



J/ Ψ 's at RHIC

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With thanks to my PHENIX (and other RHIC) colleagues,
and special thanks to

V. Cianciolò, M. Leitch, Hugo Pereira Da Costa, V.N. Tram, W. Xie

See especially [Quarkonia Production at RHIC](#), M. Leitch, SQM 2006.



Outline

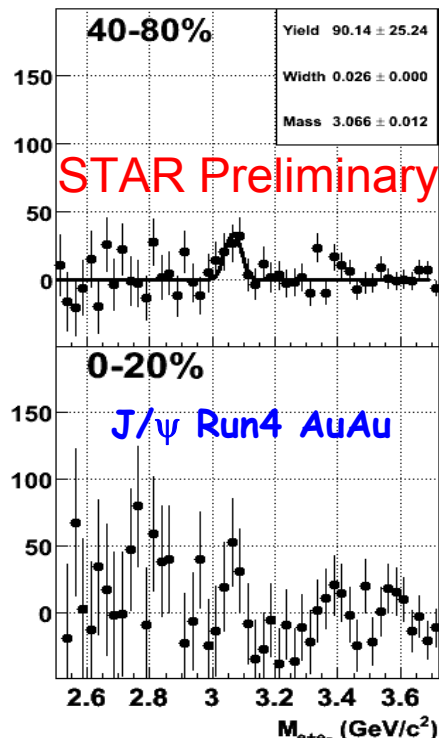
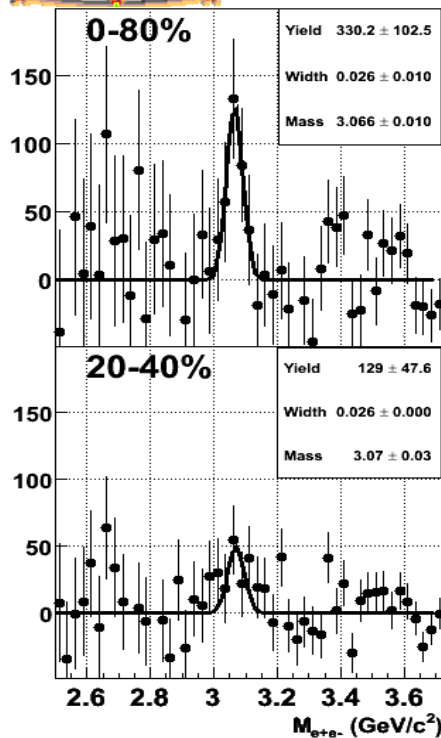
- *Experiments*
- *Yields*
- *Ratios*
- *Models*
- *Future*
- *Helmut*



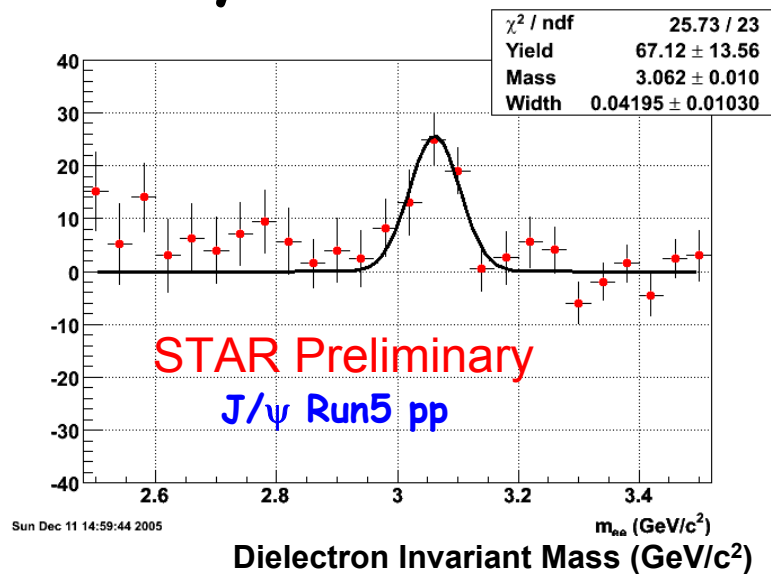
Experiments



Charmonium and Beyond in STAR

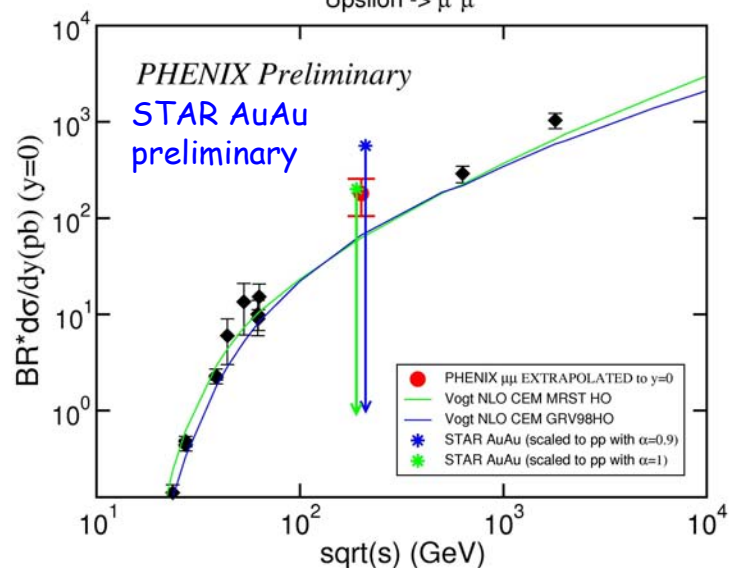


J. Gonzalez, SQM



Dielectron Invariant Mass (GeV/c^2)

PHENIX 200 GeV p-p
Upsilon $\rightarrow \mu^+ \mu^-$



Signal	RHIC Exp. (Au+Au)	RHIC I (>2008)	RHIC II	LHC ALICE ⁺
J/ψ $\rightarrow e^+e^-$	PHENIX	3,300	45,000	9,500
J/ψ $\rightarrow \mu^+\mu^-$		29,000	395,000	740,000
Υ $\rightarrow e^+e^-$	STAR	830	11,200	2,600
Υ $\rightarrow \mu^+\mu^-$	PHENIX	80	1,040	8,400

(Compilation courtesy of M. Leitch, Strange Quark Matter, 2006)



The PHENIX detector

Central arms:

hadrons, photons, electrons

$$J/\Psi \rightarrow e^+e^-$$

$$p > 0.2 \text{ GeV}/c$$

$$|y| < 0.35$$

$$\Delta\phi = \pi$$

Muon arms:

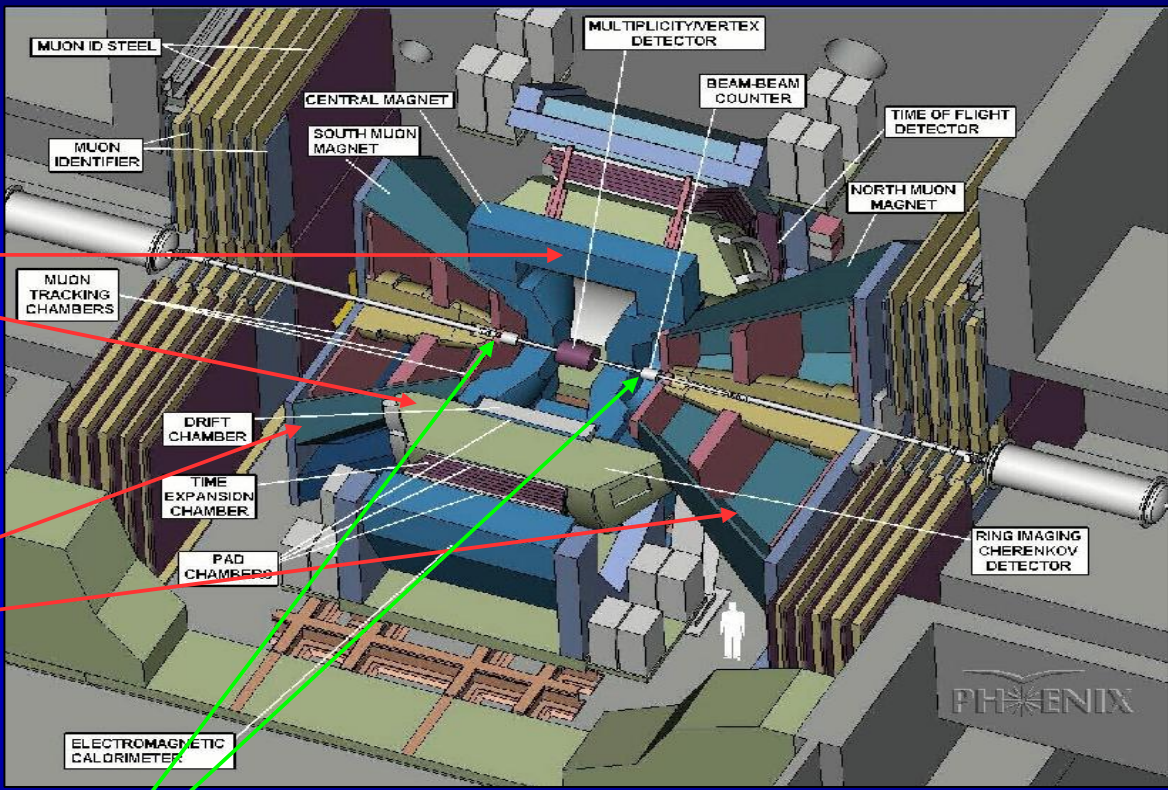
muons at forward rapidity

$$J/\Psi \rightarrow \mu^+\mu^-$$

$$p > 2\text{GeV}/c$$

$$1.2 < |y| < 2.4$$

$$\Delta\phi = 2\pi$$



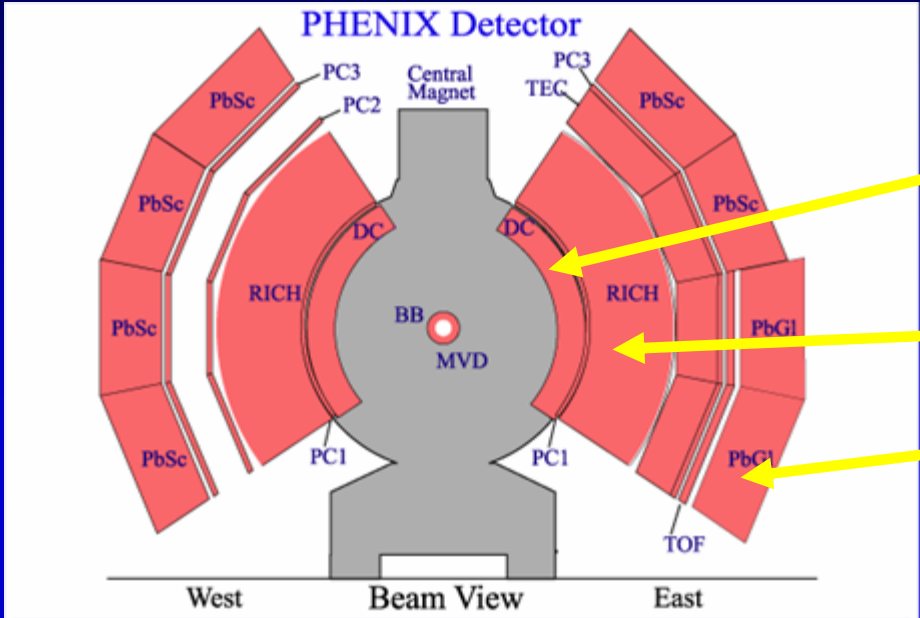
Centrality measurement:

We use beam beam counters together with zero degree calorimeters

Centrality is mapped to N_{part} (N_{col}) using Glauber model



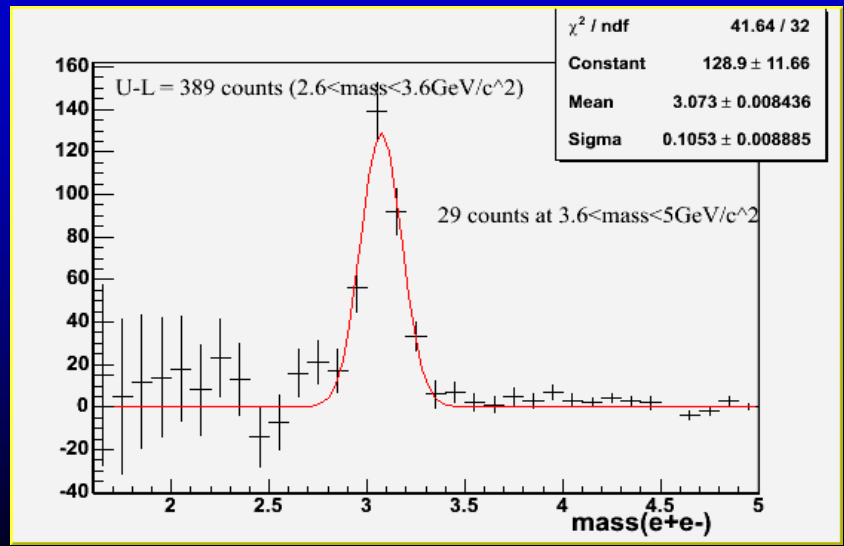
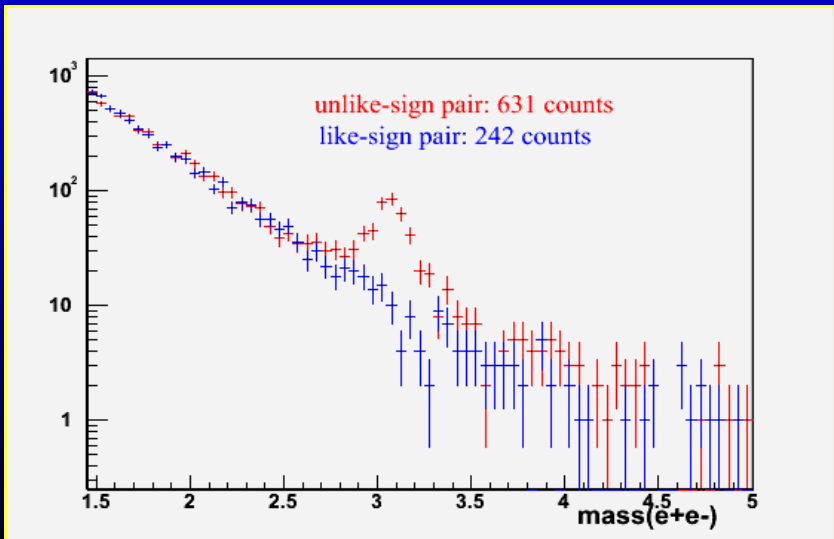
How PHENIX Measures J/ψ's



- High resolution tracking and momentum measurement from Drift chamber.

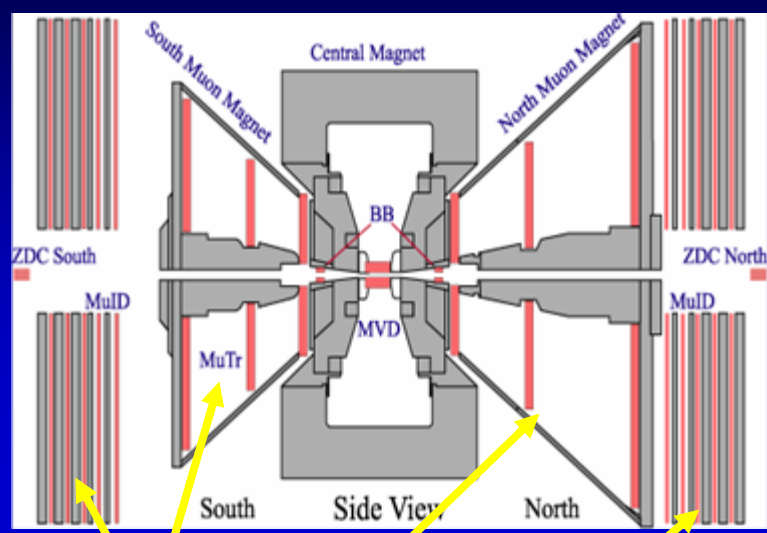
Good electron identification from Ring Imaging Cherenkov detector (RICH) and Electromagnetic Calorimeter (EMCal).

High rate capability: powerful level-1 electron trigger



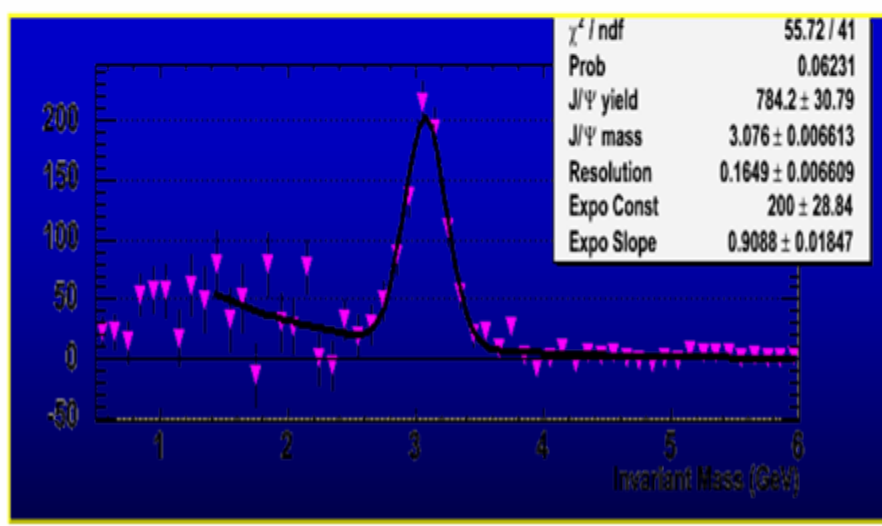
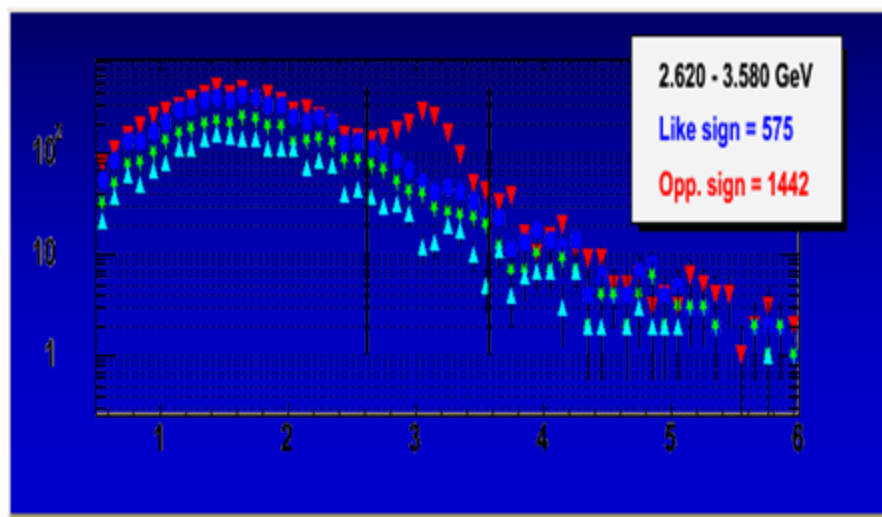


How PHENIX Measures J/ψ 's



Good momentum resolution and muon identification from μ Trk and μ ID.

High rate capability: powerful level-1 dimuon trigger





PHENIX Data Sets

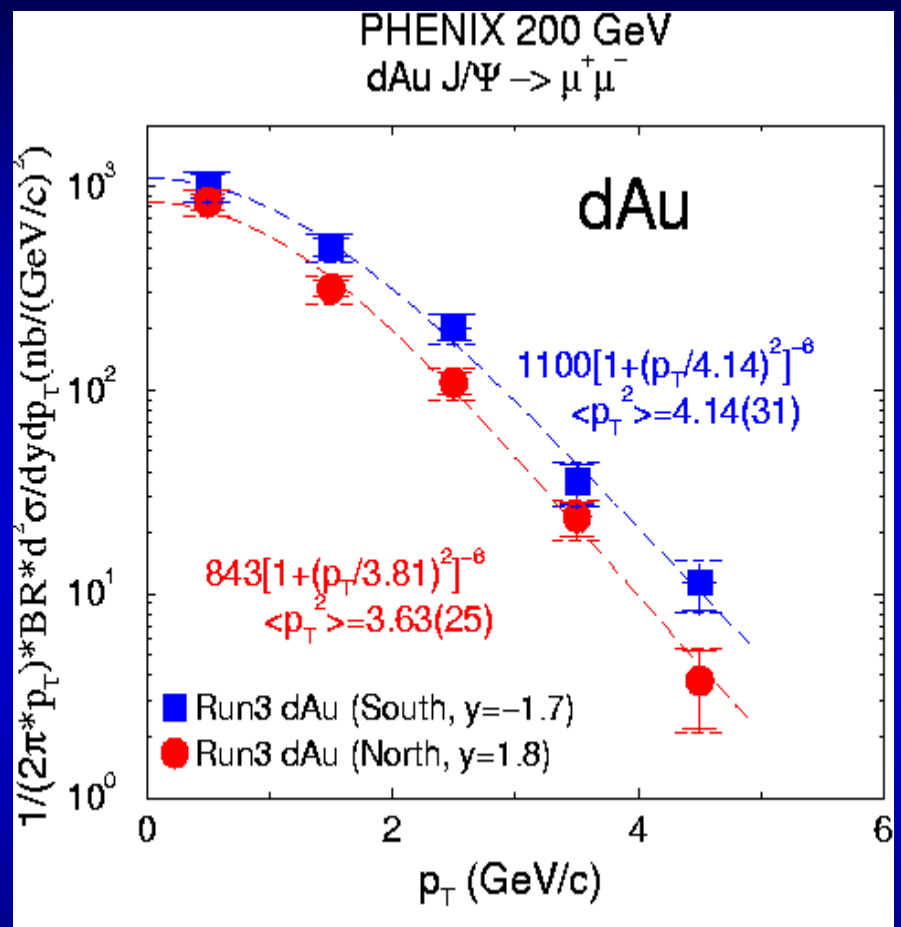
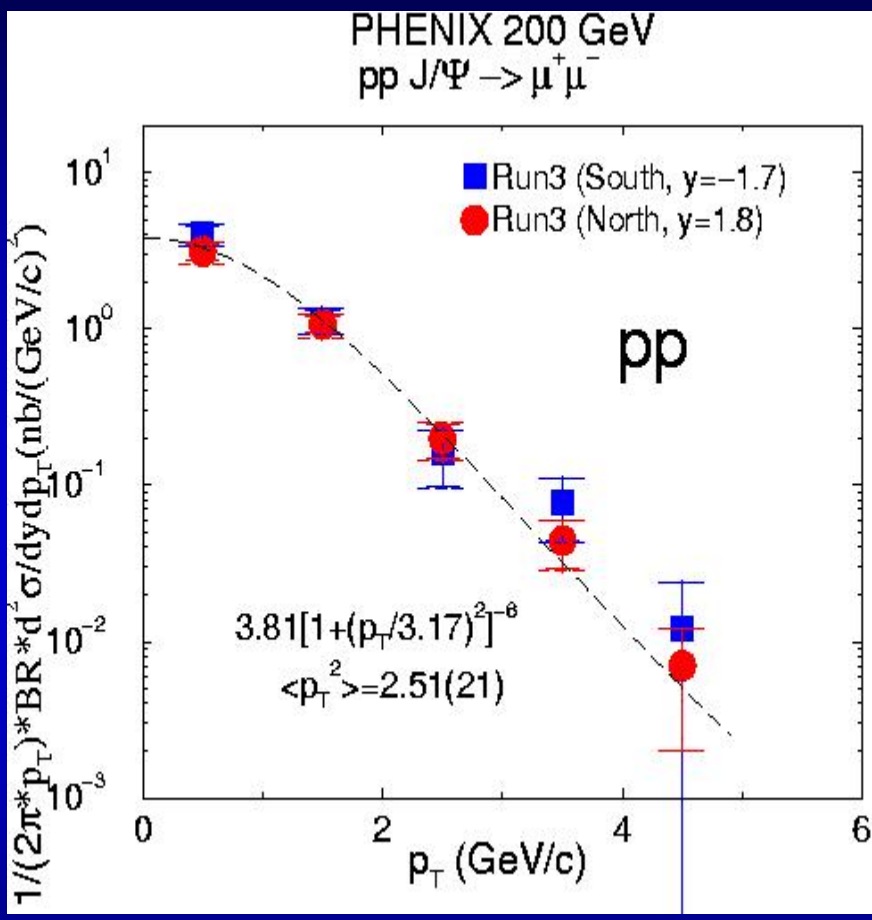
Run	Year	Species	$s^{1/2}$ [GeV]	$\int Ldt$	N_{Tot}	p-p Equivalent	Data Size
01	2000	Au+Au	130	$1 \mu b^{-1}$	10M	$0.04 pb^{-1}$	3 TB
02	2001/2002	Au+Au	200	$24 \mu b^{-1}$	170M	$1.0 pb^{-1}$	10 TB
		p+p	200	$0.15 pb^{-1}$	3.7G	$0.15 pb^{-1}$	20 TB
03	2002/2003	d+Au	200	$2.74 nb^{-1}$	5.5G	$1.1 pb^{-1}$	46 TB
		p+p	200	$0.35 pb^{-1}$	6.6G	$0.35 pb^{-1}$	35 TB
04	2003/2004	Au+Au	200	$241 \mu b^{-1}$	1.5G	$10.0 pb^{-1}$	270 TB
		Au+Au	62	$9 \mu b^{-1}$	58M	$0.36 pb^{-1}$	10 TB
05	2004/2005	Cu+Cu	200	$3 nb^{-1}$	8.6G	$11.9 pb^{-1}$	173 TB
		Cu+Cu	62	$0.19 nb^{-1}$	0.4G	$0.8 pb^{-1}$	48 TB
		Cu+Cu	22.5	$2.7 \mu b^{-1}$	9M	$0.01 pb^{-1}$	1 TB
		p+p	200	$3.8 pb^{-1}$	85B	$3.8 pb^{-1}$	262 TB



Fields



Cross sections vs p_T (p+p, d+Au)

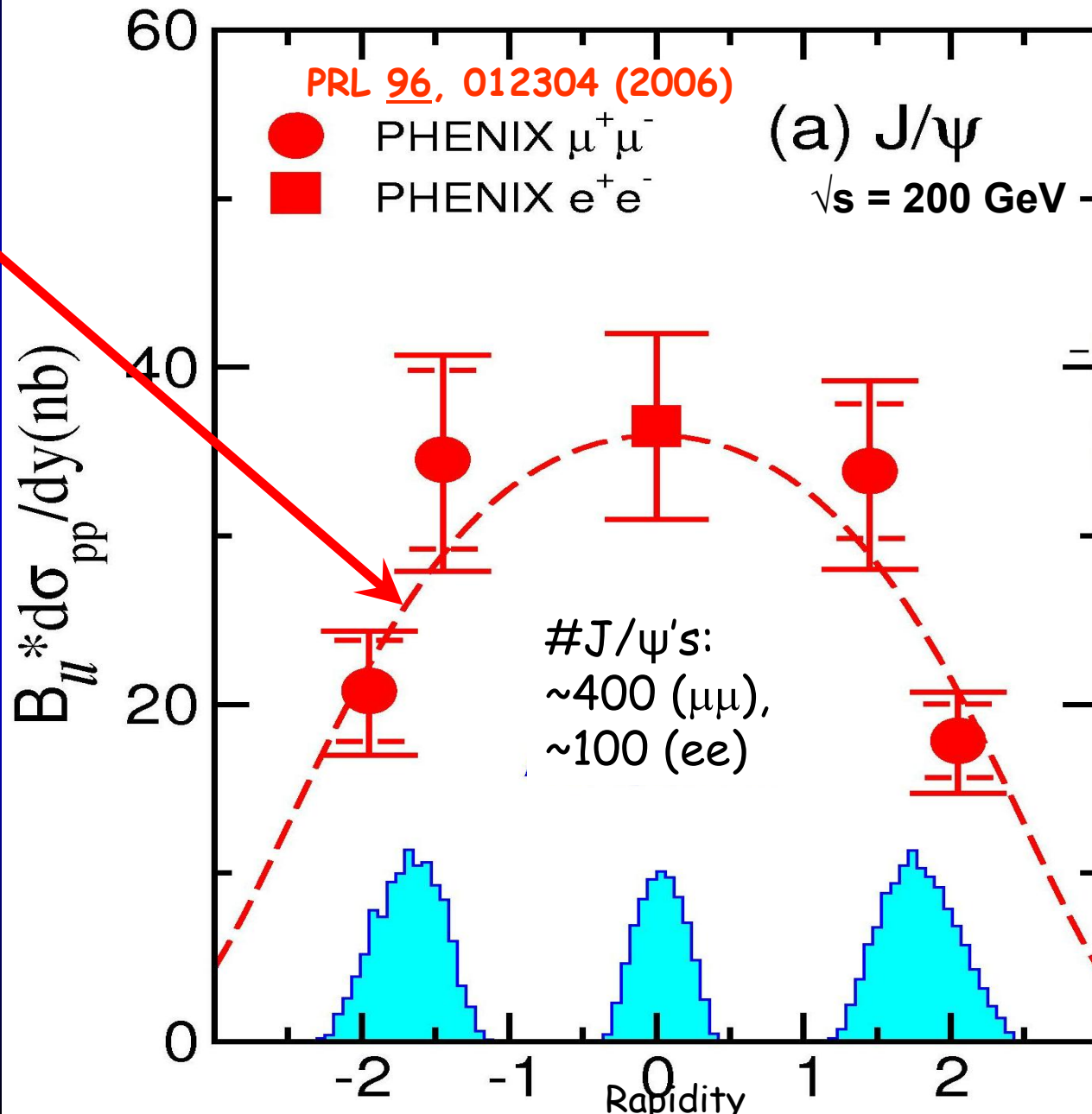


Fits to $\frac{dn}{p_T dp_T} \sim A[1+(p_T/B)^2]^{-6}$ used to extract $\langle p_T \rangle^2$



Cross section vs Rapidity (p+p)

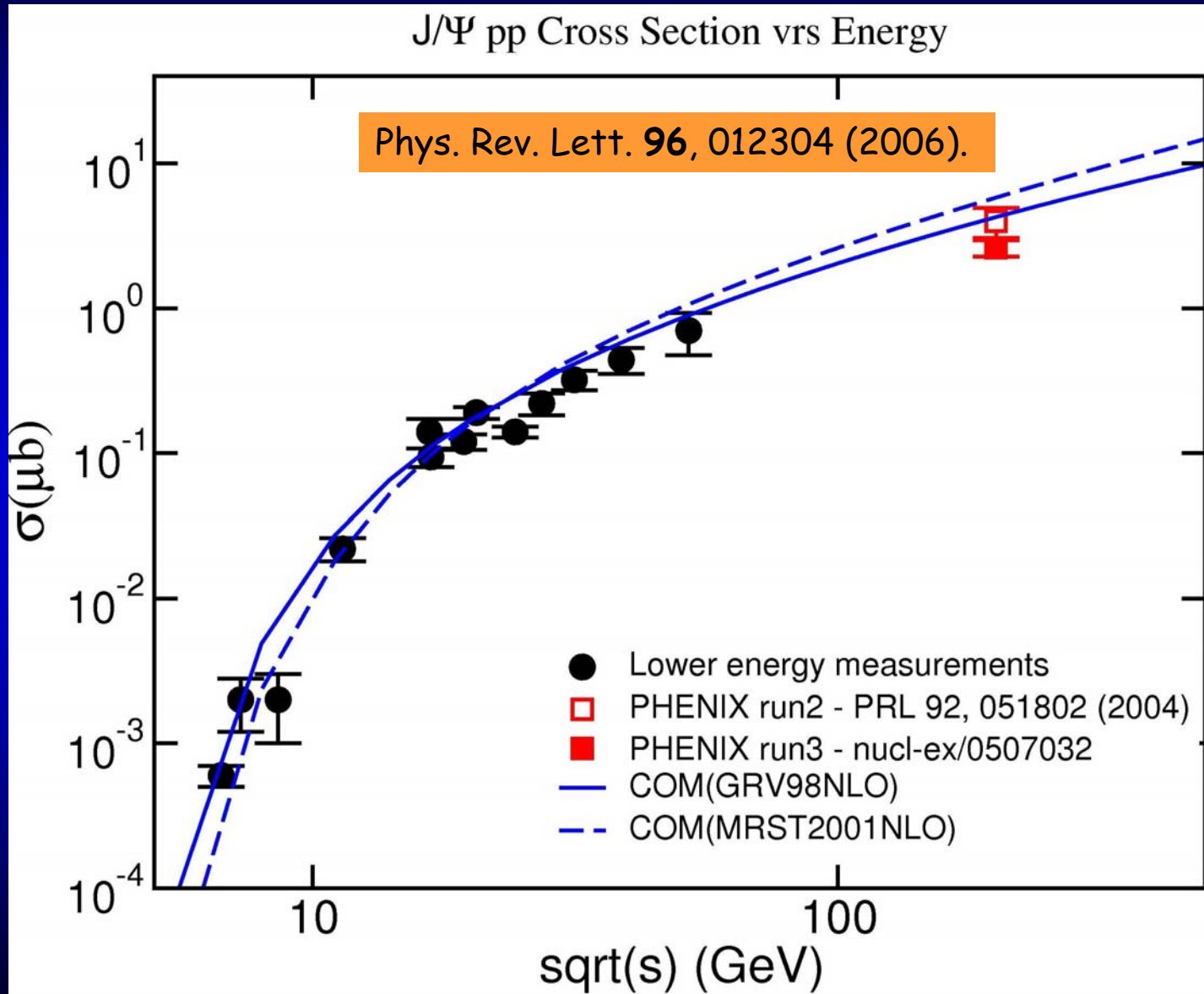
- Good agreement with PYTHIA shape
- These are Run-3 data
- Run-5: $\sim \times 10$
- Run-6: \sim Run-5 today, expected to double





p+p Reference

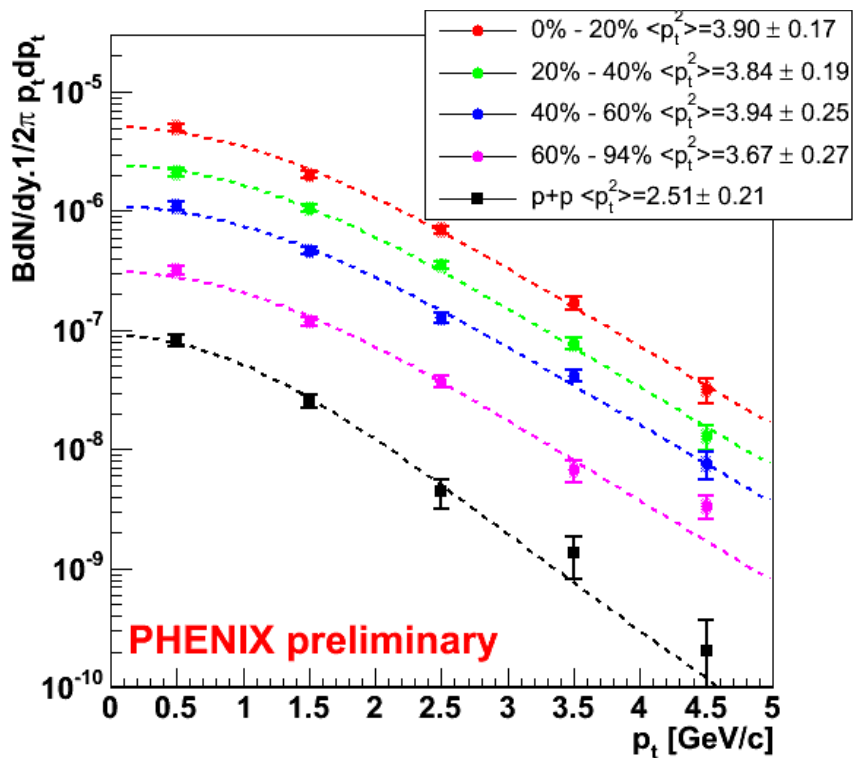
- ⇒ Consistent with trend of world's data
- ⇒ ~Consistent with at least one COM (Color Octet Model) calculation



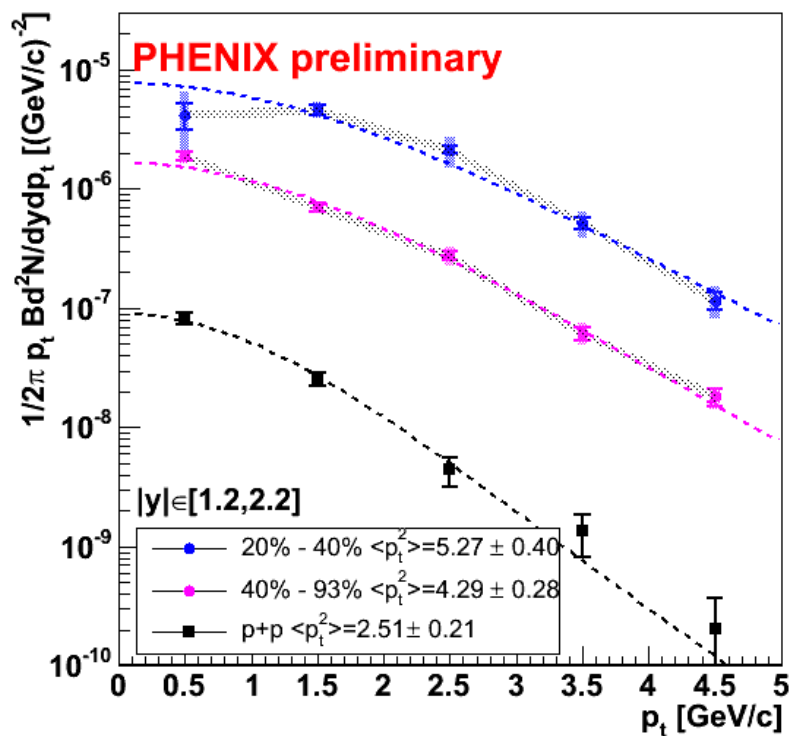


Invariant yields vs p_T (A+A)

Cu+Cu ($|y| \in [1.2, 2.2]$)



Au+Au ($|y| \in [1.2, 2.2]$)



Fits to $\frac{dn}{p_T dp_T} \sim A[1 + (p_T/B)^2]^{-6}$ used to extract $\langle p_T \rangle^2$



Ratios

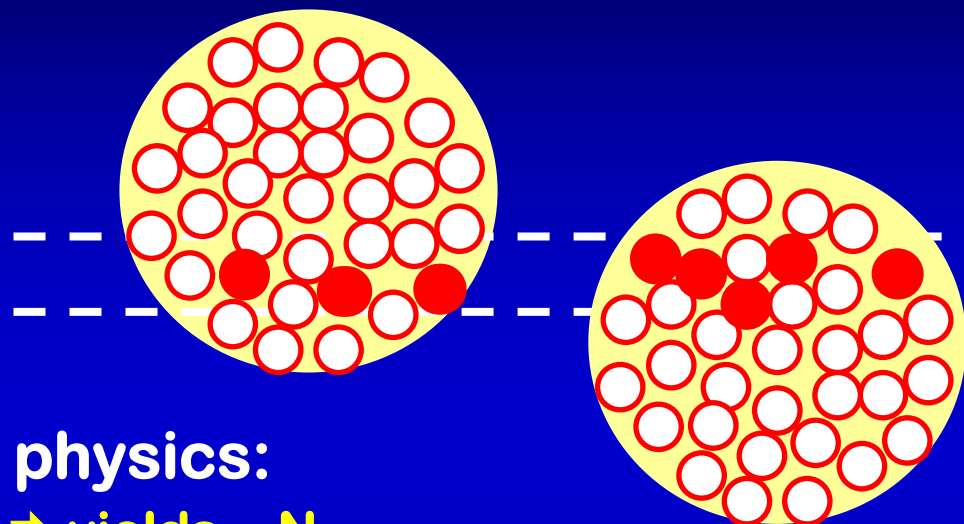
Basis for Comparison

Focus on some slice of collision:

- Assume 3 nucleons struck in A, and 5 in B
- Could weight contribution as

◆ $N_{\text{part}} \equiv 3 + 5$

◆ $N_{\text{coll}} \equiv 3 \times 5$



Choice driven by underlying physics:

- Low $p_T \Rightarrow$ large cross sections \Rightarrow yields $\sim N_{\text{part}}$
 - ◆ Soft, non-perturbative, “wounded nucleons”, ...
- High $p_T \Rightarrow$ small cross sections \Rightarrow yields $\sim N_{\text{coll}}$
 - ◆ Hard, perturbative, “binary scaling”, point-like, $A \cdot B$, ...

J/Ψ 's are complicated ...

- (This a general theme of what follows)
- Will use N_{coll} to form ratios R_{dA} and R_{AA}
- Will present same as functions of both N_{part} and N_{coll}

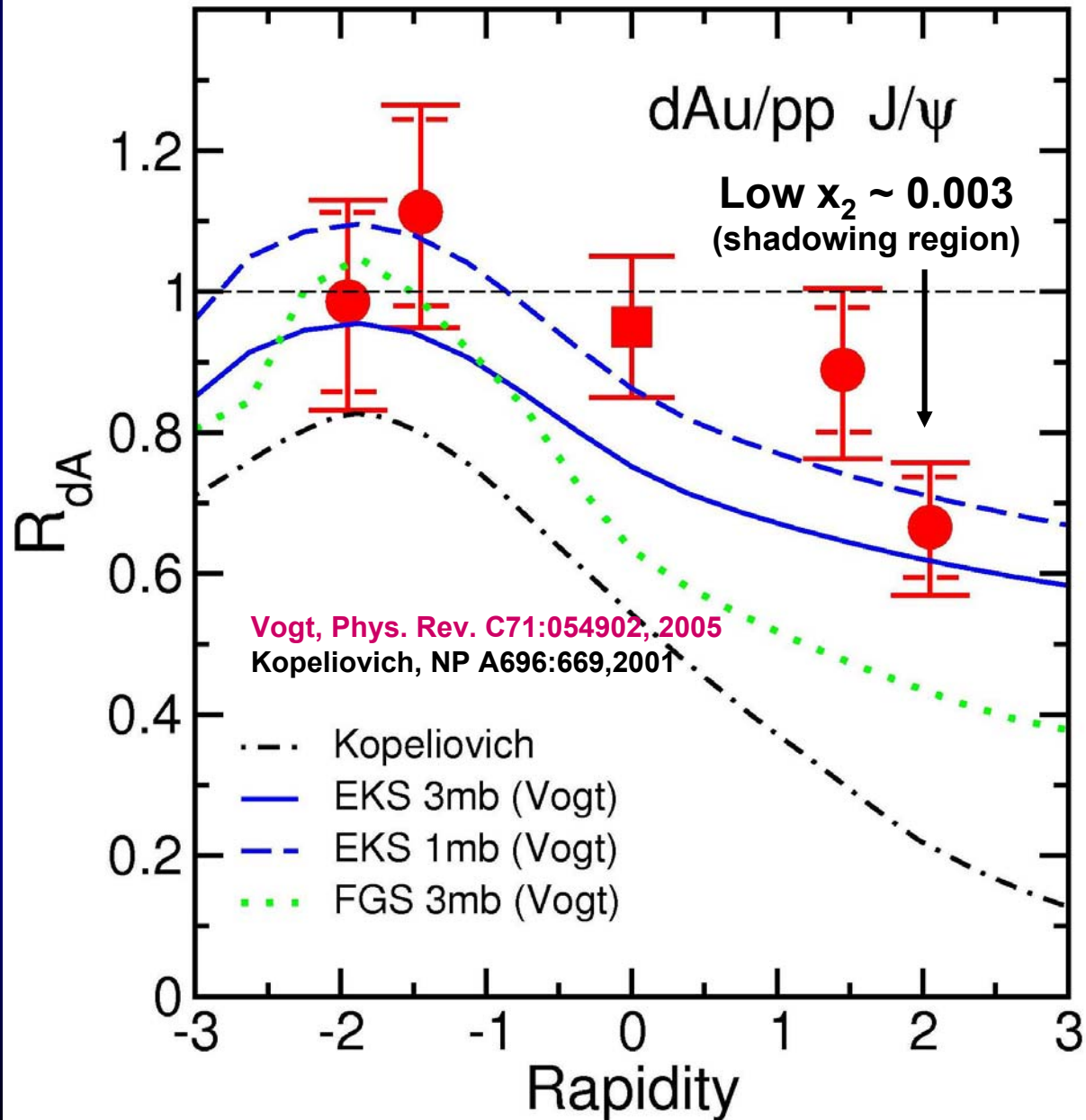


Cold Nuclear Matter (CNM)

Modest nuclear effects as a function of rapidity

- Suppression for low $x_2 \sim 0.003$ (x_{Au})
- No suppression for higher $x_2 \sim 0.02-0.09$

Modest agreement with models incorporating shadowing and absorption





Hot Nuclear Matter

- Good agreement between

- Au+Au (muons, red)
- Cu+Cu (muons, blue)

in overlap region

- Central arm data

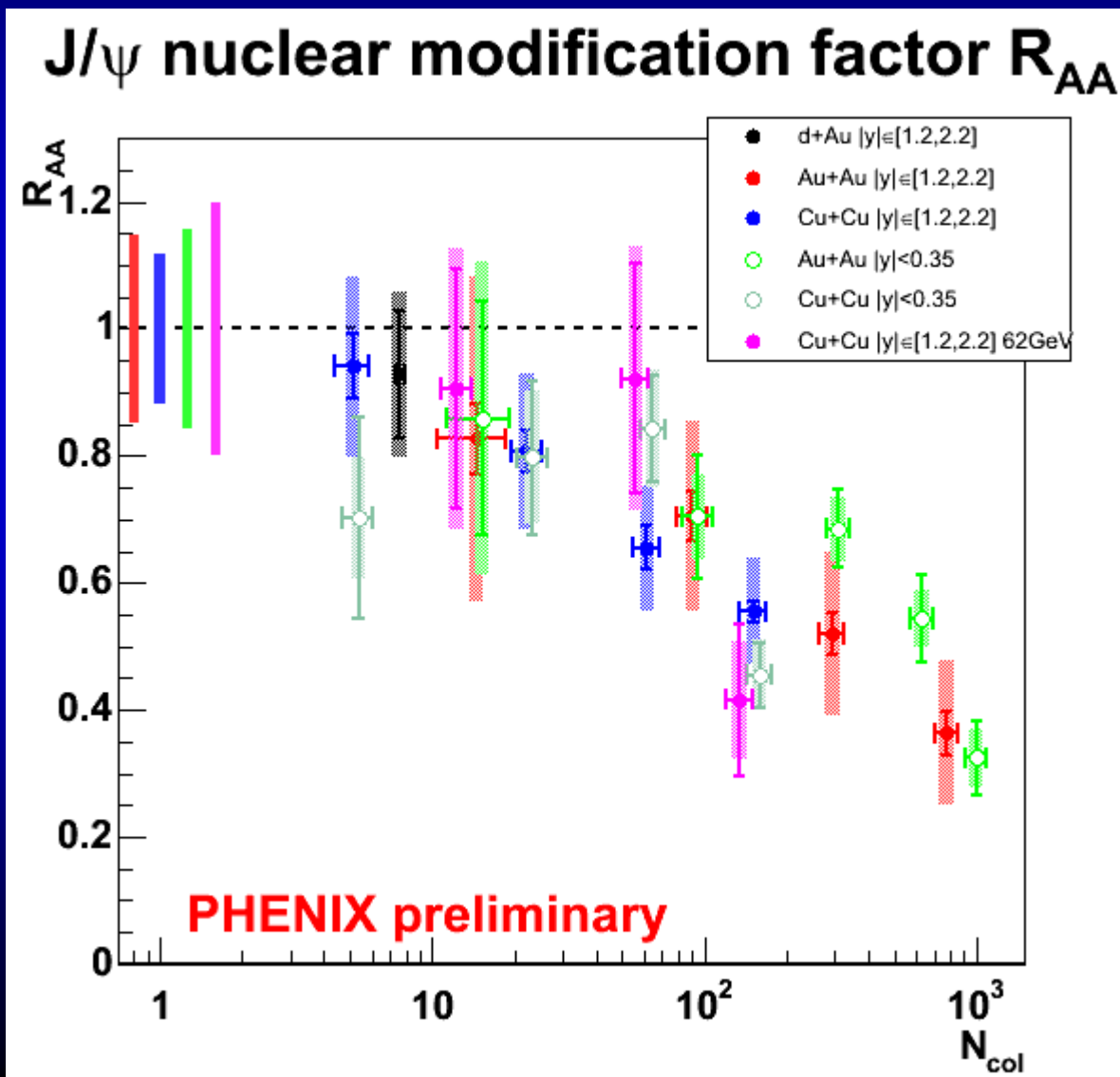
- Au+Au (green)

slightly higher(?)
than muon arms

- Note different
rapidity intervals

- ◆ Central $|y| < 0.35$
- ◆ Muons $|y| \in [1.2, 2.2]$

- “Large” suppression
effects in A+A

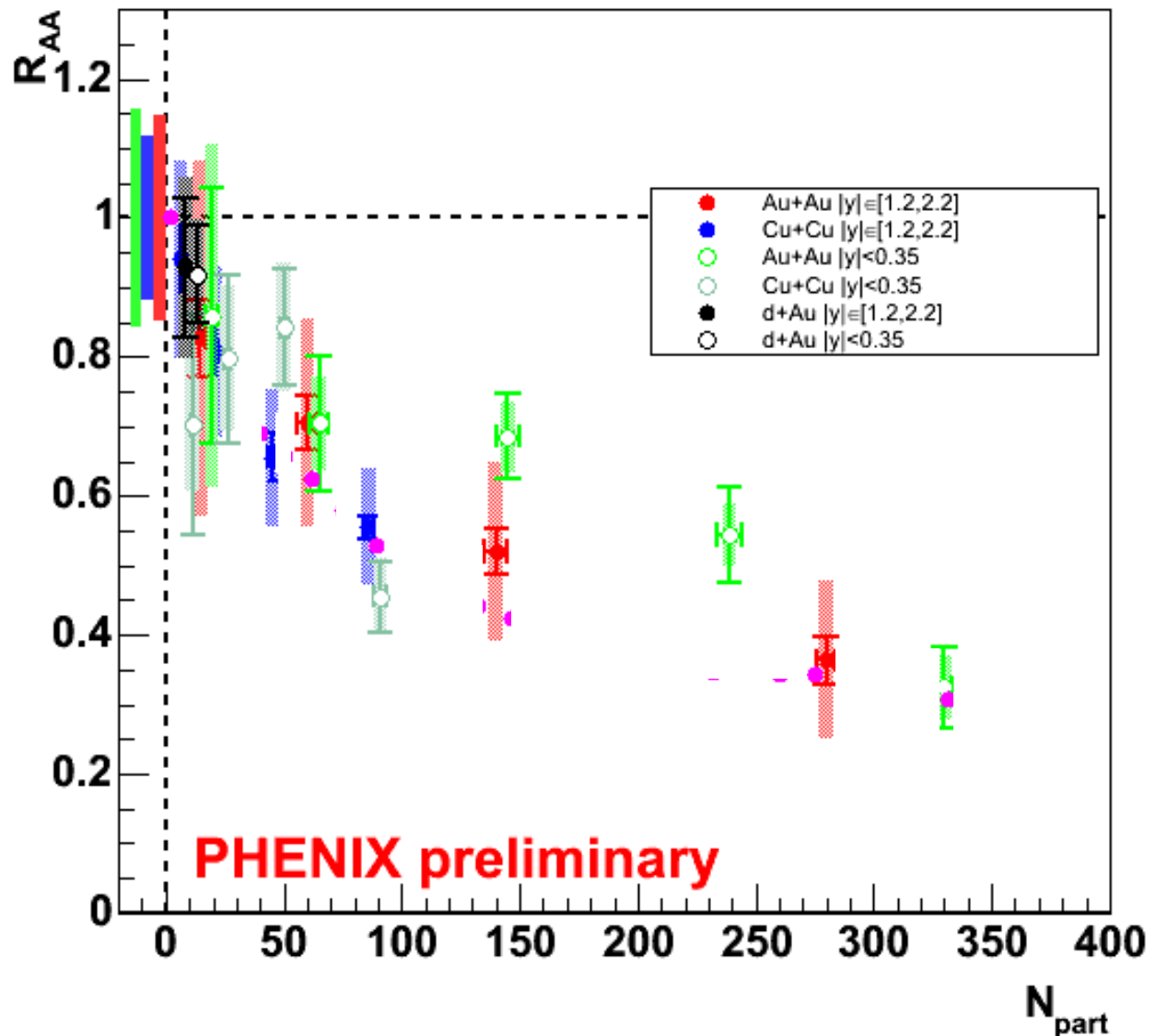




Large Compared to What?

- “Large”, of course, with respect to the d+A reference data
- Comparable to the “large” suppression observed by NA50
- (Note change of horizontal axis to N_{part})

J/ψ nuclear modification factor R_{AA}

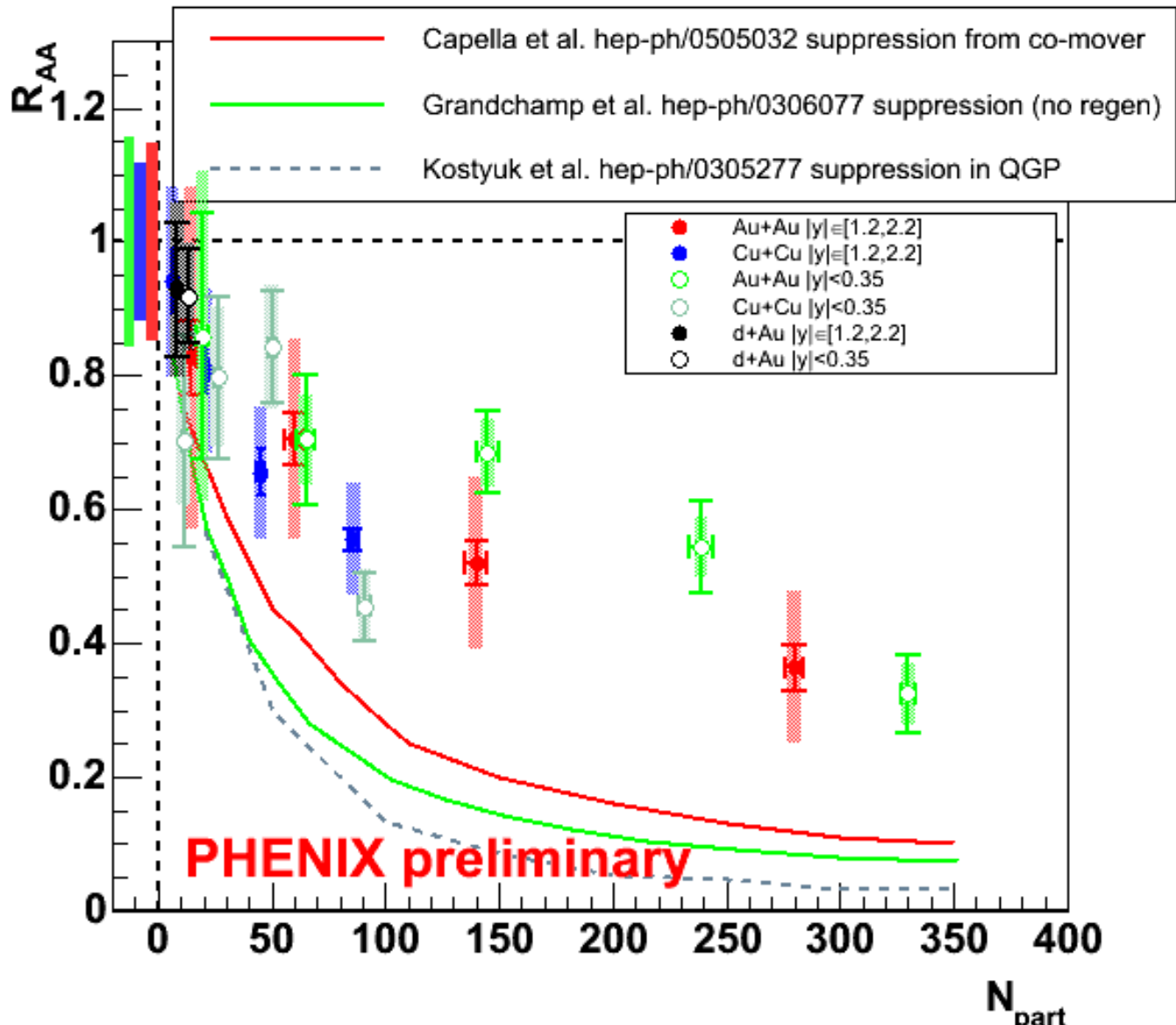




But Not Truly "Large"

- That is:
Much smaller suppression effects than expected based on scaling from SPS to RHIC energies
- Without "new" mechanisms

J/ψ nuclear modification factor R_{AA}





Models



Modeling "normal" suppression

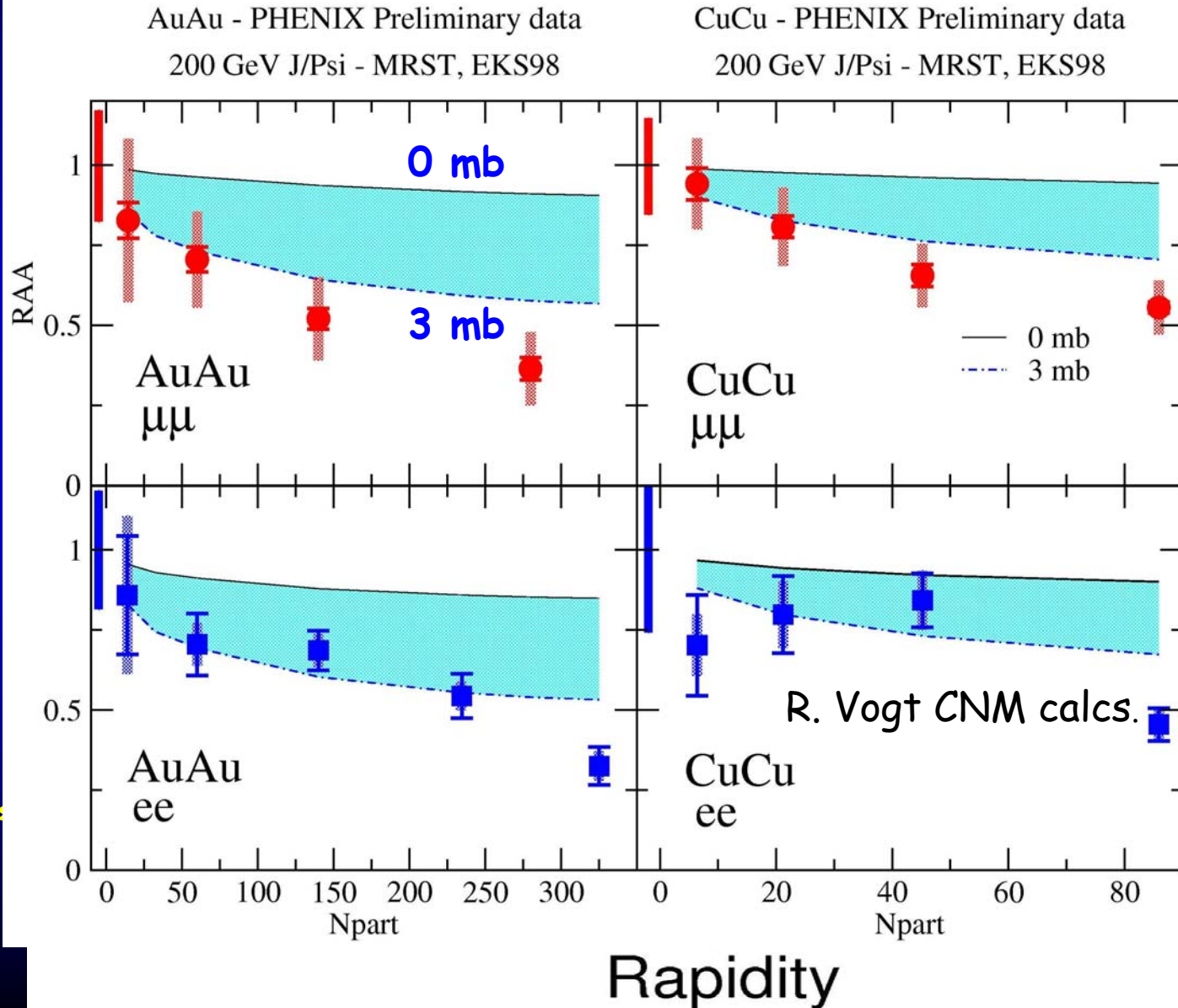
Input:

- "EKS" shadowing
- Absorption cross-section σ_{abs}

Output:

- Only modest constraints on value of σ_{abs}
- Improvement needed as "background" other mechanisms

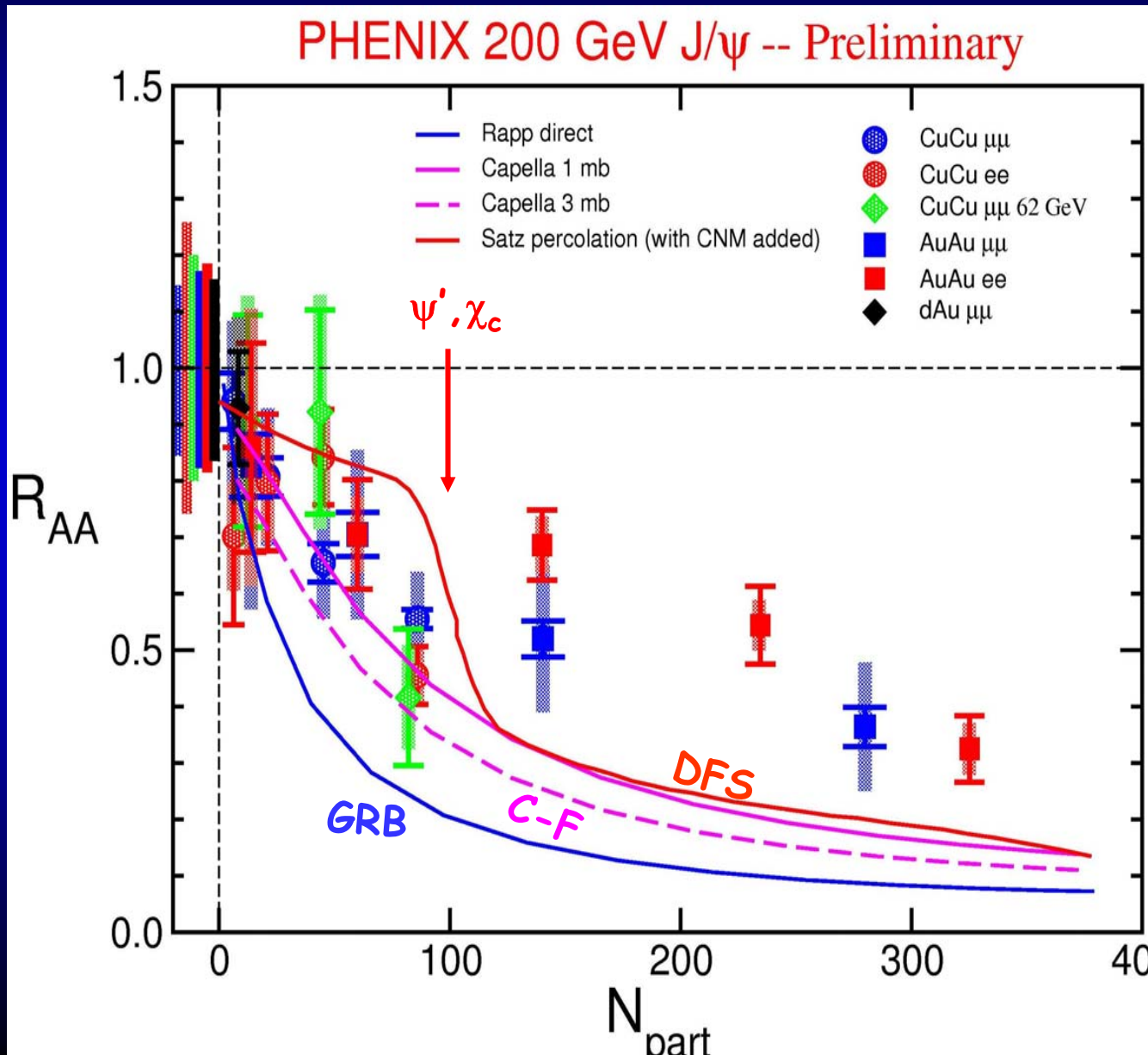
200 GeV d+Au -> J/Psi





“Normal” Anomalous Suppression

- Grandchamp, Rapp, Brown; [PRL 92, 212301 \(2003\)](#)
 - In-media dissolution
- Cappella, Ferreira; [Eur. Phys. J C42, 419 \(2005\)](#)
 - Absorption, shadowing, comovers
- Digal, Furtunato, Satz; [Eur. Phys. J C32, 547 \(2004\)](#)
 - Absorption, shadowing, percolation, color screening

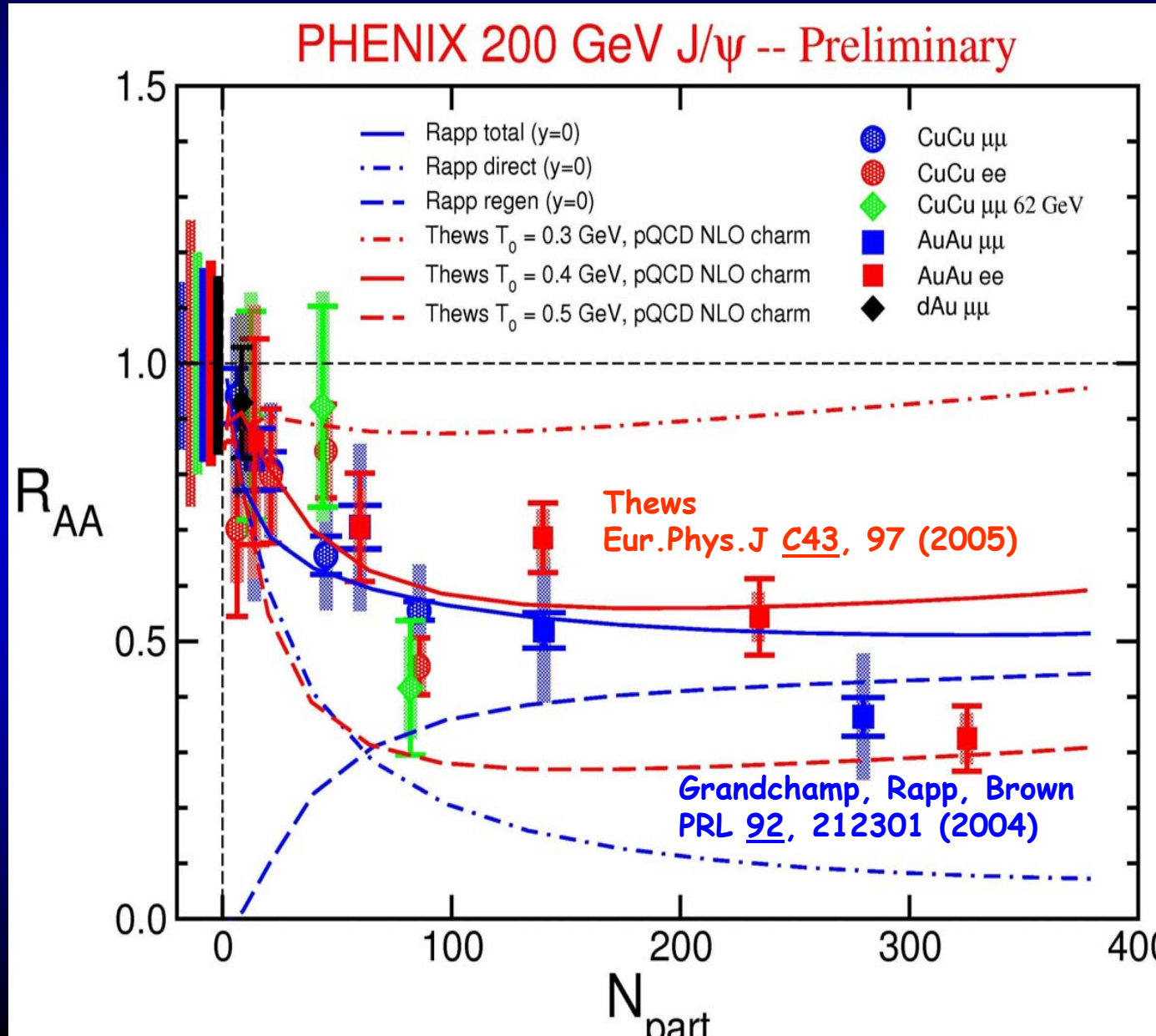




“Anomalous” Anomalous Suppression

- That is, regeneration.
- One example: Grandchamp, Rapp, Brown; [PRL 92, 212301 \(2003\)](#)

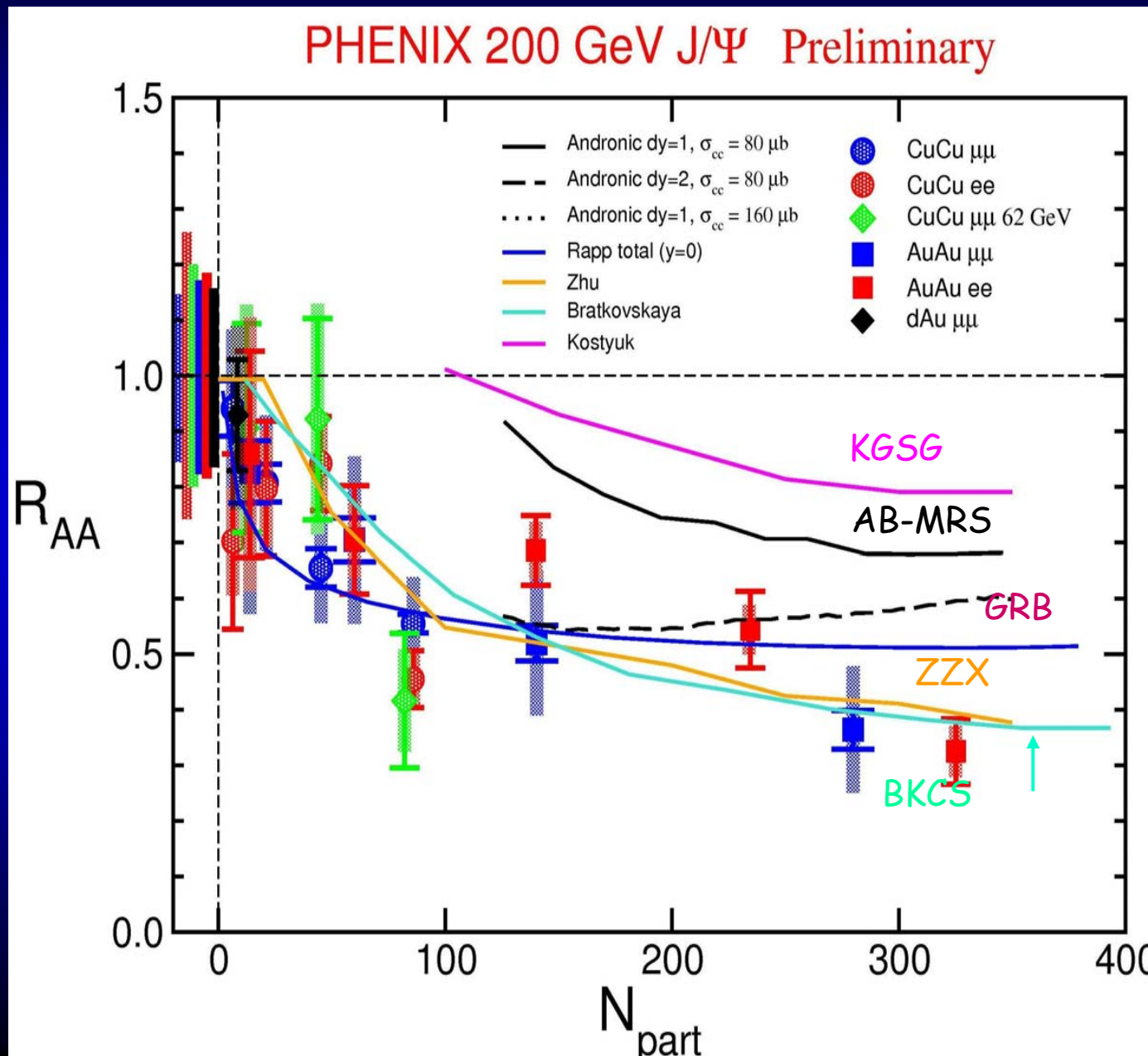
□ In-media dissolution
 ➔ Plus regeneration from “off-diagonal” c-cbar pairs





No Suppression of Models...

- GRB – see previous slide
- Thews- previous slide, next talk
- Andronic, Braun-Munzinger, Redlich, Stachel; [Phys. Lett. B571, 36 \(2003\)](#)
 - Color screening, recombination
- Bratkovskaya, Kostyuk, Cassing Stoecker; [Phys. Rev. C69, 054903 \(2004\)](#)
 - Hadron-string dynamics
- Kostyuk, Gorenstein, Stoecker, Greiner; [Phys. Rev. C68, 041902 \(2003\)](#)
 - Color screening, recombination
- Zhu, Zhuang, Xu; [Phys. Lett. B607, 107 \(2005\)](#)
 - Co-movers, gluon break-up, no regeneration

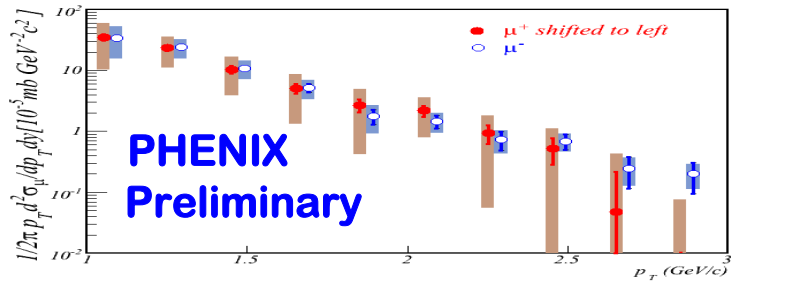
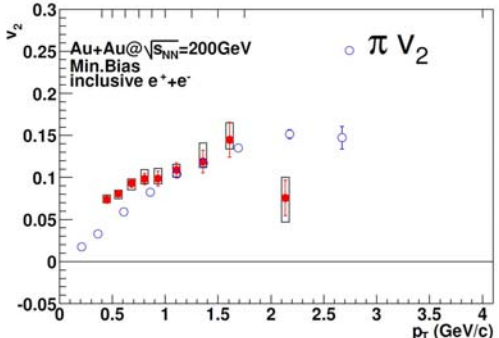
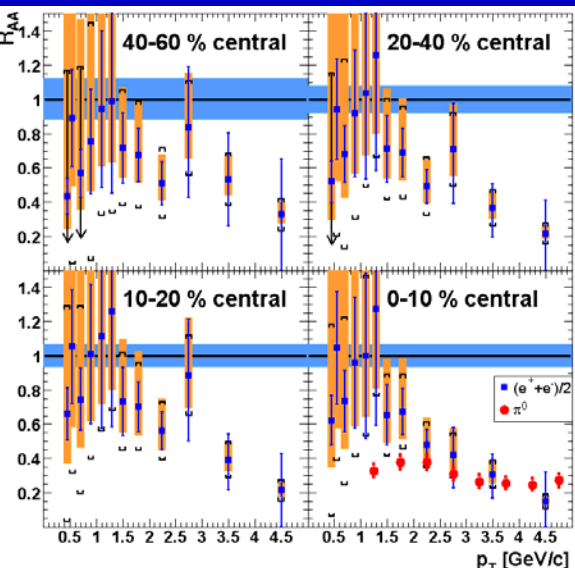
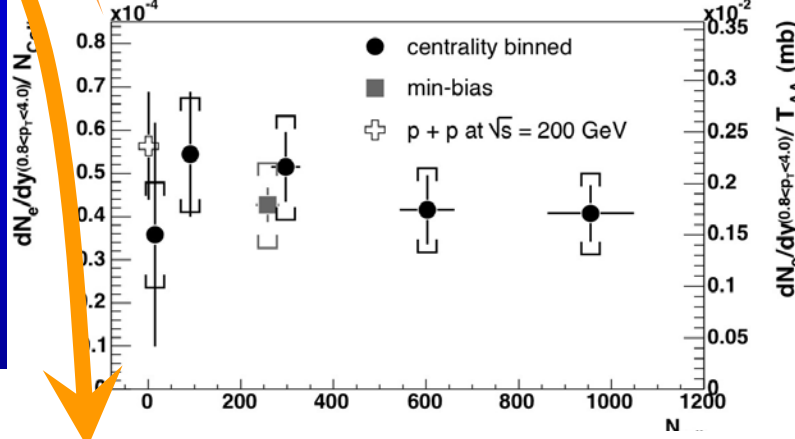
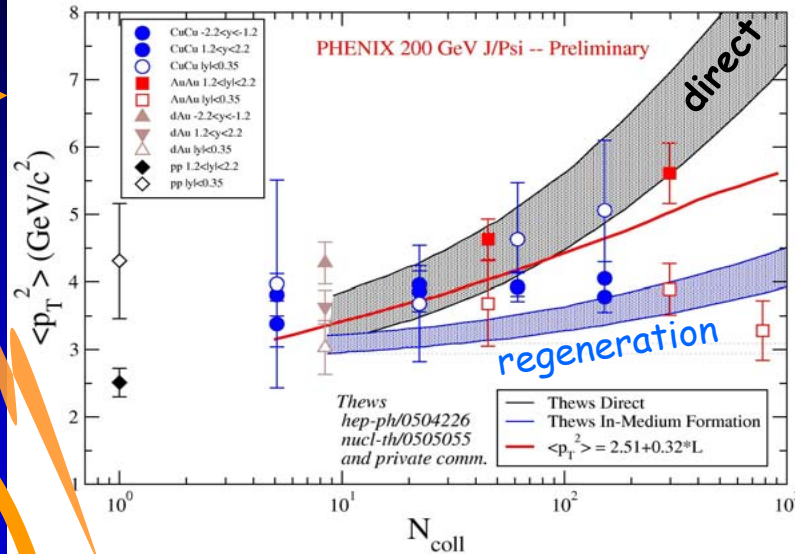




There Is So Much More To The Story

- (A Sampling)

- $\langle p_T^2 \rangle$ behavior
- Open charm yields
[Phys. Rev. Lett. 94, 082301 \(2005\).](#)
- Open charm distributions
- Charm energy loss
[Phys. Rev. Lett. 96, 032301 \(2006\).](#)
- Charm flow
[Phys. Rev. C72, 024901 \(2005\)](#)

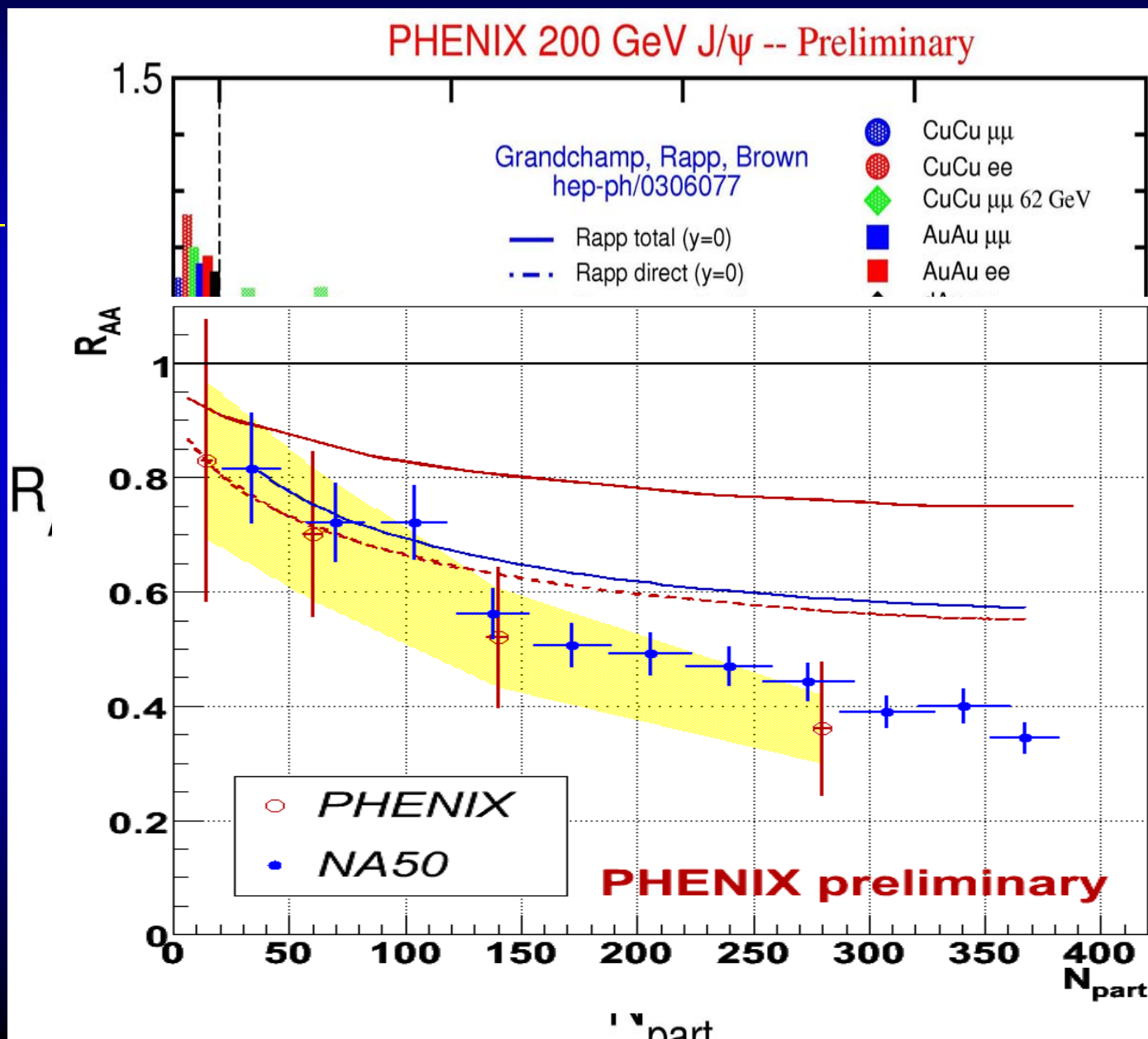




Not Suppressing the J/ψ (!)

- Karsch, Kharzeev, Satz;
hep-ph/0512239

- Based on LQCD results suggesting $T_{J/\psi} \sim 2 T_C$
- Suppression (only) of Ψ' and χ_c
- See talk(s) to follow



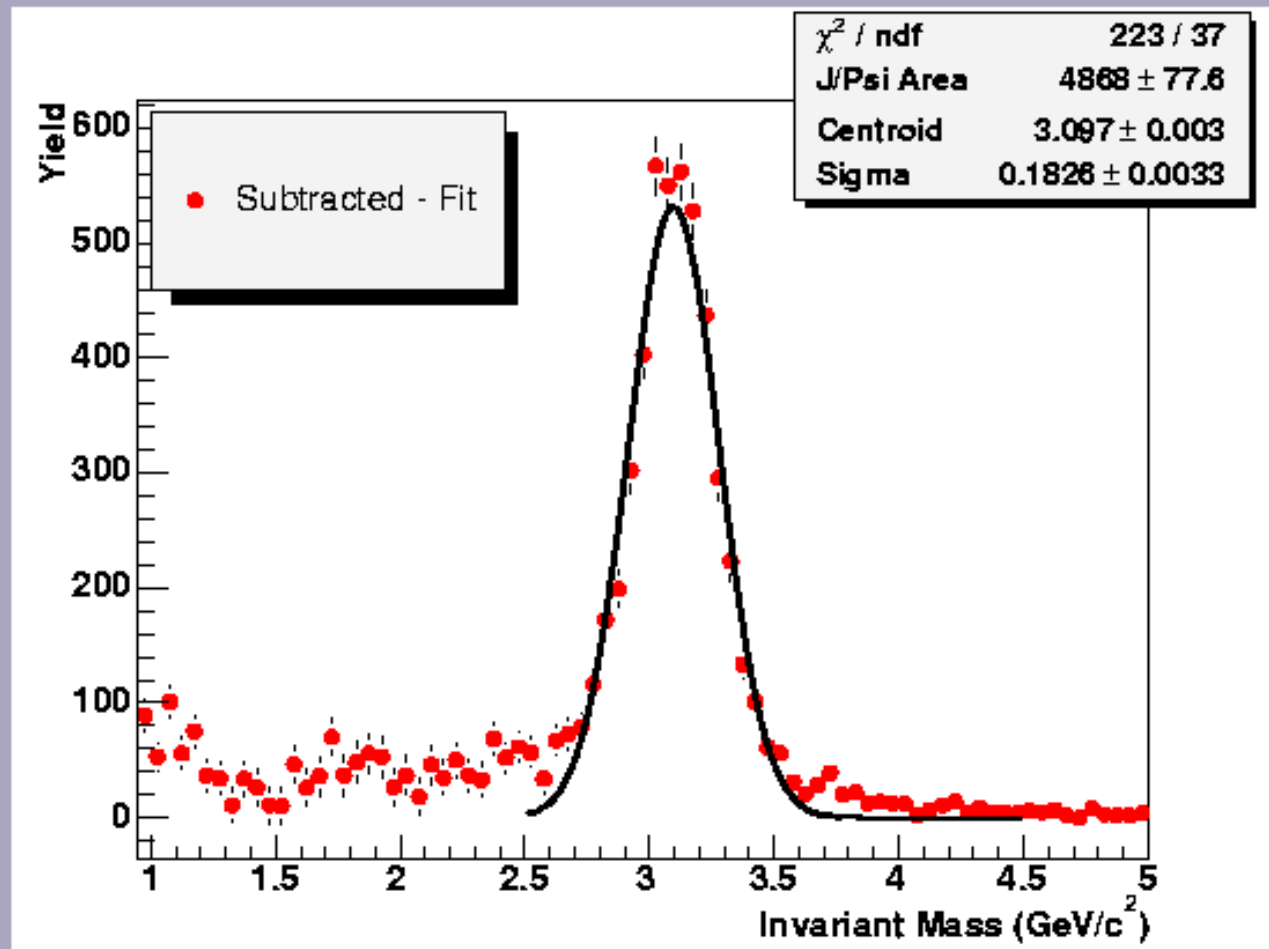


Future

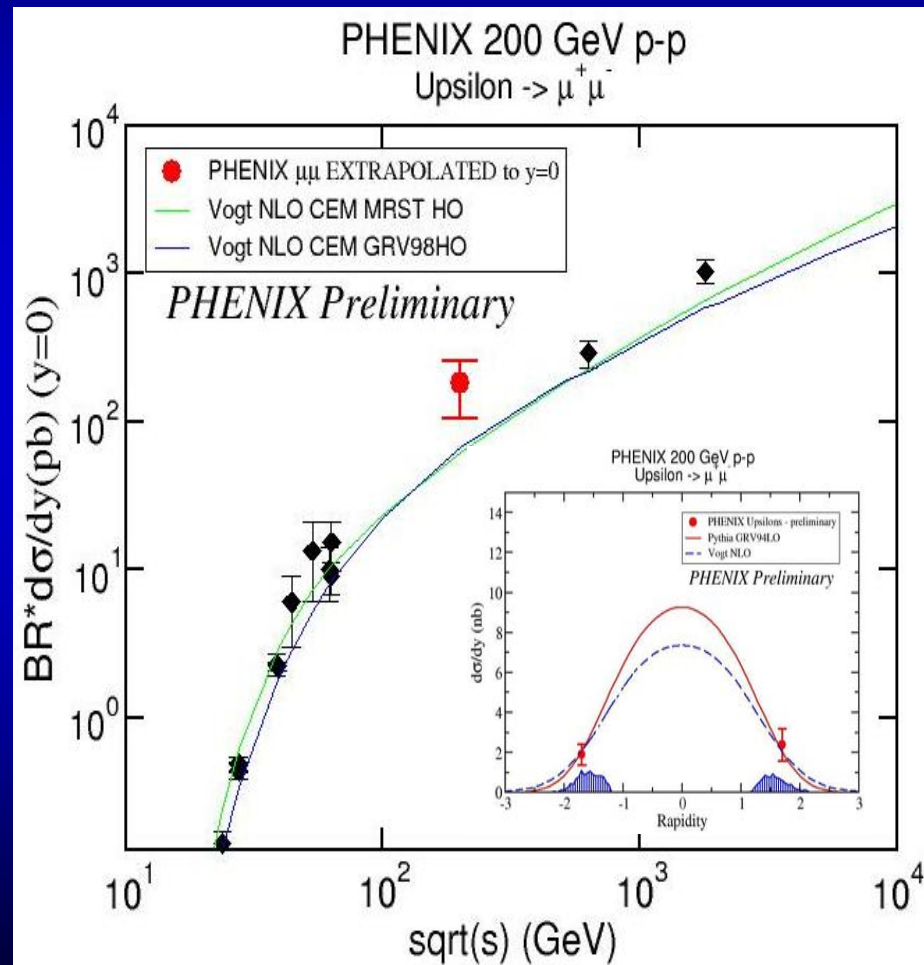
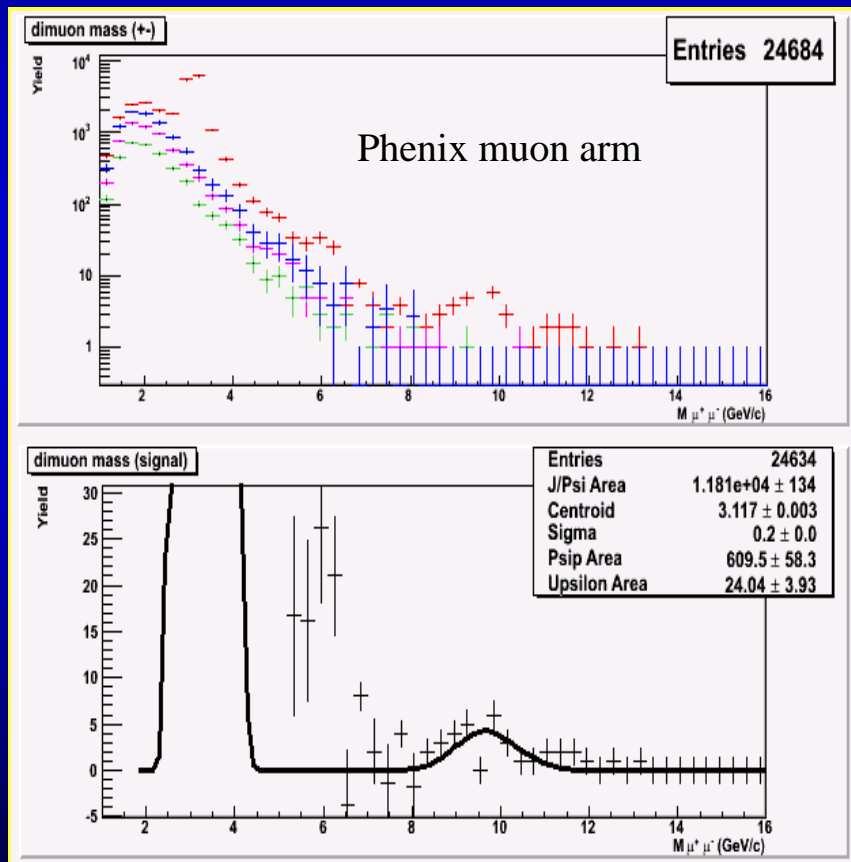


Run-6 (as of today)

- 200 GeV polarized p+p collisions
- Results from ~ real-time production of Level-2 triggers
- Already equals or exceeds all of Run-5
- Expect ~ 3x this by end of proton running (~ 10 pb⁻¹)



- From Run-5 p+p: $\sim 3 \text{ pb}^{-1}$
- Current Run-6 : $\sim 4 \text{ pb}^{-1}$

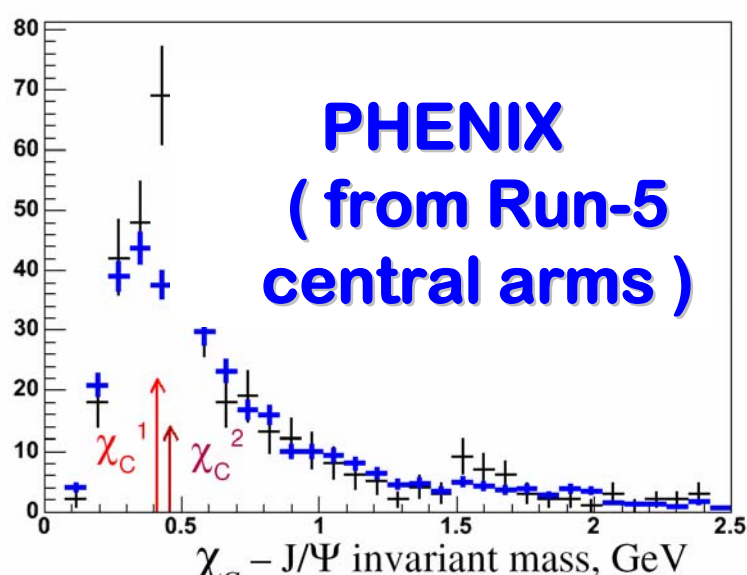
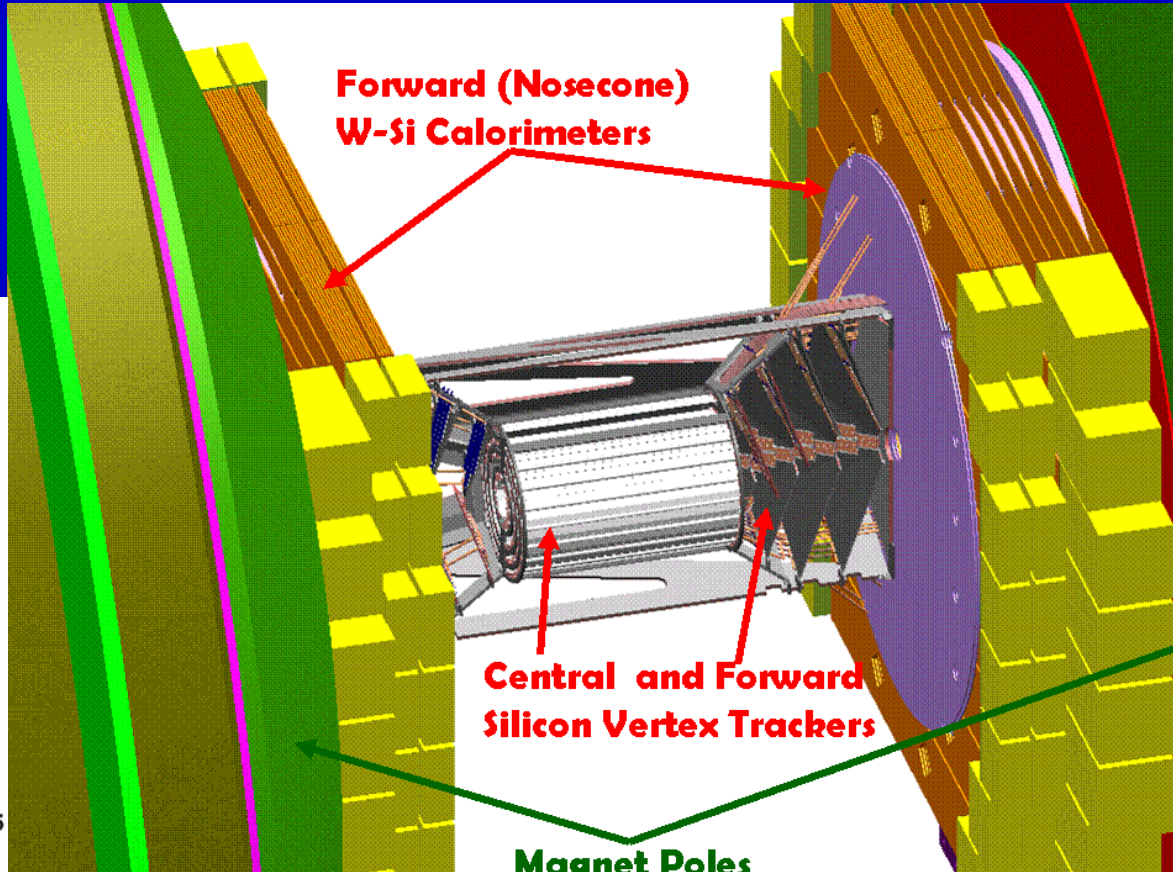


First Upsilon at RHIC



Future: Nose Cone Calorimeter

- Replace central arm magnet nosecones (Cu) w/ tungsten-silicon calorimeters
- Coverage at forward/backward rapidity: $0.9 < |\eta| < 3.5$
 - γ/π^0 separation for $p_T < 30$ GeV/c
 - Jet identification
- γ identification gives good acceptance for $\chi_c \rightarrow J/\Psi + \gamma$





RHIC-II - Heavy Flavor Yields

All numbers are first rough estimates
(including trigger and reconstruction efficiencies)
for 12 weeks Au+Au run ($\int L_{\text{eff}} dt \sim 18 \text{ nb}^{-1}$)

Signal	RHIC Exp.	Obtained	RHIC I (>2008)	RHIC II	LHC/ALICE ⁺
$J/\psi \rightarrow e^+e^-$	PHENIX	~800	3,300	45,000	9,500
$J/\psi \rightarrow \mu^+\mu^-$		~7000	29,000	395,000	740,000
$\gamma \rightarrow e^+e^-$	STAR	-	830	11,200	2,600
$\gamma \rightarrow \mu^+\mu^-$	PHENIX	-	80	1,040	8,400
$B \rightarrow J/\psi \rightarrow e^+e^-$	PHENIX	-	40	570	N/A
$B \rightarrow J/\psi \rightarrow \mu^+\mu^-$		-	420	5,700	N/A
$\chi_c \rightarrow e^+e^- \gamma$	PHENIX	-	220	2,900*	N/A
$\chi_c \rightarrow \mu^+\mu^- \gamma$		-	8,600	117,000*	N/A
$D \rightarrow K\pi$	STAR	$\sim 0.4 \times 10^6$ (S/B $\sim 1/600$)	30,000**	30,000**	8000

* Large backgrounds, quality uncertain as yet

** Running at 100 Hz min bias

+ 1 month (= year), P. Crochet, EPJdirect A1, a (2005) and private comm.

T. Frawley, PANIC'05,
RHIC-II Satellite Meeting



Helmut

Lepton pair production in hadron collisions

K. Kinoshita,* H. Satz, and D. Schildknecht

Department of Theoretical Physics, University of Bielefeld,
(Received 12 October 1977)

We construct a model for lepton pair production in the low-mass region and compare its predictions with the production characteristics of inclusive hadron production at small p_T . This model is then applied to the pair data up to 3.1 GeV. We then apply our model at higher masses as a pointlike-parton annihilation.

I. INTRODUCTION

The aim of this work is to examine a hadronlike model for lepton pair production in hadron-hadron collisions. By "hadronlike" we mean that the functional behavior of the transverse- and longitudinal-momentum spectra, their dependence on the incident energy, and all relevant parameters of these distributions, are taken from low- p_T hadron physics.

also show that

$$f(E_T) = e^{-\lambda E_T}$$

with $\lambda = 6 \text{ GeV}^{-1}$ quite well.

For the p_T like production function $f(E_T)$ from

$$\left. \frac{d\sigma}{dp_T^2} \right|_{y=0}$$

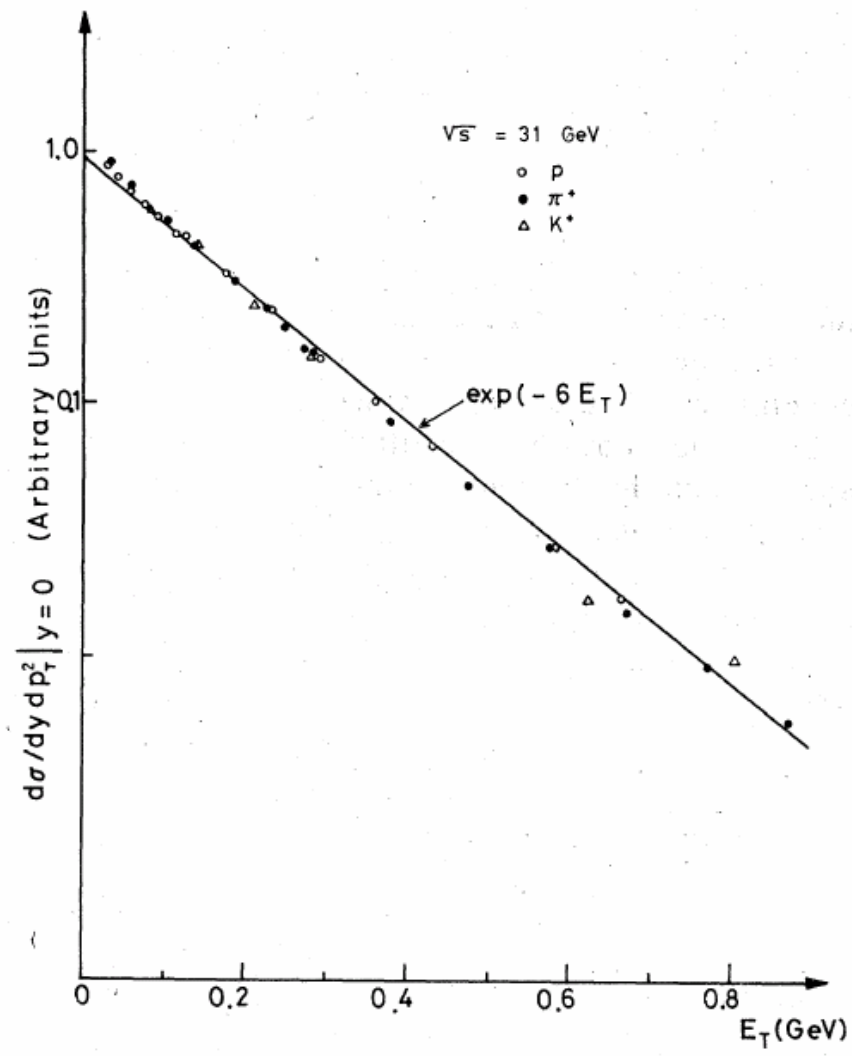


FIG. 1. Data (Ref. 3) on $pp \rightarrow (p, \pi^+, K^+) + X$ at 90° compared with $\exp(-6E_T)$.

• First "m_T scaling" plot? →



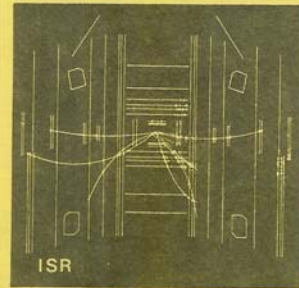
My First Helmut Sighting

- *Deconfinement in Dense Matter*, by H. Satz
- Abstract:
 - We survey recent developments in the study of strongly interacting matter by means of quantum chromodynamics on the lattice. The Monte Carlo evaluation of SU(2) Yang-Mills systems is shown to provide a unified thermodynamic description from the hadronic region through the deconfinement transition to a non-interacting gluon gas. We ... discuss possible signals for the formation of a quark-gluon plasma in relativistic heavy ion collisions.

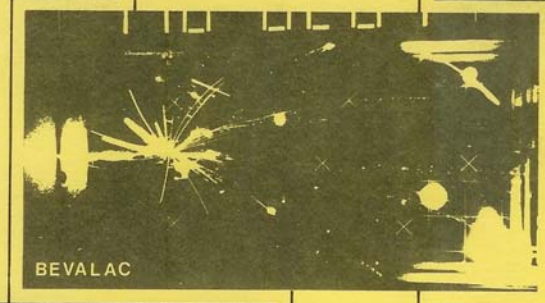
LBL-12652
UC-34
CONF-8105104

5TH HIGH ENERGY HEAVY ION STUDY MAY 18-22, 1981

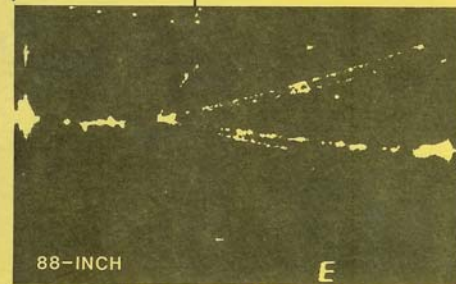
PROCEEDINGS



ISR



BEVALAC



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LAWRENCE BERKELEY LABORATORY
UNIVERSITY OF CALIFORNIA
BERKELEY, CALIFORNIA 94720

PREPARED FOR THE U.S. DEPARTMENT OF ENERGY UNDER CONTRACT W-7405-ENG-48

OCTOBER 1981



A Furtive Note-Taker...

TALK SATZ

Phase Transition From Nuclear Matter to Quark Matter

H. Satz

1. High temperature - free q 's and g 's

$$\epsilon(T) = \frac{\pi^2}{30} T^4 \left\{ 2 N_c^6 + \frac{7}{2} N_c^2 N_f \right\} \quad (\text{Stefan-Boltzmann})$$

$$p = \frac{1}{3} \epsilon \quad c_v(T) = \left(\frac{\partial \epsilon}{\partial T} \right)_V$$

2. Hadronic resonance gas

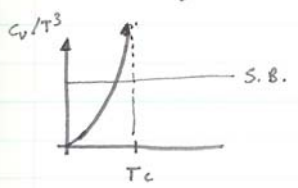
$$\tau(m) = d \delta(m - m_0) + c \delta(m - 2m_0) m^{-a} e^{-b m} \quad m_0 = 1$$

$$\text{Then } \epsilon(T) = \epsilon_0(T) + \frac{c T^{3/2}}{(2\pi)^{3/2}} \int_{m_0}^{\infty} dm m^{-a + \frac{3}{2}} e^{-m(4-b)}$$

\downarrow ($\sim T^4$)

$$\text{and } c_v(T) = c_v^0(T) + \frac{3(\epsilon - \epsilon_0)}{2T} + \int_{m_0}^{\infty} \dots m^{-a + \frac{7}{2}} \dots$$

So (Taking $a=4, c=d=1$) c_v diverges @ $T=T_c = \frac{1}{b}$; $\epsilon_n(T_c)$ finite



My notes from Helmut's 1981 talk

3. Work in (restricted) model with

- a.) $SU(2)$ instead of $SU(3)$ - because fermions require much more comp. time.
- b.) pure g 's, no q 's

DECONFINEMENT IN DENSE MATTER

How does strongly interacting matter behave at high temperature and/or density?

Transition from nuclear to quark matter

[Zeldovich 1957; Zeldovich, Okun, Pikelner 1965; Ivanenko, Mandelstam 1967; Pacini 1966; Itoh 1970; ...]

Three approaches to critical strong interaction physics:

Helmut's 1981 talk

- CRITICAL HADRON PHYSICS
- TWO-PHASE MODELS
- QCD

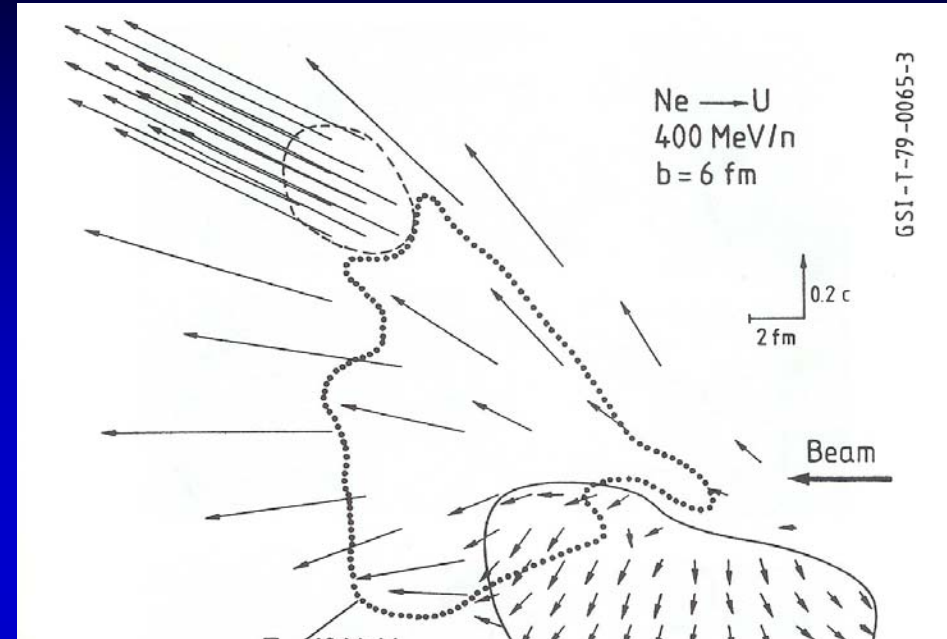
Brief comments only on first two, then

TRANSITION TO QUARK MATTER IN QCD

Perturbative QCD, INSTANTON GAS, STRONG COUPLING EXPANSION:

High or low temperature limits only

- State of the art in graphics:



- State of the art in simulations:





The End is the Beginning

C. SIGNALS FOR QUARK MATTER (and for PHASE TRANSITION)

E. L. Feinberg 1976, E. Shuryak 1980, G. Domokos & J. B. Goldmann 1981, I. Montvay 1980, K. Kajantie & H. Mäffinen 1981, J. Rafelski 1981 (and more...)

Three possibilities:

- photons or lepton pairs from "hot" interior emitted during quark stage
- discontinuity in photon or lepton pair spectra at parameters corresponding to T_c
- quantum number information from hadrons as signals for "chemical" equilibrium

Troubles:

- rate measurements; calculate dN/dM^2 for lepton pairs from quark matter - these pairs are also produced from $A \approx p-p$ collisions without quark matter.
Better: functional form $e^{-M/T}$ vs $\frac{1}{M^4}$
- hadronic vs. quark matter: clear quantum

CONCLUSIONS:

- (1) QCD PREDICTS TRANSITION FROM NUCLEAR MATTER TO QUARK-GLUON PLASMA
- (2) ESTIMATES OF TRANSITION PARAMETERS INDICATE THAT RELATIVISTIC HEAVY ION COLLISIONS SHOULD PROVIDE SUFFICIENT ENERGY DENSITY
- (3) ON SIGNATURES, STILL MUCH WORK TO DO...

Little bang ~ big bang

when you look, quark matter has vanished



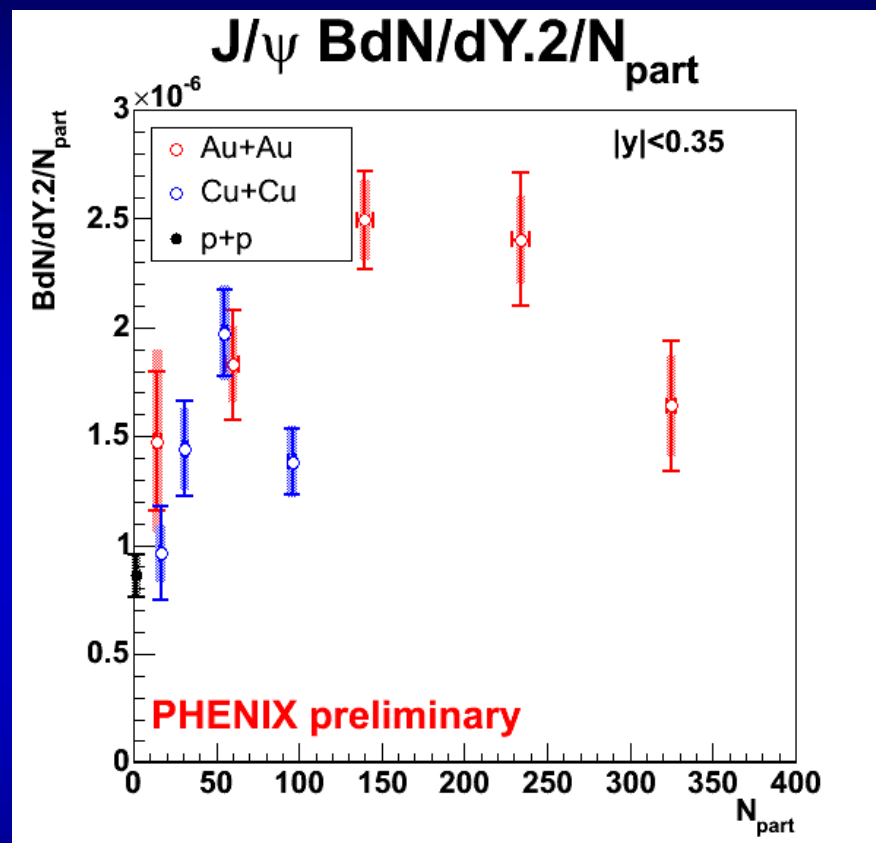
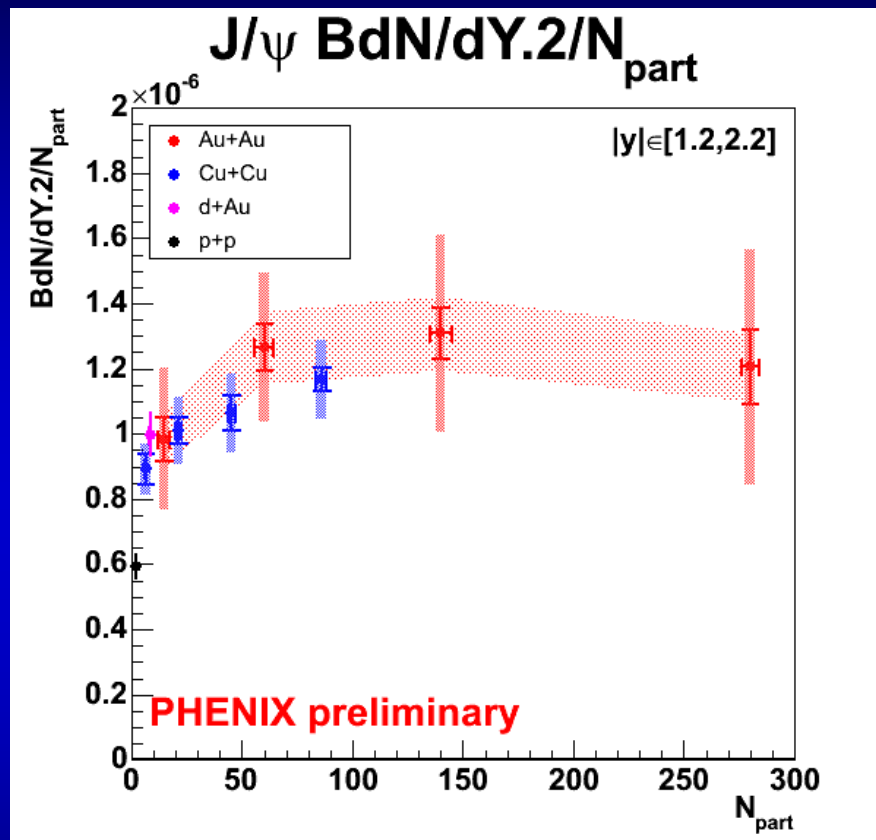
*Thanks Helmut,
for the continuing
intellectual
journey!*



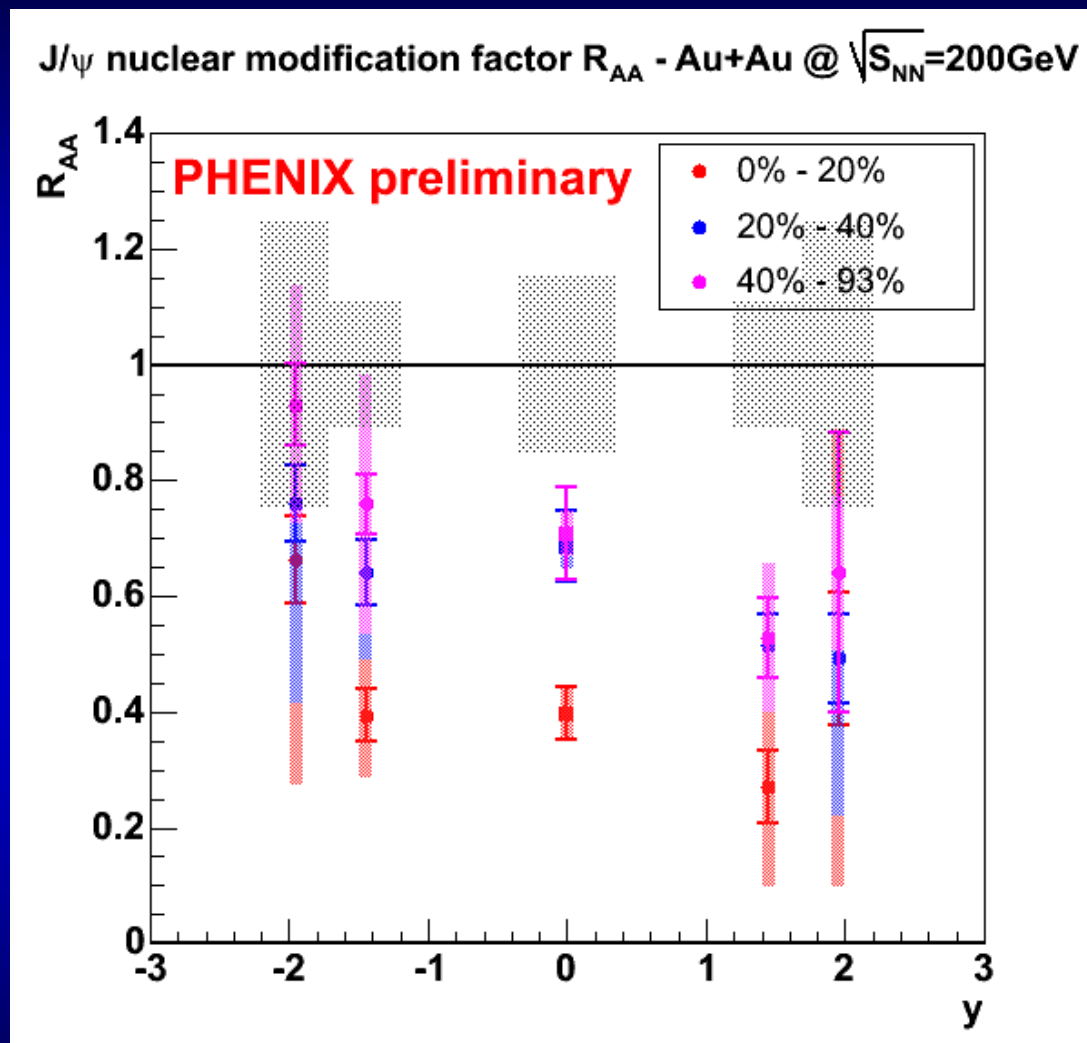
Backup



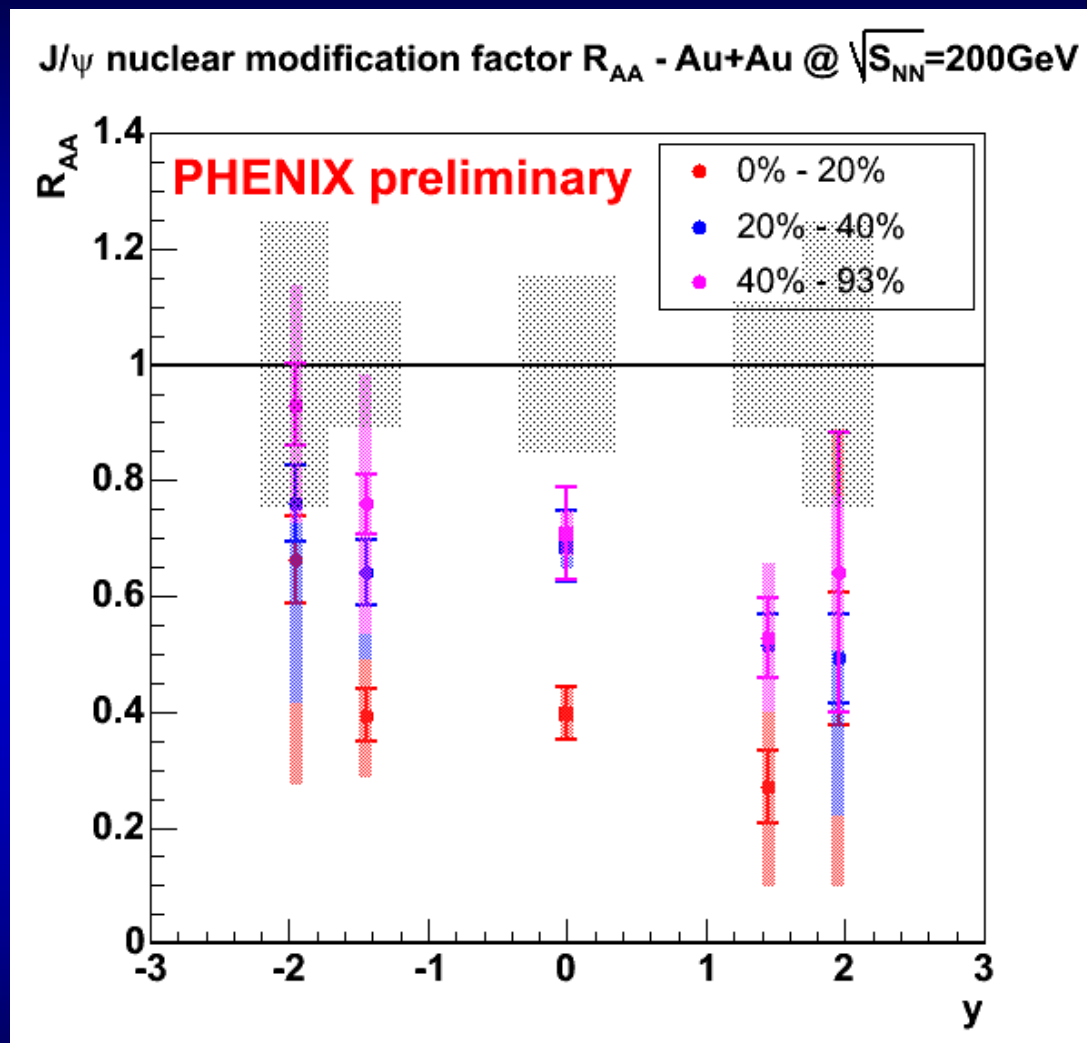
Test of N_{part} scaling



R_{AA} vs rapidity



R_{AA} vs rapidity





J/Ψ production in d+Au vs centrality

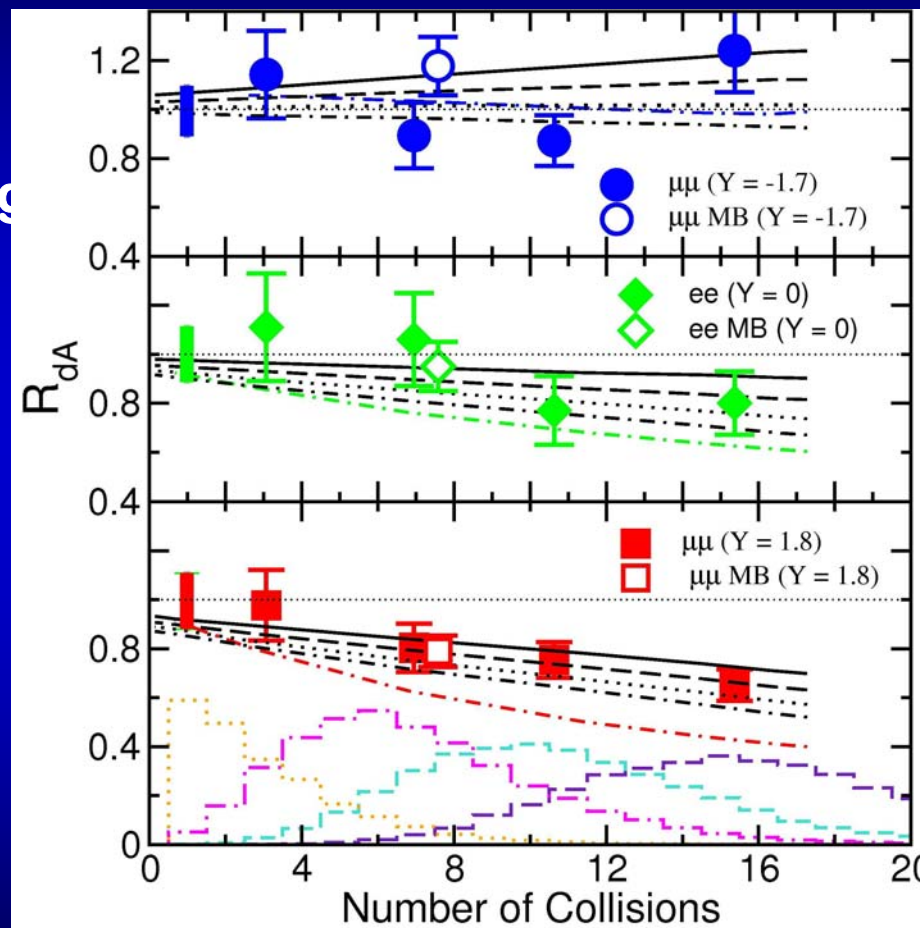
Small centrality dependence

Models with absorption + shadowing

- shadowing EKS98
- $\sigma^{abs} = 0$ to 3 mb

$\sigma^{abs} = 1$ mb good agreement

$\sigma^{abs} = 3$ mb is an upper limit



Weak shadowing and weak nuclear absorption observed