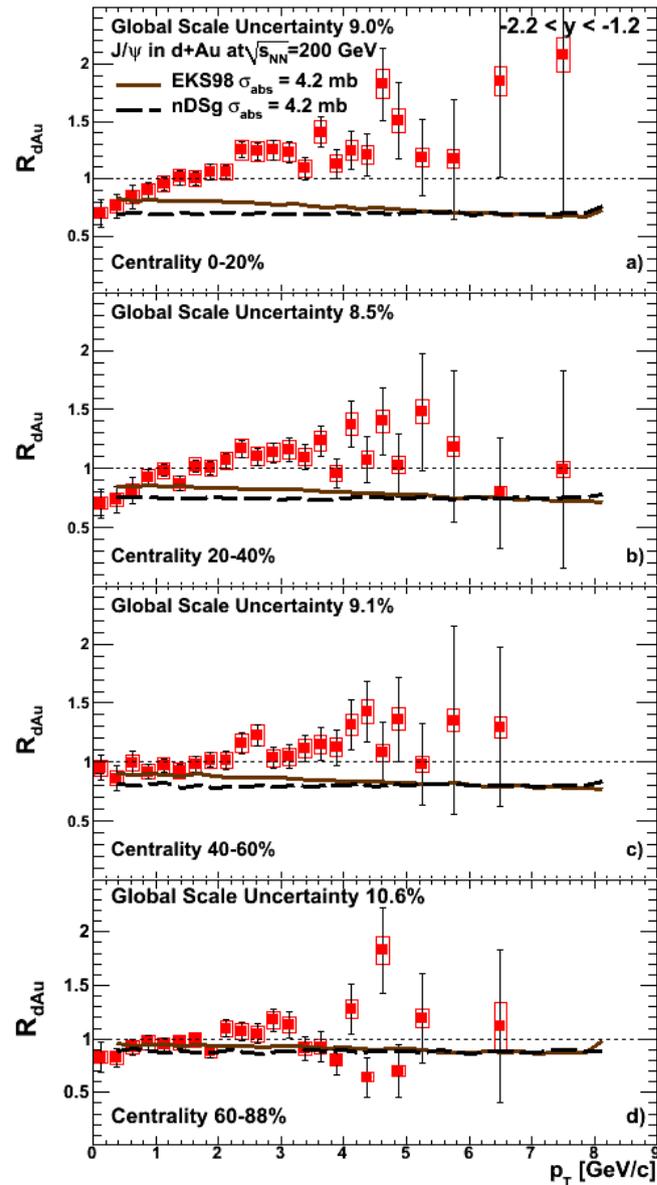


- PHENIX has made new measurements of the p_T dependence of R_{dAu} .
 - J/ψ measured over p_T , rapidity, and centrality.
- Rapidity dependence is inconsistent with modifications which are linearly or exponentially dependent on the nuclear thickness.
- p_T dependence of R_{dAu} is similar at mid & forward rapidities with a more rapid decrease in suppression with increasing p_T at backward rapidity.
- Average R_{dAu} at high- $p_T > 1$ at backward rapidity.
- Need a detailed understanding of CNM effects and their propagation to A+A to understand hot nuclear matter effects.

Backup

- J/ψ invariant yields vs p_T for each centrality.





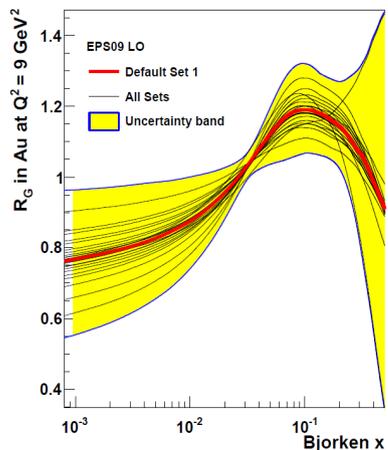
1) 1st Calculation includes two components.

1) Gluon modification (shadowing) from EPS09 nPDF – parametrization of DIS+pA data.

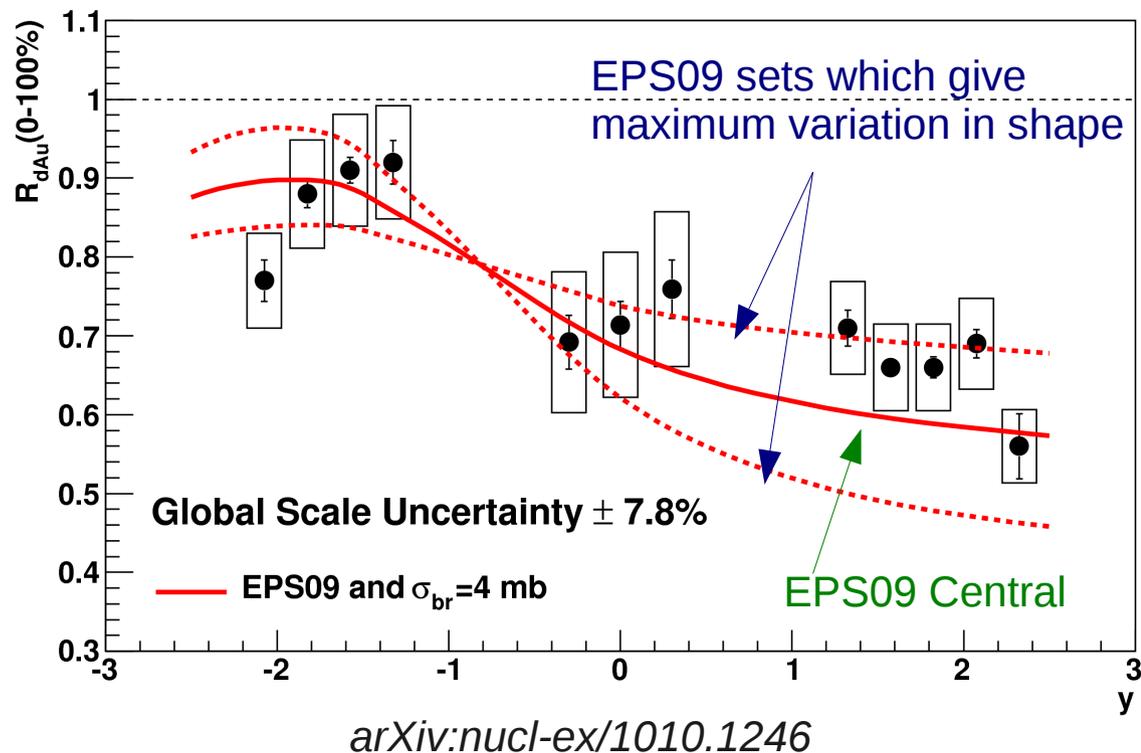
- Calculations are modification vs. nucleon impact parameter (r_T) in the Au nucleus.
- Fold r_T distribution with PHENIX centrality distributions calculated from Glauber MC.

2) Nuclear Break-up cross section σ_{br} – due to collisions of J/ψ with nucleons

- $\sigma_{br} = 4$ mb chosen to match backward rapidity data.
- Shows reasonable agreement over all rapidity, as expected for MB data.



LO EPS09 R_G vs. Bjorken x for Au
 ← no intrinsic centrality dependence



1) 1st Calculation includes two components.

1) Gluon modification (shadowing) from EPS09 nPDF – parametrization of DIS+pA data.

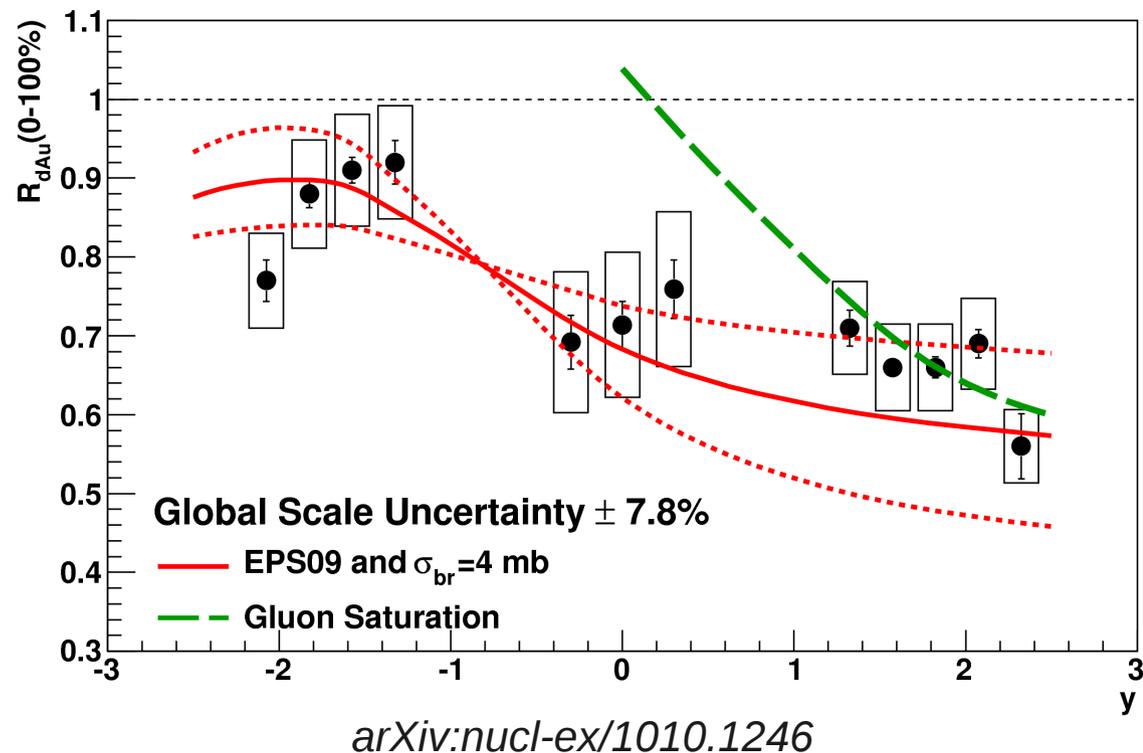
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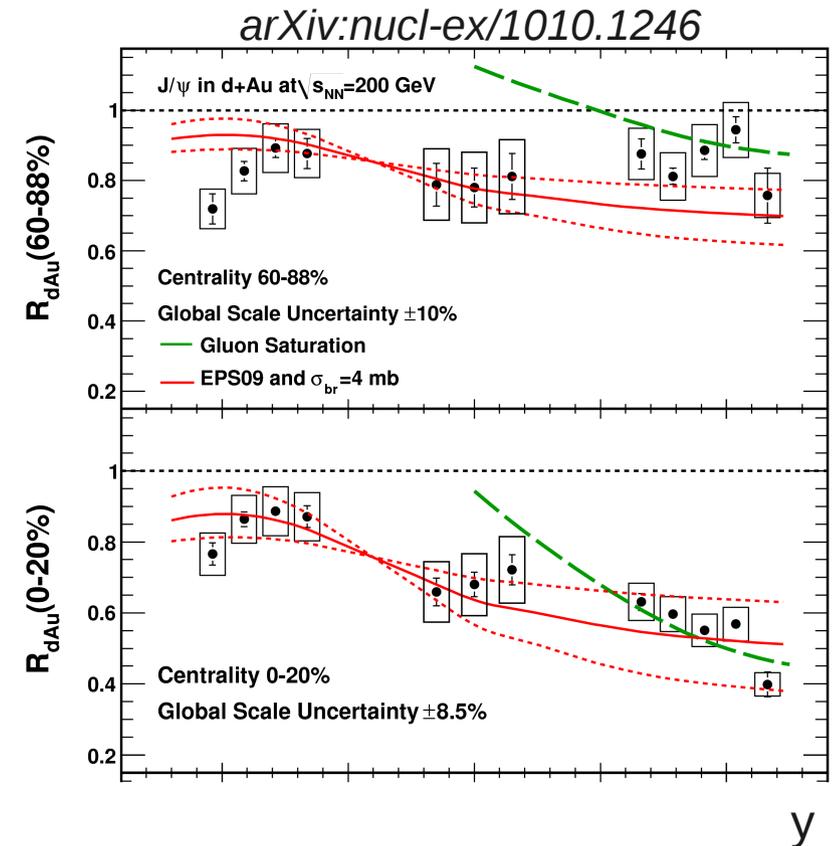
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- Shows reasonable agreement over all rapidity. As expected for MB data.

2) Calculation by Kharzeev and Tuchin (*Nucl. Phys. A 770 (2006)40*)

- Includes Gluon Saturation at low x
- Shows good agreement @ +y.
- Unrealistic at backward and mid rapidity.
- Validity uncertain for peripheral events?

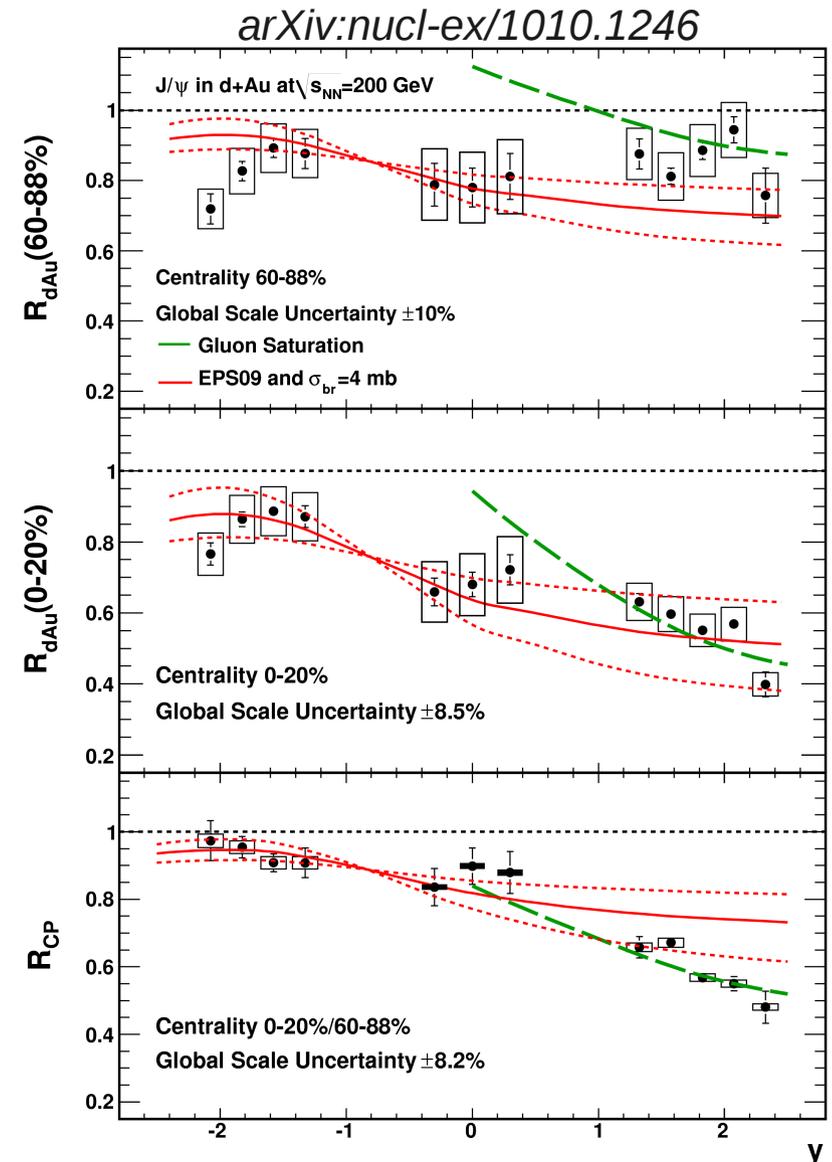


- Want to investigate centrality (impact parameter) dependence.
 - Divide into percentage bins based on BBC charge (0% - most central, 100% - most peripheral).
- Calculations by the same models as detailed on previous slide.
 - Must introduce centrality dependence into EPS09 – **arbitrarily** choose linear dependence on nuclear thickness (common assumption).
 - Shadowing + break-up does not describe forward rapidity data for peripheral collisions.
 - Gluon saturation model still describes data well at +y.

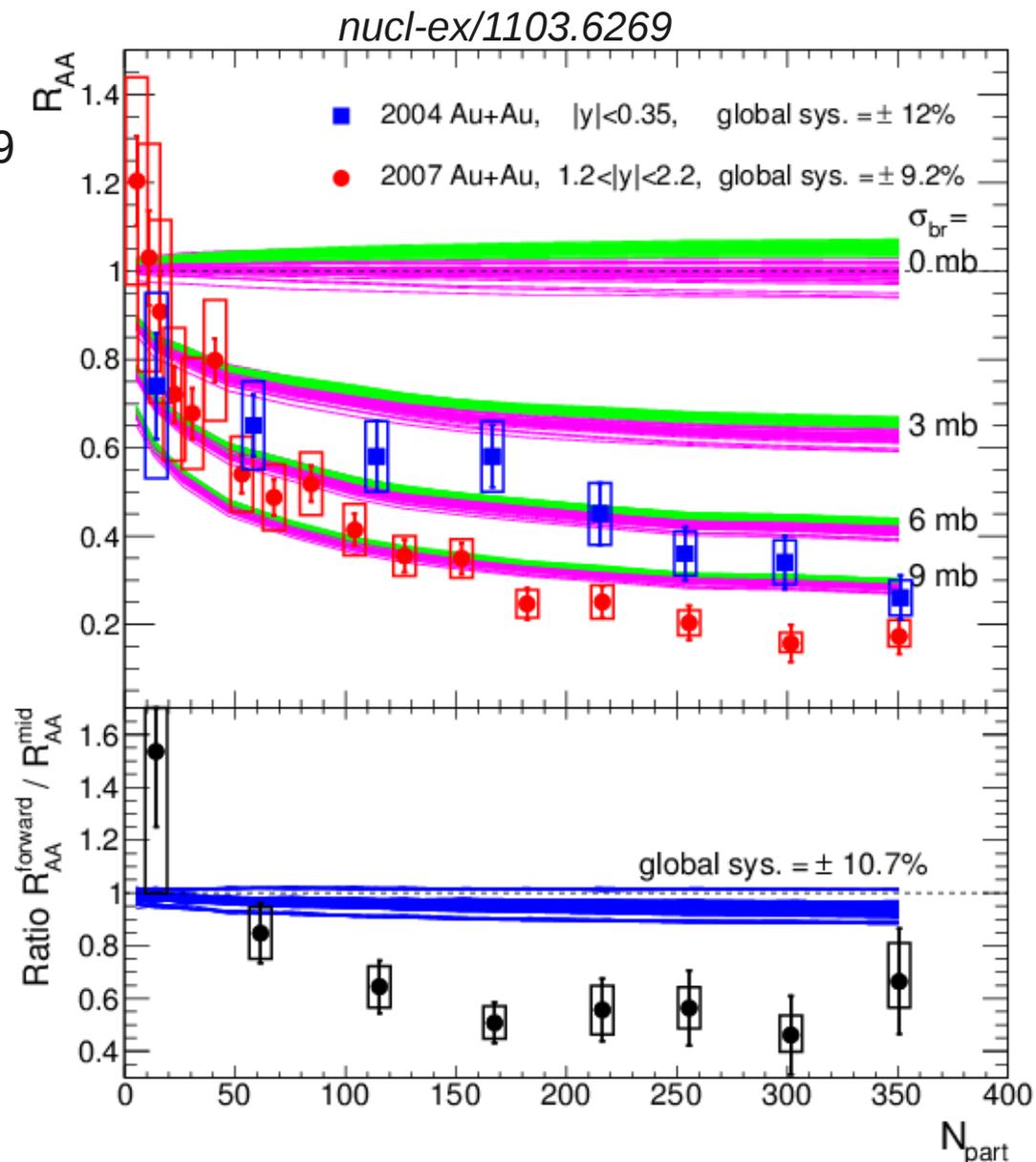


- Take the ratio of central R_{dAu} to peripheral $R_{dAu} \rightarrow R_{cp}$
 - Significant reduction of systematic errors
 - Shadowing + σ_{br} describes backward & midrapidity well.
 - Failure of Shadowing + σ_{br} to describe R_{cp} at large y seems to be due to poor description of centrality dependence.
 - Gluon saturation model appears to provide a better description of the centrality dependence, although it is not clear how reliable it is for peripheral collisions where there should be less coherent effects.

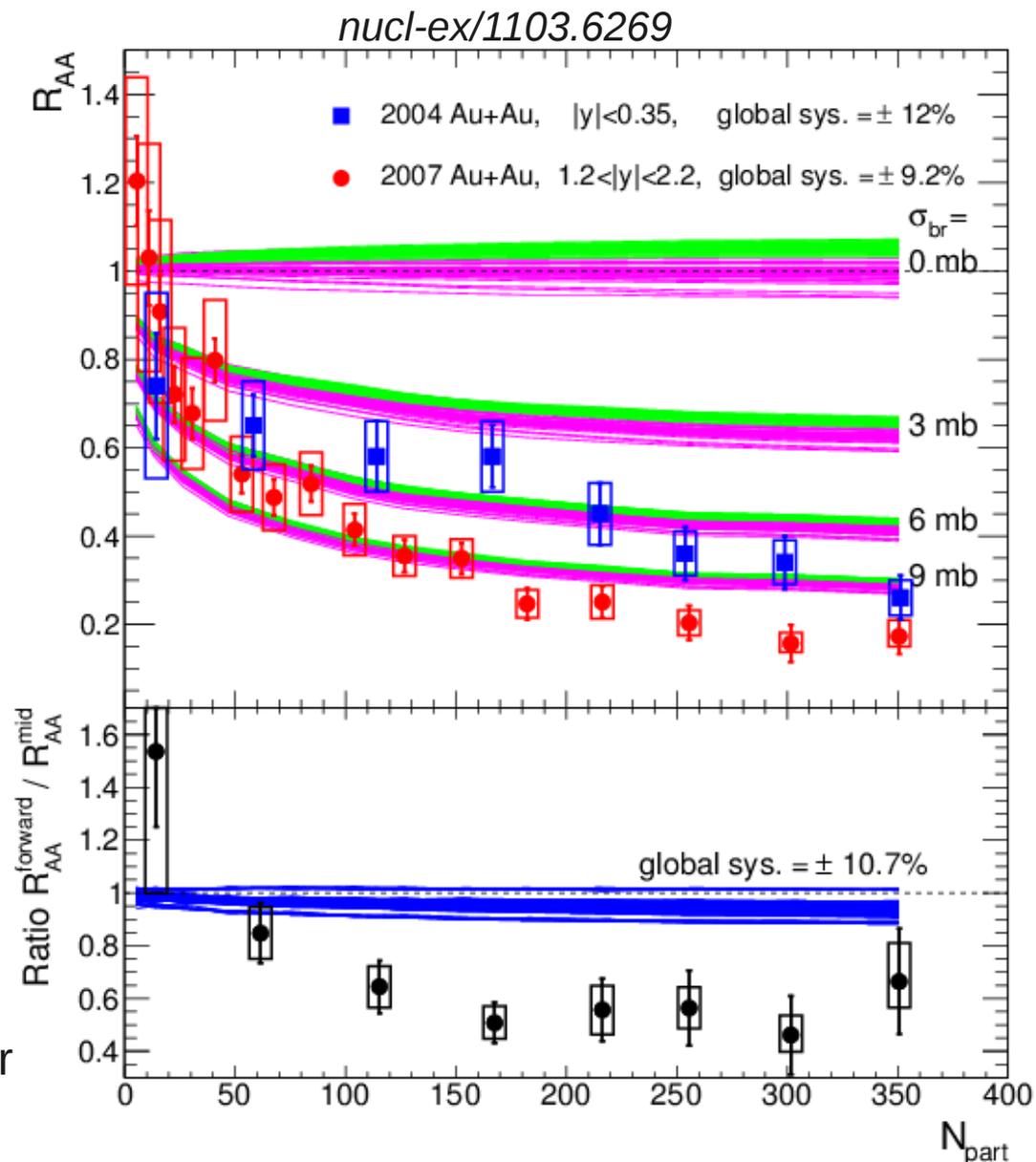
$$R_{cp}(0 - 20\%) = \frac{R_{dAu}(0 - 20\%)}{R_{dAu}(60 - 88\%)}$$



- Calculations show CNM effects extrapolated to Au+Au for (**linear**) EPS09 + (**exponential**) nuclear breakup.
- **Green curves** – predictions for $|y| < 0.35$ for each of 31 EPS09 sets for 0, 3, 6, 9 mb breakup cross sections
- **Magenta curves** – same as green curves, but for $1.2 < |y| < 2.2$
- Bottom panel shows the ratio of data & CNM predictions.

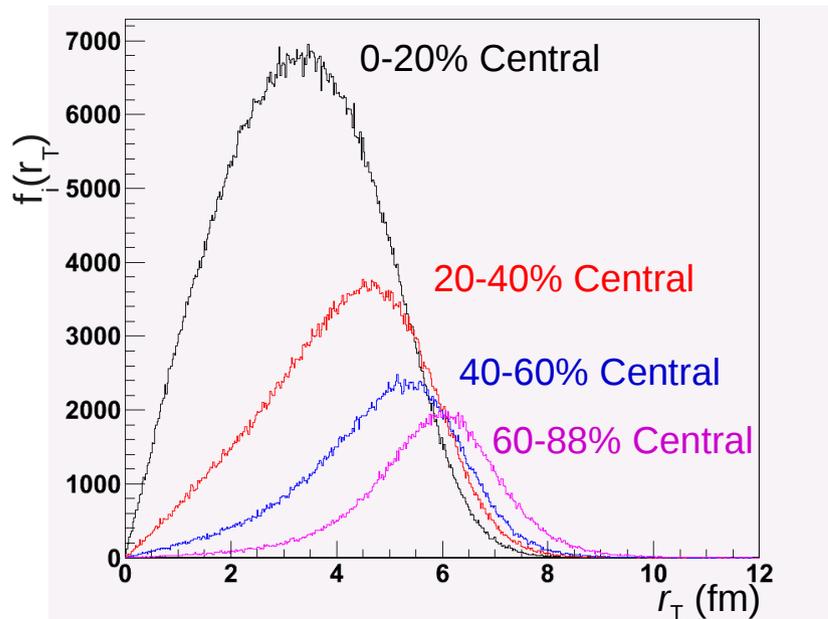
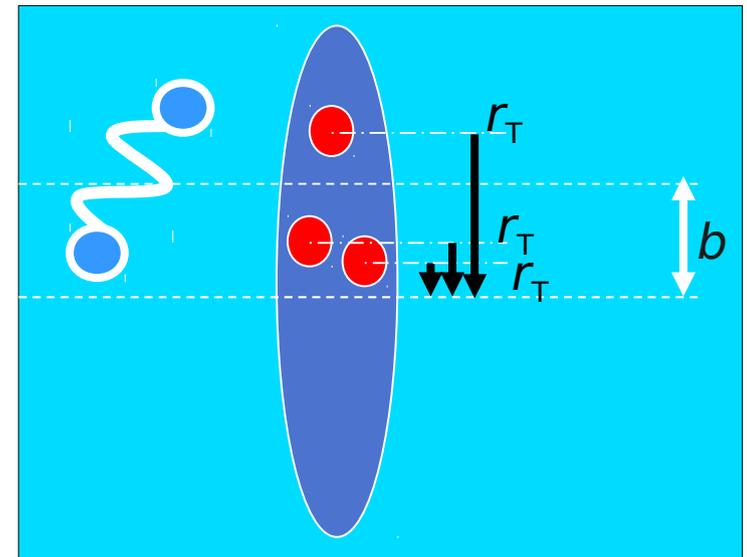


- What do we learn from d+Au:
 - Linear EPS09 w/ a 4mb breakup cs adequately describes backward & midrapidity d+Au data, not sufficient at forward rapidity.
- What does this imply in Au+Au:
 - Suppression at **midrapidity** stronger than expected from d+Au alone.
 - Can not make a similar statement about the forward rapidity data until we understand d+Au at forward y.
 - Linear EPS09 does not explain difference of suppression with rapidity.
- Still, clear evidence of hot nuclear matter effects, although not yet quantifiable.



Simple Geometrical Model

- Would like to understand how the suppression depends on centrality.
- In d+Au relevant parameter is transverse position of the struck nucleon in each N-N collision $\rightarrow r_T$
- Use Glauber MC of N-N hit positions in d+Au events to generate r_T distributions.



- Use a simple parametrization of the nuclear modification based on the density weighted longitudinal thickness in the Au nucleus $\rightarrow \Lambda(r_T)$.

$$\Lambda(r_T) = \frac{1}{\rho_0} \int dz \rho(z, r_T)$$

Woods-Saxon



Resulting r_T distributions from MC for PHENIX centrality bins.

Consider, for example, three functional forms for the nuclear modification vs nuclear thickness at r_T , $\Lambda(r_T)$, with one free strength parameter a

$$M(r_T, a) = 1 - a\Lambda(r_T)$$

$$M(r_T, a) = 1 - a\Lambda(r_T)^2$$

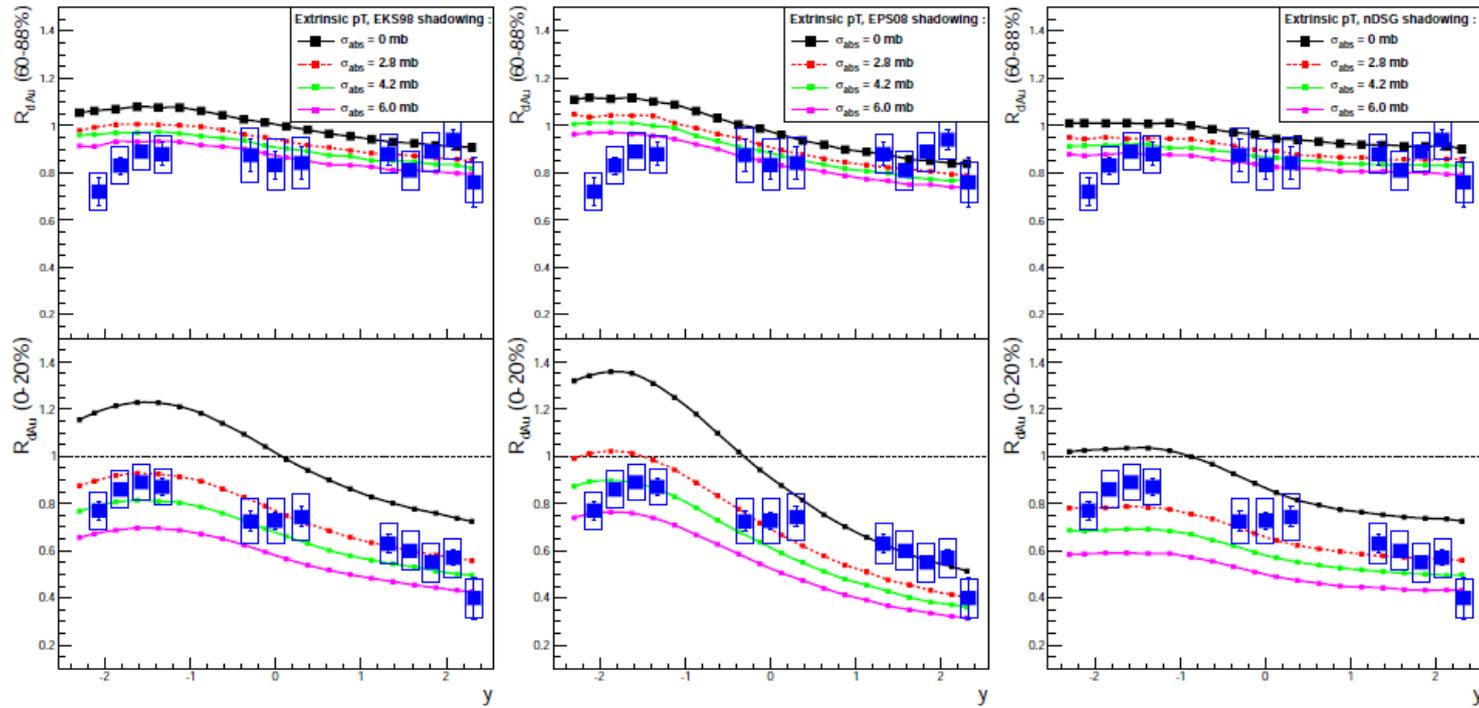
$$M(r_T, a) = e^{-a\Lambda(r_T)}$$

The modification factor R_{dAu} for a given centrality bin (i) is then given by

$$R_{dAu,i}(a) = \int f_i(r_T) M(r_T, a) dr_T$$

r_T distributions
from PHENIX MC

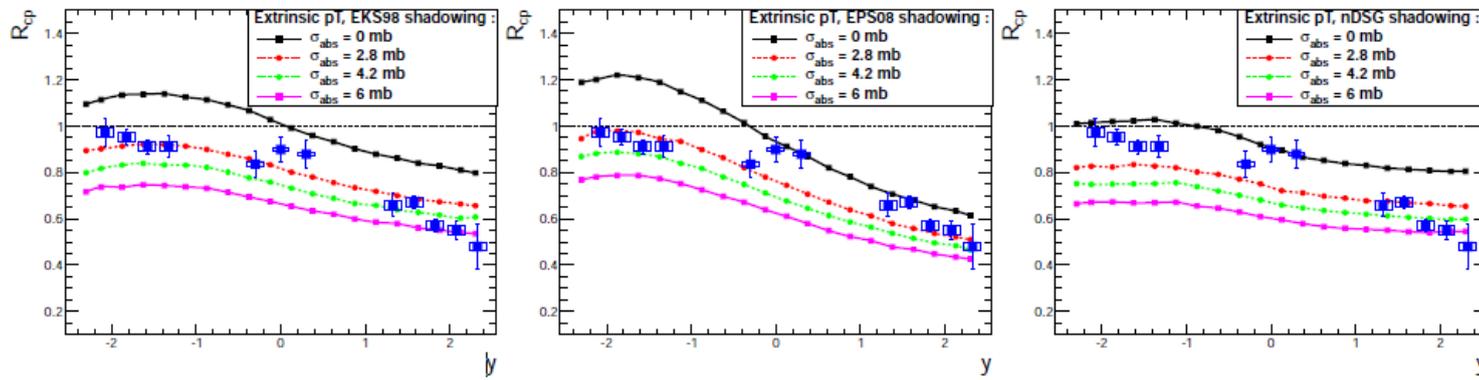
Modification vs $\Lambda(r_T)$ and
free parameter a



(a) EKS98

(b) EPS08

(c) nDSg



(a) EKS98

(b) EPS08

(c) nDSg

nPDF Comparison

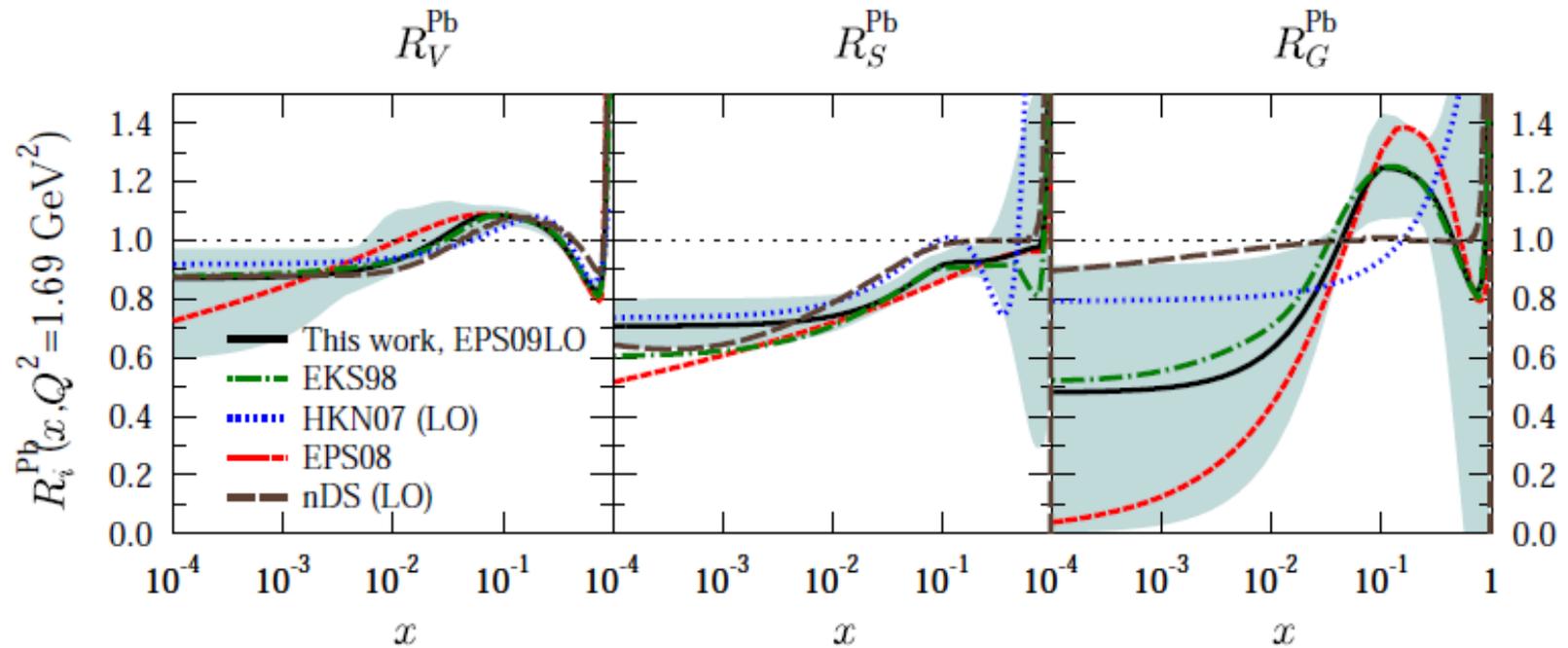


Fig. 12. Comparison of the nuclear parton distribution functions and shadowing effects.