

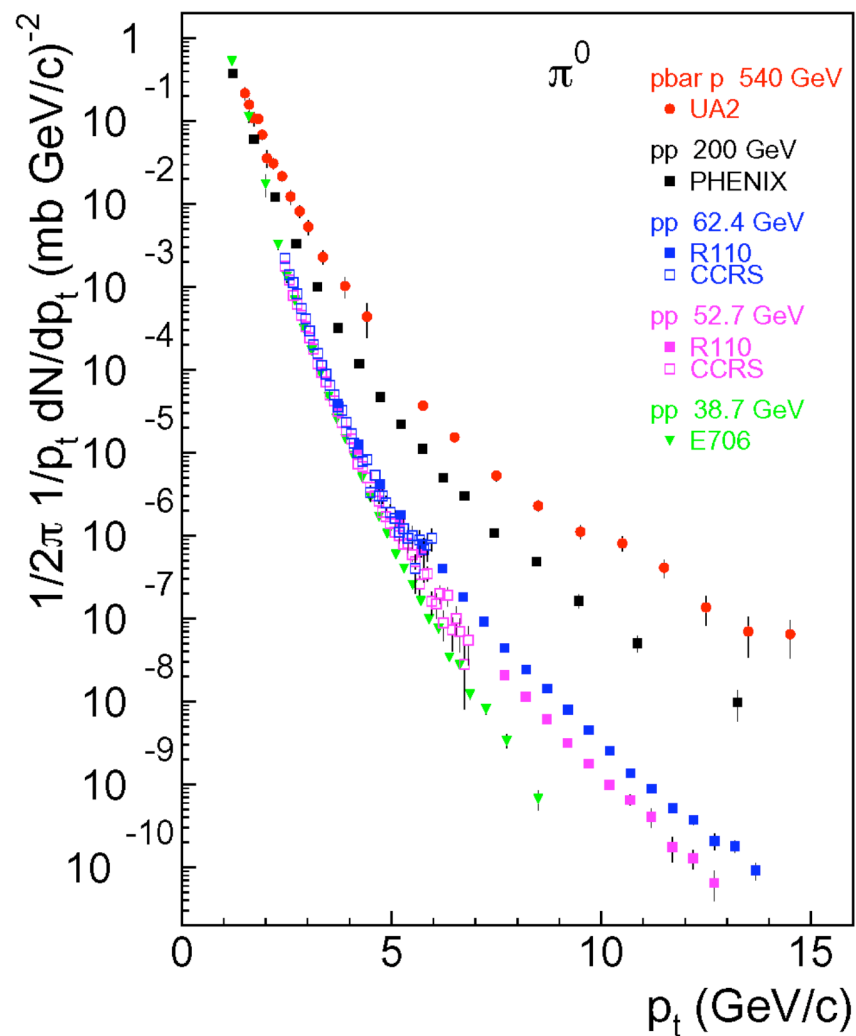
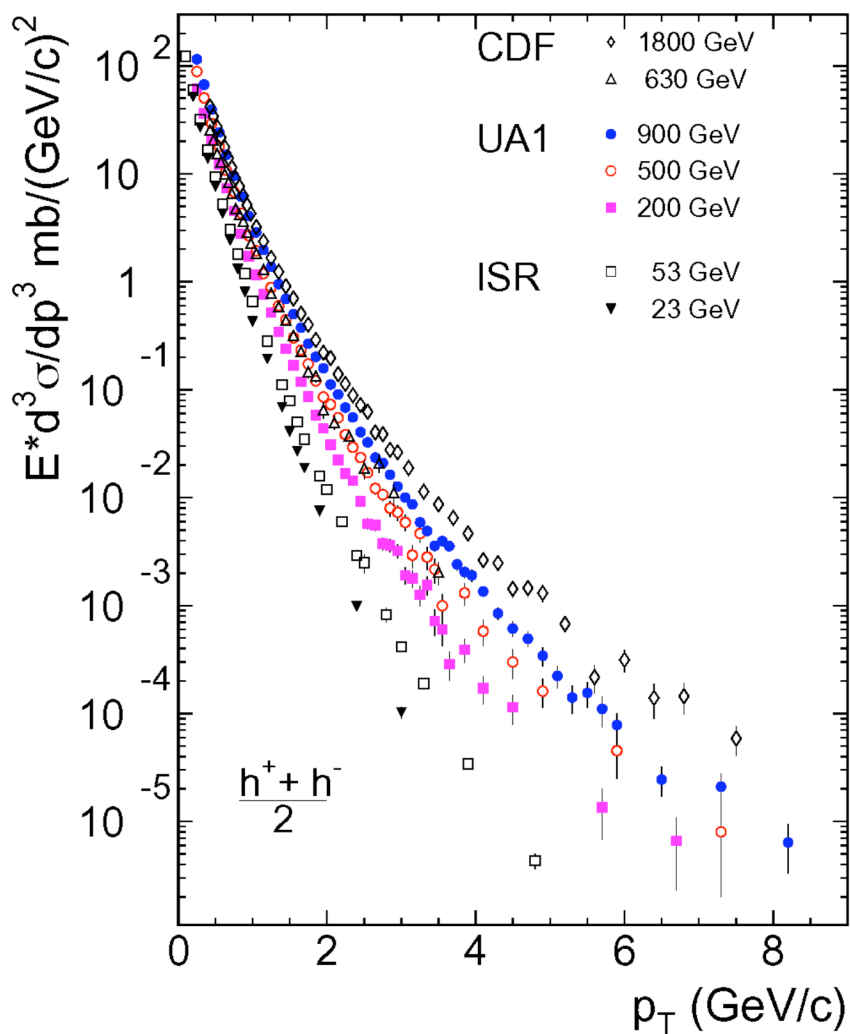
x_T scaling in Au+Au collisions at $s_{NN}=130$ and 200 GeV for π^0 and charged hadrons from PHENIX

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Inclusive single hadron high p_T spectra in p-p collisions



The invariant cross section for the single-particle inclusive reaction $p + p \rightarrow C + X$ where particle C has transverse momentum p_T near mid-rapidity, was given by the general scaling form [54]:

$$E \frac{d^3\sigma}{dp^3} = \frac{1}{p_T^n} F\left(\frac{2p_T}{\sqrt{s}}\right) \quad \text{where} \quad x_T = 2p_T/\sqrt{s}$$

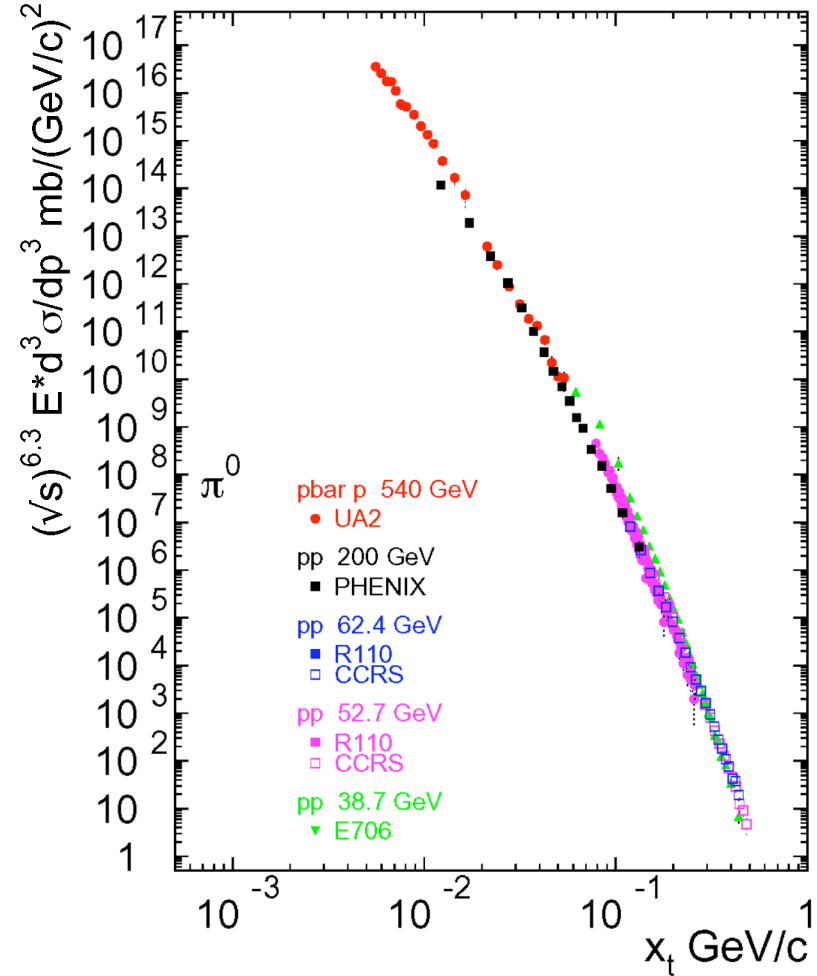
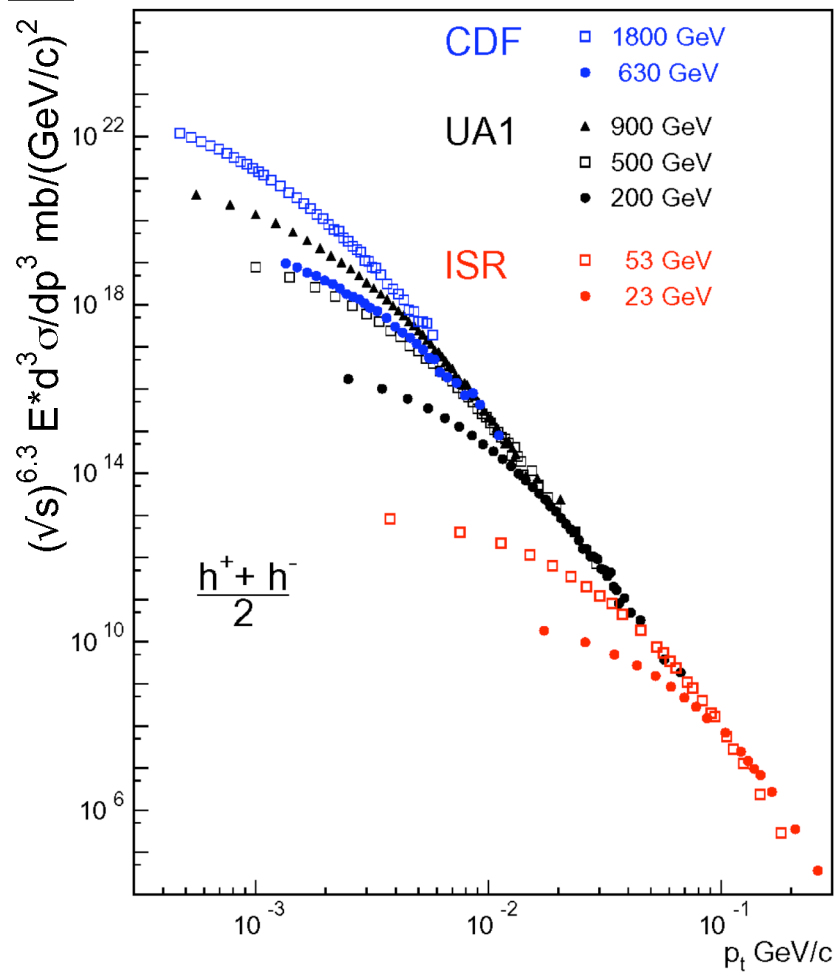
There are 2 factors: a function F which depends only on the ratio of momenta, and a dimensioned factor, p_T^{-n} , where n depends on the quantum exchanged in the hard-scattering. For QED or Vector Gluon exchange [53], $n = 4$. For the case of quark-meson scattering by the exchange of a quark [54], $n=8$.

Inclusion of QCD [58] into the scaling form led to the x_T -scaling law

$$E \frac{d^3\sigma}{dp^3} = \frac{1}{\sqrt{s}^{n(x_T, \sqrt{s})}} G(x_T)$$

where the “ x_T -scaling power” $n(x_T, \sqrt{s})$ should equal 4 in lowest order (LO) calculations, analogous to the $1/q^4$ form of Rutherford Scattering in QED. The structure and fragmentation functions, which scale as the ratios of momenta are all in the $G(x_T)$ term. Due to higher order effects such as the running of the coupling constant, $\alpha_s(Q^2)$, the evolution of the structure and fragmentation functions, and the initial state k_T , measured values of $n(x_T, \sqrt{s})$ in $p + p$ collisions are in the range from 5 to 8.

x_T scaling in p-p collisions $x \sim 0.05-0.10$

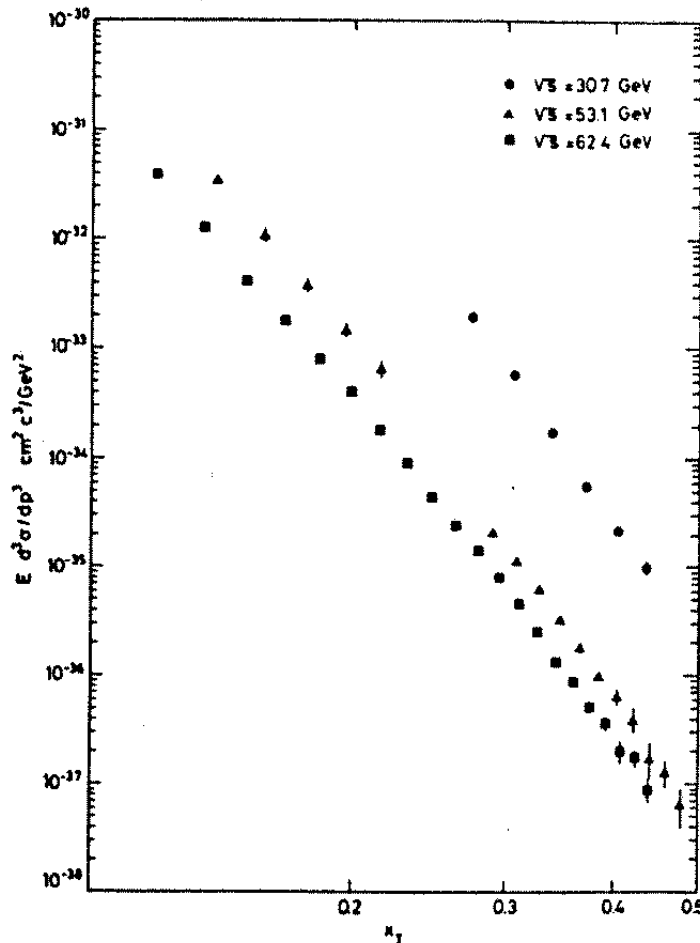


The x_T -scaling of the single particle inclusive data is nicely illustrated by a plot as a function of x_T , with $n(x_T, \sqrt{s}) = 6.3$ of:

$$\sqrt{s}^{n(x_T, \sqrt{s})} \times E \frac{d^3\sigma}{dp^3} = G(x_T)$$

$n(x_T, \sqrt{s} \sim 60 \text{ GeV})$ WORKS, $n \approx 5.1$

Invariant cross section plotted vs $x_T = 2p_T/\sqrt{s}$ p-p collisions



Calculate n from ratio of Invariant cross sections at fixed x_T

$n = 5.1 \pm 0.4$
including all systematic errors

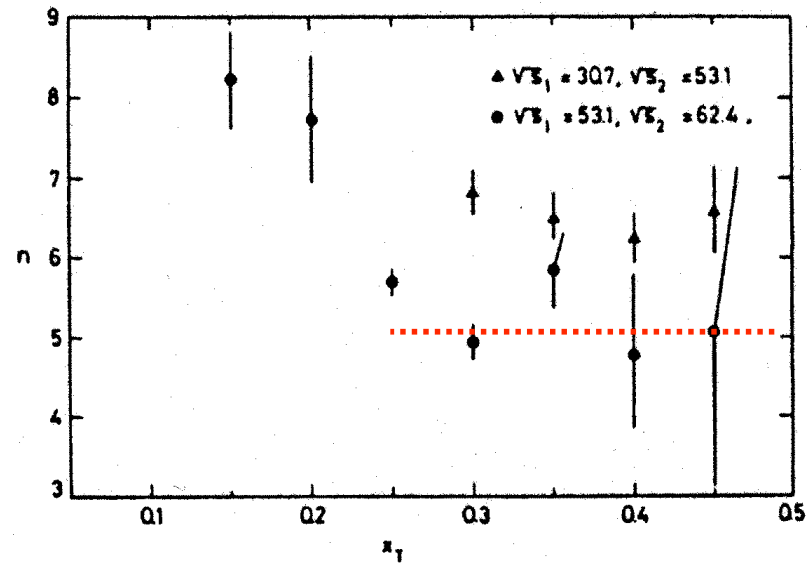
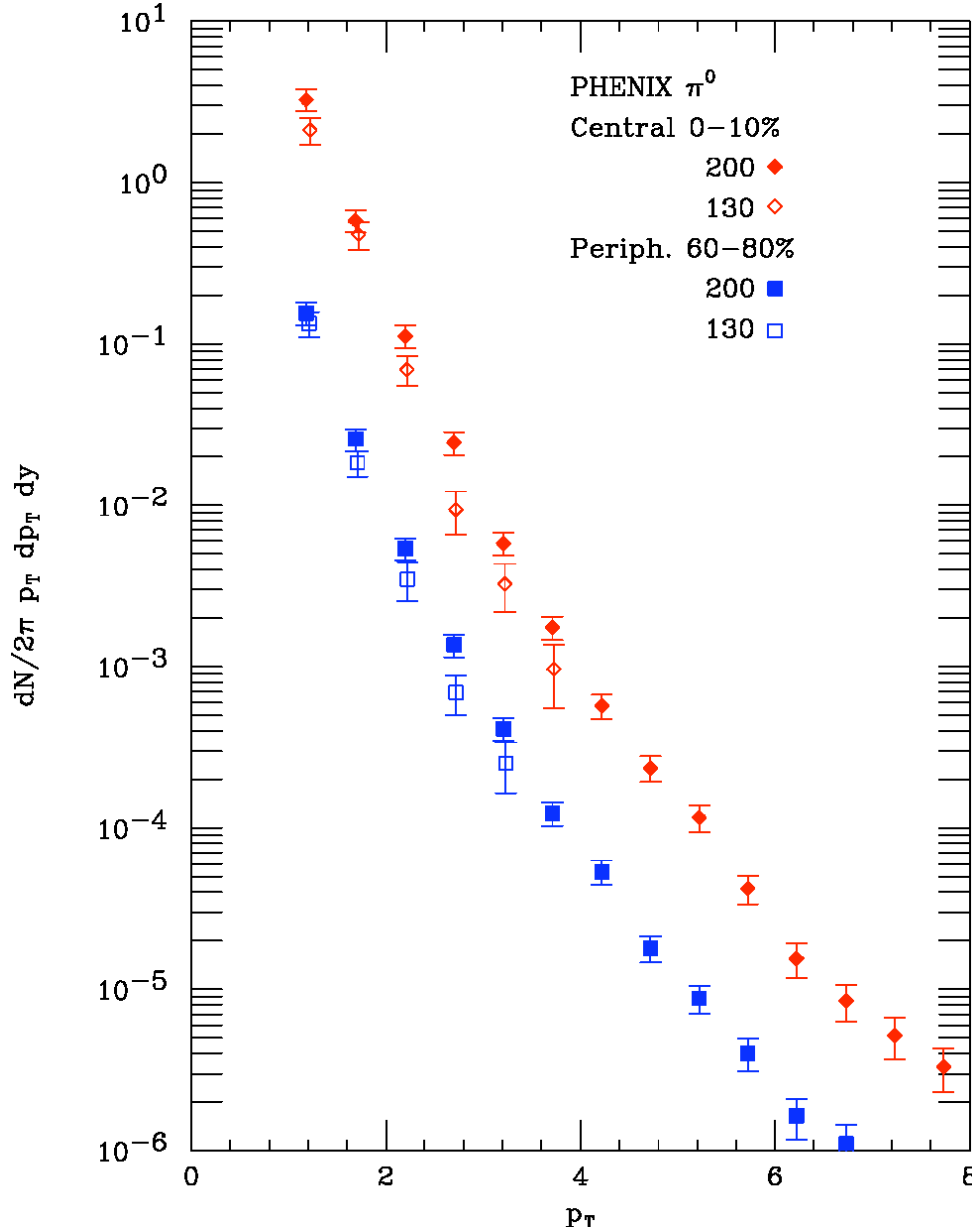


Figure 5: Left: CCOR invariant cross section vs $x_T = 2p_T/\sqrt{s}$. Right: $n(x_T, \sqrt{s})$ derived from the combinations indicated. *The systematic normalization at $\sqrt{s} = 30.6$ has been added in quadrature. Note that the absolute scale uncertainty cancels!*

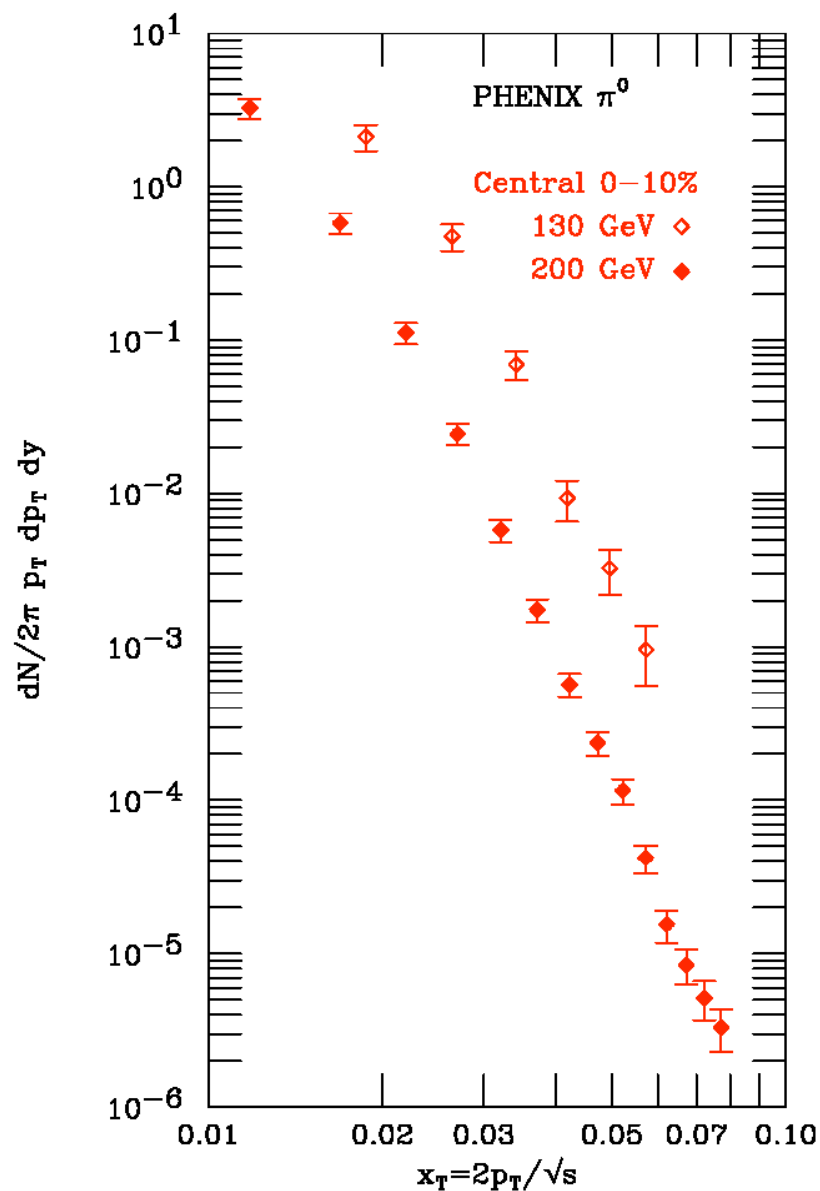
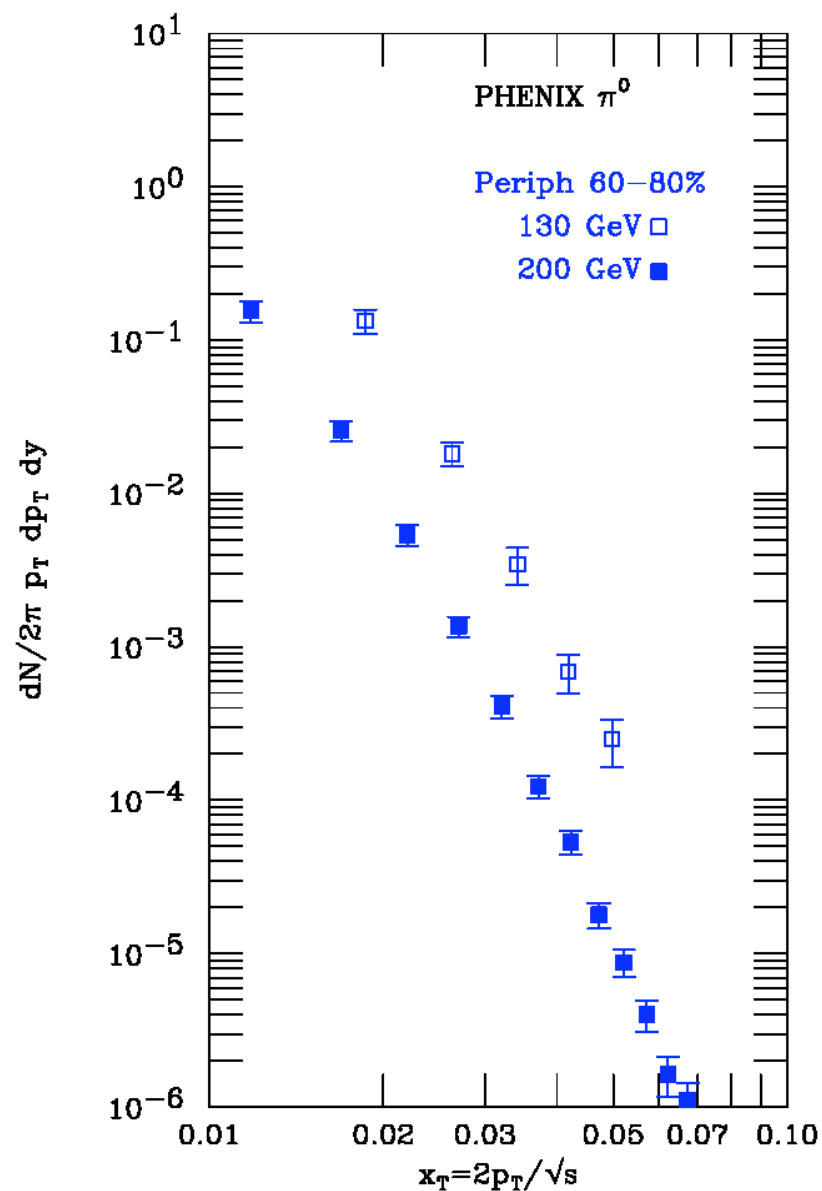
PHENIX
Semi-
Inclusive π^0
Au+Au
 $s_{NN}=130$ and
200 GeV
vs p_T

Peripheral 60 -- 80%

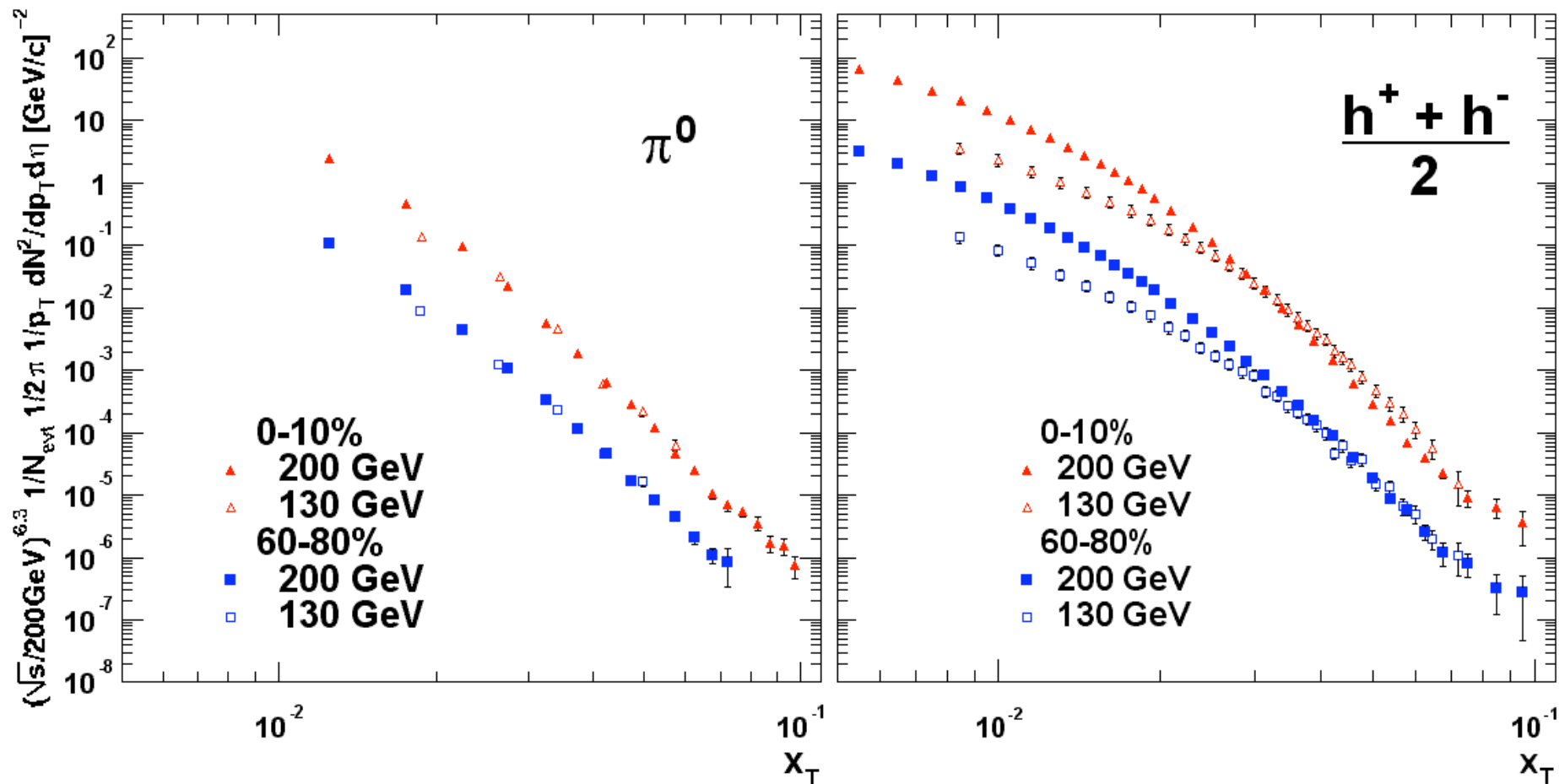
Central 0-10%



Same data vs x_T on log-log plot



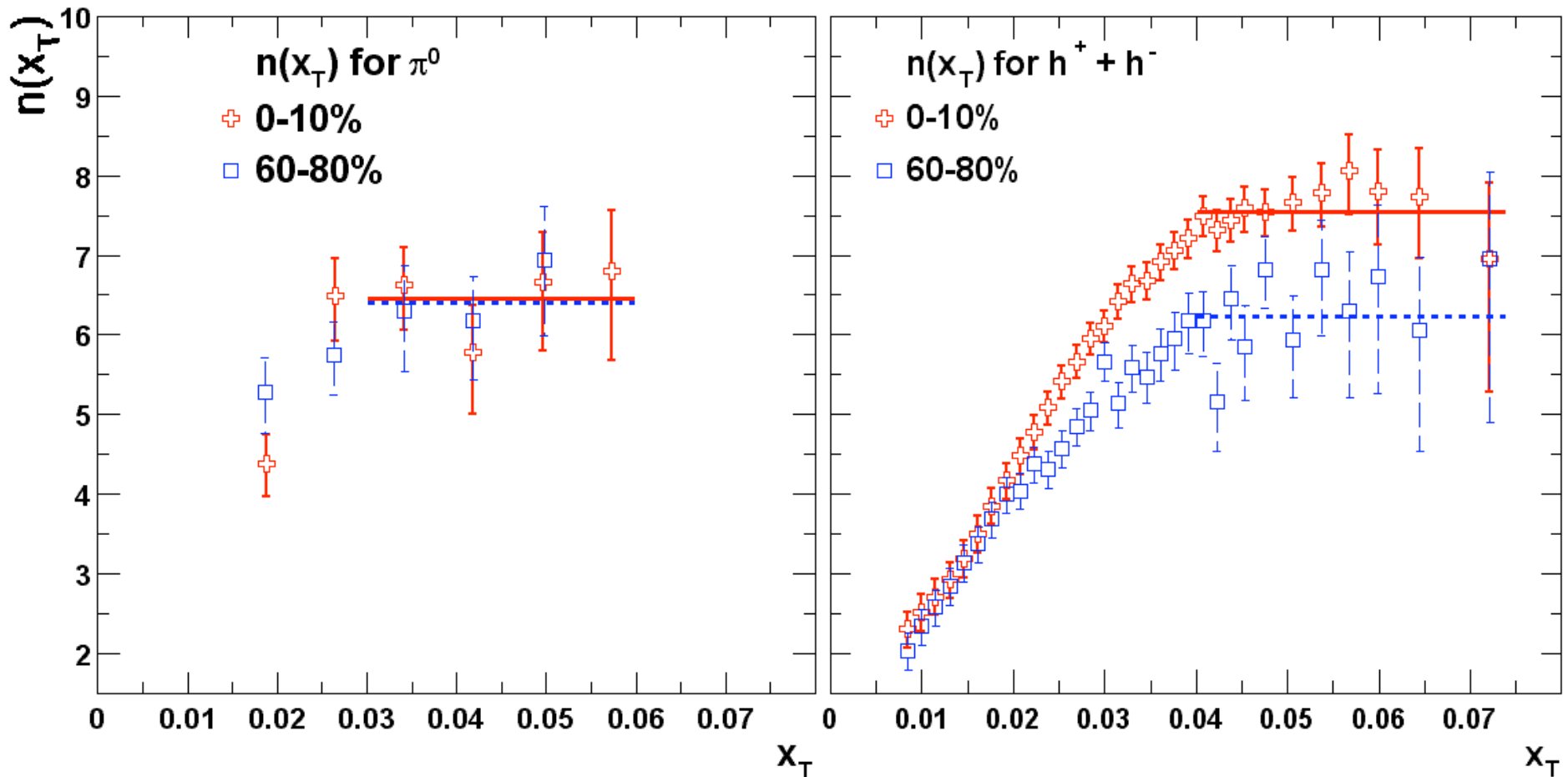
π^0 and $(h^+ + h^-)/2$ x_T scaled $n=6.3$



- Can calculate $n(x_T)$ point-by-point by the ratio of π_{inv} at fixed x_T for 2 different s

$$\left(\frac{\sqrt{s_1}}{\sqrt{s_2}} \right)^{n(x_T, \sqrt{s})} = \frac{\frac{Ed^3\sigma}{dp^3}(x_T, \sqrt{s_2})}{\frac{Ed^3\sigma}{dp^3}(x_T, \sqrt{s_1})}$$

$n(x_T)$ point-by-point 200/130



- π^0 x_T scales in both peripheral and central Au+Au with same value of $n=6.3$ as in p-p
- $(h^+ + h^-)/2$ x_T scales in peripheral same as p-p but difference between central and peripheral is significant

Precision values of $n(x_T)$

For a more quantitative analysis, the $Au + Au$ data for a given centrality and hadron selection are fitted simultaneously for $\sqrt{s_{NN}} = 130$ and 200 GeV to the form,

$$E \frac{d^3\sigma}{dp^3}(x_T, \sqrt{s}) = \left(\frac{A}{\sqrt{s}} \right)^n (x_T)^{-m}$$

Fitting results for π^0 over $0.03 < x_T < 0.06$		
parameters	0-10% centrality bin	60-80% centrality bin
A	0.973 ± 0.232	0.843 ± 0.3
m	8.48 ± 0.17	7.78 ± 0.22
n	$6.41 \pm 0.25(stat)$ $\pm 0.49(sys)$	$6.33 \pm 0.39(stat)$ $\pm 0.37(sys)$
Fitting results for $h^+ + h^-$ over $0.04 < x_T < 0.074$		
A	2.30 ± 0.44	0.62 ± 0.27
m	8.74 ± 0.28	8.40 ± 0.43
n	$7.53 \pm 0.18(stat)$ $\pm 0.40(sys)$	$6.12 \pm 0.33(stat)$ $\pm 0.36(sys)$

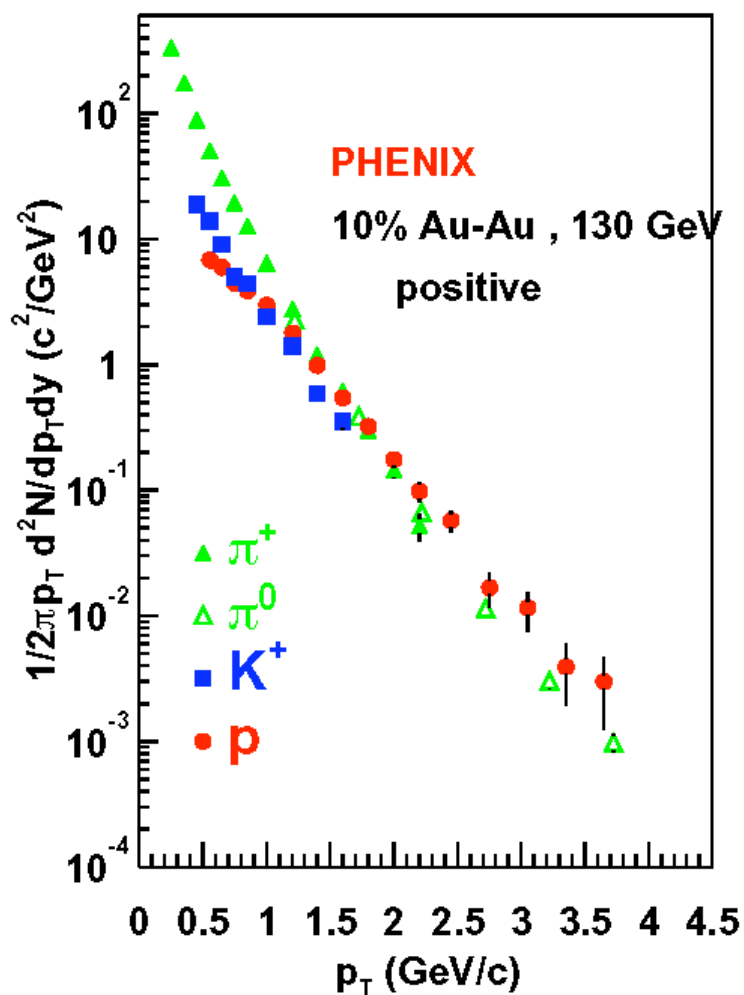
- $\Delta n = n_{cent} - n_{periph} = 1.41 \pm 0.43$ for $(h^+ + h^-)/2$ *significant*
- $\Delta n = n_{cent} - n_{periph} = 0.09 \pm 0.47$ for π^0

Conclusions

- x_T scaling in peripheral Au+Au collisions at $\sqrt{s_{NN}}=130$ and 200 GeV for both π^0 and $(h^+ + h^-)/2$ with the same value of $n \sim 6.3$ as in p-p collisions in this x_T range indicates that hard-scattering is the dominant production mechanism for high p_T particles in Au+Au collisions.
- π^0 production exhibits x_T scaling with the same value of $n \sim 6.3$ in both central and peripheral Au+Au collisions
- This implies that the dynamics of suppressed high p_T π^0 production in Au+Au collisions is consistent with hard-scattering according to pQCD with scaling structure and fragmentation functions as in p-p collisions.
- Perhaps a little puzzling since this indicates that the energy loss of the parton scales with its energy, if energy loss. However, scaling is consistent with gluon-saturation.

Conclusions, continued

- For $(h^+ + h^-)/2$ the difference in n between central and peripheral collisions is significant: $\Delta n = 1.41 \pm 0.43$ and is consistent with the large proton and anti-proton enhancement compared to charged pions, which appears to violate x_T scaling from 130 to 200 GeV:
- » The range $0.04 < x_T < 0.074$ corresponds to $2.6 < p_T < 4.8 \text{ GeV}/c$ at 130 GeV and $4 < p_T < 7 \text{ GeV}/c$ at 200 GeV
But protons are enhanced for the same p_T range $2 < p_T < 4.5 \text{ GeV}/c$ for both 130 and 200 GeV



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