

A Summary of Earlier p+A Experiments' Physics & Results

Workshop on the Physics of p+A Collisions at RHIC

BNL, Jan 7-9, 2013

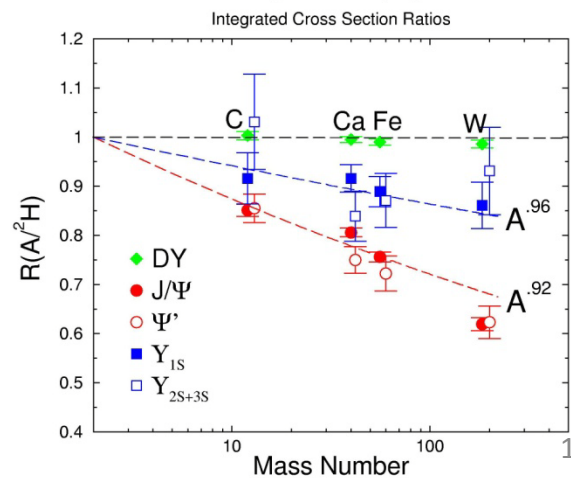
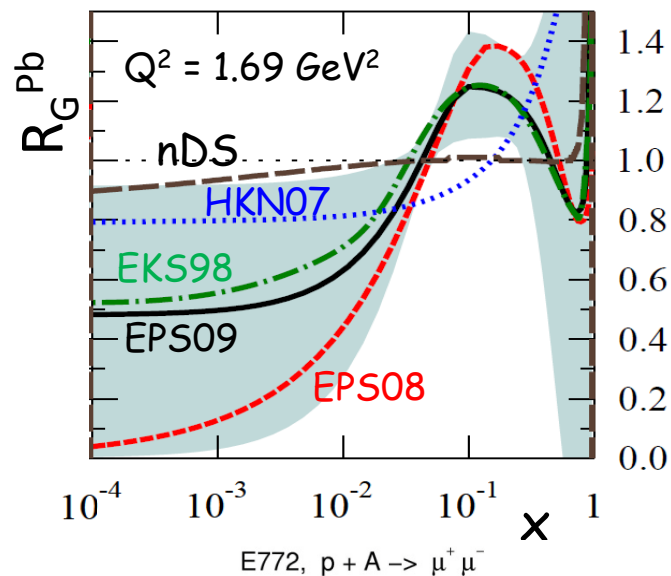
Mike Leitch, LANL Retired Fellow

Cold Nuclear Matter (CNM) physics

Lessons & Clues from the data:

- shadowing
- initial-state vs final-state
- absorption
- dE/dx from Drell-Yan
- p_T broadening
- Intrinsic Charm
- very small- x shadowing (E665)

*FNAL E772/789/866; SPS NA38/50/60;
FNAL E665 (DIS); Cronin expt, ...*



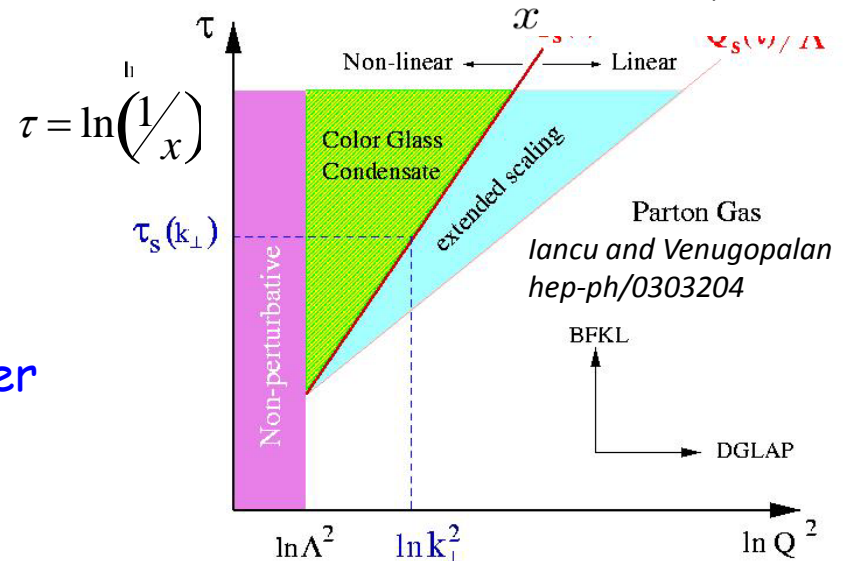
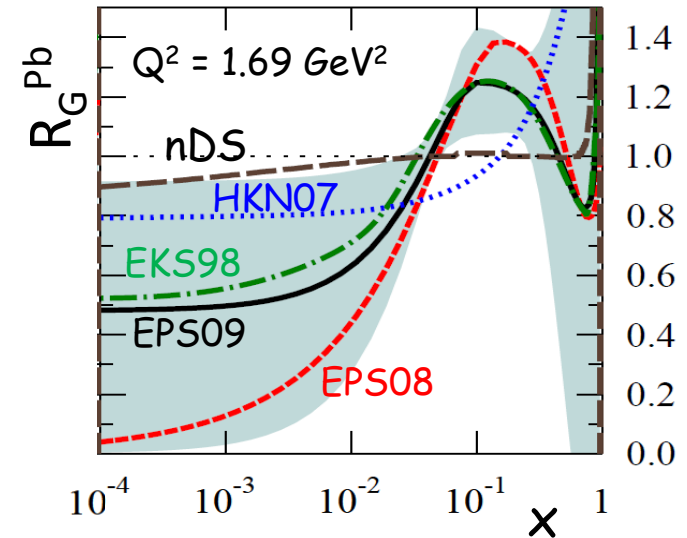
Saturation of Small-x Gluons or Shadowing

Leading twist gluon shadowing, e.g.:

- EPS09 - phenomenological fit to DIS & DY data with large uncertainties, Eskola, Paukkunen, Salgado, JHEP 0904:065,2009
- Also coherence models & higher-twist (HT) shadowing, e.g. Vitev

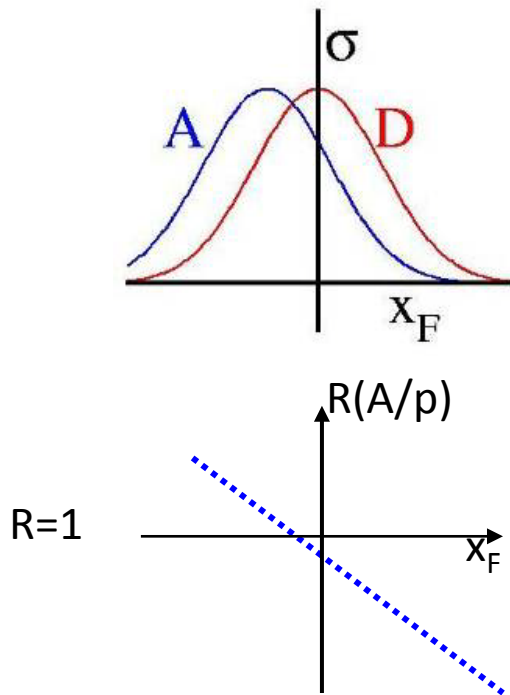
Small-x gluon saturation or Color Glass Condensate (CGC)

- At low-x there are so many gluons that $2 \rightarrow 1$ diagrams become important and deplete the low-x region
- Nuclear amplification: $x_A G(x_A) = A^{1/3} x_p G(x_p)$, i.e. gluon density is $\sim 6x$ higher in Gold than the nucleon

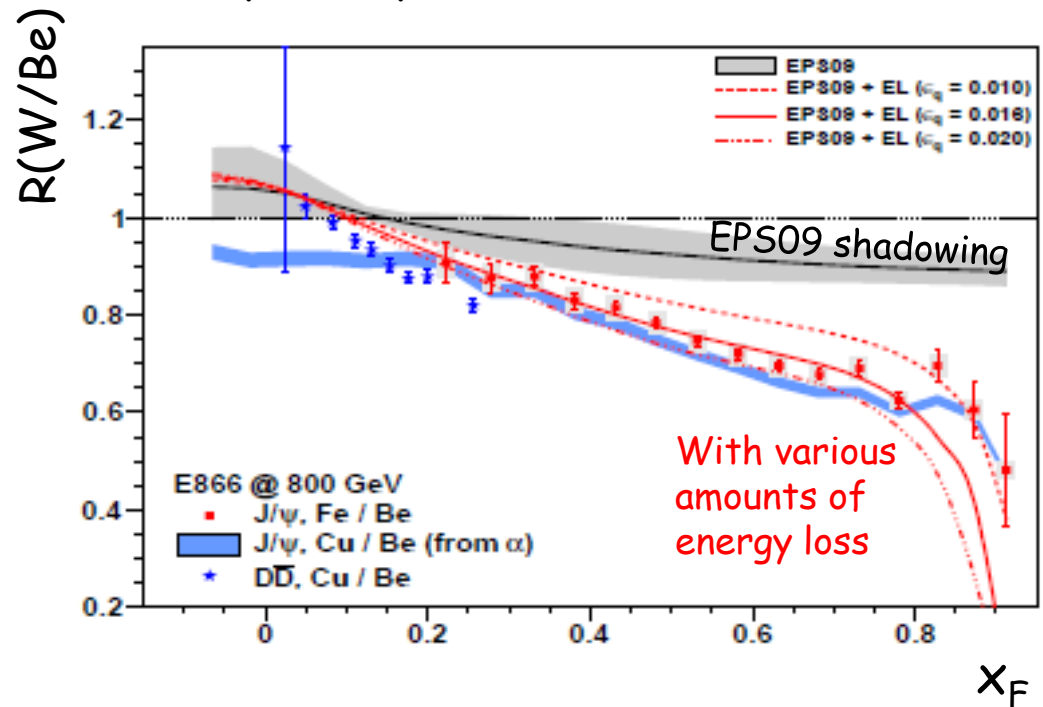


Energy Loss of Partons in Nuclei

Initial-state energy loss



E866/NuSea, 800 GeV, J/ Ψ
p+W / p+Be ratio



R. Vogt - Jan. 2011

Cronin effect, or p_T broadening from soft initial-state pre-hard-interaction scatterings

Cronin et al., PRD 11, 3105(1975)

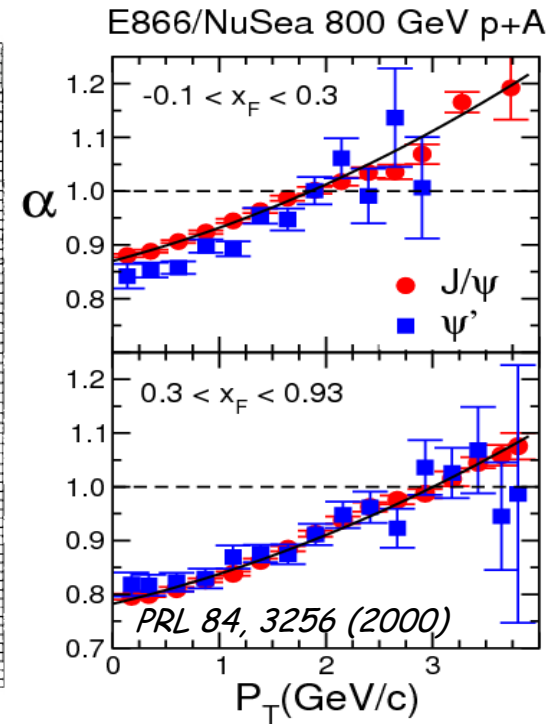
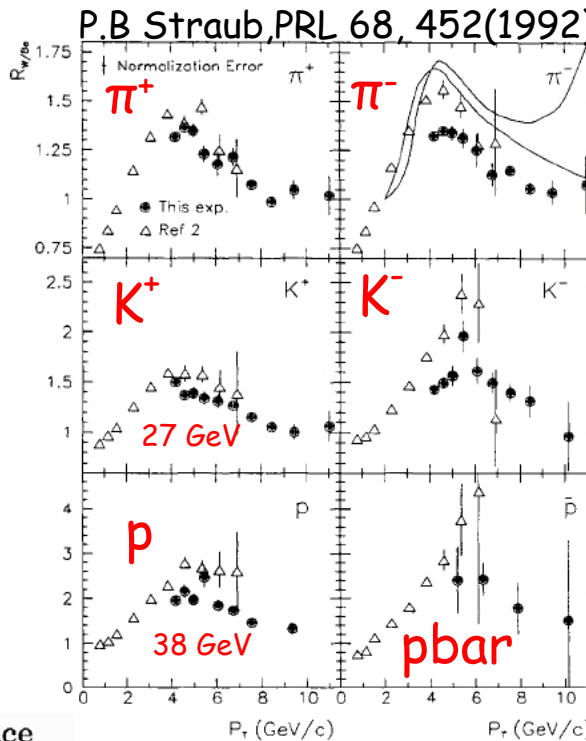
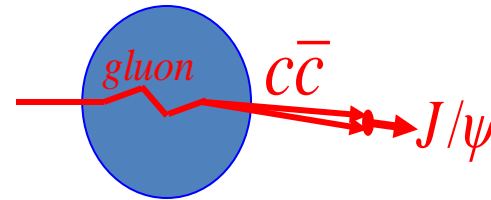
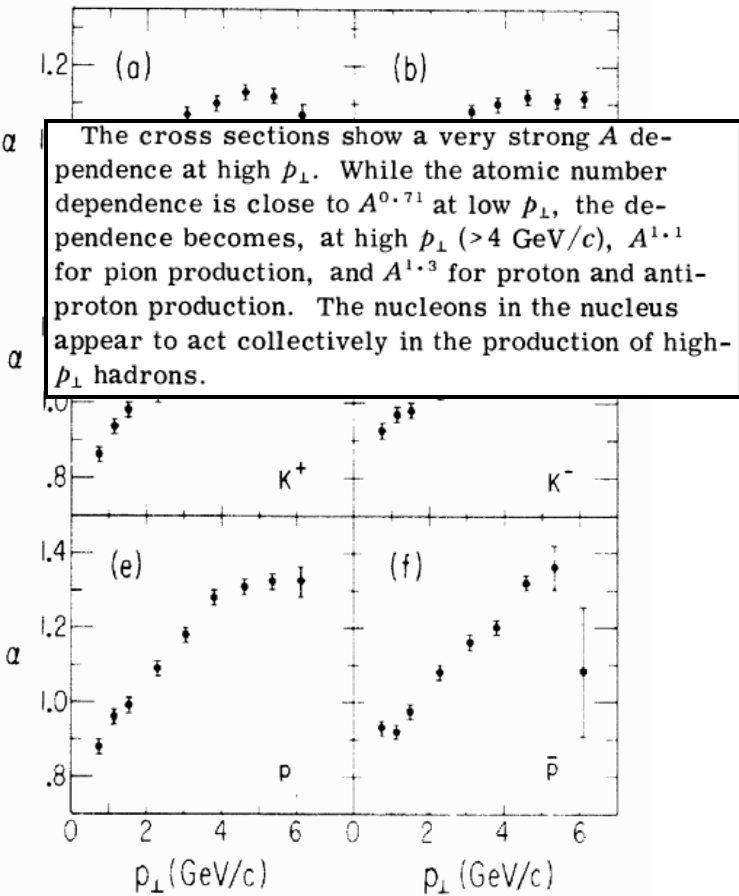
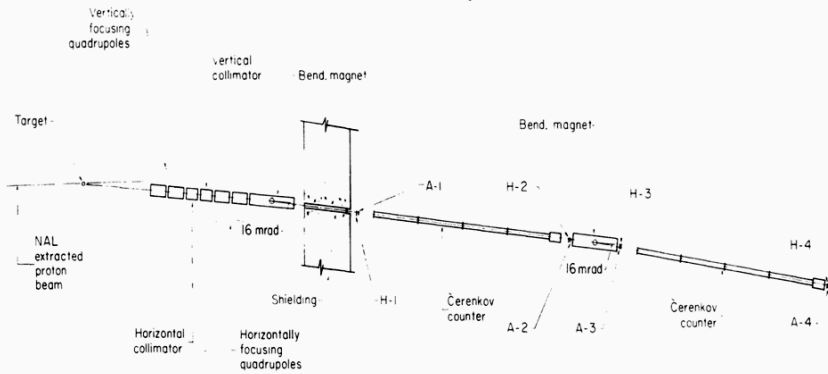


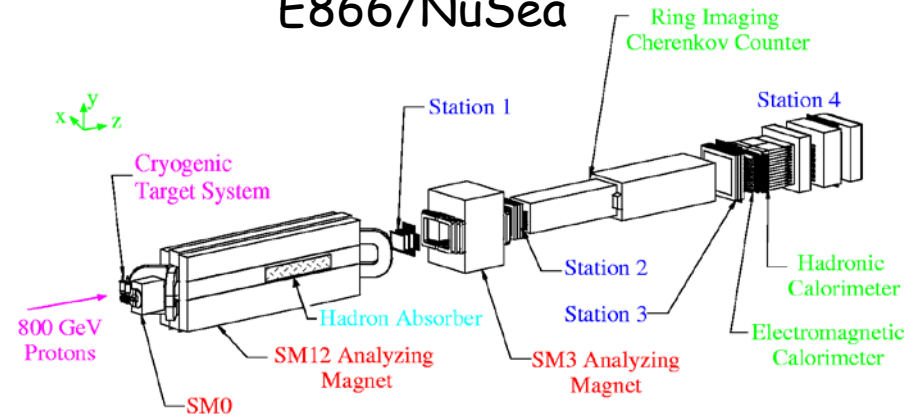
FIG. 17. Plots of the power α of the A dependence versus p_{\perp} for the production of hadrons by 300-GeV protons; (a) π^+ , (b) π^- , (c) K^+ , (d) K^- , (e) p , and (f) \bar{p} .

Lessons and Clues from the Data

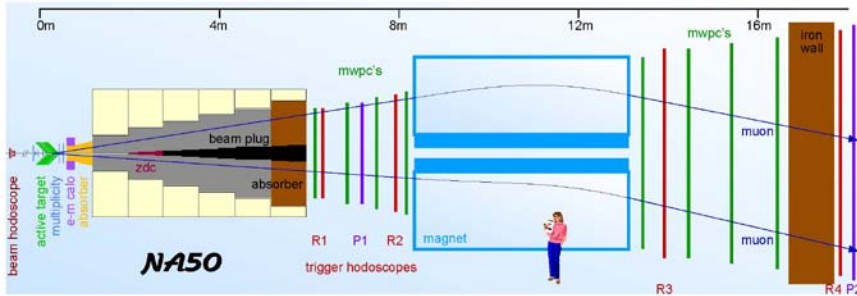
Cronin expt.



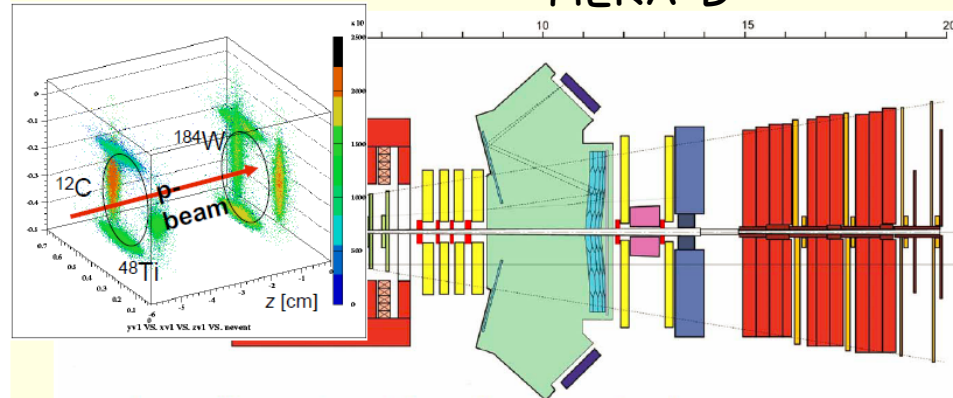
E866/NuSea



NA50



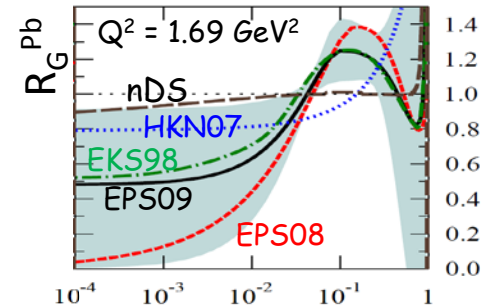
HERA-B



Suppression Not Universal vs x_2 as Expected for Shadowing

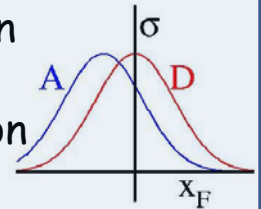
PHENIX, E866, NA3 J/ ψ Comparison

$$\sigma_{pA} = \sigma_{pp} A^\alpha$$



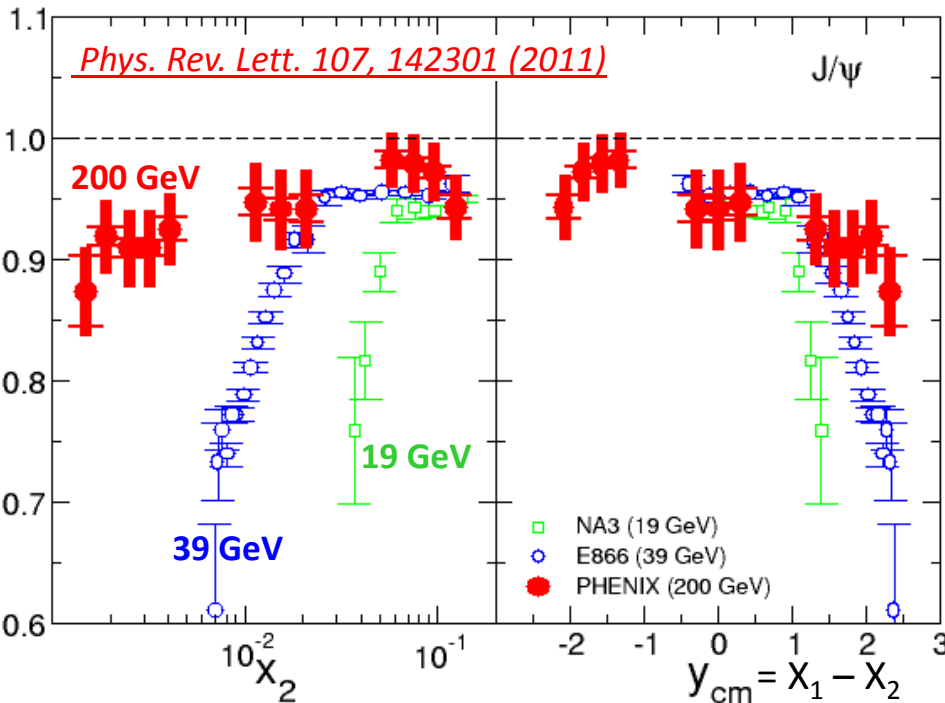
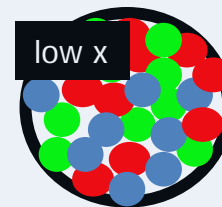
Closer to scaling with x_F or rapidity
 • initial-state gluon energy loss?

Energy loss of incident gluon shifts effective x_F and produces nuclear suppression which increases with x_F



• or gluon saturation?

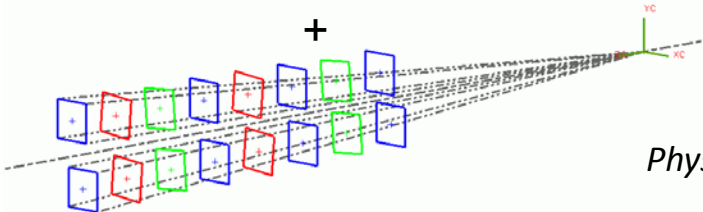
Gluon saturation from non-linear gluon interactions for the high density at small x ; amplified in a nucleus.



(x_2 is x in the nucleus)

Fermilab E789: $D^0 \rightarrow K\pi$ and $B \rightarrow J/\psi X$ (charm & beauty using silicon vertex detector)

Dimuon spectrometer

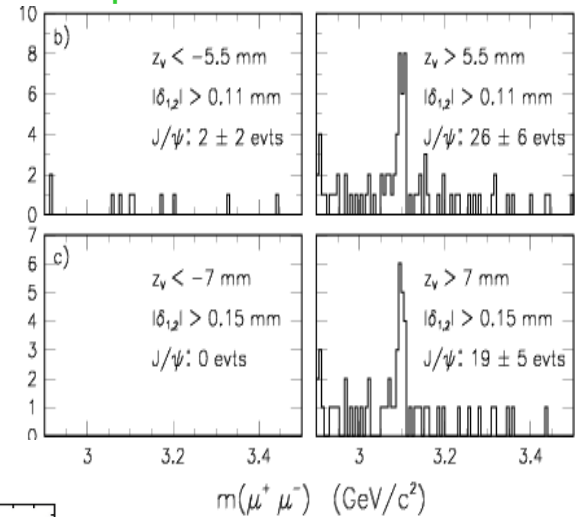


16-plane, 50 μ m pitch/8.5k
strip silicon vertex detector

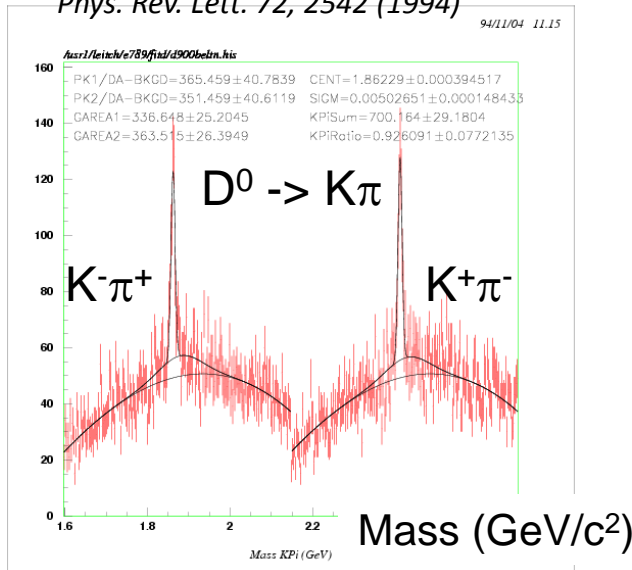
$B \rightarrow J/\psi + X$

Phys. Rev. Lett. 74, 3118 (1995)

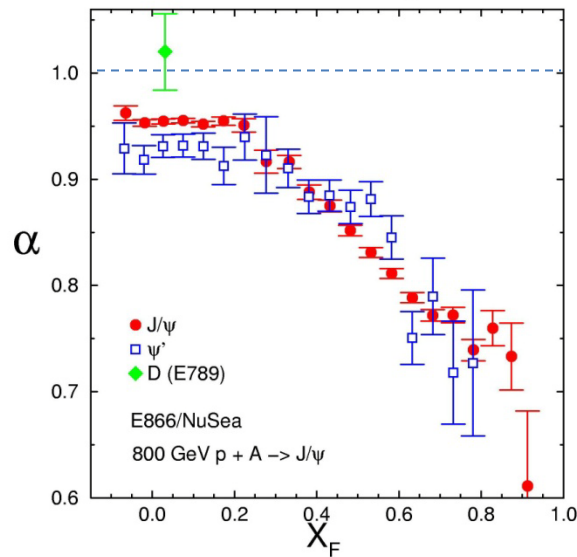
upstream downstream



Phys. Rev. Lett. 72, 2542 (1994)



E866/NuSea, $\sigma = \sigma_N * A^\alpha$



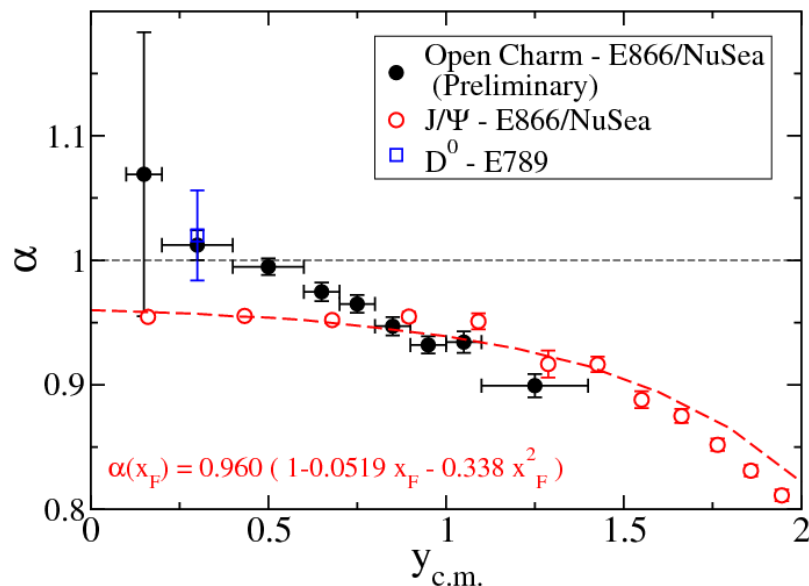
Compare closed and open charm to isolate initial-state CNM effects

Comparing Open & Closed Charm Isolating Initial-state Effects

E866/NuSea 800 GeV p+A

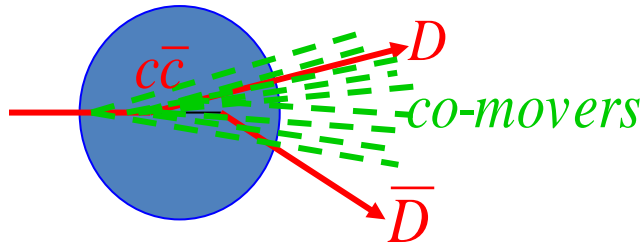
Open-charm p+A nuclear dependence (single- μ $p_T > 1$ GeV/c) - very similar to that of J/ψ :

- implies that dominant effects are in the initial state
 - e.g. dE/dx, CGC (since shadowing disfavored by lack of x_2 scaling)
- weaker open-charm suppression at $y=0$ attributed to lack of absorption for open charm



follow this example at RHIC

Absorption at mid-rapidity, J/ψ & ψ' in Fixed Target Measurements



Absorption (or dissociation) of J/ψ or pre- J/ψ ($c\bar{c}$) as they exit nucleus into two D mesons, by nucleus or co-movers

Power law parameterization $\sigma = \sigma_N * A^\alpha$

$$\alpha = 0.954 \pm 0.003 \quad \text{E866/NuSea @ } x_F=0$$

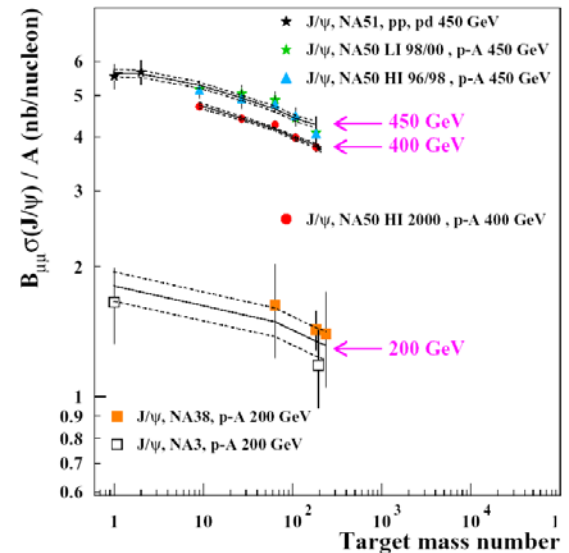
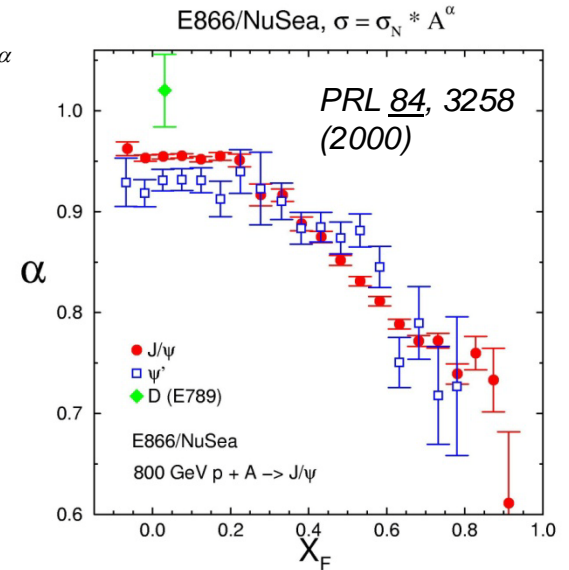
$$\alpha = 0.941 \pm 0.004 \quad \text{NA50, QM04}$$

Absorption model parameterization (from pA)

$$\sigma_{J/\psi} = 4.18 \pm 0.35 \text{ mb} \quad \text{NA50 QM05}$$

$$\sigma_{\psi'} = 7.6 \pm 1.1 \text{ mb}$$

$$\sigma_{pA} = \sigma_{pp} A^\alpha$$



Absorption or Dissociation of J/ψ after correcting for shadowing

Energy Dependence; after corrections for EKS98 shadowing

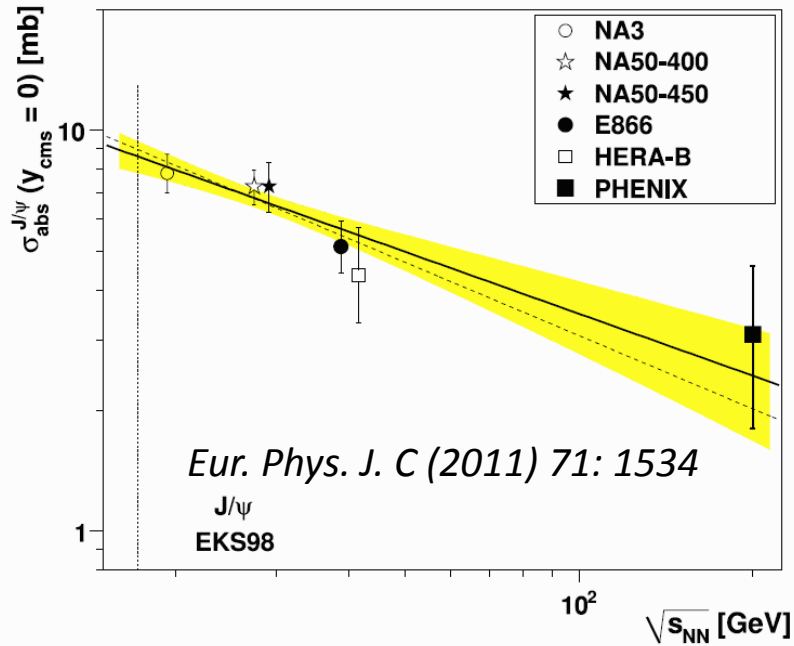
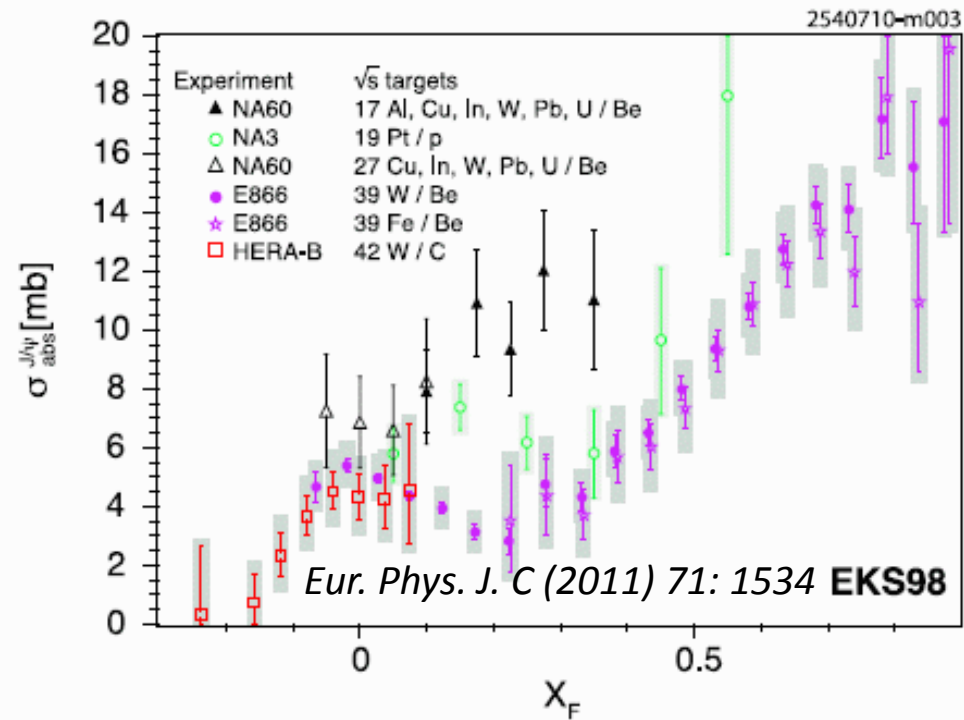


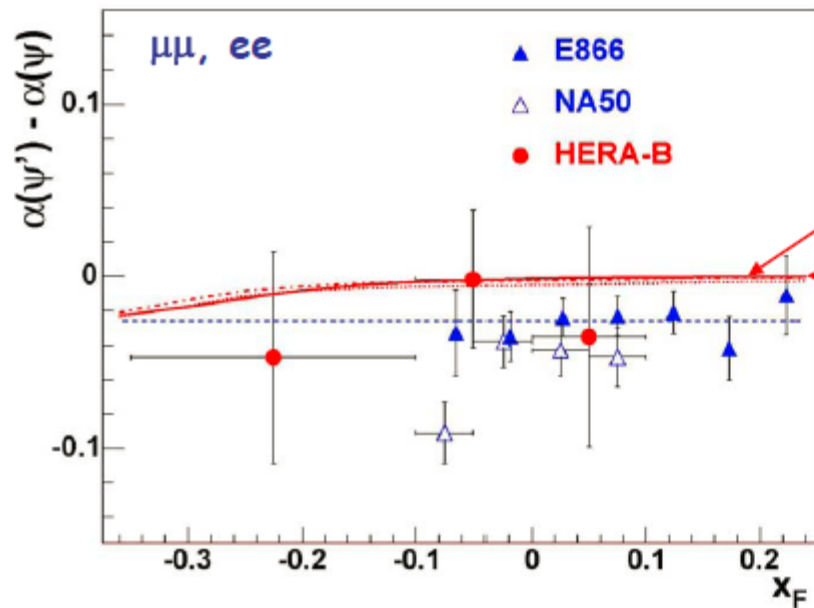
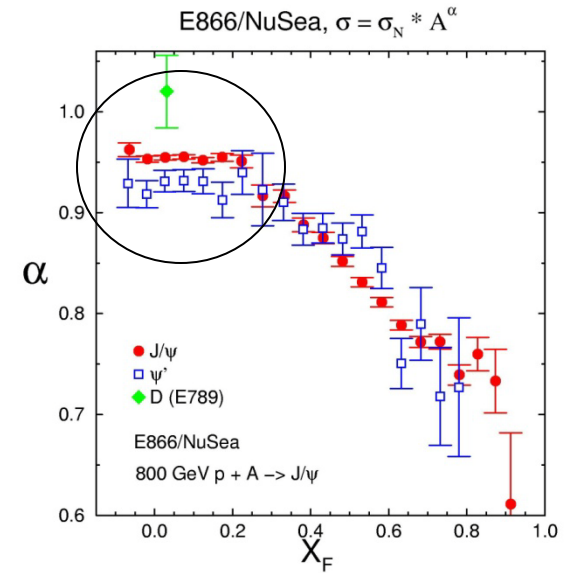
Fig. 77 The extracted energy dependence of $\sigma_{\text{abs}}^{J/\psi}$ at midrapidity. The *solid line* is a power-law approximation to $\sigma_{\text{abs}}^{J/\psi}(y = 0, \sqrt{s_{NN}})$ using the EKS98 [876, 877] shadowing parametrization with the CTEQ61L parton densities [878, 879]. The *band* indicates the uncertainty in the extracted cross sections. The *dashed curve* shows an exponential fit for comparison. The data at $y_{\text{cms}} \sim 0$ from NA3 [880], NA50 at 400 GeV [872] and 450 GeV [873], E866 [874], HERA-B [881], and PHENIX [663] are also shown. The *vertical dotted line* indicates the energy of the Pb+Pb and In+In collisions at the CERN SPS. Adapted from [875] with kind permission, copyright (2009) Springer

Apparent absorption vs x_F ; after EKS98 shadowing is removed



ψ' absorbed more than J/ψ near $x_F = 0$

Meson	M(GeV)	R(Fm)	BE (MeV)
J/ψ	3.1	.45	~640
Ψ'	3.7	.88	~52



A-dependence: $\Delta\alpha(x_F) = \alpha' - \alpha$

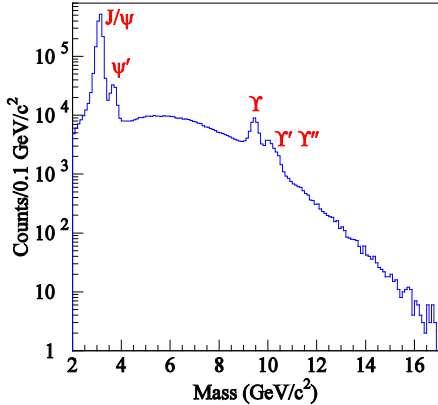
CEM for color-1 nuclear absorption

NRQCD for color-1 & -8 nuclear absorption

fit with const: $\Delta\alpha(\text{E866}) = -0.026 \pm 0.005$

All consistent with no x_F dependence.
 We use average of E866 and NA50:
 $\Delta\alpha(\text{E866, NA50}) = -0.030 \pm 0.004$

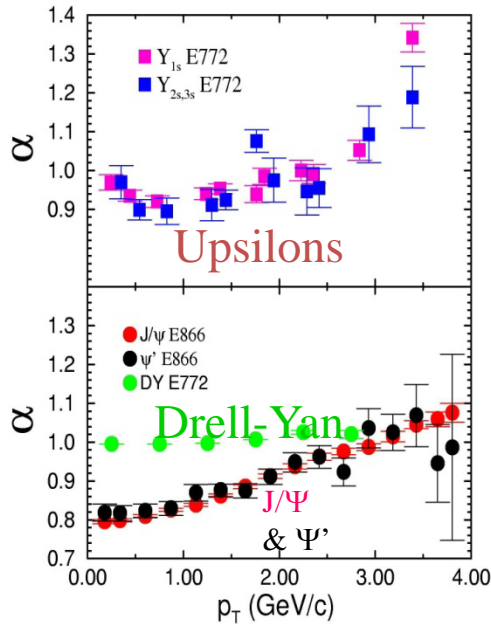
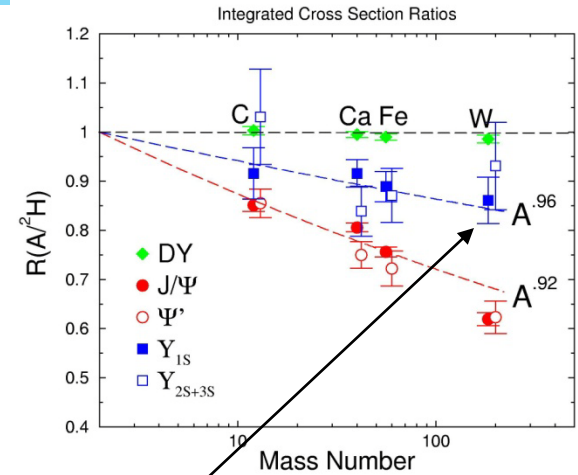
Contrasting Υ 's with J/ψ 's



$\sqrt{s} = 39 \text{ GeV}$ (E772 & E866)

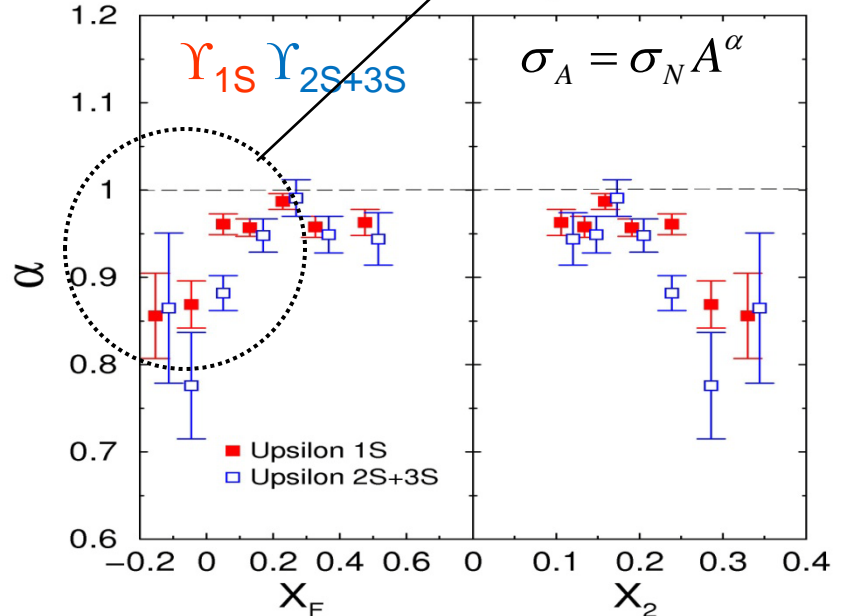
- less absorption
- not in shadowing region (large x_2)
- similar p_T broadening

E772, $p + A \rightarrow \mu^+ \mu^-$



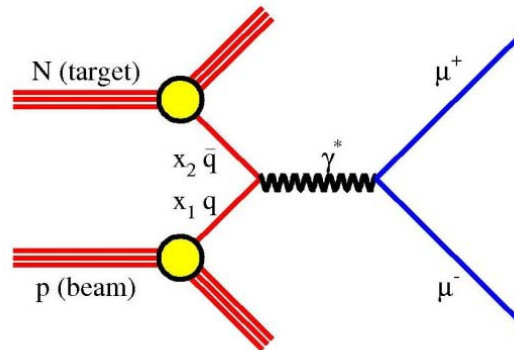
But careful: Υ suppression is from data for $x_F < 0$ or $x_2 > 0.2$ (in the EMC region)

E772 800 GeV p+A

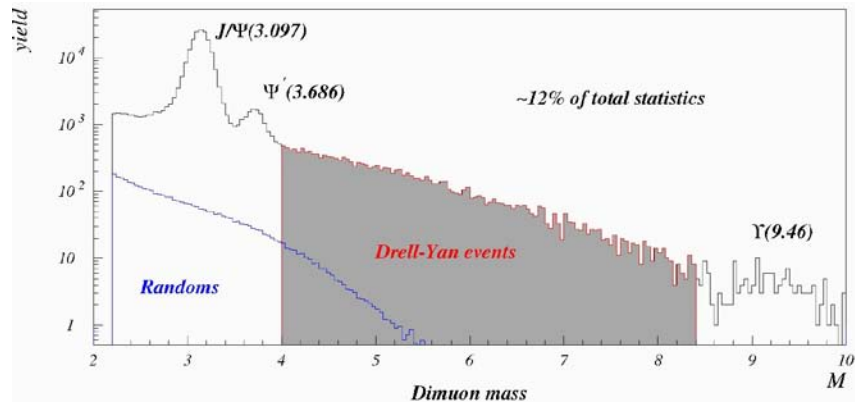
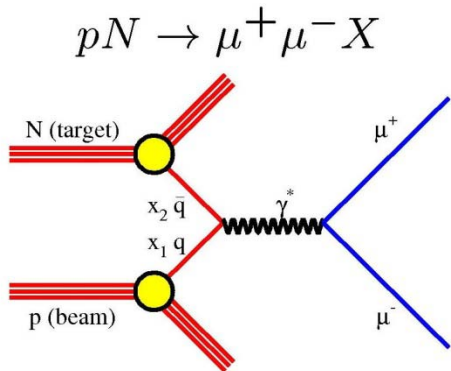


Drell-Yan

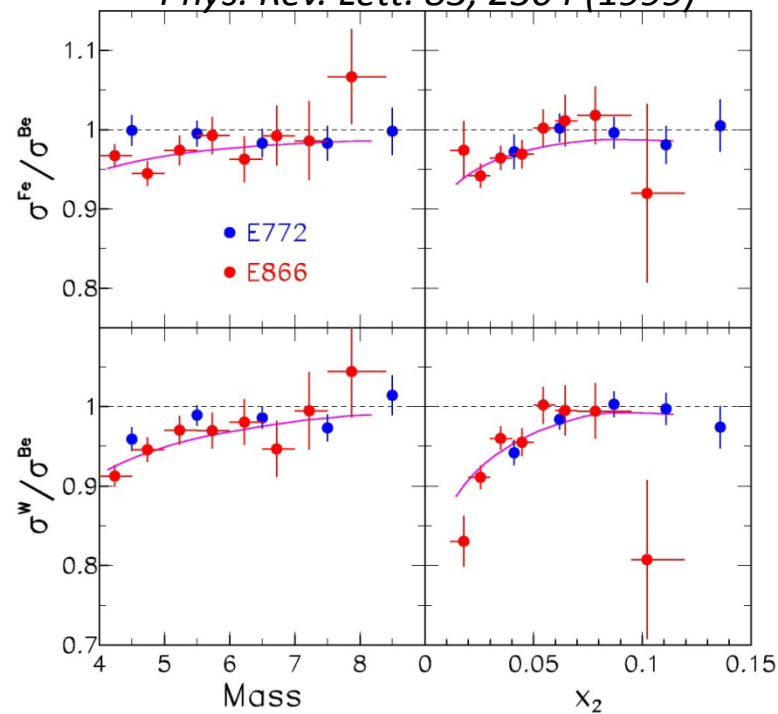
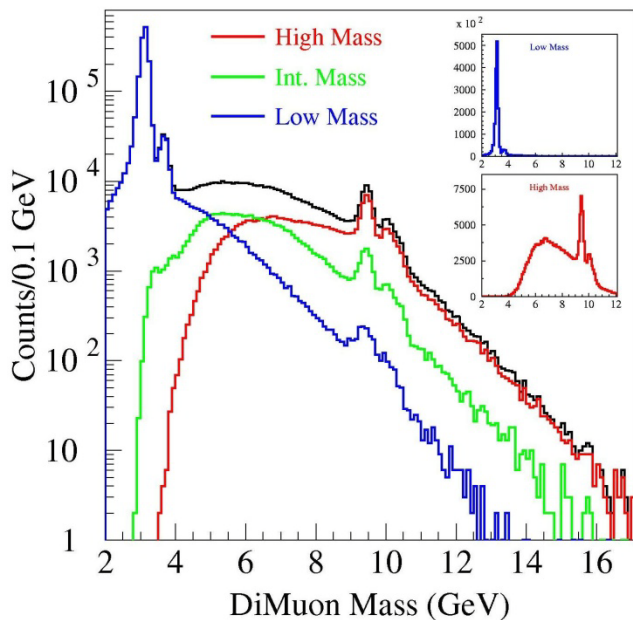
$$pN \rightarrow \mu^+ \mu^- X$$



Drell-Yan from E772/E789/E866 at Fermilab



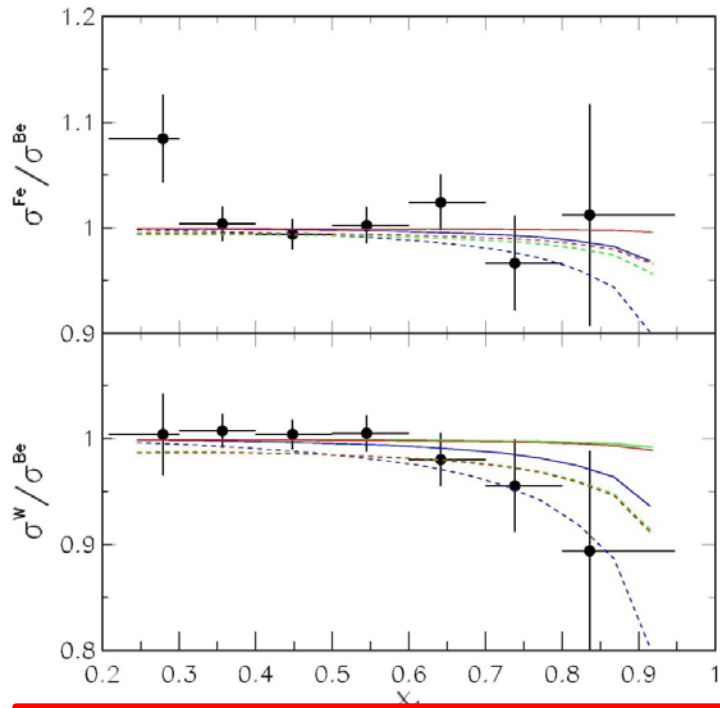
Phys. Rev. Lett. 83, 2304 (1999)



Quark Energy Loss from Drell-Yan data Two ways

assuming EKS Shadowing

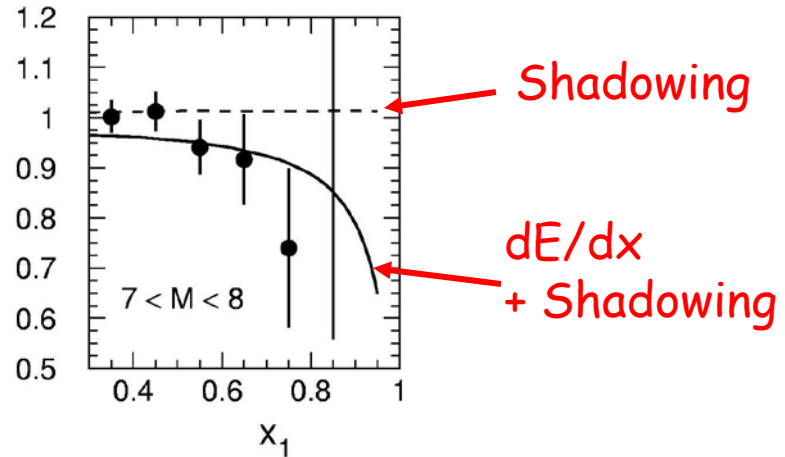
Phys. Rev. Lett. 83, 2304 (1999)



- Gavin and Milana $\Delta E < 0.14 \text{ \%}/\text{fm}$
- Brodsky and Hoyer $\Delta E < 0.44 \text{ GeV}/\text{fm}$
- Baier et al. $\Delta E < (0.046 \text{ GeV}/\text{fm}^2) \times L^2$

assuming Kopeliovich Shadowing

Johnson, Kopeliovich et al., PRL 86, 4483 (2001)



Analysis of our p-A Drell-Yan data using the Kopeliovich model.

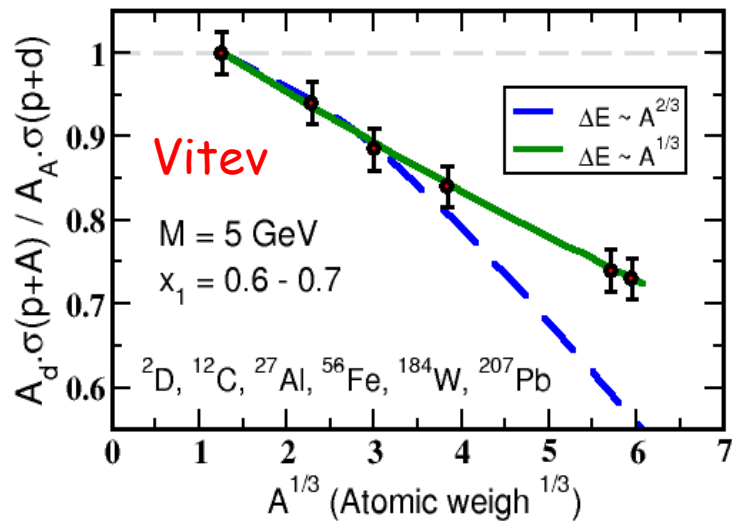
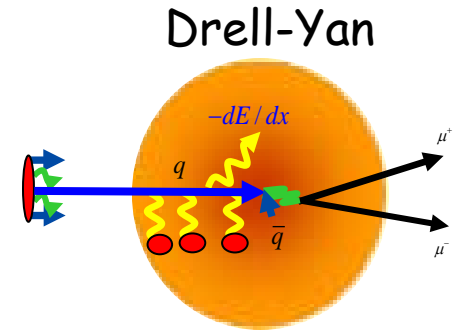
$$\boxed{dE/dx = 2.32 \pm 0.52 \pm 0.5 \text{ GeV}/\text{fm}}$$

Confusion due to having both shadowing and dE/dx contributing to Drell-Yan suppression for measurements at this energy

Drell-Yan to constrain parton energy loss in CNM

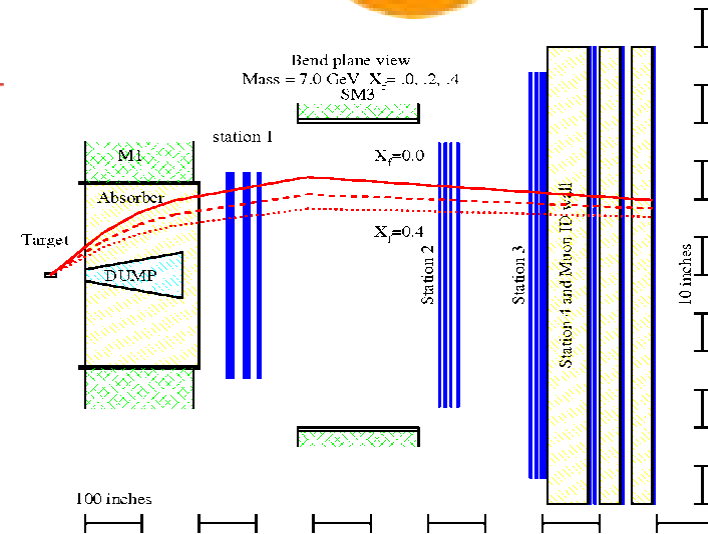
FNAL E906/Seaquest - Drell-Yan

In E906 at 120 GeV, nuclear suppression in Drell-Yan should only be from dE/dx ($x_2 > 0.1$)



$$\frac{\Delta E}{E} \propto \frac{\mu^2 L^2 \ln E / Q_0}{\lambda_g E}$$

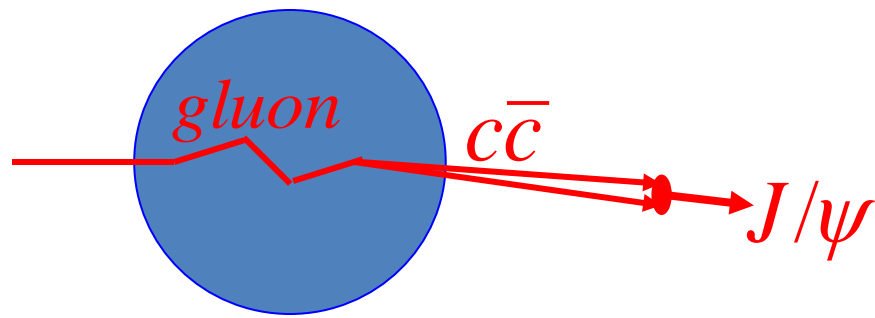
$$\frac{\Delta E}{E} \propto \frac{L}{\lambda_g} \ln \frac{E}{Q_0}$$



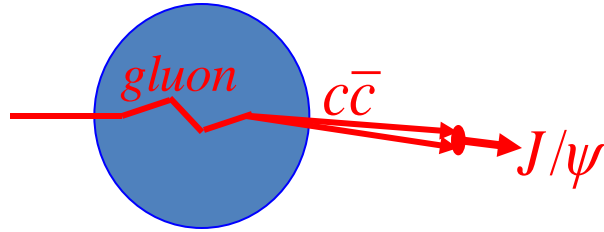
Distinguish radiative from collisional (L^2 vs L)

See Paul Reimer's talk

Transverse Momentum Broadening

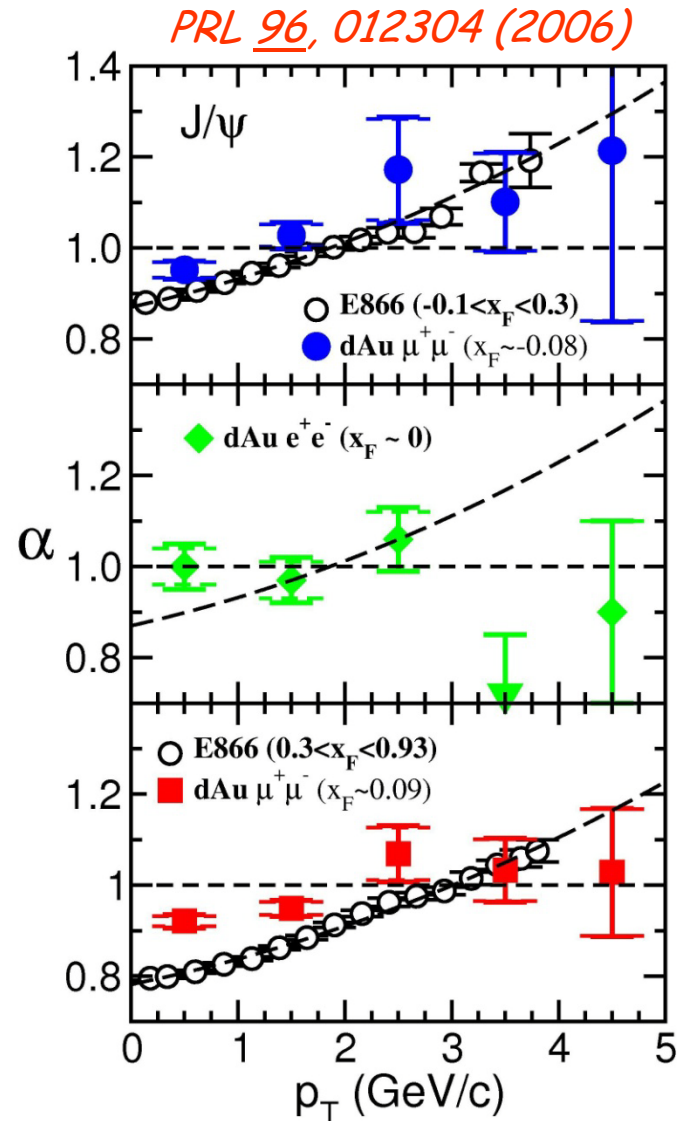
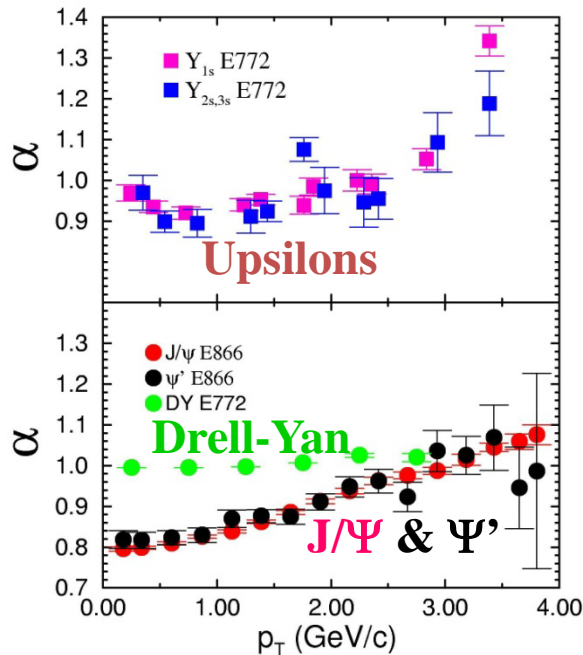


Transverse Momentum Broadening - DY, J/ψ, Upsilon



Initial-state gluon multiple scattering causes p_T broadening (or Cronin effect)

$$\sigma_A = \sigma_N A^\alpha$$



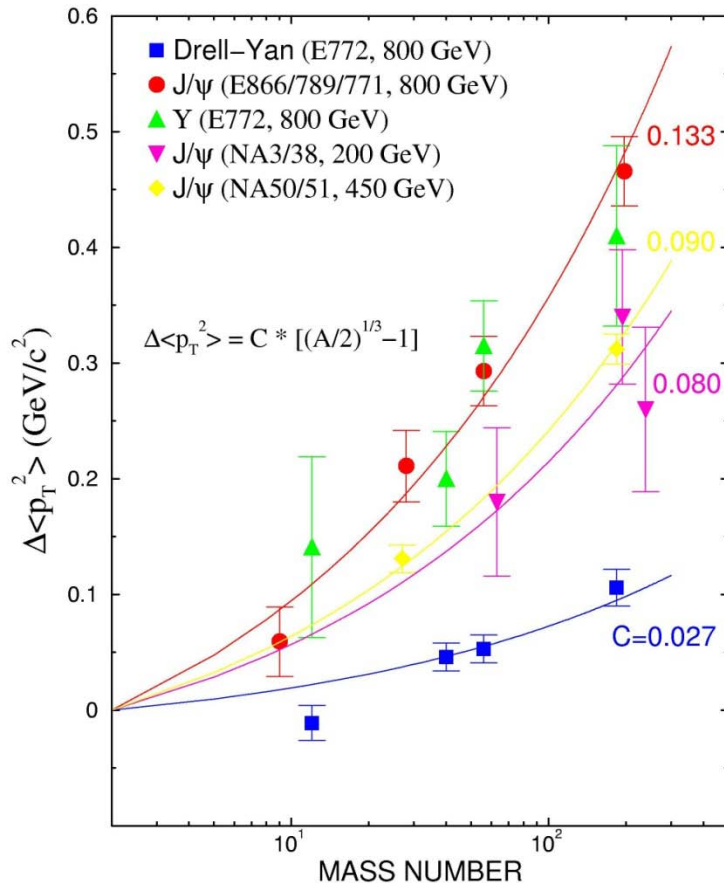
High x_2
 ~ 0.09

PHENIX 200 GeV results show p_T broadening comparable to that at lower energy ($\sqrt{s}=39$ GeV in E866/NuSea)

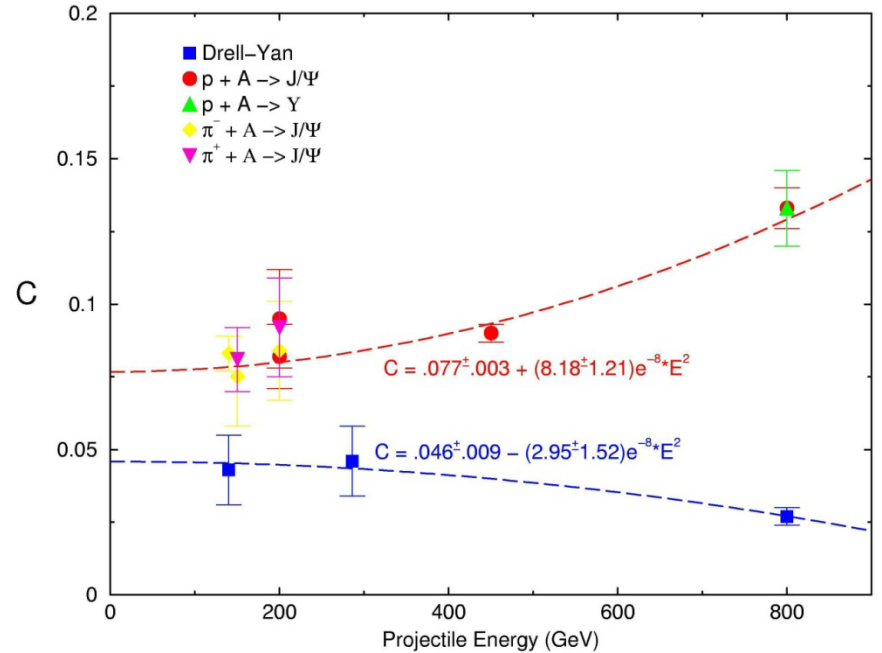
Low x_2
 ~ 0.003

Systematics of P_T Broadening E866/NuSea

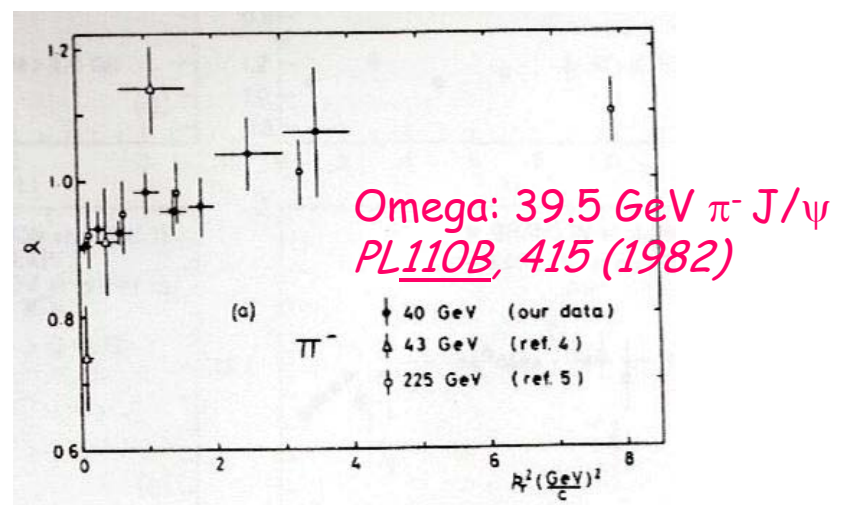
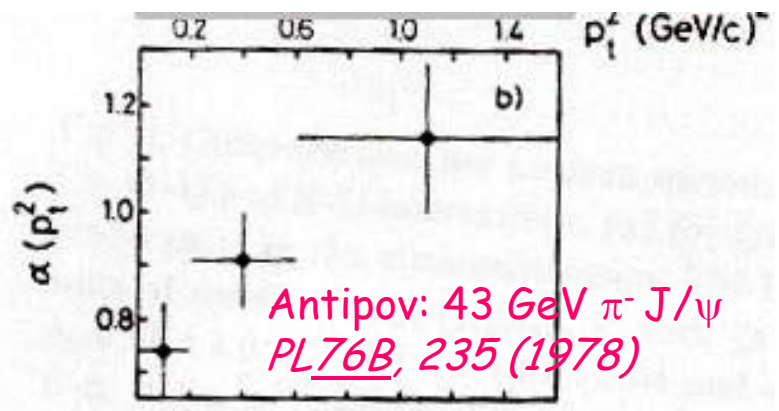
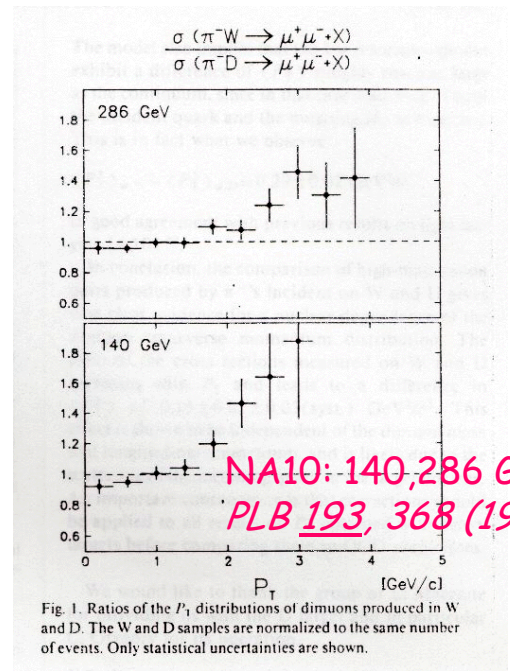
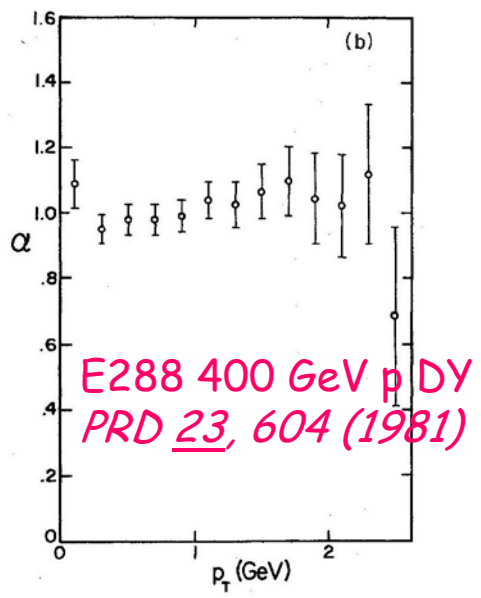
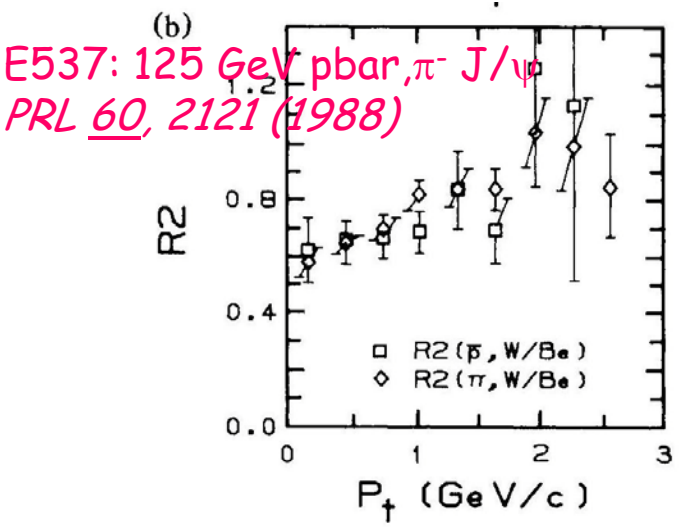
$$\Delta \langle p_T^2 \rangle = \Delta \langle p_T^2 \rangle_A - \Delta \langle p_T^2 \rangle_D$$



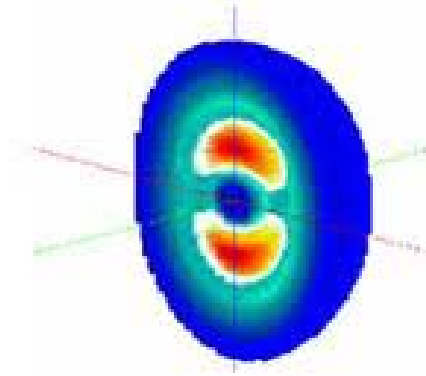
$$\Delta \langle p_T^2 \rangle = C * \left[\left(\frac{A}{2} \right)^{1/3} - 1 \right]$$



P_T Broadening for different energies and probes - other data



Deuteron vs the Proton



Deuteron is very weakly bound & not a typical nucleus in terms of CNM effects

- Binding energies:
 - D - 2.22 MeV
 - Alpha - 28.3 MeV

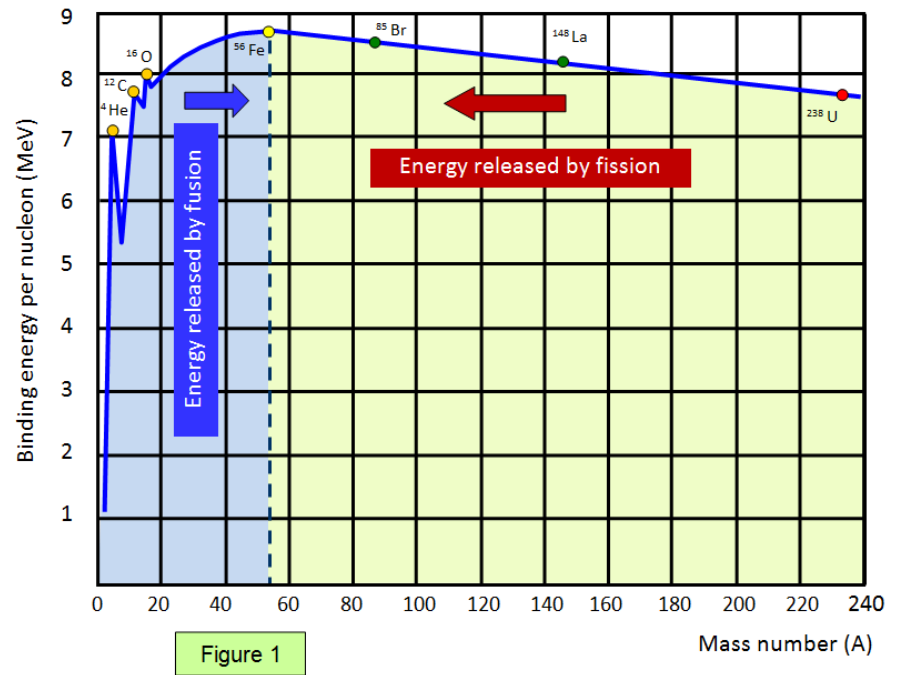
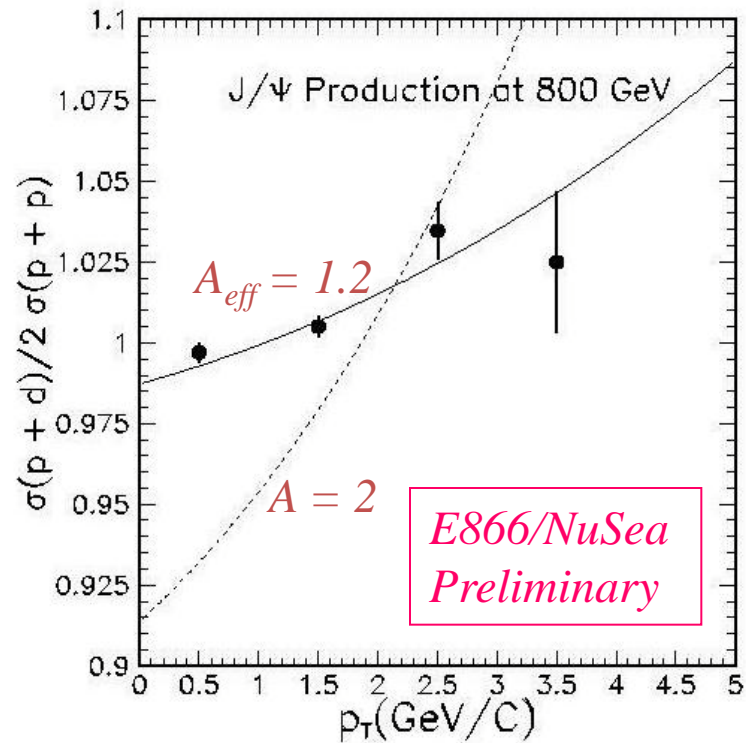
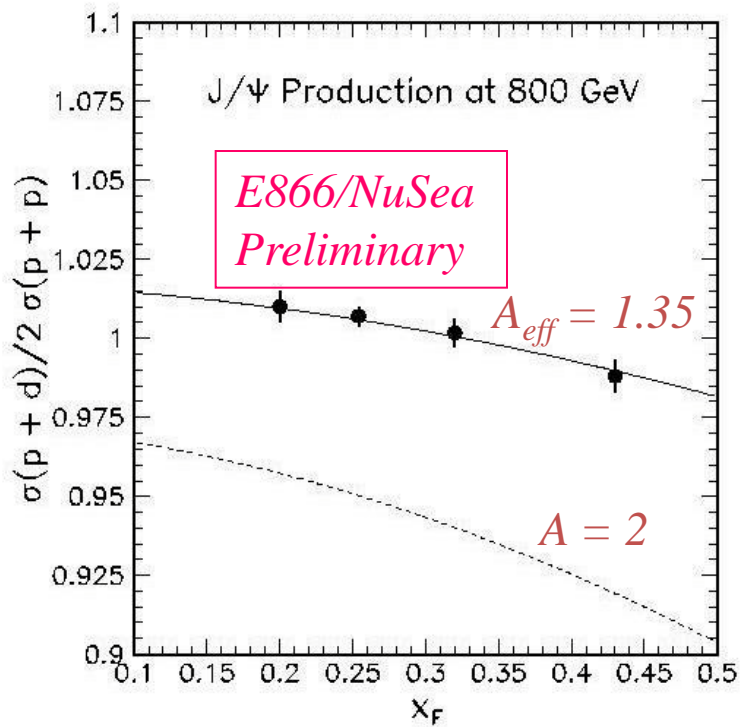


Figure 1

J/ψ Nuclear dependence weaker than expected for Deuterium/Hydrogen !



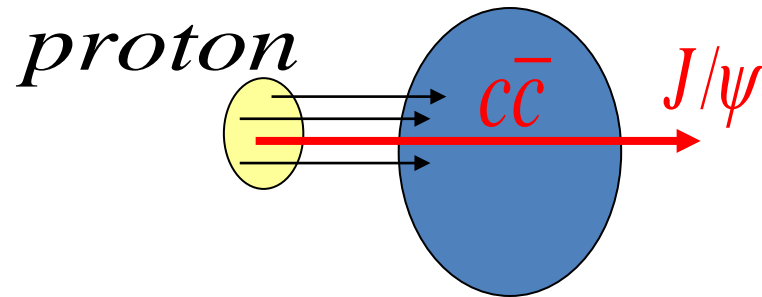
Nuclear dependence in deuterium seems to follow the systematics of larger nuclei, but with an effective, A_{eff} , smaller than the $A=2$ of deuterium

From fits to E866/NuSea
 p + Be, Fe, W data: $\sigma_{pA} \sim \sigma_{pp} A^\alpha$

$$\alpha(x_F) \propto 1 - 0.052x_F - 0.034x_F^2$$

$$\alpha(p_T) \propto 0.06p_T + 0.011p_T^2$$

Intrinsic Charm



Intrinsic charm components of incident proton produce J/ψ at large x_F . $A^{2/3}$ dependence from surface stripping of proton's light quarks (Brodsky)

Intrinsic charm contribution to Quarkonia

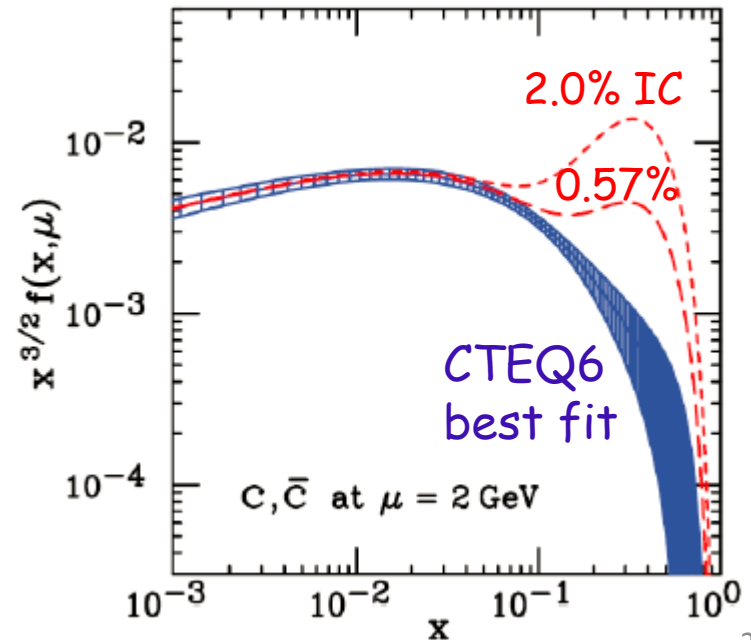
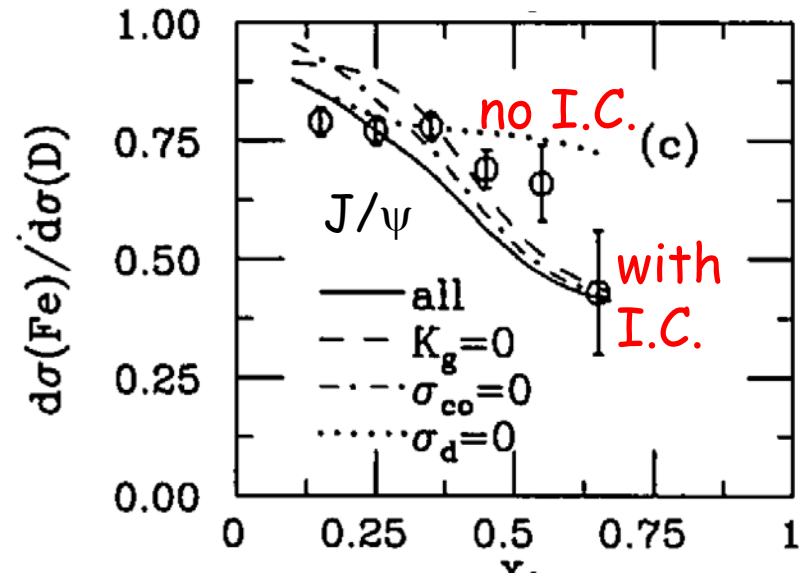
At large $x_F (\geq 0.5)$ intrinsic $c\bar{c}$ components of the projectile proton can dominate the production of charm pairs

- $A^{2/3}$ dependence via surface stripping of light quarks to free charm pair component

Vogt, Brodsky, Hoyer, NP B360, 67 (1991)
(also includes absorption and shadowing)

No conclusive evidence for IC:
"IC constrained to be from zero (no IC) to a level 2-3 times larger than previous model estimates"

Pumplin, Lai, Tung, PRD 75, 054029 (2007)



Flattening of shadowing at very small x?

F_2 structure functions (q & \bar{q}) may show leveling out of shadowing at smallest x values probed by DIS on nuclear targets (E665)

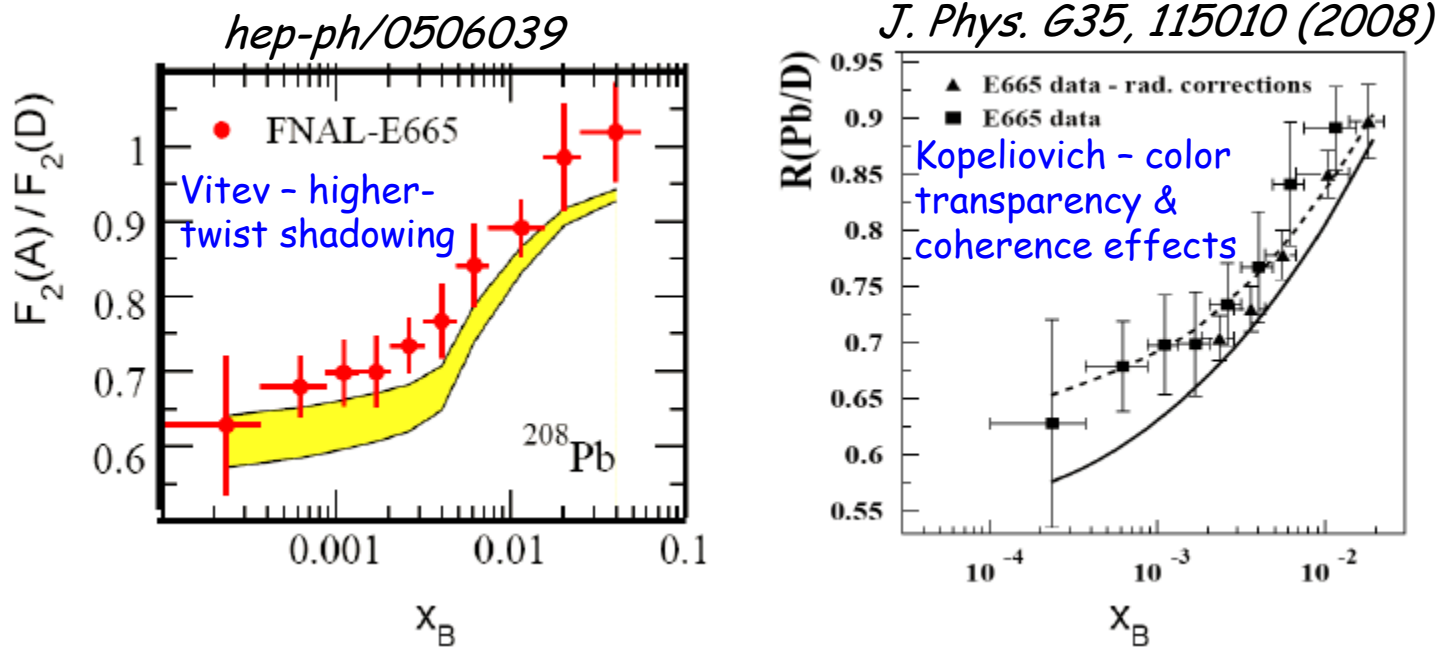
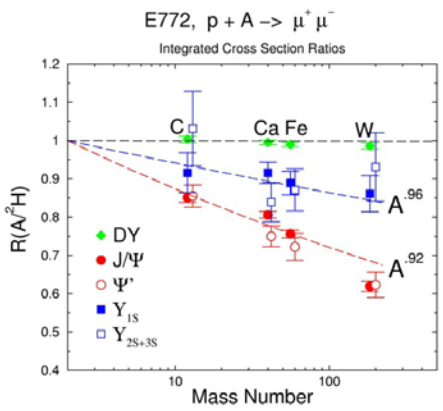


Figure 3.34: Data from E665 compared with (left) a shadowing calculation from Vitev [307] and with (right) color transparency and coherence calculations from Kopeliovich [223]. In the right panel, the solid (dashed) curve is with (without) the contribution from gluon shadowing.

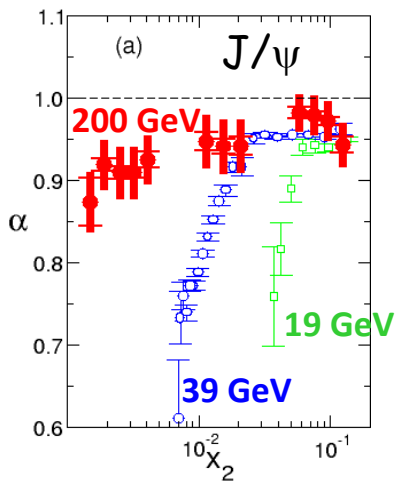
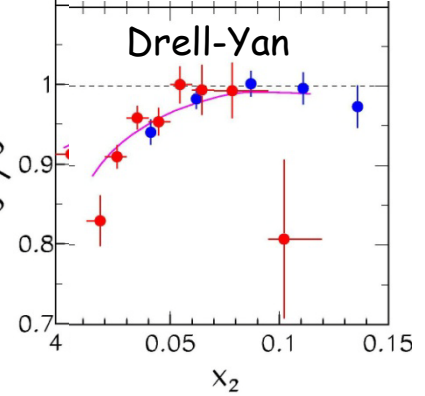
e.g. probe \bar{q} with very-forward Drell-Yan on nuclear targets



suppression of vector mesons

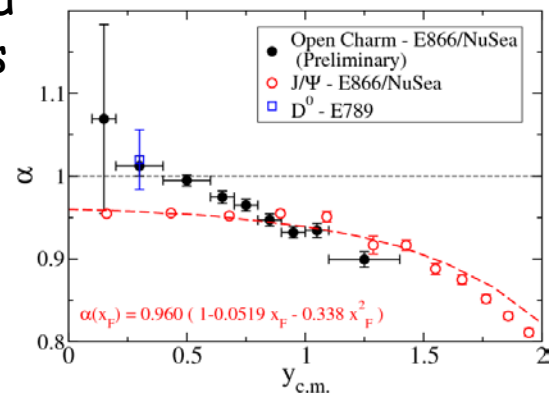
Summary p-A data lessons

DY shows shadowing or dE/dx

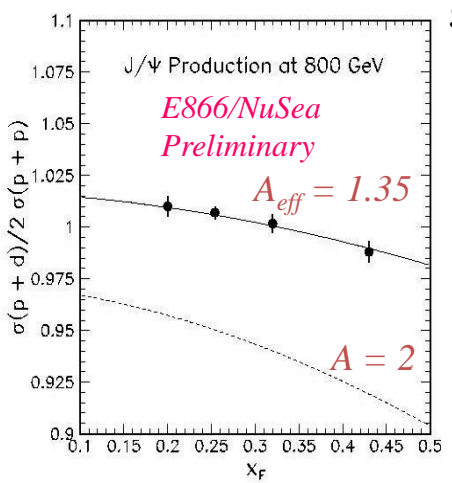


J/Psi suppression doesn't scale like shadowing

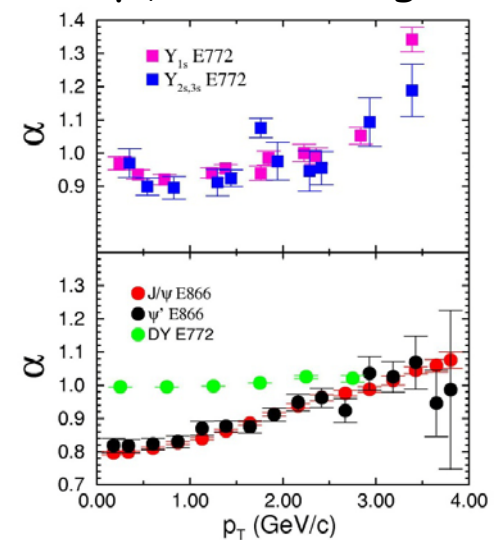
Open vs closed charm isolates initial state



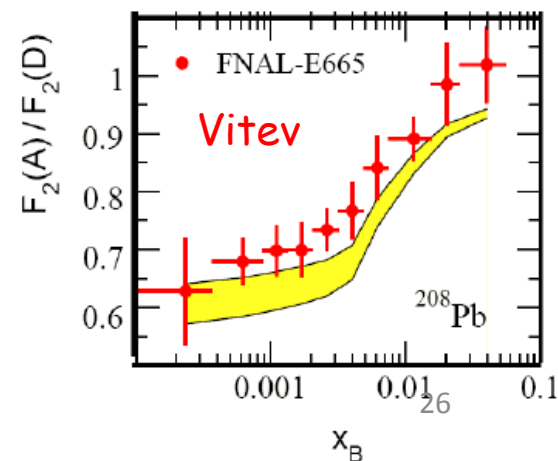
p+d J/Psi suppression is weak



everything has p_T broadening



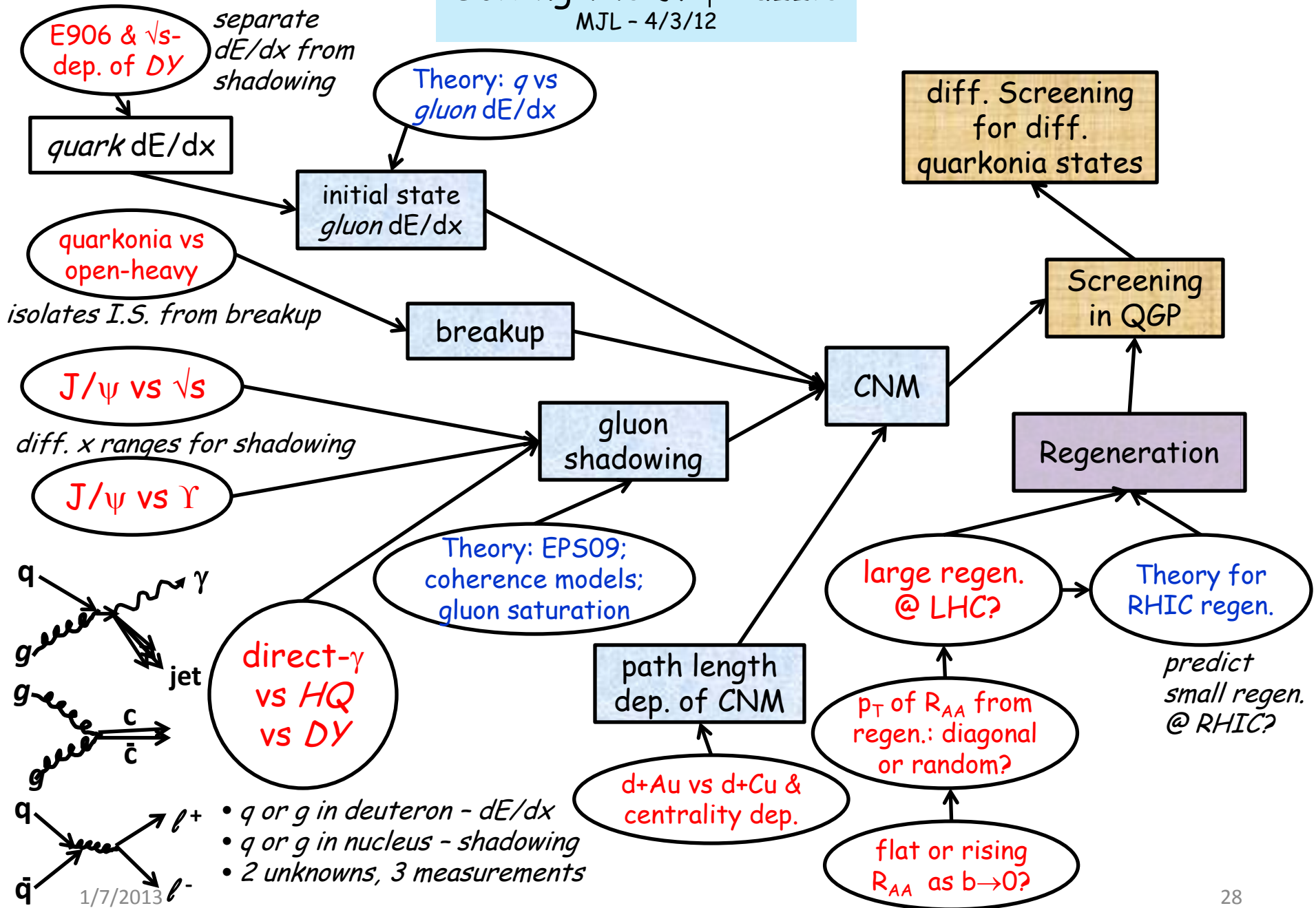
Shadowing may flatten for x < 10^-3?



Backups

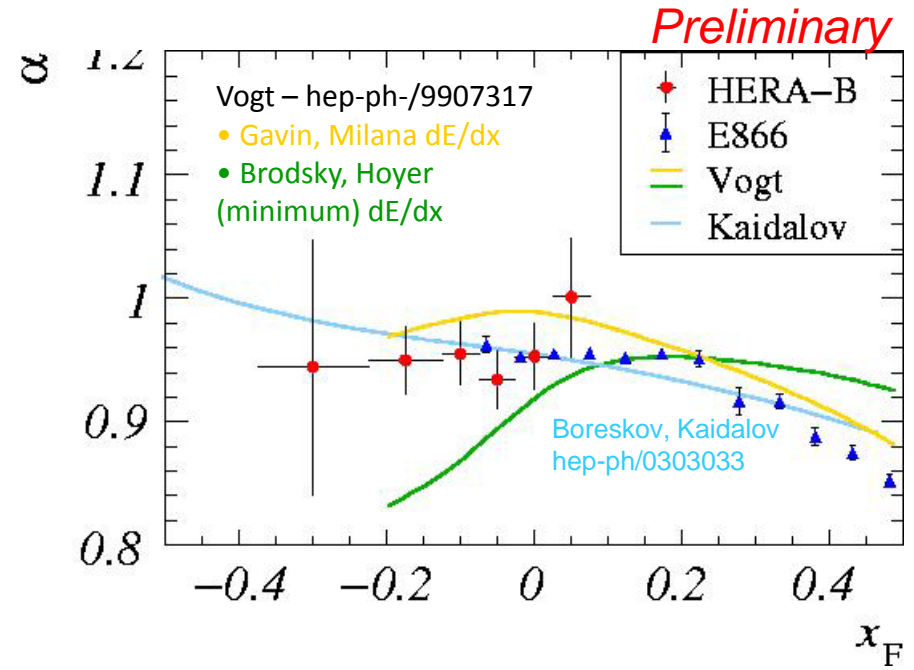
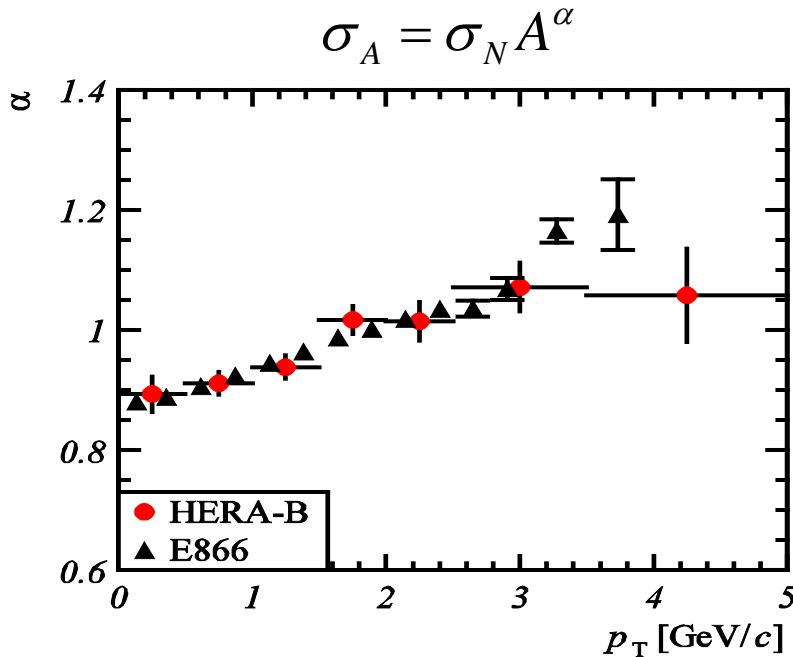
Solving the J/ ψ Puzzle

MJL - 4/3/12



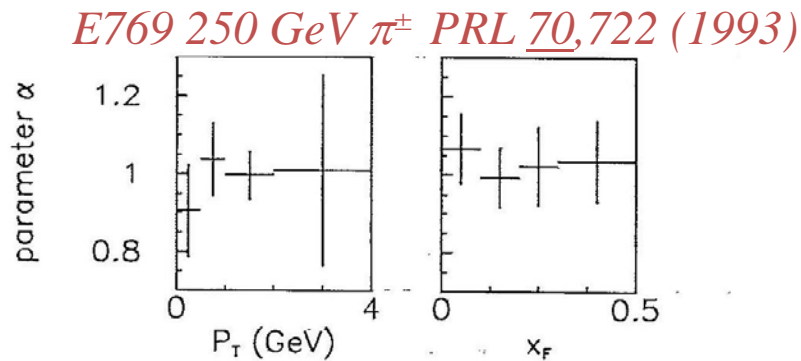
HERA-B - J/ψ A dependence

A. Zoccoli (HERA-B) – talk @ Hard Probes 2004



- Previous result of FNAL E866 extended to $x_F = -0.35$
- Result from 15% of full $\mu^+ \mu^-$ sample, statistical uncertainties only, similar results for e^+e^-
- Work on systematics ongoing. Complete the analysis on the full data sample.

Open Charm Nuclear Dependence : x_F Dependence?



WA82 340 GeV π PRB 284,453 (1992)

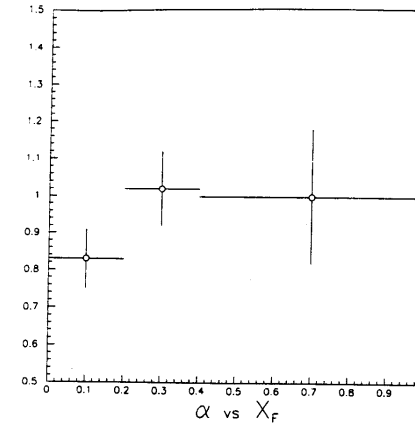


FIG. 4. Dependence of the parameter α on P_T and x_F for D^+ and D^0 .

Fig. 4. The parameter α for charm production as measured in three x_F intervals.

Vogt et al., NP 383,643 (1992)

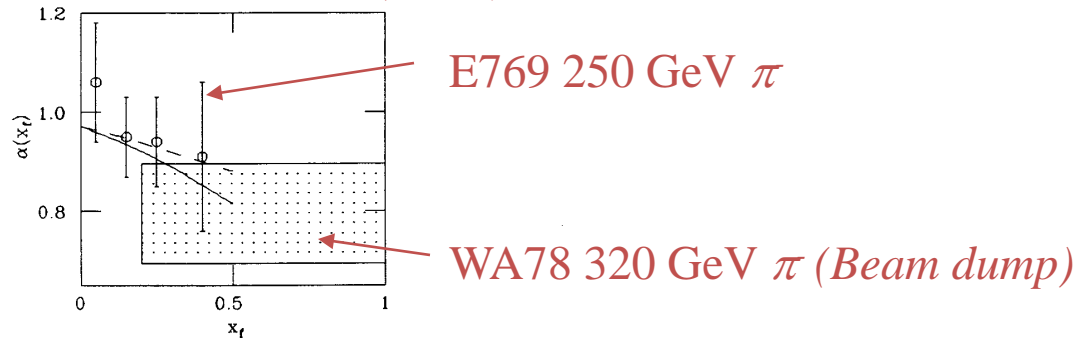
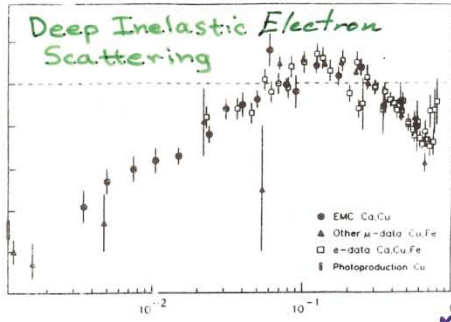


Fig. 14. We show $\alpha(x_f)$ for π^-A interactions at 300 GeV as calculated in our model. The dashed curve shows delta function fragmentation, the solid curve, the Peterson function. These results are compared to those for D^\pm mesons produced by 250 GeV π^-A interactions [3] and the effective α found by the WA78 beam dump experiment [23], indicated by the band with $\langle x_f \rangle = 0.31$.

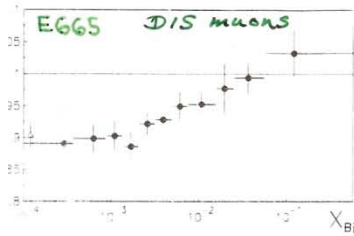
What About A-dependence of Drell-Yan??



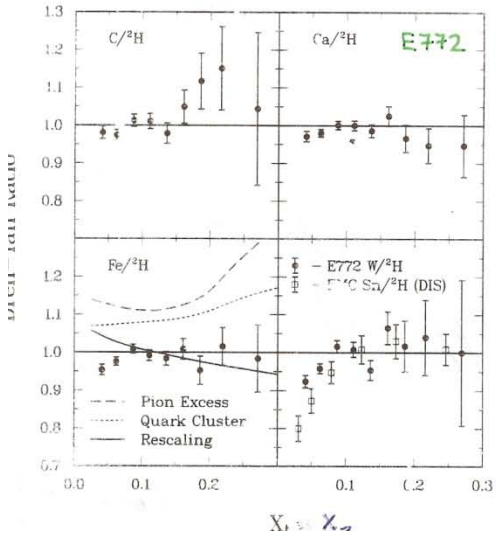
SHADOWING

$x < 0.1$ quarks suppressed in nuclei.

Gluons probably also suppressed.

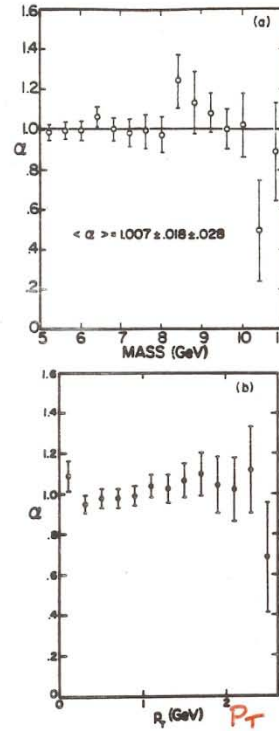


E665
Deep Inelastic Scattering
(470 GeV/c μ)
Z. Phys. C67, 403 (1995)



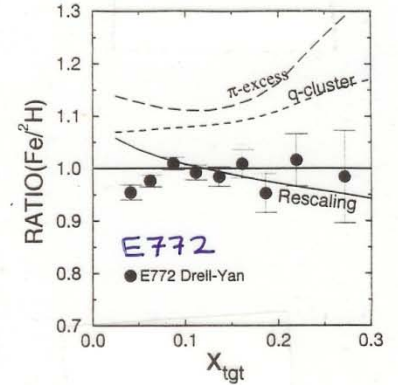
Explains low-x Drell-Yan

E772
 $P+A \rightarrow \mu^+\mu^-$
@ 800 GeV/c
PRL 64, 2479 (1990)

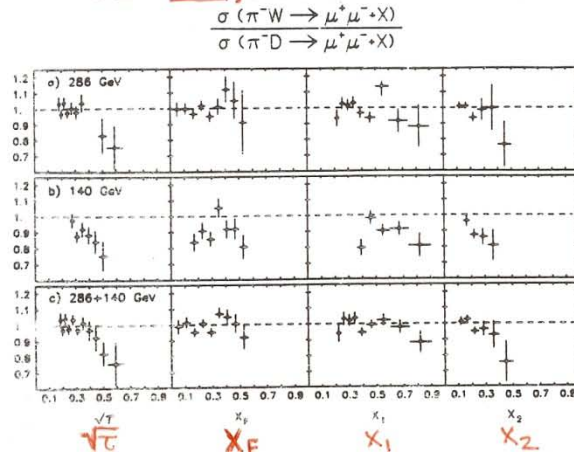


E288 - 400 GeV protons
PRD 23, 604 (1981).

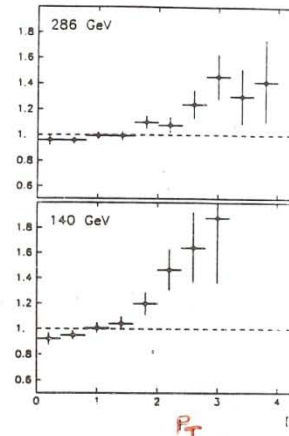
NUCLEAR DEPENDENCE OF DRELL-YAN.



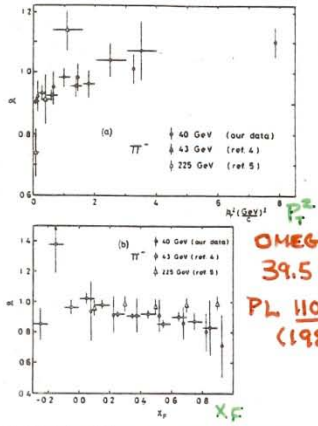
NA10 - 140 GeV/c + 288 GeV/c π^-
PL B 193, 368 (1987).



$$\frac{\sigma(\pi^- W \rightarrow \mu^+ \mu^- X)}{\sigma(\pi^- D \rightarrow \mu^+ \mu^- X)}$$



**MORE J/ψ
NUCLEAR DEP.**



**OMEGA Spectr.
39.5 GeV/c π⁻
PL 110B, 415
(1982).**

**NA3
Z. Phys.
20, 101
(1983).
150 ±
200 ±
280 GeV/c**

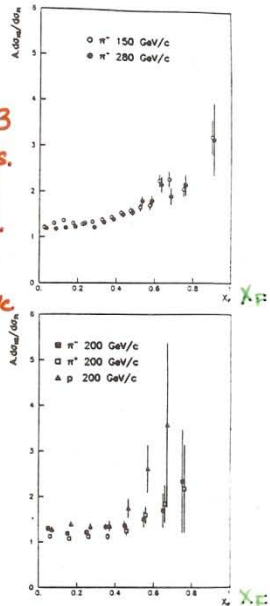
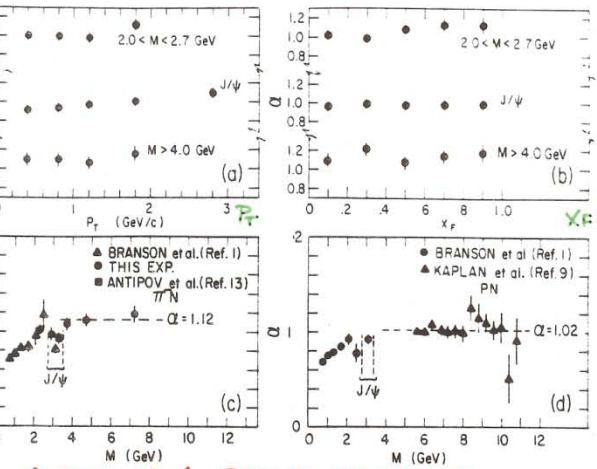
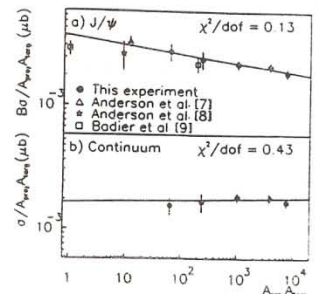


Fig. 2. Ratio of J/ψ cross sections on hydrogen and platinum $R(x_F) = A \frac{d\sigma(H)}{dx_F} / \frac{d\sigma(Pt)}{dx_F}$ for the various π⁻, π⁺, p data of the experiment

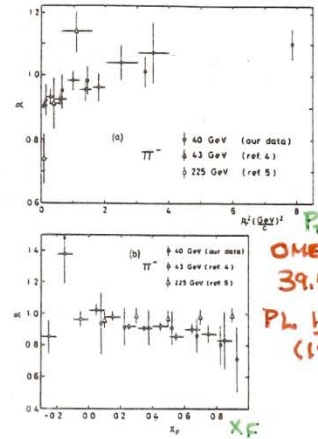


**Anderson et al. PRL 42, 944 (1979).
225 GeV/c π[±], K[±], p[±]**

**NA38
PL 270, 105
(1991).**



**MORE J/ψ
NUCLEAR DEP.**



**OMEGA Spectr.
39.5 GeV/c π⁻
PL 110B, 415
(1982).**

Fig. 3. Variation of A-dependence parameter $\alpha, A^\alpha = \sigma(W)/\sigma(H_2)$, for π⁻ induced J/ψ as a function of p_T² and x_F. No corrections have been made for the effects of Fermi motion.

**NA3
Z. Phys.
20, 101
(1983).
150 ±
200 ±
280 GeV/c**

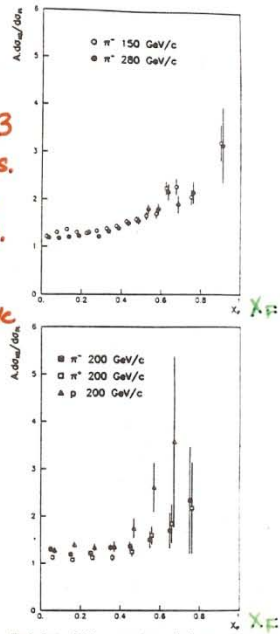
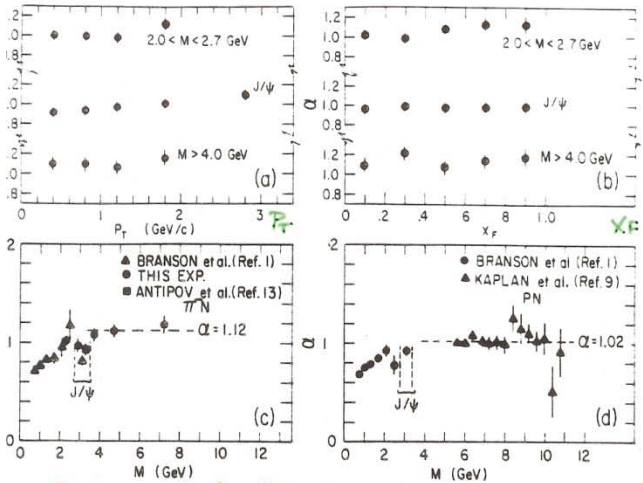
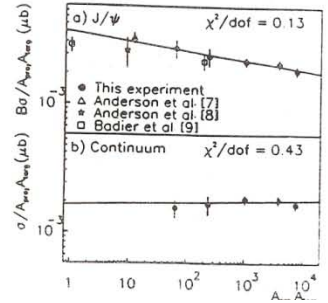


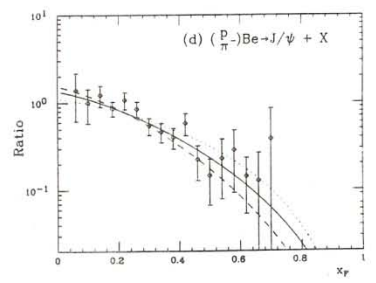
Fig. 2. Ratio of J/ψ cross sections on hydrogen and platinum $R(x_F) = A \frac{d\sigma(H)}{dx_F} / \frac{d\sigma(Pt)}{dx_F}$ for the various π⁻, π⁺, p data of the experiment



**Anderson et al. PRL 42, 944 (1979).
225 GeV/c π[±], K[±], p[±]**

**NA38
PL 270, 105
(1991).**





NUCLEAR DEPENDENCE OF J/ψ PRODUCTION

E672/706
530 GeV/c π^+, p
(FNAL-Pub-91/62-E,
FNAL-Conf-89/151-E)

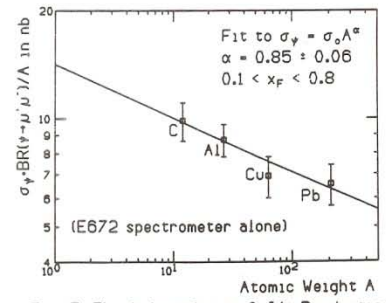
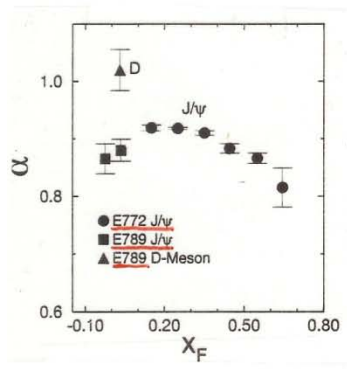
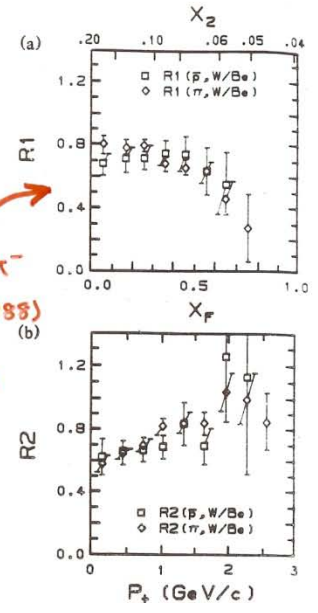


Fig. 7. The A-dependence of J/ψ Production



E537
125 GeV/c \bar{p}, π^-
PRL 60, 2121 (1988)



EMC - Theoretical Models

Nuclear Binding & Fermi Motion

Excess pions (nuclear binding)

- Loss of valence quark momentum & enhanced sea (pions)
- Pions have valence (large x) anti-quarks so anti-quark momentum is enhanced

Fermi motion

- Spreads quark momentum near x=1

Rescaling (Close et al.; Jaffe)

$$F_2^{Fe}(x, Q^2) \cong F_2^D(x, \xi Q^2), \quad \xi \approx 2$$

~15% phenomenological increase in confinement scale

Overlapping nucleons (6-quark clusters, Close, Jaffe,...)

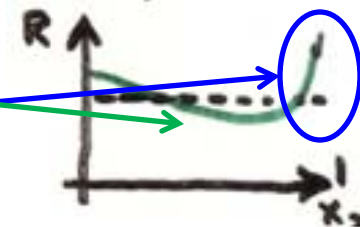
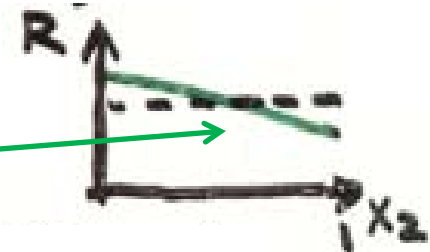
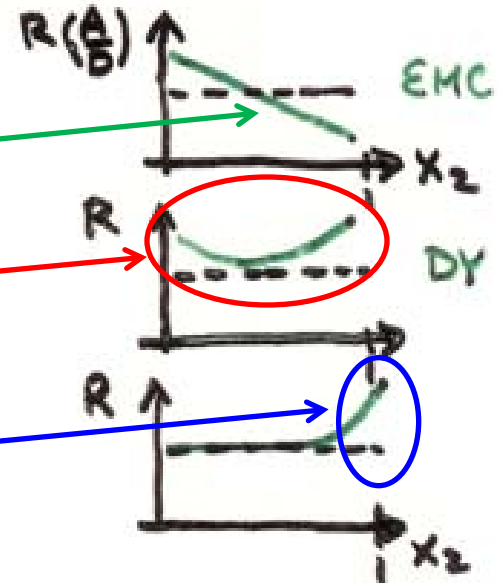
- Loss of valence quark momentum

Multi-quark Clusters (Pirner & Vary; Carlson & Havens)

Loss of valence quark momentum

relative enhancement at x~1 (one quark can have > nucleon momentum)

(Ericson & Thomas
Bickerstaff & Thomas
Berger & Coester)



Sudakov suppression

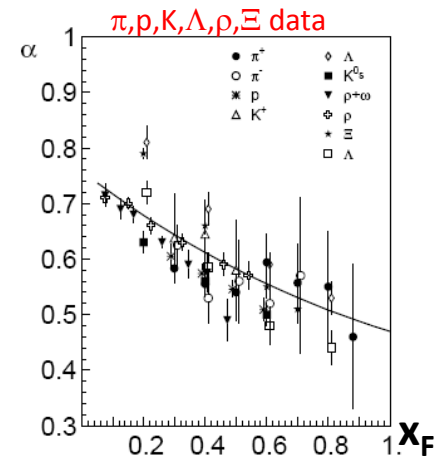
Kopeliovich hep-ph/0501260

Universal suppression at large x_F seen in data for various reactions:

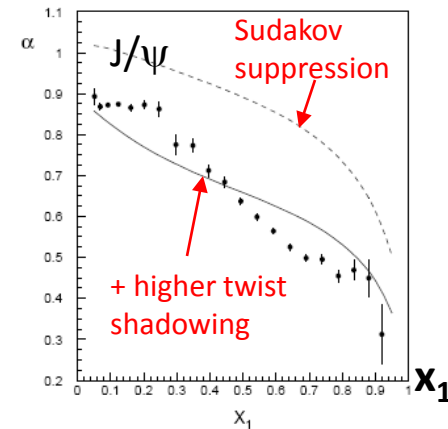
- forward light hadrons, Drell-Yan, heavy flavor
- often attributed to shadowing since x_2 is small

But this common suppression mechanism can also be viewed as Sudakov suppression :

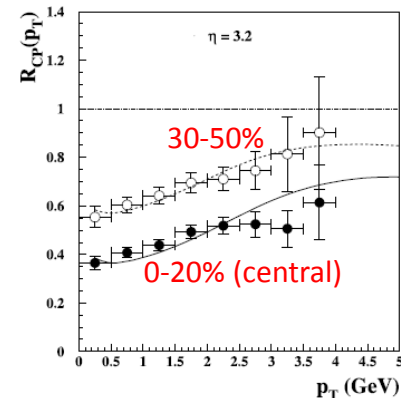
- no particles produced as $x_F \rightarrow 1$ due to energy conservation
- more multiple interactions make the effect larger in nuclei



describes universal suppression vs x_F for low energy data

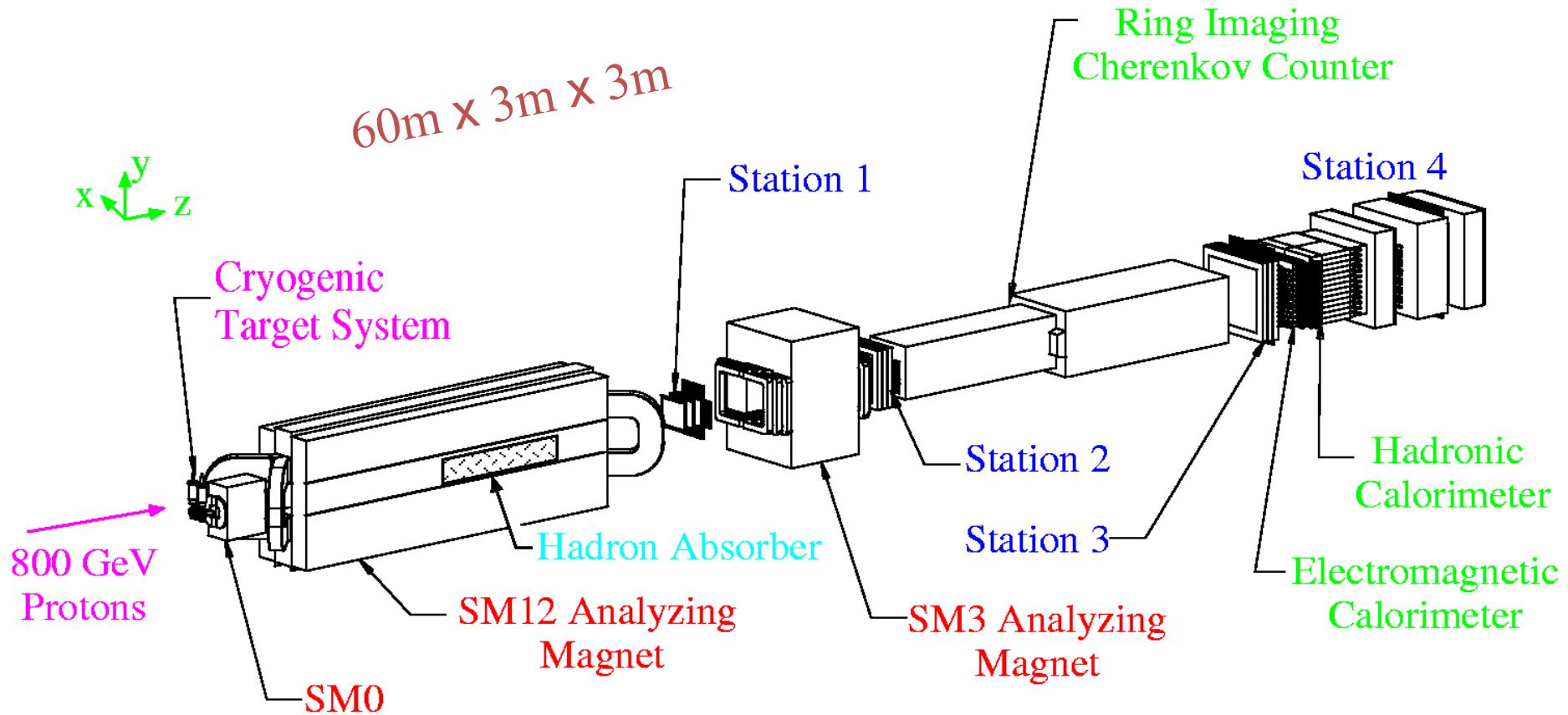


Close to 800 GeV J/ψ suppression vs x_1 ($x_1 \approx x_F$)



and Brahm's forward rapidity ($\eta = 3.2$) hadrons

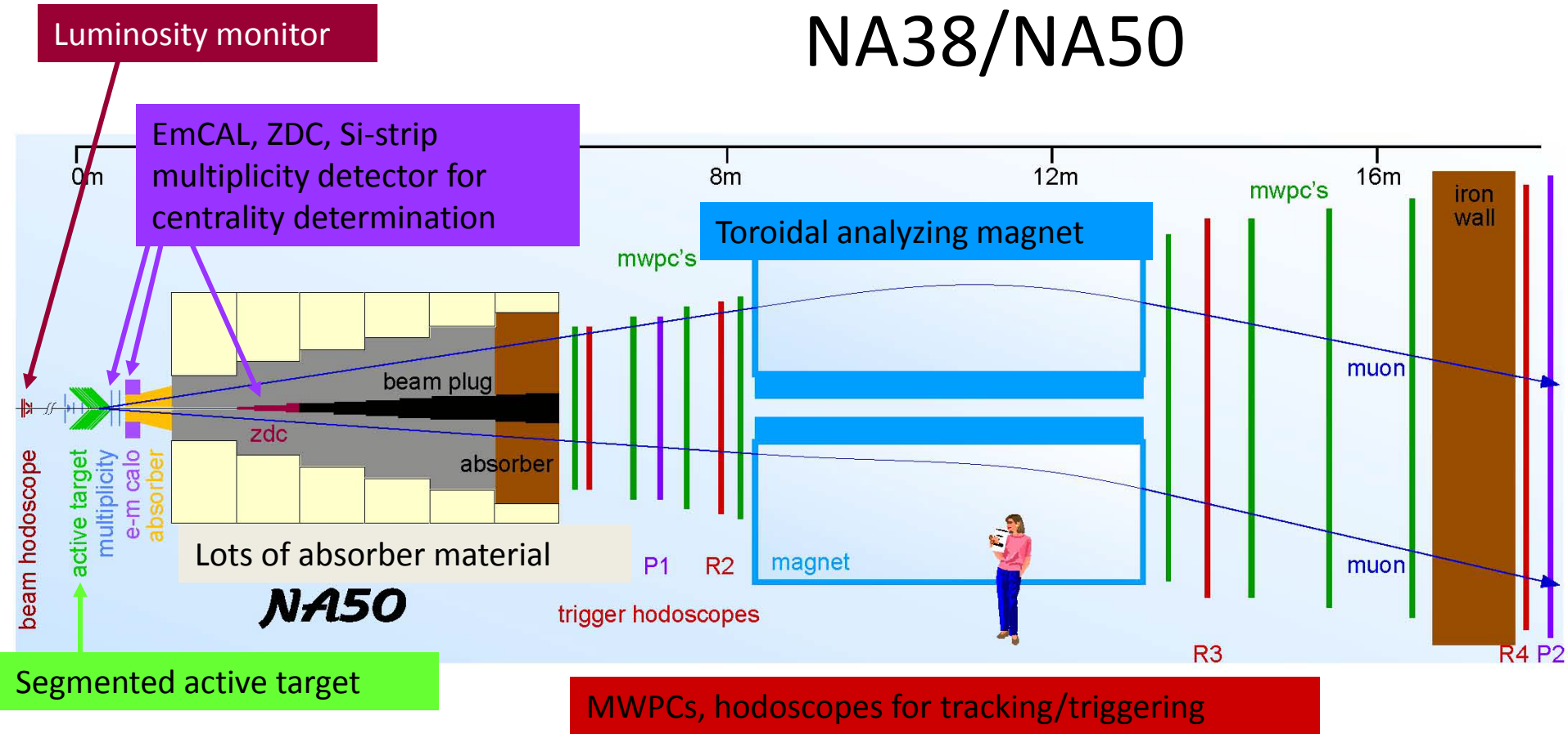
Fermilab E866/NuSea Detector



- Forward x_F high mass μ -pair spectrometer
- Liquid hydrogen and deuterium targets
- Two acceptance defining magnets (SM0, SM12)

- Beam dump (4.3m Cu)
- Hadronic absorber (13.4 I_0 -Cu, C, CH₂)
- Momentum analyzing magnet (SM3)
- Three tracking stations
- Muon identifier wall & 4th tracking

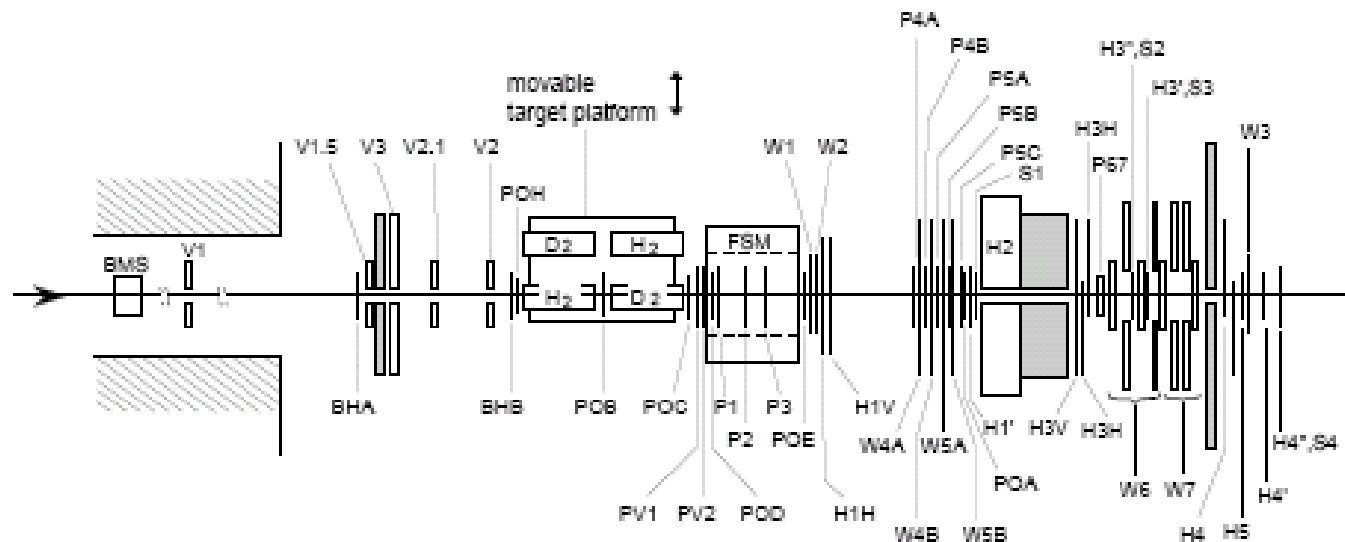
NA38/NA50



Note benefits of a fixed target experiment:

- Lots of luminosity.
- Boost in lab frame gives very high momentum particles, so resolution is improved and background can be greatly reduced by many meters of absorber.

1989 NMC SPECTROMETER (Top view)



- BMS
 - V1, V1.5, V3, V2.1, V2
 - BHA, BHB
 - POA-E, POH, PV1-2, P1-3, P4A-SC, P67
 - FSM
 - W1-3, W4A-5B, W6-7
 - H1H, H1V, H3V, H3H, H4, H5
 - H1', H3', H4'
 - S1-4, H3'', H4''
 - H2
- Beam momentum spectrometer
 - Veto counters
 - Beam hodoscopes
 - Proportional chambers
 - Forward spectrometer magnet
 - Drift chambers
 - Large angle trigger (T1) hodoscopes
 - Small angle trigger (T2) hodoscopes
 - Small x trigger (T14) hodoscopes
 - Hadron calorimeter
 - Iron absorbers

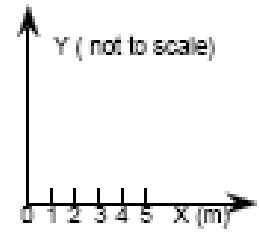
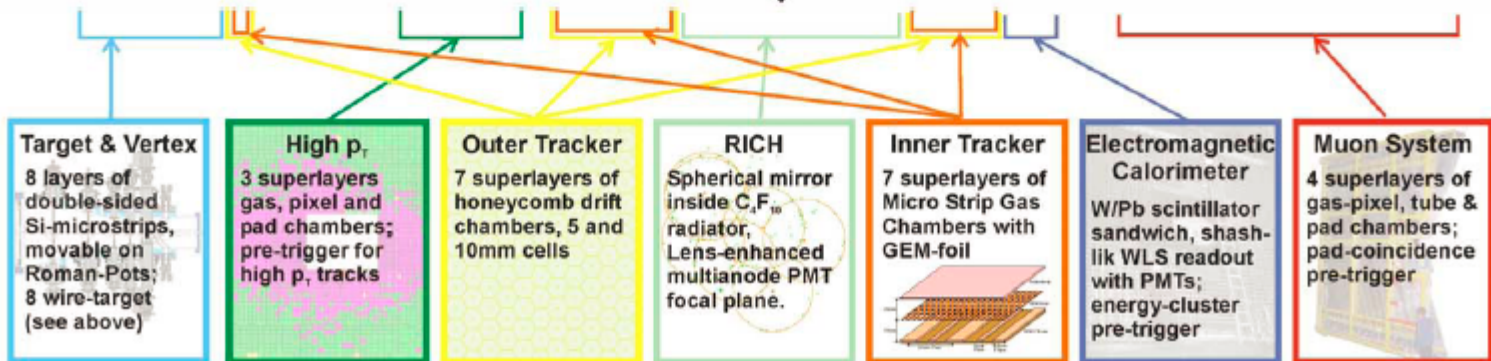
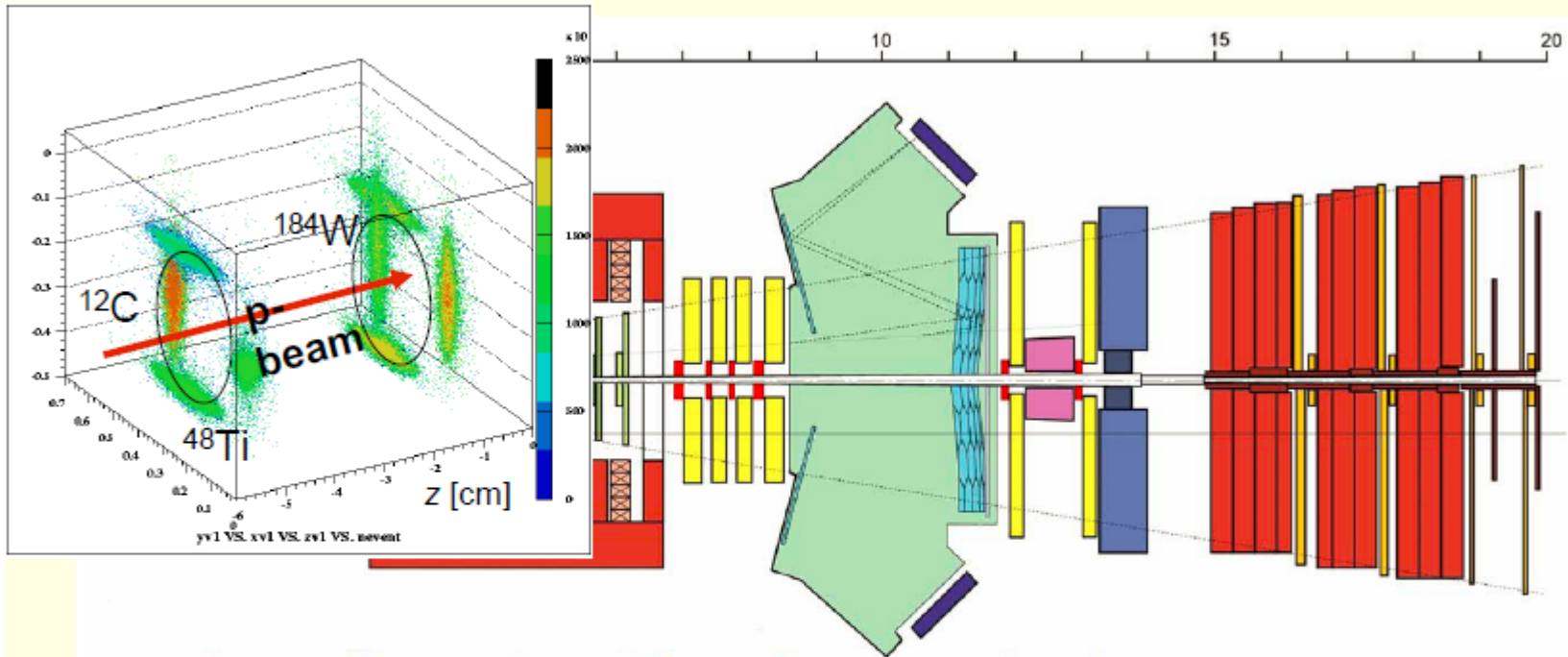
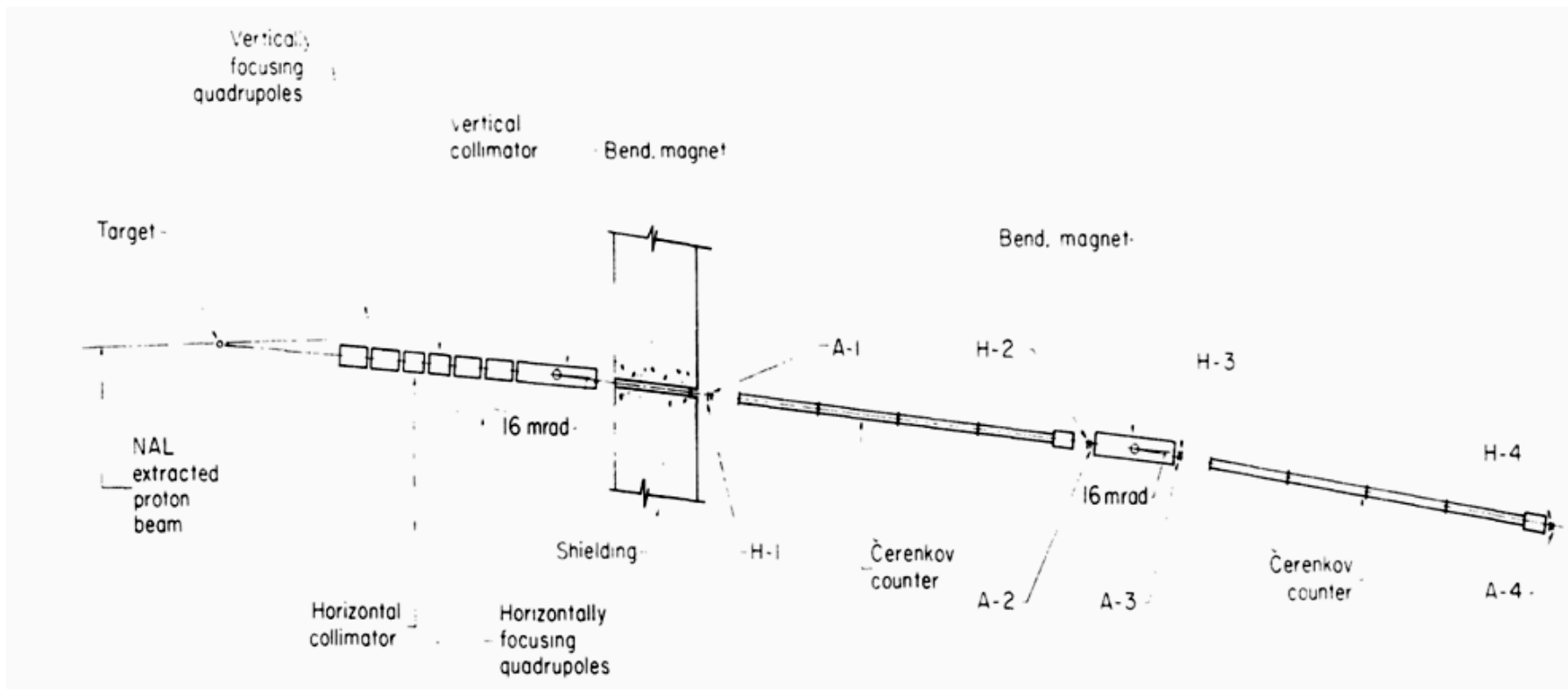


Figure 1: The NMC spectrometer for the 1989 data taking. The beam calibration spectrometer is located downstream and not shown.

The HERA-B Detector



Production of hadrons at large transverse momentum at 200, 300 and 400 GeV at Fermilab J.W. Cronin et al., PRD 11, 3105 (1975)



A Dependence of J/ψ and ψ' Not Identical: Size Matters

Color octet mechanism suggested that J/ψ and ψ' A dependence should be identical — Supported by large uncertainties of early data

More extensive data sets (NA50 at SPS, E866 at FNAL) show clear difference at midrapidity [NA50 ρ_L fit gives $\Delta\sigma = \sigma_{\text{abs}}^{\psi'} - \sigma_{\text{abs}}^{J/\psi} = 4.2 \pm 1.0$ mb at 400 GeV, 2.8 ± 0.5 mb at 450 GeV for absolute cross sections]

Suggests we need to include formation time effects

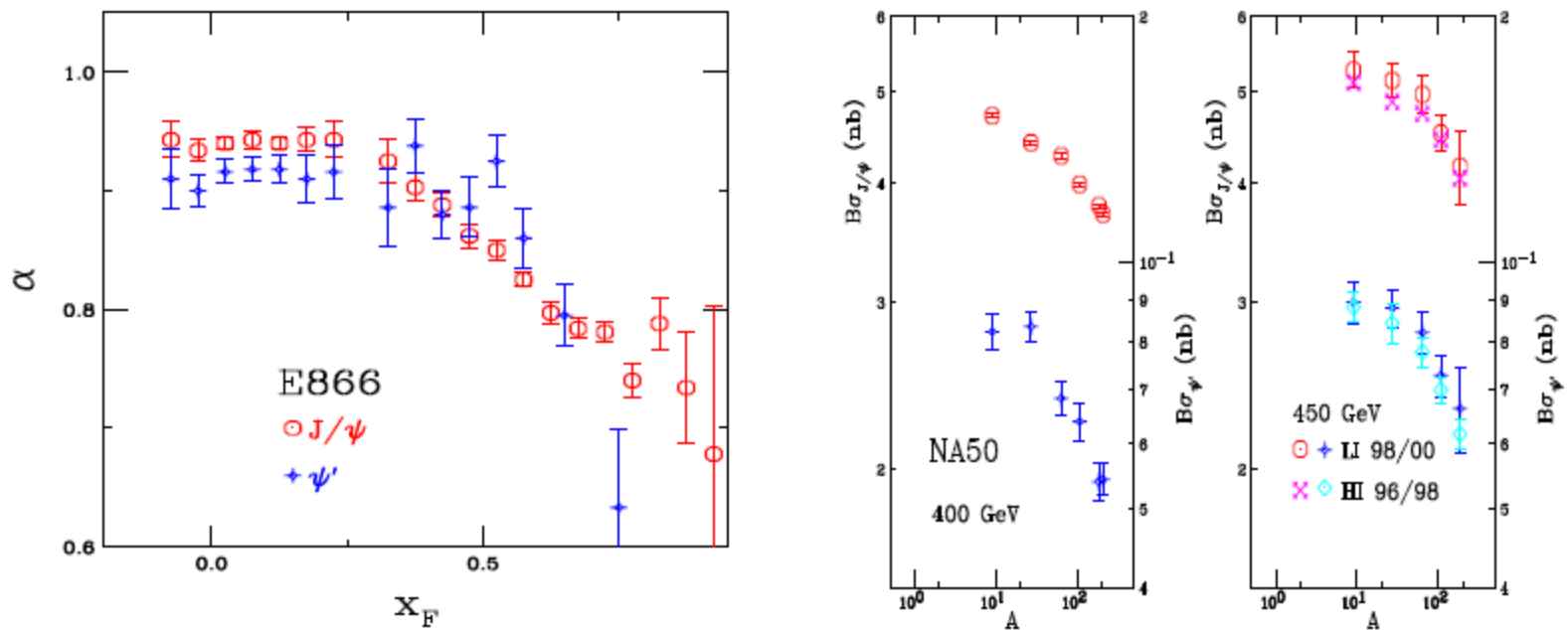
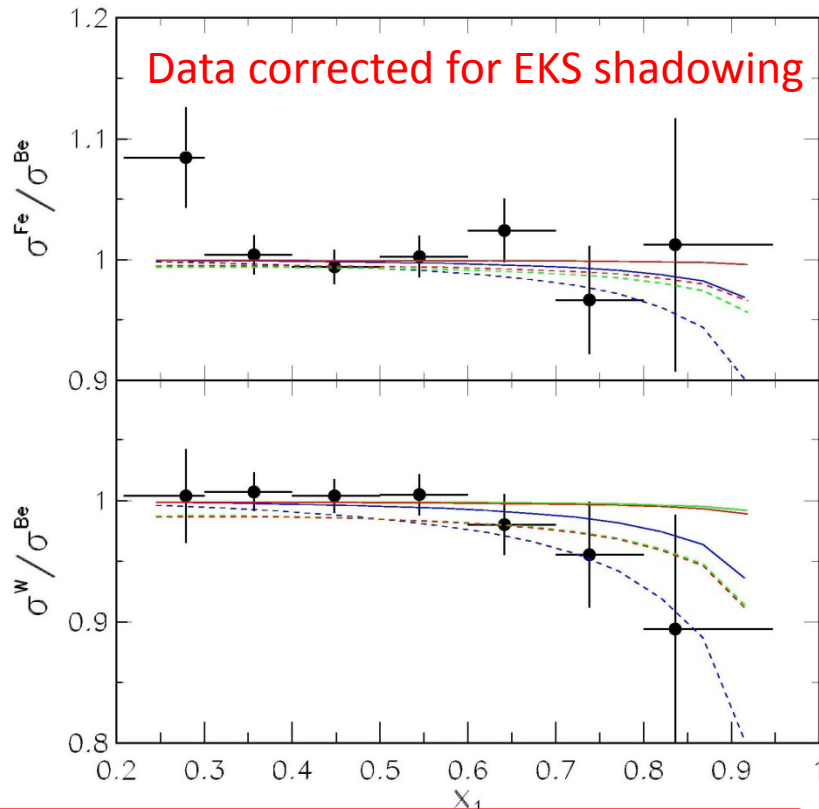


Figure 15: The J/ψ A dependence (left) as a function of x_F at FNAL ($\sqrt{s_{NN}} = 38.8$ GeV) and (right) and a function of A at the SPS (NA50 at $p_{\text{lab}} = 400$ and 450 GeV) for J/ψ and ψ' production.

Quark Energy Loss from Drell-Yan data assuming EKS Shadowing

Parton energy loss limits



- Gavin and Milana $\Delta E < 0.14 \text{ \%}/\text{fm}$
- Brodsky and Hoyer $\Delta E < 0.44 \text{ GeV}/\text{fm}$
- Baier et al. $\Delta E < (0.046 \text{ GeV}/\text{fm}^2) \times L^2$

Phys. Rev. Lett. 83, 2304 (1999)

- S.Gavin, J.Milana, Phys.Rev.Lett 68(1992)1834

$$\Delta x_1 = kC x_1 \left(\frac{Q_0}{Q} \right)^n A^{\frac{1}{3}}$$

$C = 4/3$ for quarks, $C = 3$ for gluons, $n =$ for "soft physics" Experimental fit:

$$\Delta x_1 = -k_1 x_1 A^{\frac{1}{3}}$$

- S.Brodsky, R.Hoyer, Phys.Lett. B298(1993)165

$$|\Delta x_1| \leq \frac{\langle k_{\perp} \rangle L_a}{2E}$$

Experimental fit: $\Delta x_1 = -\frac{k_2}{s} A^{\frac{1}{3}}$

- R.Baier et al., Nucl.Phys. B484(1997)265
R.Baier et al., Nucl.Phys. B531(1998)403

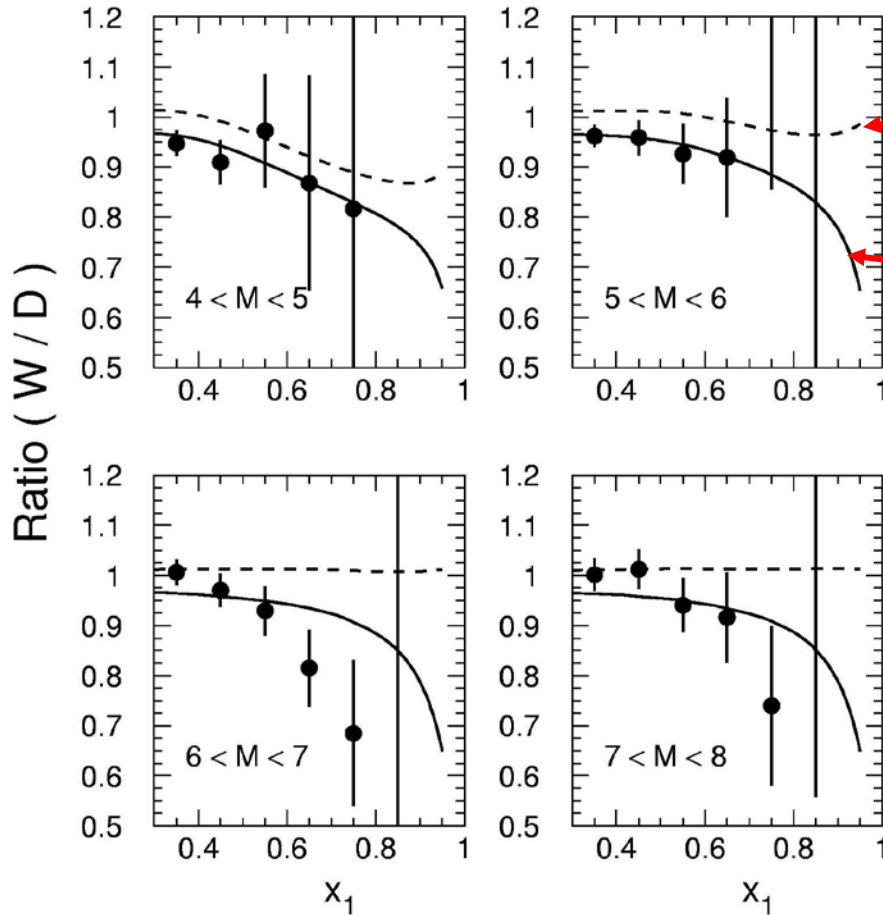
$$|\Delta x_1| = \frac{3\alpha_s}{2} \frac{m_p}{s} L_a \langle p_{\perp}^2 \rangle$$

Experimental fit: $\Delta x_1 = -\frac{k_3}{s} A^{\frac{2}{3}}$

Quark Energy Loss from Drell-Yan data assuming Kopeliovich Shadowing

Johnson, Kopeliovich et al., PRL 86, 4483 (2001)

data from
E772 - PRL 64, 2479 (1990)



Shadowing
dE/dx + Shadowing

Analysis of our p-A Drell-Yan data using the Kopeliovich model. Dashed lines with shadowing only; solid lines with parton energy loss of

$$dE/dx = 2.32 \pm 0.52 \pm 0.5 \text{ GeV/fm}$$

Confusion due to having both shadowing and dE/dx contributing to Drell-Yan suppression for measurements at this energy

Intrinsic charm contribution to Quarkonia

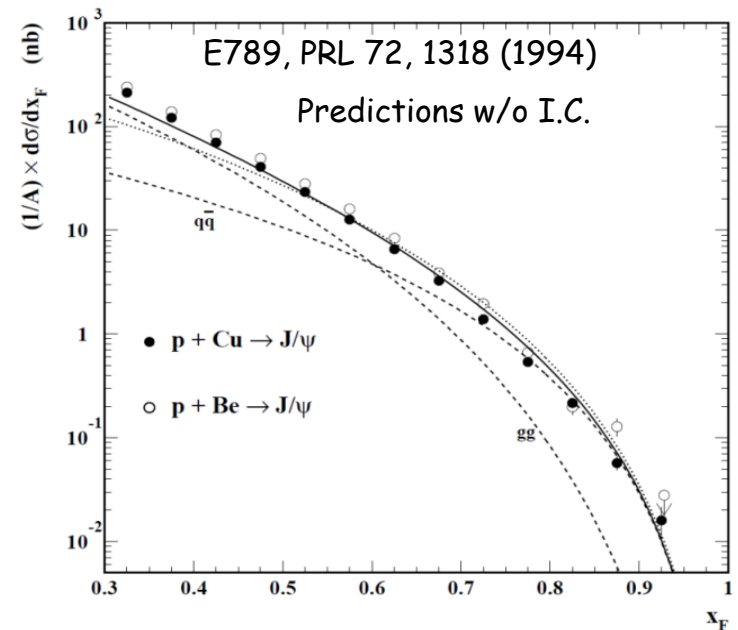
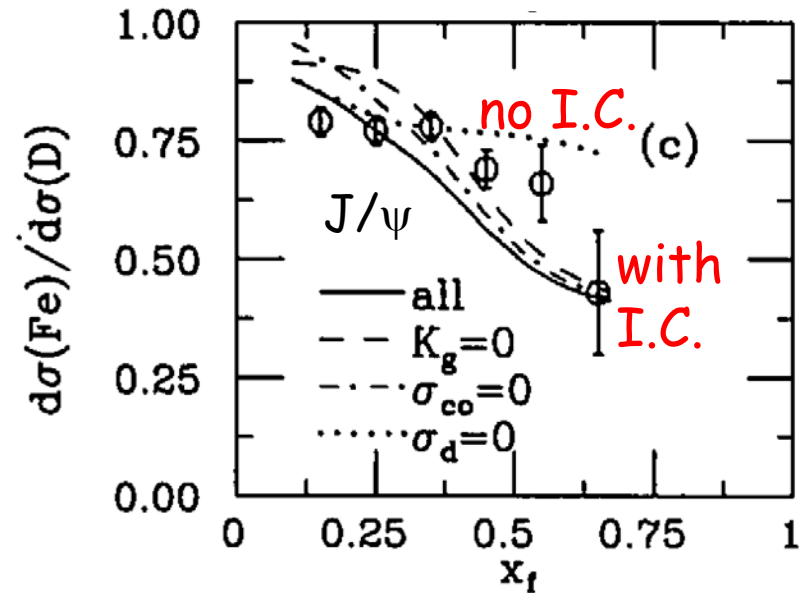
At large $x_F (\geq 0.5)$ intrinsic $c\bar{c}$ components of the projectile proton can dominate the production of charm pairs

- $A^{2/3}$ dependence via surface stripping of light quarks to free charm pair component

Vogt, Brodsky, Hoyer, NP B360, 67 (1991)
(also includes absorption and shadowing)

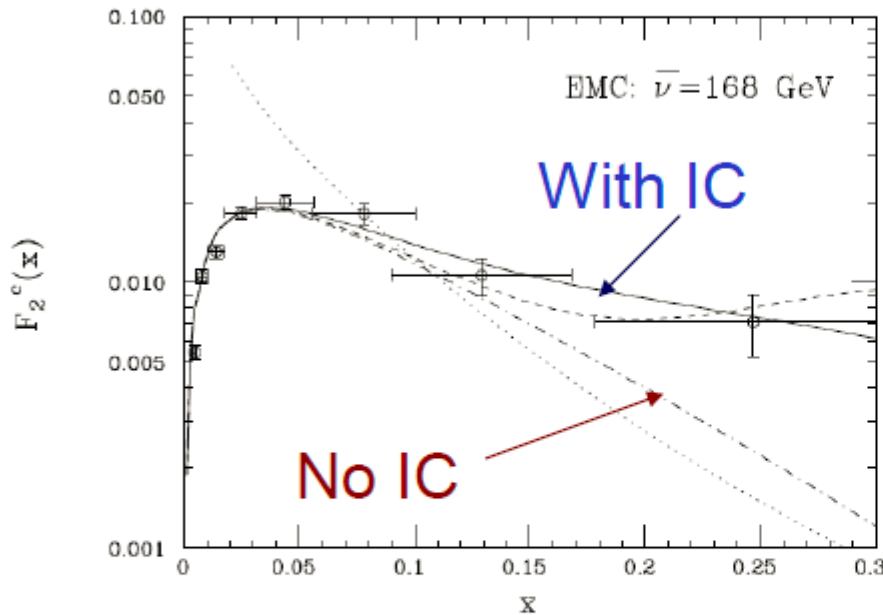
But E789 set limit on I.C. contribution via shape of cross section vs x_F

- $< 2.3 \times 10^{-3}$ nb/nucleon (1.8 nb/nucleon predicted)

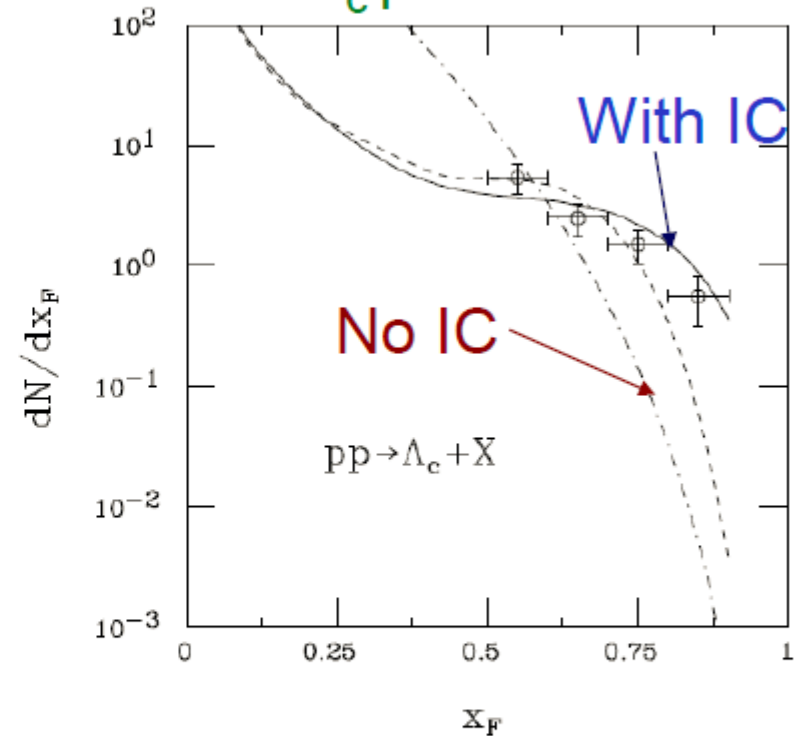


“Evidence” for the “intrinsic” charm (IC)

DIS data



Λ_c production



Gunion and Vogt (hep-ph/9706252)

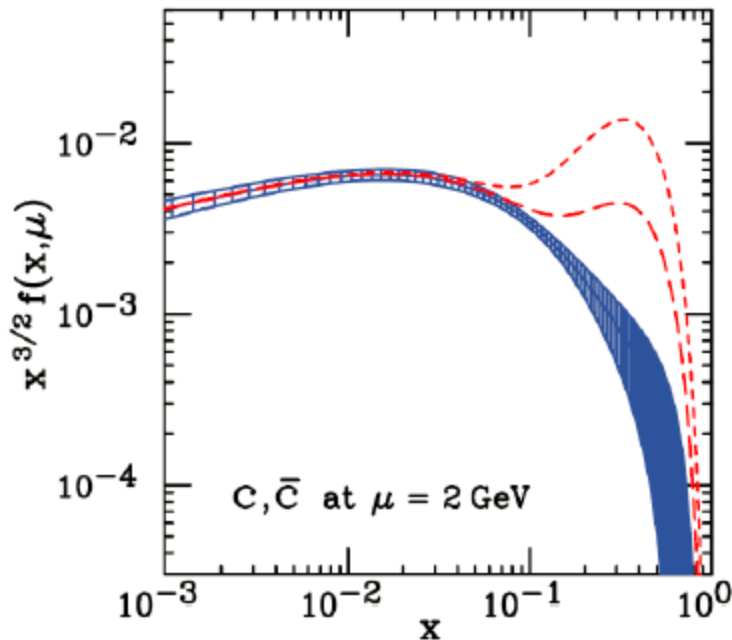
“Evidence” appears to be rather weak

No conclusive experimental evidence for intrinsic-charm

PHYSICAL REVIEW D 75, 054029 (2007)

Charm parton content of the nucleon

J. Pumplin,^{1,*} H. L. Lai,^{1,2,3} and W. K. Tung^{1,2}



Blue band corresponds to CTEQ6 best fit, including uncertainty

Red curves include intrinsic charm of 0.57% and 2.0%

We find that the range of IC is constrained to be from zero (no IC) to a level 2–3 times larger than previous model estimates. The behaviors of typical charm distributions within this range are described, and their implications for hadron collider phenomenology are briefly discussed.