

PHENIX **for Beginners**

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Kurchatov Institute, Moscow, Russia

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- What is the QGP phase transition?
- What is the best strategy for observing it?
- How has PHENIX implemented that strategy?
 - ❑ An introduction to PHENIX physics results
 - ❑ A (very) brief introduction to PHENIX detector technologies
- Goals:
 - ❑ Provide overview of the many observables already measured by PHENIX in RHIC Run-1
 - ❑ Understand the trade-offs made by experimenters in optimizing a detector
 - ❑ Understand vast potential of PHENIX in future RHIC runs



Q. How to liberate quarks and gluons from
~1 fm confinement scale?

A. Create an energy density

$$e > \sim (\hbar / 1 \text{ fm})^4 \sim 0.2 \text{ GeV} / \text{fm}^3 \sim \text{Normal nuclear density} ??$$

➡ Need better control of dimensional analysis:

$$e = g \frac{p^2}{30} T^4$$

Energy density for “g” massless d.o.f

$$= \left(\frac{1}{1} \times 2 \times 8_g + \frac{7}{8} \times 2_s \times 2_a \times 2_f \times 3_c \right) \frac{p^2}{30} T^4$$

8 gluons, 2 spins;
2 quark flavors, anti-quarks,
2 spins, 3 colors

$$= 37 \times \frac{p^2}{30} T^4$$

37 (!)

$$\gg 12 \times T^4 \gg 12 \times \left(\frac{\hbar}{1 \text{ fm}} \right)^4 \gg 2.4 \text{ GeV} / \text{fm}^3$$

“Reasonable”
estimate



- **Compare**

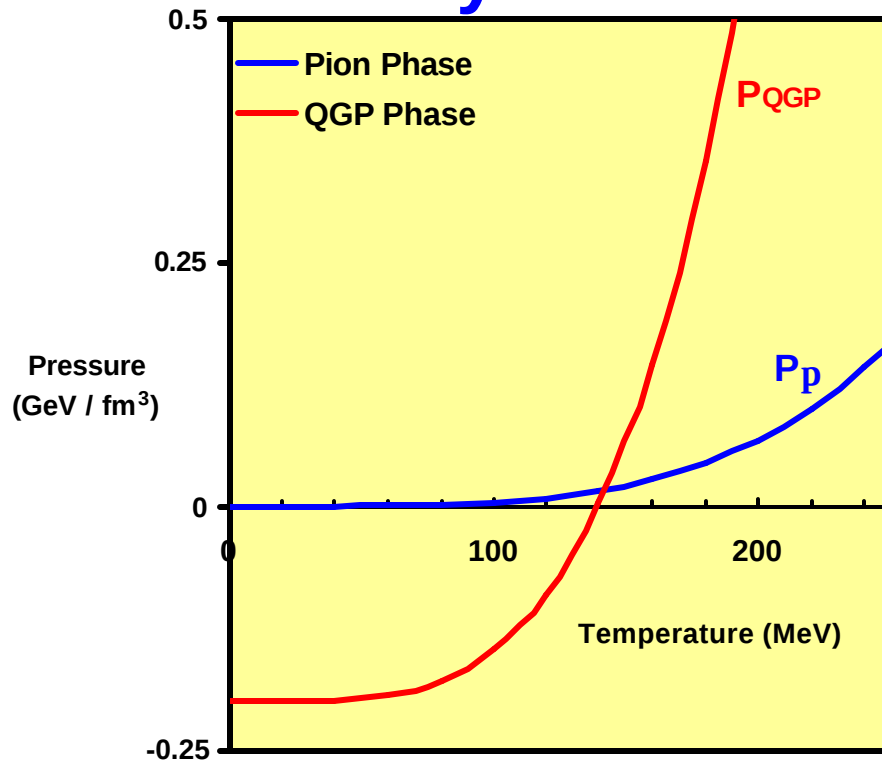
$$P_p = 3 \frac{p^2}{90} T^4$$

Pressure of “pure” pion gas at temperature T

$$P_{QGP} = g \frac{p^2}{90} T^4 - B, \quad g = 37$$

Pressure in plasma phase with
“Bag constant” $B \sim 0.2 \text{ GeV} / \text{fm}^3$

- **Select system with higher pressure:**



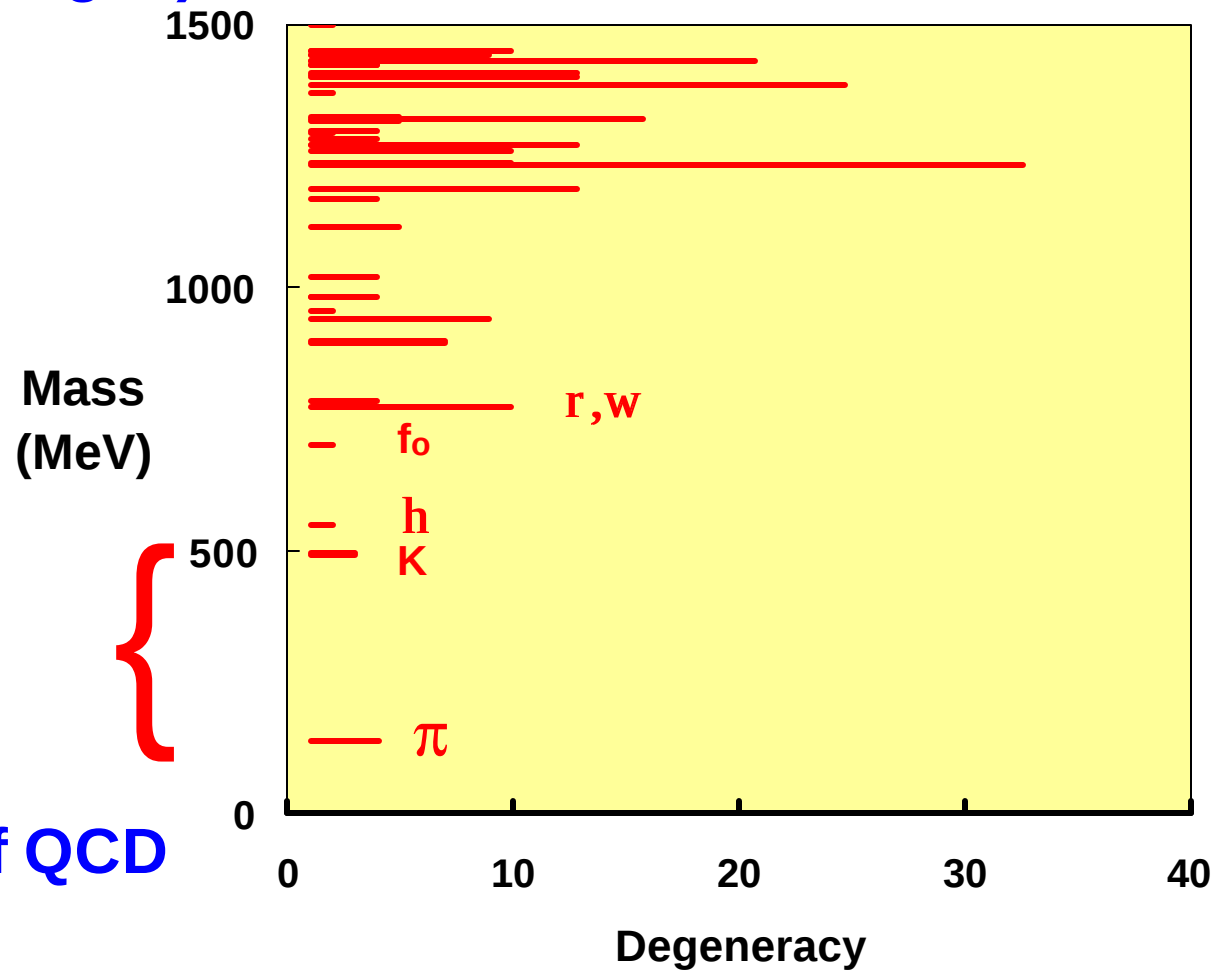
► Phase transition at $T \sim 140 \text{ MeV}$
with latent heat $\sim 0.8 \text{ GeV} / \text{fm}^3$

Compare to best estimates (Karsch, QM01)
from lattice calculations:
 $T \sim 150\text{-}170 \text{ MeV}$
latent heat $\sim 0.7 \pm 0.3 \text{ GeV} / \text{fm}^3$



- Previous approach (using a “pure” pion gas) works because QCD is a theory with a mass gap
- This gap is a manifestation of the approximate $SU(2)_R \times SU(2)_L$ chiral symmetry of QCD with pions as the Nambu-Goldstone bosons

Hadron 'level' diagram



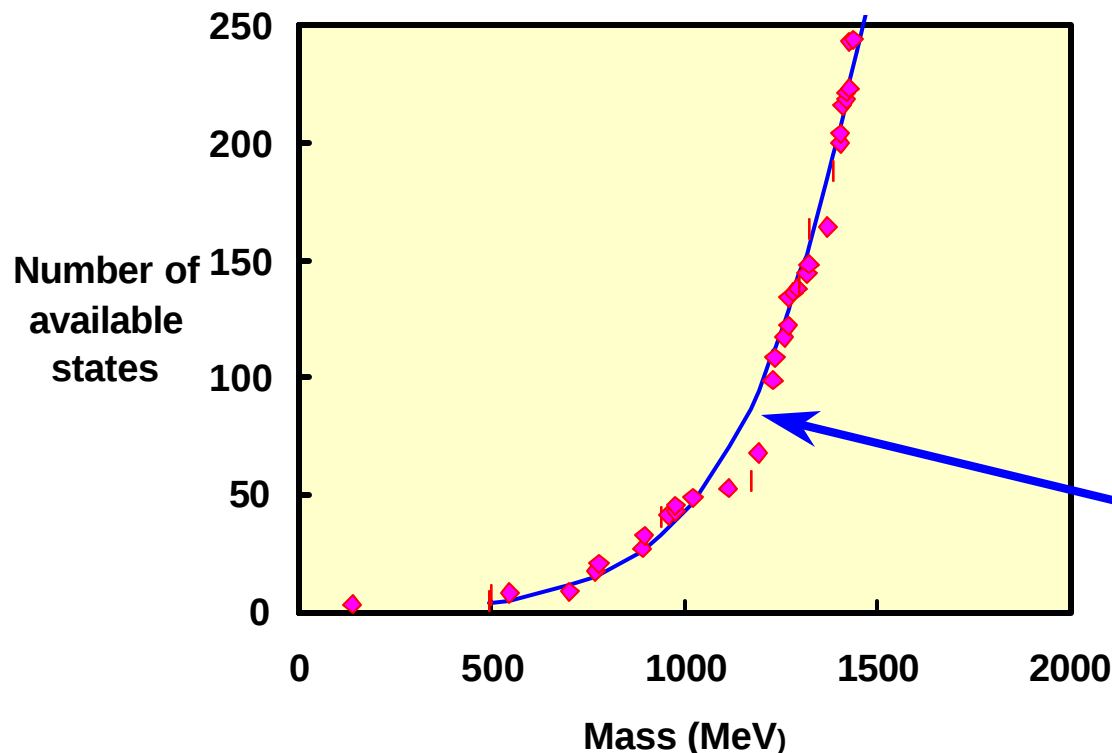


The frightening density of hadronic levels led to concepts of

- A “limiting temperature” T_H (Hagedorn, 1965)
- A phase transition(?) in hadronic matter

before quarks were understood as underlying constituents

Density of States vs Energy



$$r(m) \propto \frac{dn}{dm} \sim m^a e^{m/T_H}$$

$$P \propto \int_0^\infty r(m) e^{-m/T} dm$$

$$\sim \int_0^\infty m^a e^{m(\frac{1}{T_H} - \frac{1}{T})} dm$$

requires $T < T_H$

Fit to this form
with $T_H = 163 \text{ MeV}$

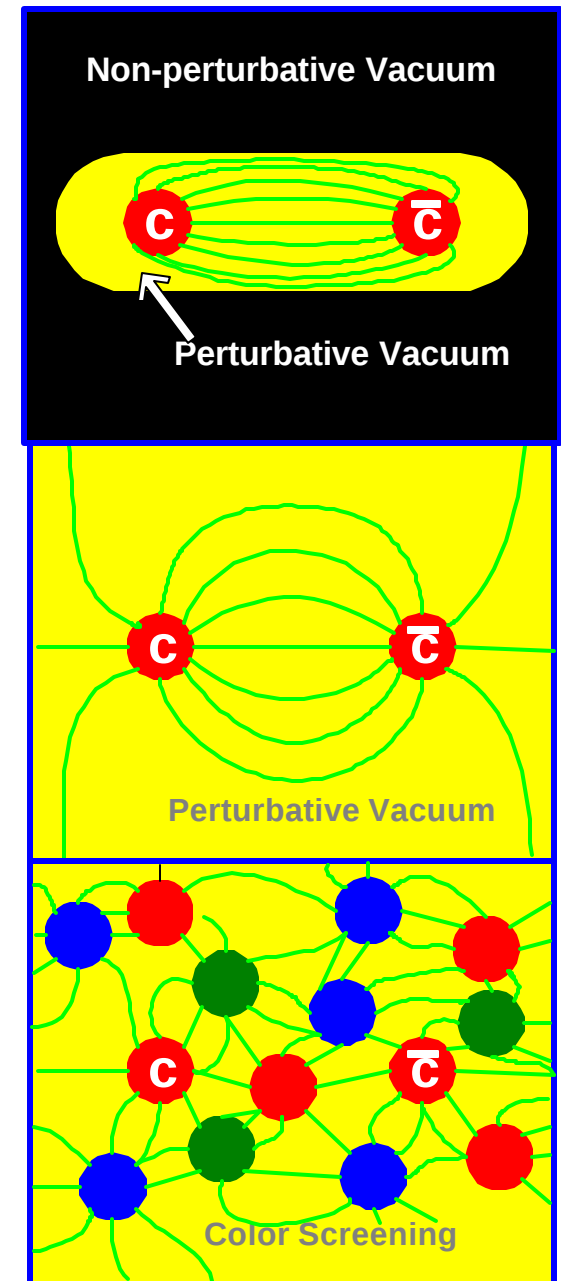


- Explore non-perturbative “vacuum” that confines color flux by melting it requires temperature $T \sim \hbar / (1 \text{ fm}) \sim 200 \text{ MeV}$

- ➔ Particle production
- ➔ Our ‘perturbative’ region is filled with
 - ◆ gluons
 - ◆ quark-antiquark pairs
- ➔ A Quark-Gluon Plasma (QGP)

- Experimental method:
Energetic collisions of heavy nuclei
- Experimental measurements:
Use probes that are
 - ❑ Auto-generated
 - ❑ Sensitive to all time/length scales

12-Jan-02





Question: How to proceed with experimental design

when $\sum_{\text{Theorists}} (\overline{\text{Theoretical Opinion}}) \approx 0$

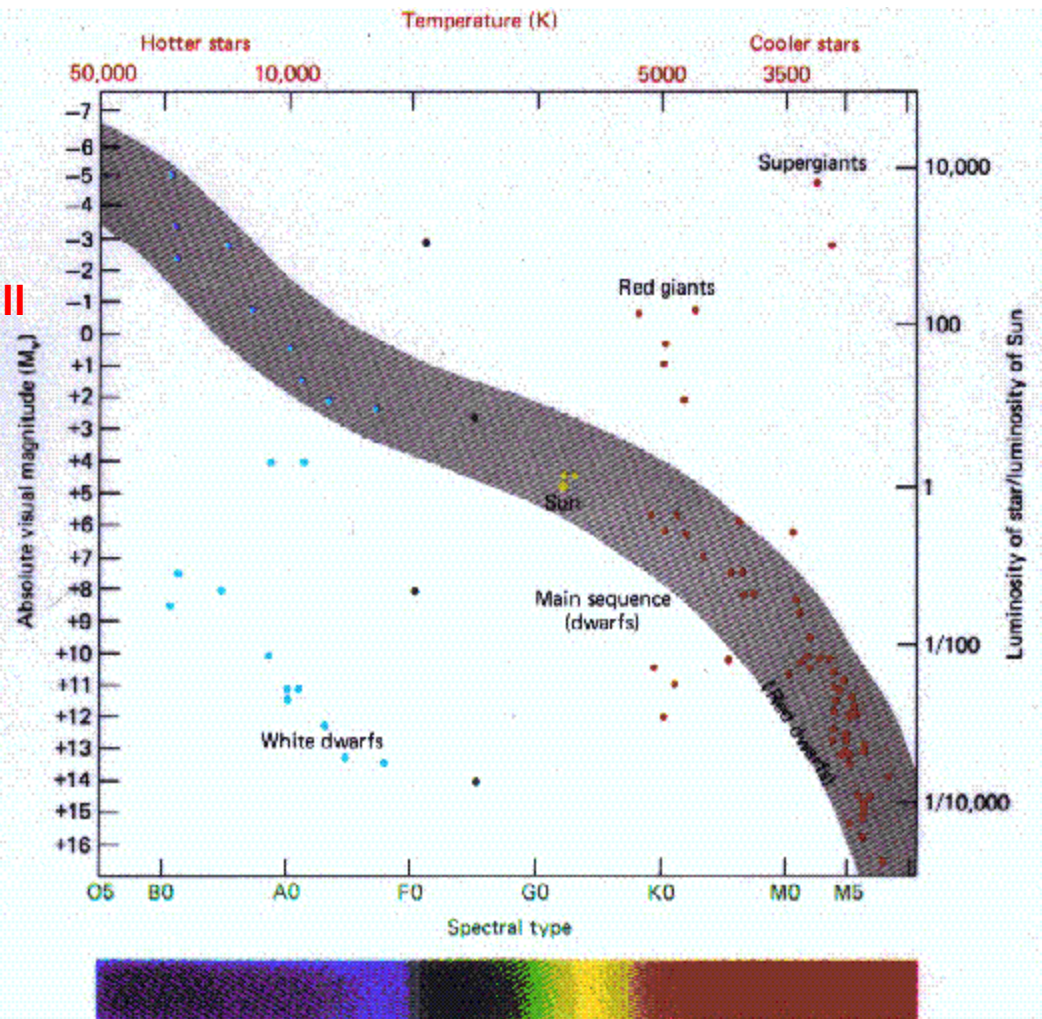
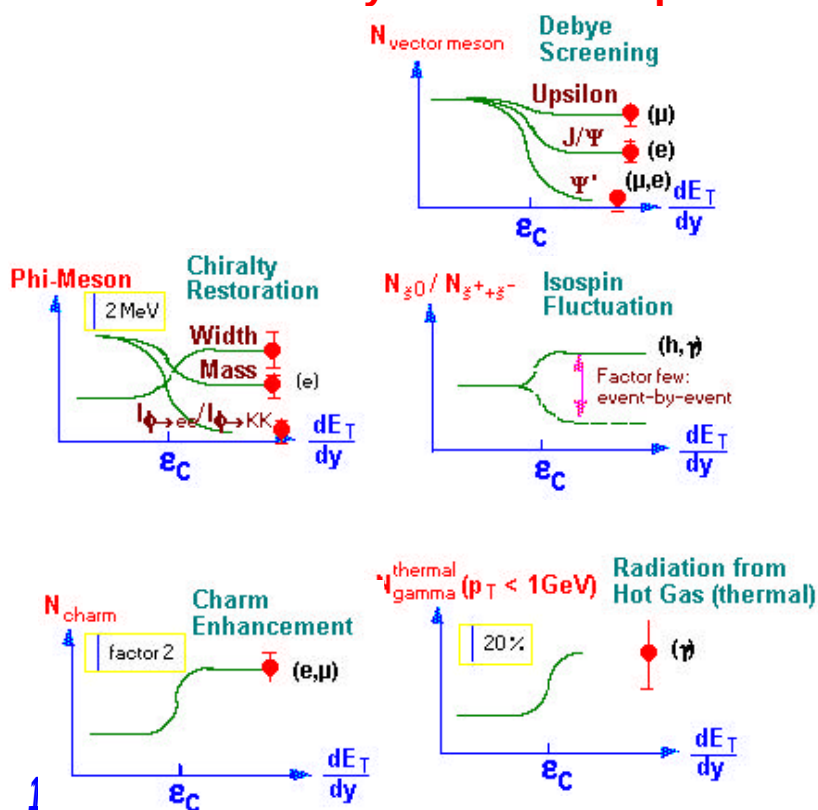
(Partial) answers:

- The QGP phase transition will not be “seen” at RHIC
 - ❑ Instead it will emerge as a consistent framework for describing the observed phenomena
 - ⇒ Avoid single-signal detectors
- There are no* cross sections at RHIC
 - * Except
 - * $S_{\text{GEOM}} \sim \text{few barns}$
 - * $S_{\text{CENTRAL}} \sim (1-10)\% S_{\text{GEOM}}$
 - * but $S_{\text{QGP}} \sim S_{\text{CENTRAL}} ??$
 - ⇒ Preserve high-rate and triggering capabilities
- Expect the unexpected
 - ❑ High gluon density ➡ production of exotics?
 - ❑ Color topology ➡ high anti-baryon production?
 - ❑ New vacuum ➡ large isospin fluctuations?
 - ⇒ Maintain flexibility as long as \$'s allow

(Guidance written in 1991)

Ample historical evidence for categorizing complex physics through correlation of observables

- For example, Hertzsprung-Russell in astronomy
- ➔ PHENIX will approach QGP via as many channels as possible



1. Deconfinement

$R(U) \sim 0.13 \text{ fm} < R(J/\psi) \sim 0.29 \text{ fm} < R(\Upsilon) \sim 0.56 \text{ fm}$

⇒ Electrons, Muons

2. Chiral Symmetry Restoration

Mass, width, branching ratio of F to e^+e^- , K^+K^- with $dM < 5 \text{ MeV}$:

⇒ Electrons, Muons, Charged Hadrons

Baryon susceptibility, color fluctuations, anti-baryon production:

⇒ Charged hadrons

DCC's, Isospin fluctuations:

⇒ Photons, Charged Hadrons

3. Thermal Radiation of Hot Gas

Prompt g , Prompt g^* to e^+e^- , m^+m^- :

⇒ Photons, Electrons, Muons

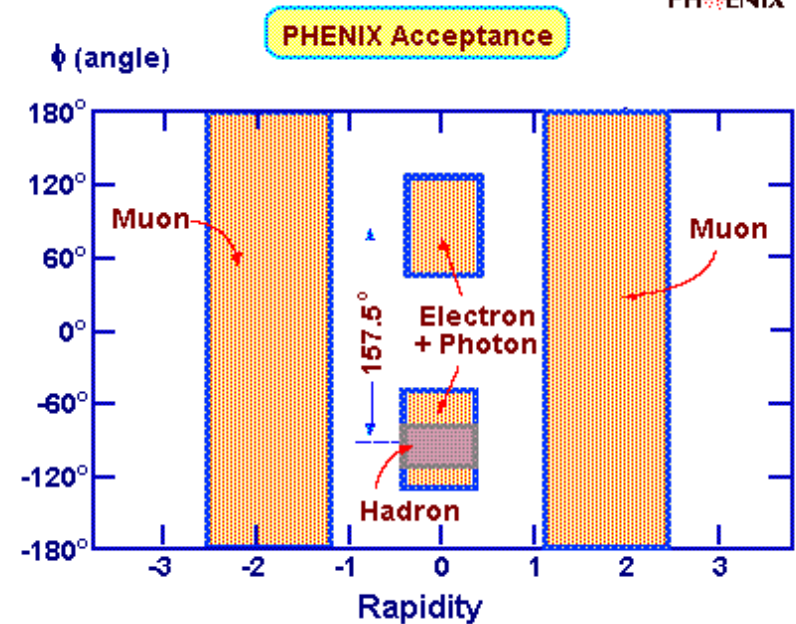
4. Strangeness and Charm Production

Production of K^+ , K^- mesons:

⇒ Hadrons

Production of F , J/ψ , D mesons:

⇒ Electrons, Muons



5. Jet Quenching

High p_T jet via leading particle spectra:

⇒ Hadrons, Photons

6. Space-Time Evolution

HBT Correlations of p^+p^+ , K^+K^+ :

⇒ Hadrons

Summary: Electrons, Muons, Photons, Charged Hadrons



- ❑ Early timescales in collision typically probed by “hard processes”.
- ❑ “Hard”
 - ➔ high momentum transfer
 - ➔ rare
 - ➔ luminosity limited

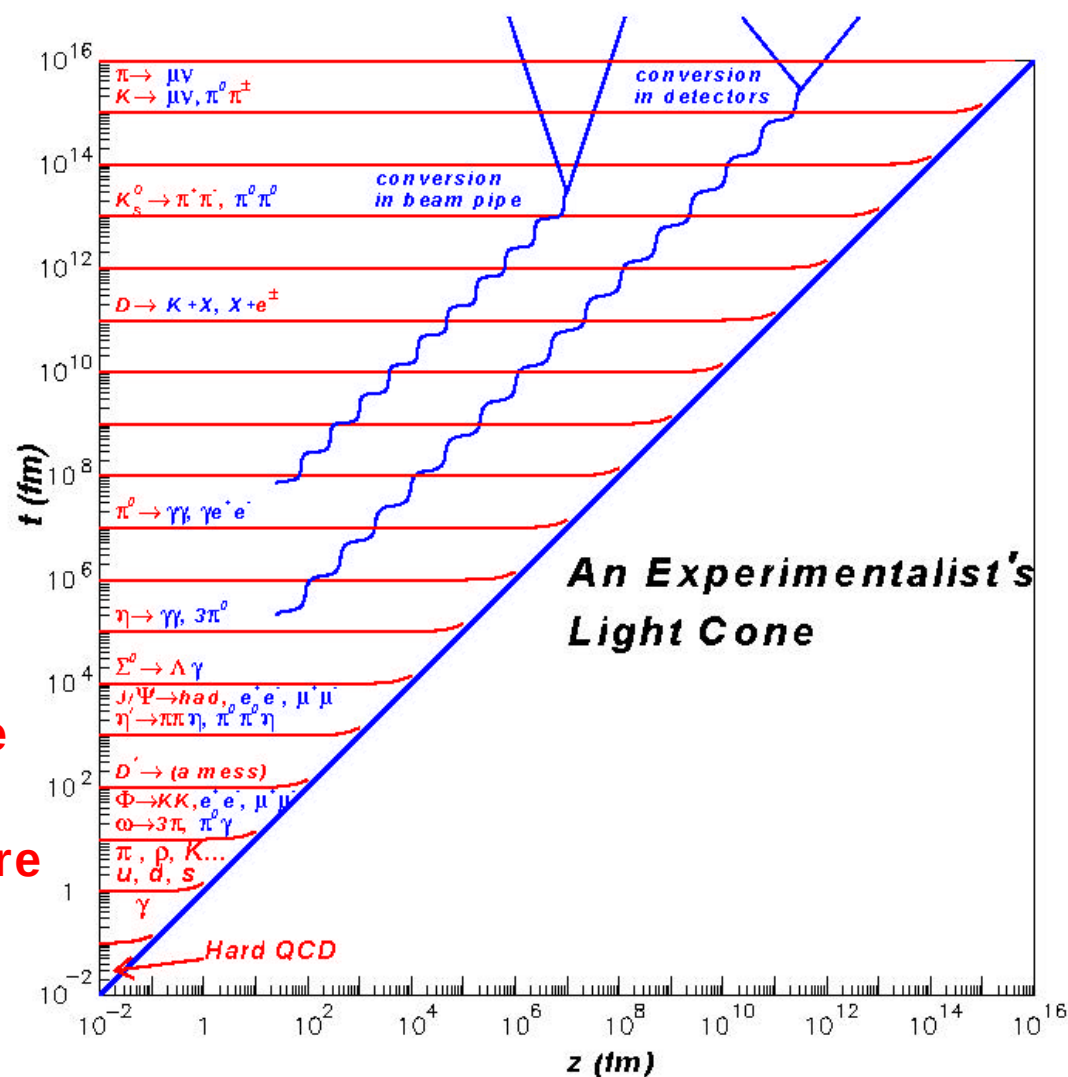
- ❑ $< \sim 1 \text{ mb}^{-1}$ recorded
- ❑ $\sim 5\text{M}$ “minimum bias” events

- ❑ $\sim 24 \text{ mb}^{-1}$ recorded
- ❑ $\sim 200\text{M}$ events “sampled”

12-Jan-02



- One must deal with hadronic and electromagnetic interactions at ~all length scales
 - (I.e., not just the first 10 fm)
 - Convenient to represent on a “logarithmic light-cone”
- More importantly, PHENIX faced severe design constraints:
 - Central arms-
minimize material in aperture
 - Muon arms-
maximize absorber in aperture

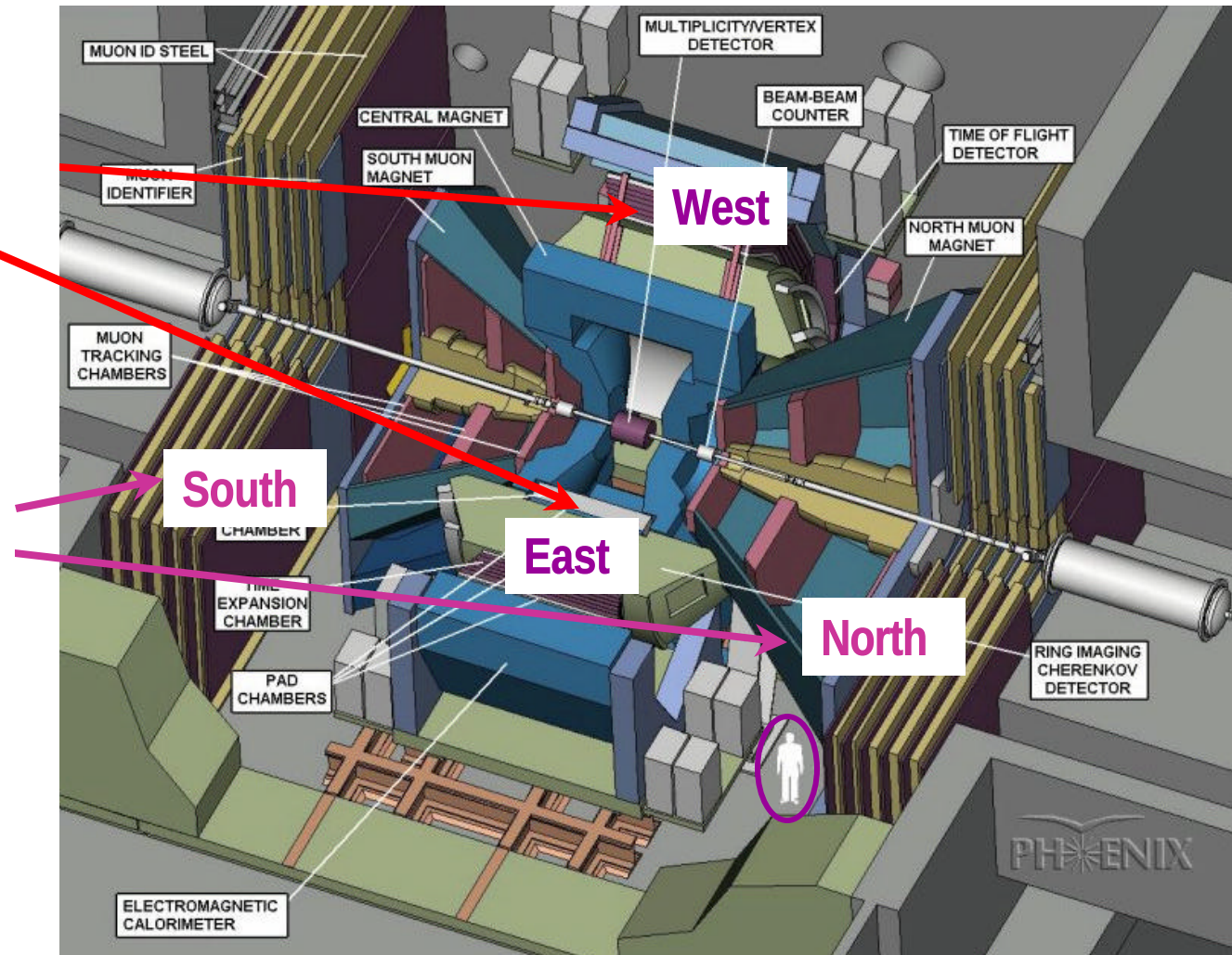


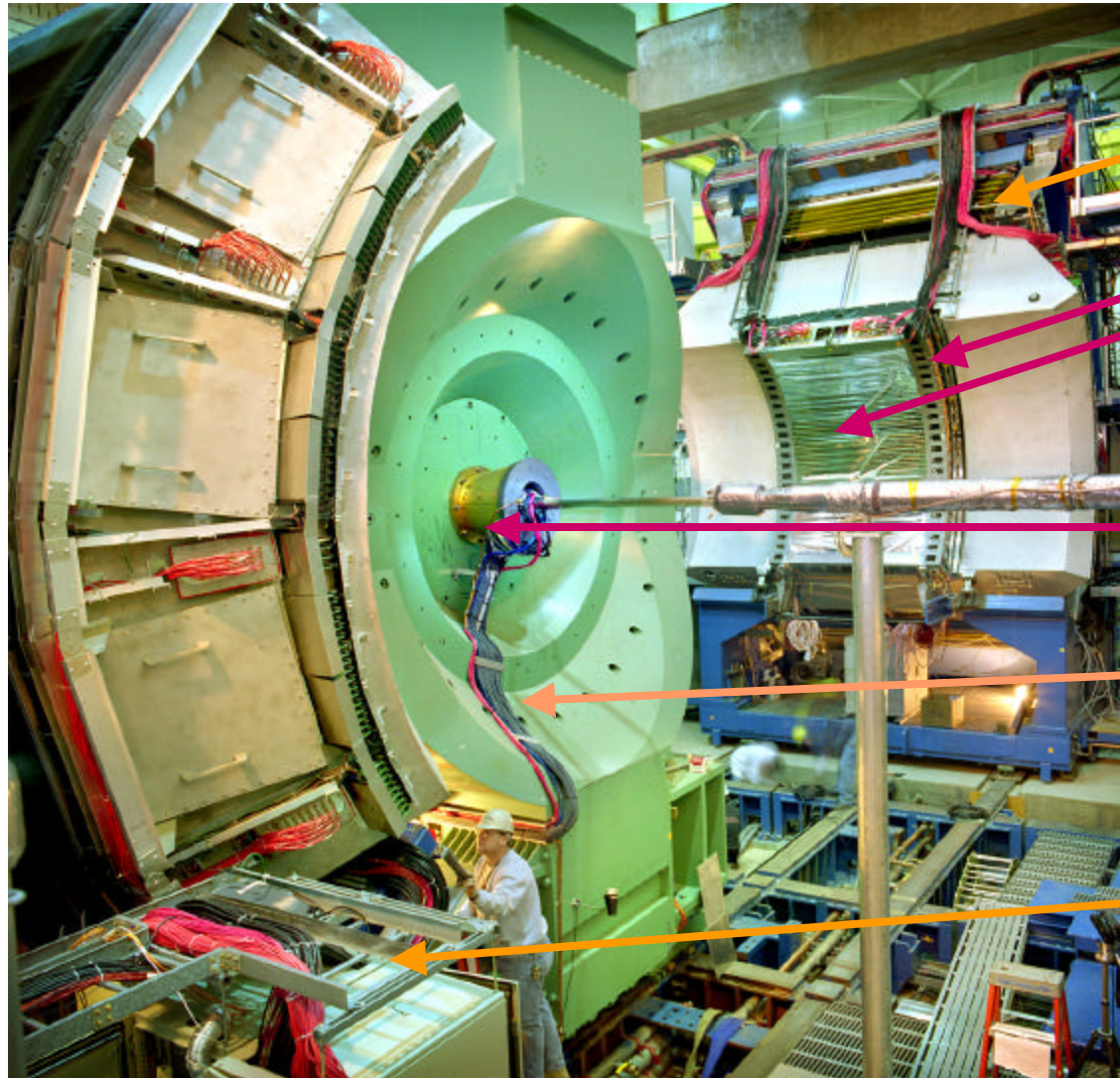


2 central spectrometers

2 forward spectrometers

3 global detectors





East Carriage

Ring Imaging Cerenkov
Drift Chamber

Beam-Beam Counter

Central Magnet

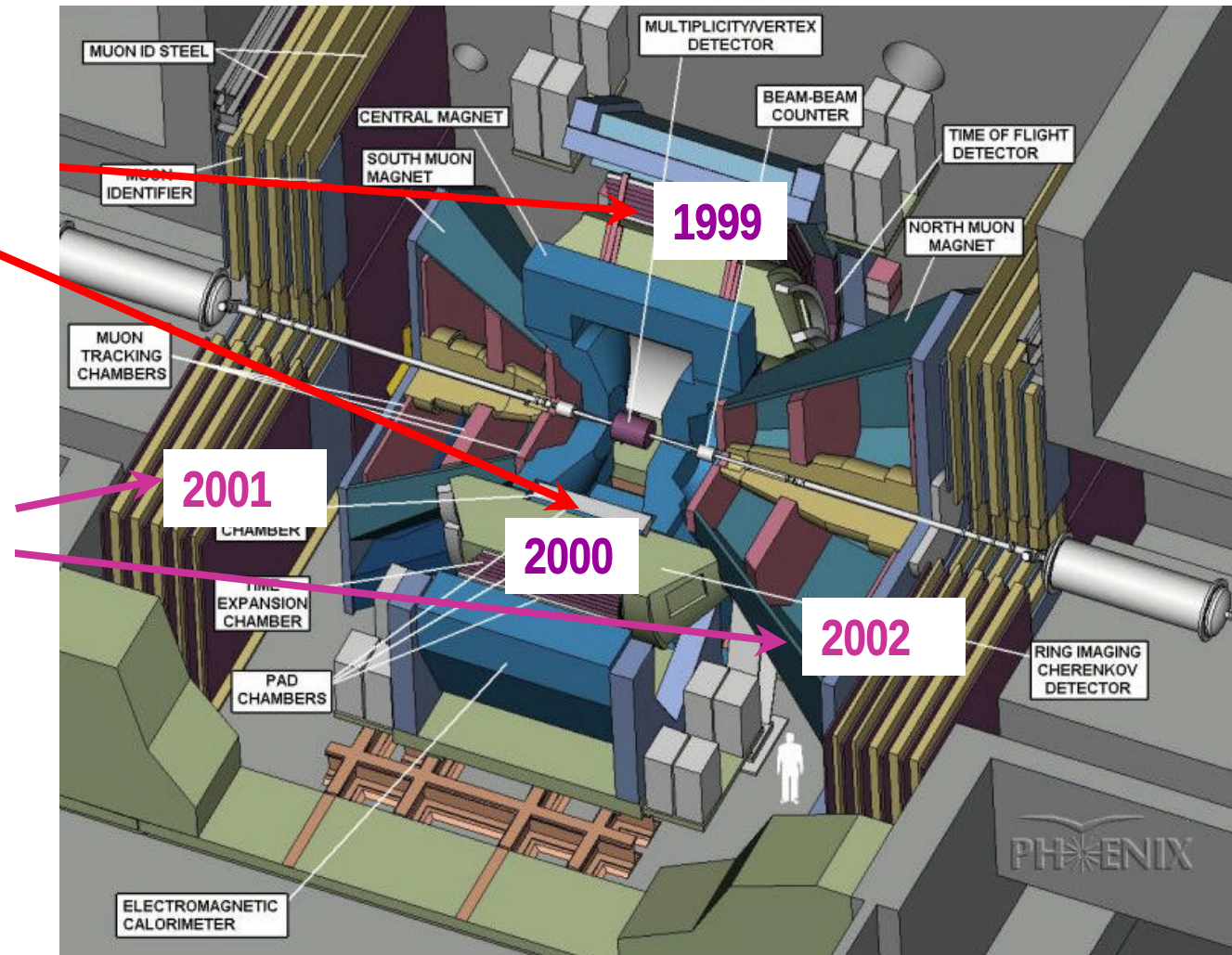
West Carriage



2 central spectrometers

2 forward spectrometers

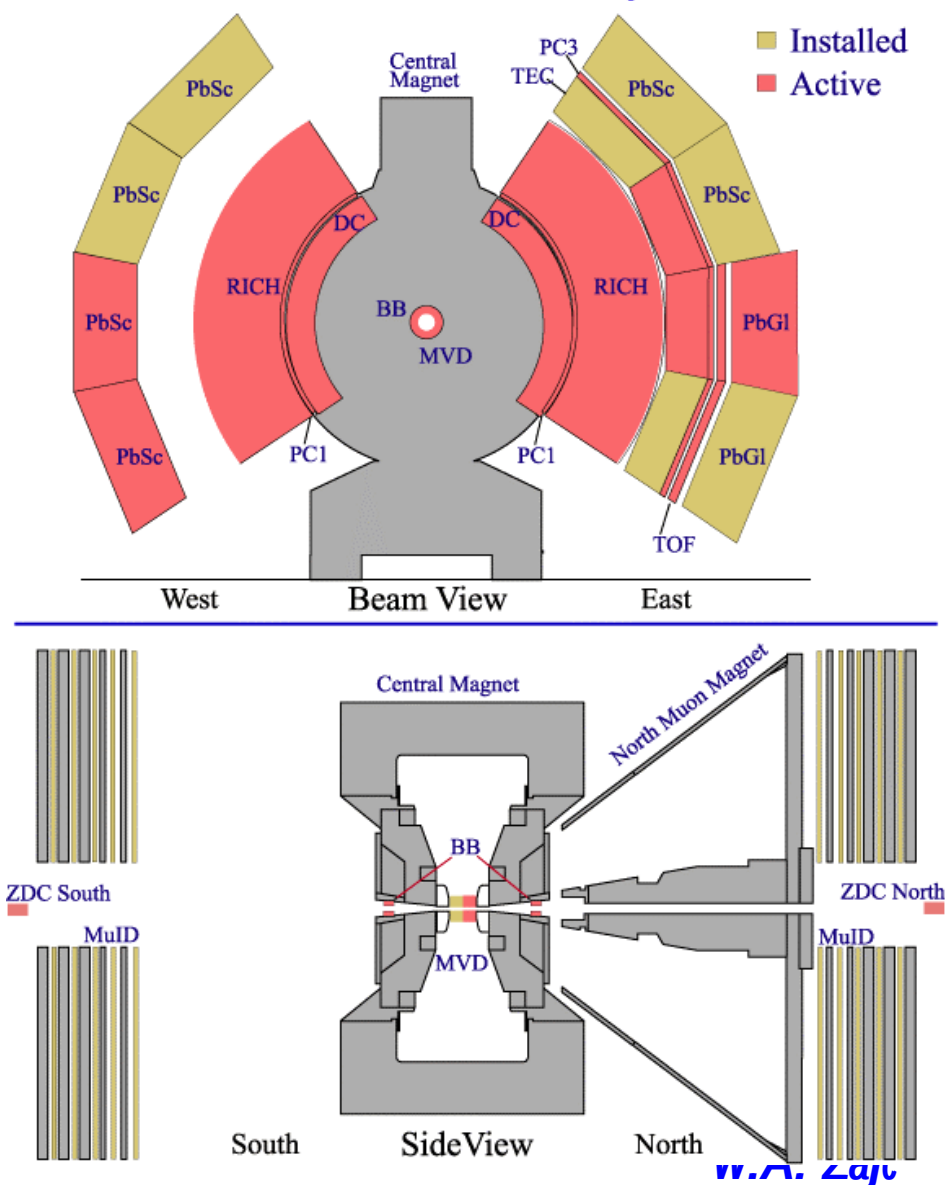
3 global detectors





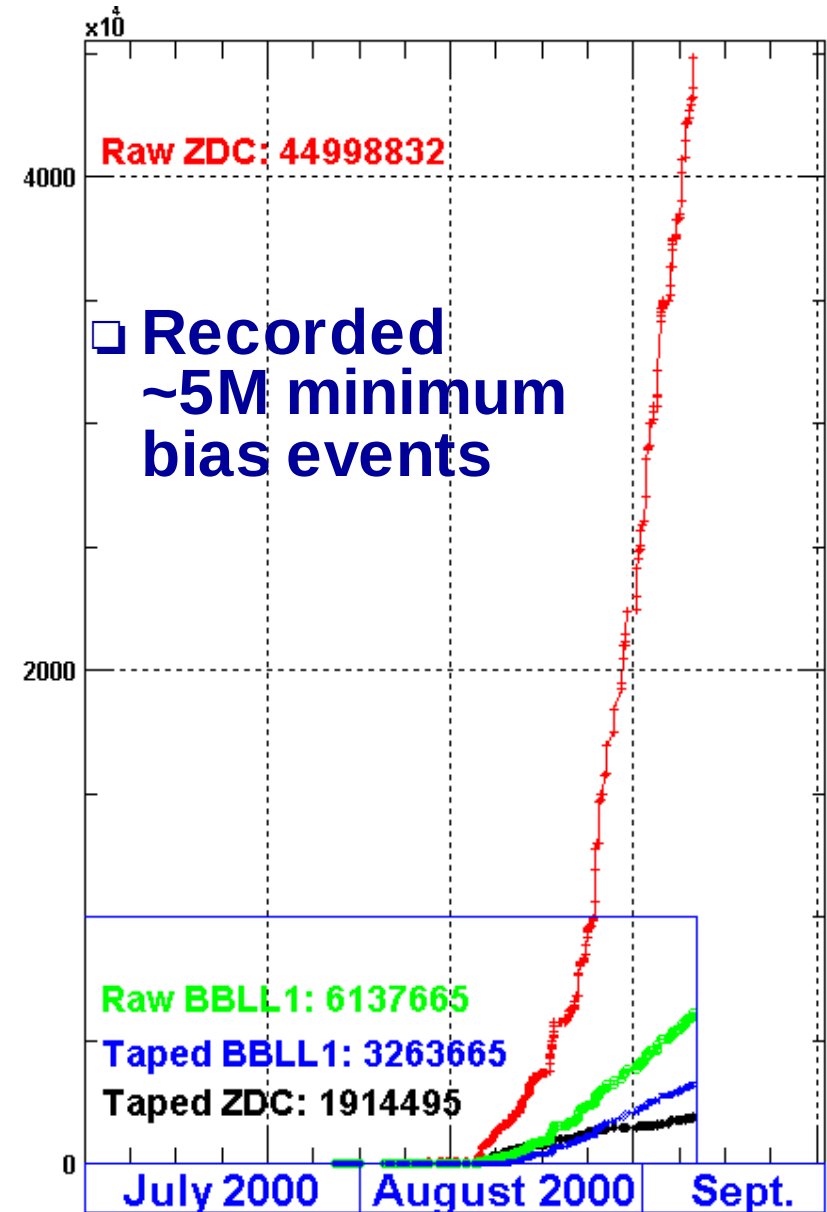
- **Two central arms**
 - ❑ Mechanically ~complete
 - ❑ Roughly half of aperture instrumented
- **Global detectors**
 - ❑ Zero-degree Calorimeters (ZDCs)
 - ❑ Beam-Beam Counters (BBCs)
 - ❑ Multiplicity and Vertex Detector (MVD, engineering run)

PHENIX Detector - First Year Physics Run





- First collisions: 15-Jun-00
- Last collisions: 04-Sep-00
- During this period:
 - Commissioned
 - ◆ Zero-Degree Calorimeters
 - ◆ Beam-Beam Counters
 - ◆ Multiplicity and Vertex Counter
 - ◆ Drift Chambers
 - ◆ Pad Chambers
 - ◆ Ring Imaging Cerenkov Counter
 - ◆ Time Expansion Chamber
 - ◆ Time-of-Flight Counters
 - ◆ Electromagnetic Calorimeter
 - ◆ Muon Identifier
 - ◆ Minimum Bias Triggers
 - ◆ Data Acquisition System



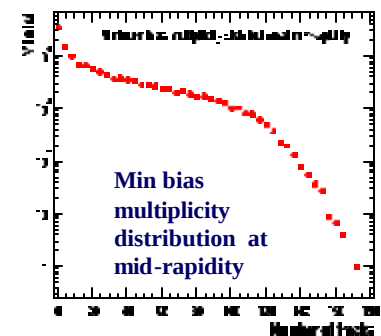
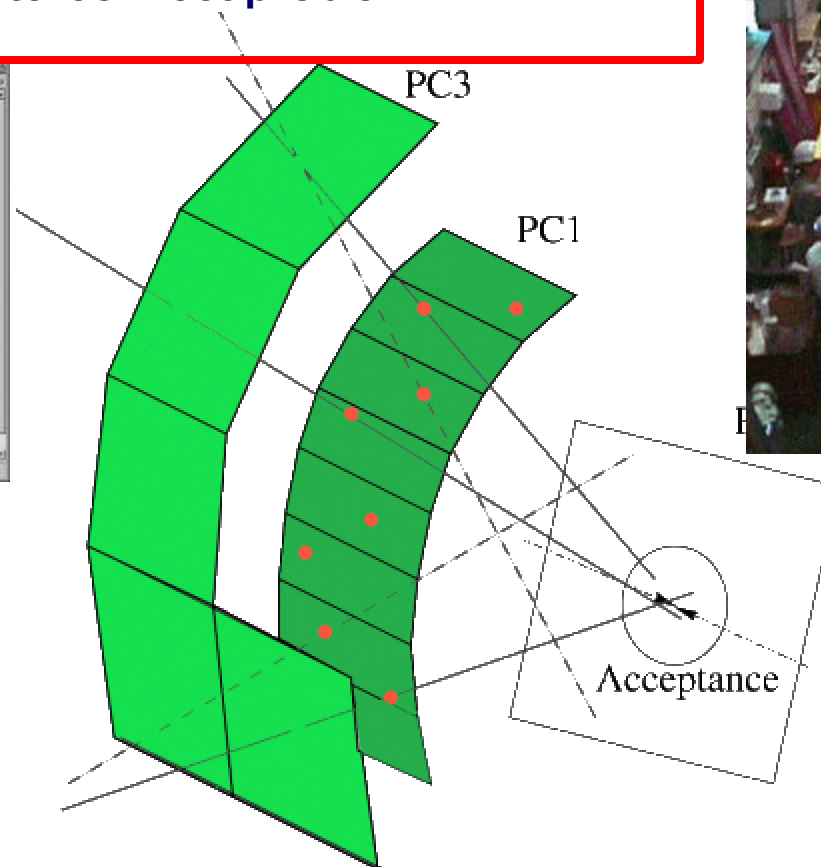


- Charged particle multiplicity
 - How many particles are produced?
 - What do they tell us about underlying reaction mechanism?
- Answered in

“Centrality Dependence of Charged Particle Multiplicity in Au-Au Collisions at $\sqrt{s} = 130$ GeV”, K. Adcox et al, PRL 86 (2001) 3500, preprint nucl-ex/0012008.



- 



Pixel Pad Cathode Pattern



Q. How many charged particles per nucleon-nucleon collision?

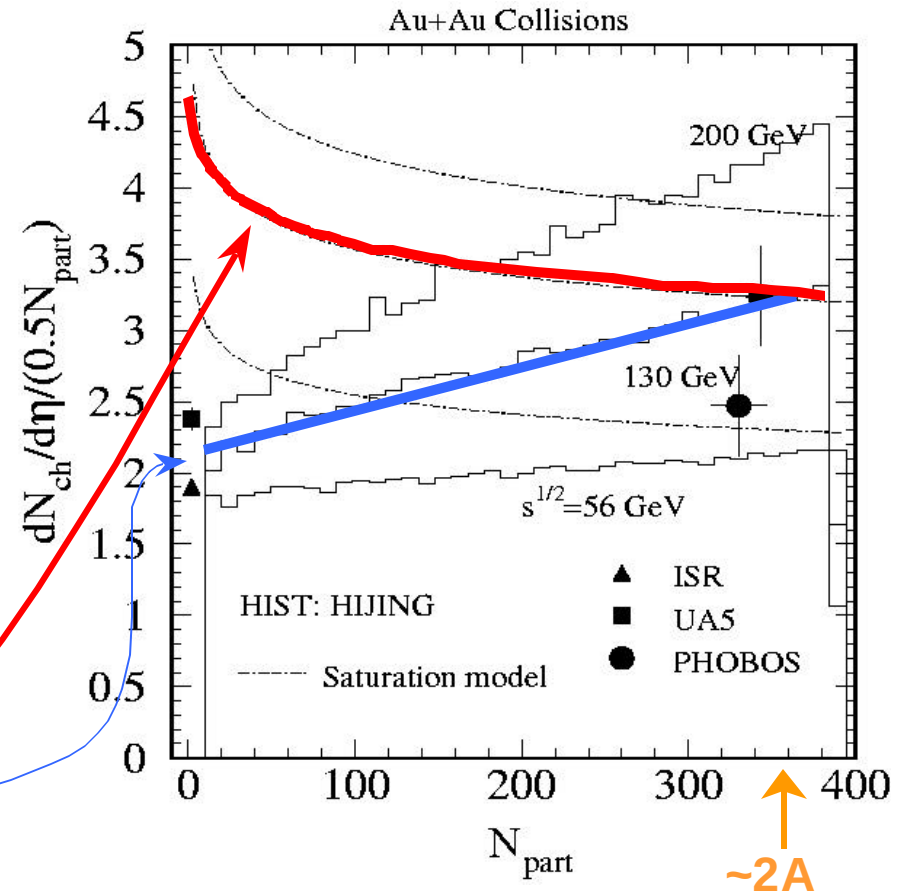
A: Express using

- $dN_{ch}/dh \sim$ “generic” dN_{ch}/dy
- N_{part}
 $\equiv$ Number of participating nucleons
 $= 2$ for p-p collision
 $\sim 2A$ for central A-A collision
- Plot multiplicity per N-N collision
 $(dN_{ch}/dh) / (N_{part}/2)$

• Does not (quite) distinguish between

- “Saturation” models, dominated by $gg \rightarrow g$
- “Cascade” models, dominated by $gg \rightarrow gg, gg \rightarrow ggg$

(X.N. Wang and M. Gyulassy, nucl-th/0008014)

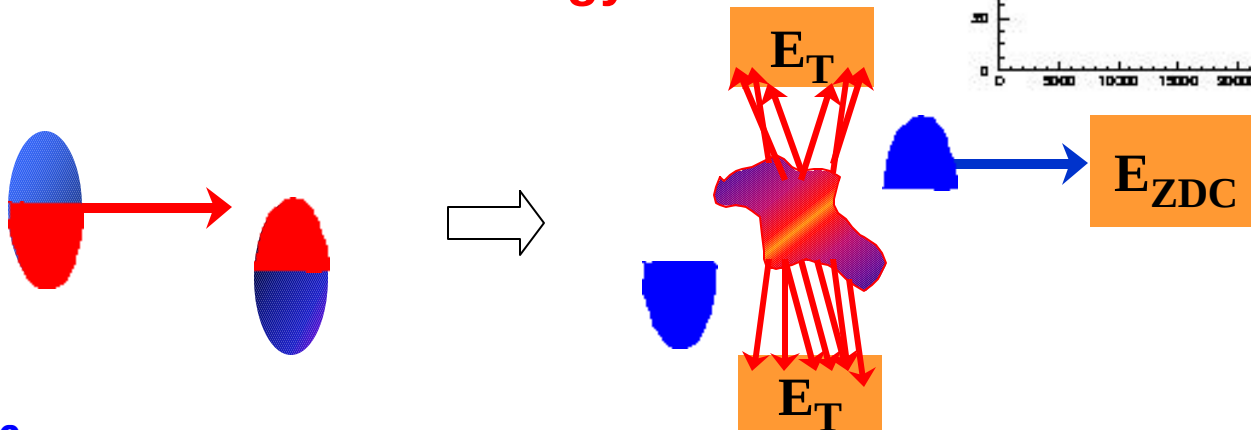
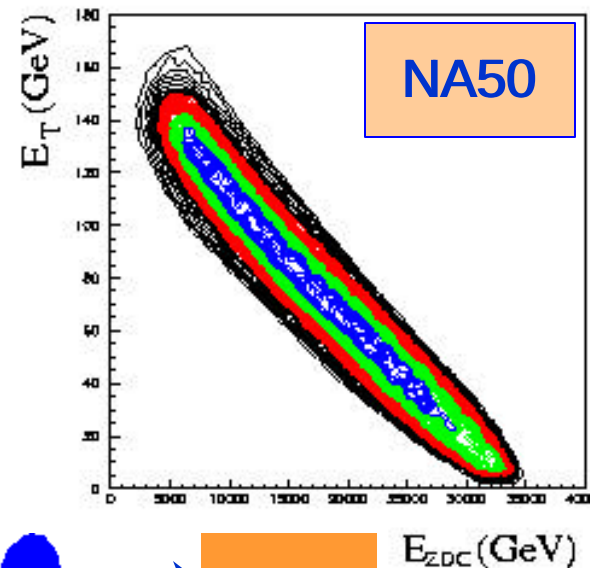


Best approach (for fixed target!):

- Directly measure in a “zero degree calorimeter”

$$N_{PART} \gg 2 \left(\frac{E_{ZDC}}{E_{PerNucleon}} \right) \quad (\text{for A+A collisions})$$

- Strongly (anti)-correlated with produced transverse energy:

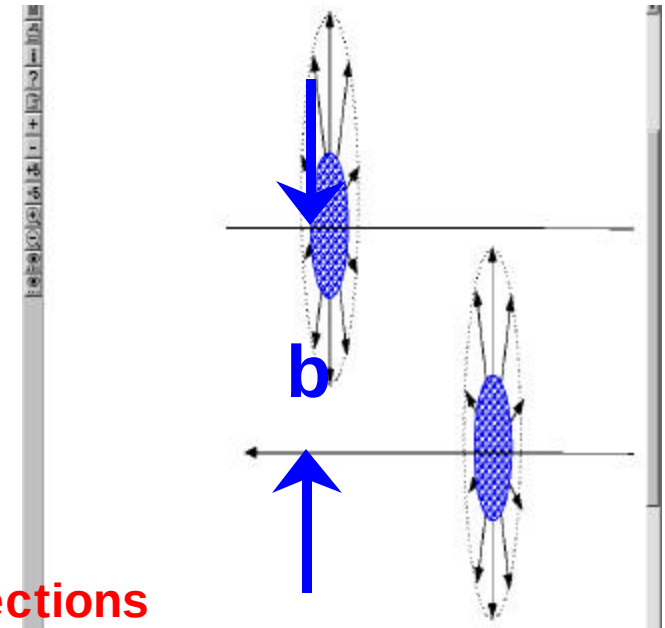
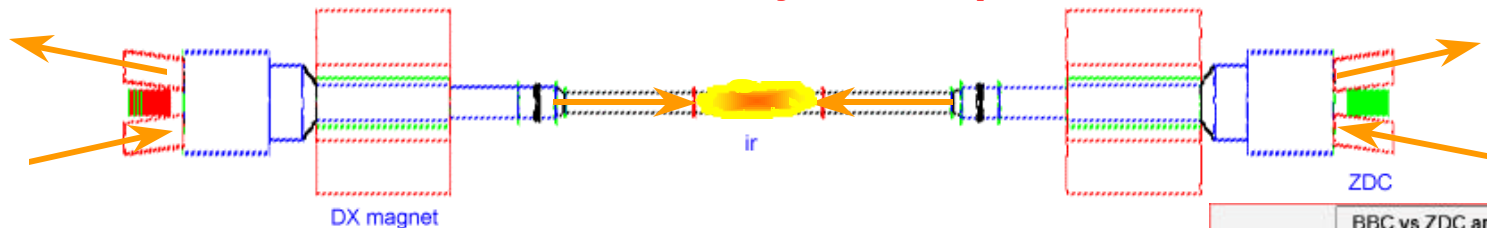


- Event characterization

- Impact parameter b is *well-defined* in heavy ion collisions
- Event multiplicity predominantly determined by *collision geometry*
- Characterize this by global measures of multiplicity and/or transverse energy correlated with
ZDC Zero Degree Calorimeter

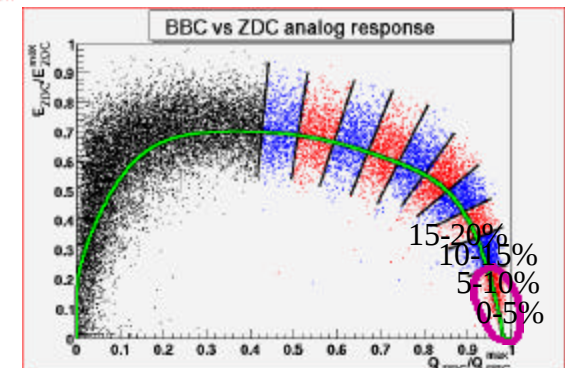
- Goals:

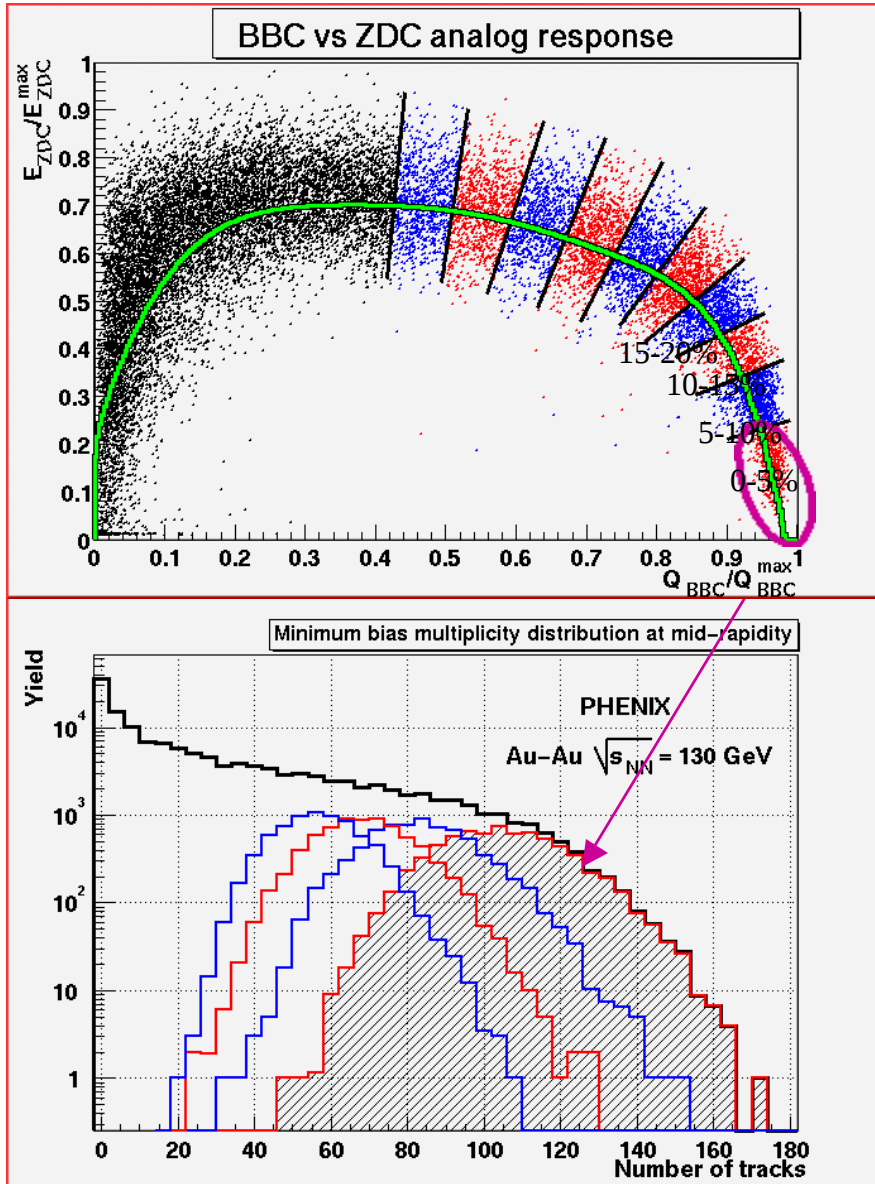
- Uniform luminosity monitoring at all 4 intersections
- Uniform event characterization by all 4 experiments



- Process:

- Correlated Forward-Backward Dissociation
- $S_{\text{tot}} = 11.0$ Barns (+/- few %)





Use combination of

- Zero Degree Calorimeters
 - Beam-Beam Counters
- to define centrality classes
- Glauber modeling
- to extract $N_{\text{participants}}$

determines Multiplicity vs. Centrality
i.e

dN_{ch}/dh vs. $N_{\text{participants}}$

which is presented as

“specific particle production”

multiplicity per N - N collision

$$\left(\frac{dN_{\text{ch}}}{dh} \right) / \left(N_{\text{part}}/2 \right)$$

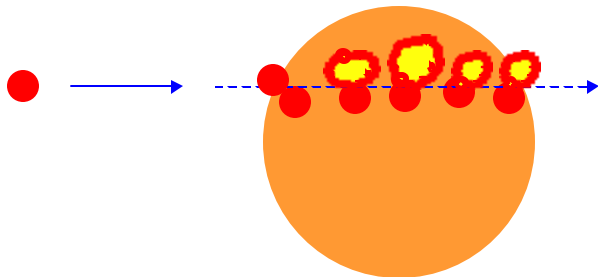
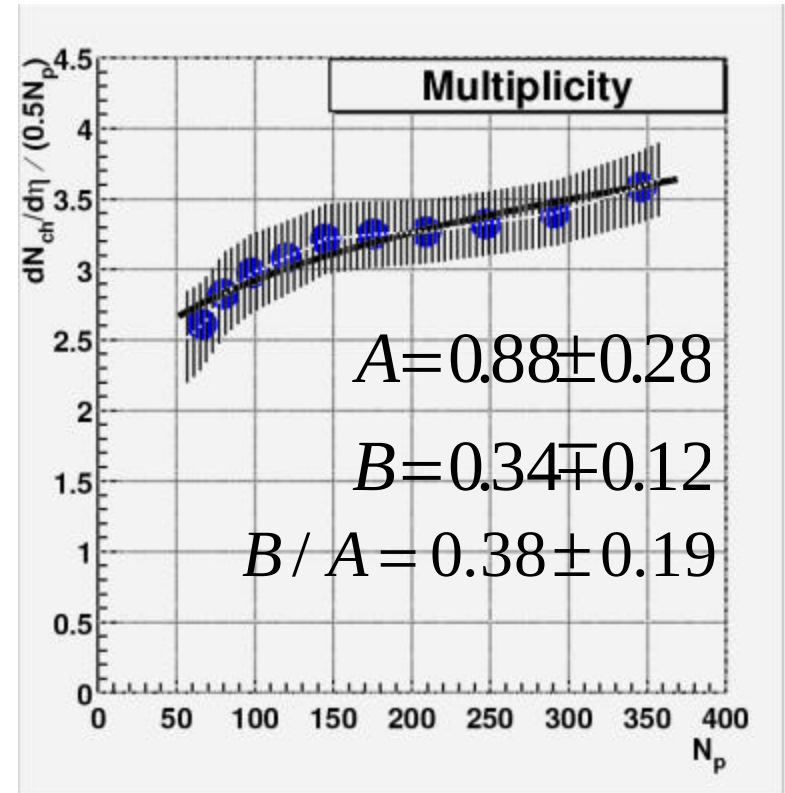


- “Specific multiplicity” is not flat

- Yields grow significantly faster than $N_{\text{participants}}$
- Evidence for term $\sim N_{\text{collisions}}$
- $$dN_{ch}/dh|_{h=0} = A' N_{part} + B' N_{coll}$$
- Qualitatively consistent with HIJING
- Inconsistent with *some* saturation models

- One interpretation:

- Evidence for both “soft” ($\sim N_{\text{participants}}$) “hard” ($\sim N_{\text{collisions}}$) processes



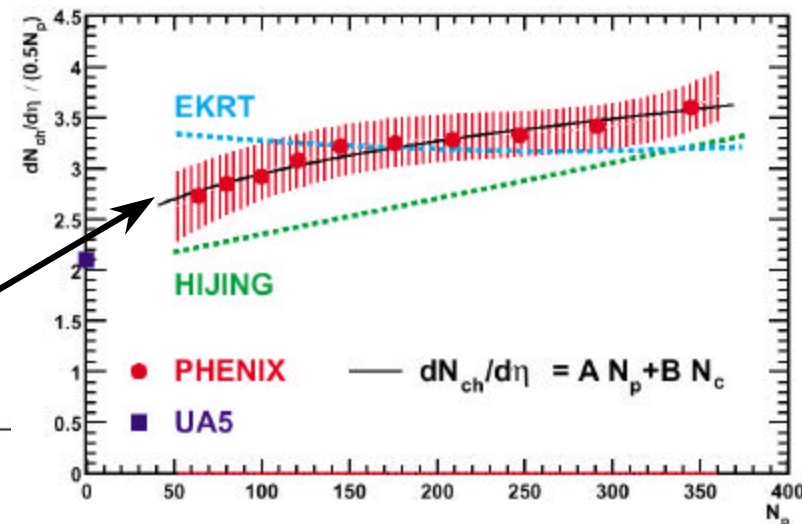
In this cartoon are there
5 N-N collisions $\Rightarrow 5 \times N_{pp}$?
OR
6 participants $\Rightarrow 3 \times N_{pp}$?

PHENIX Consistency between Experiments

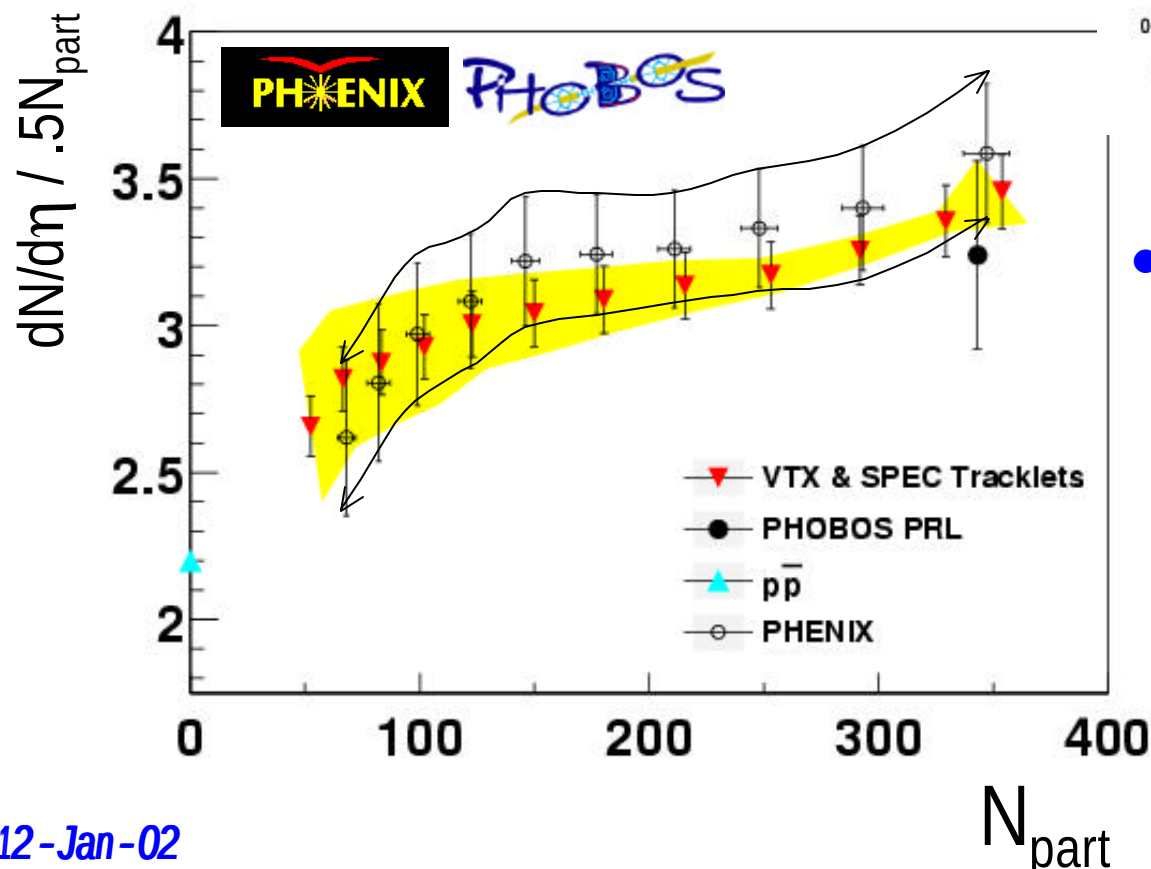


• Trend

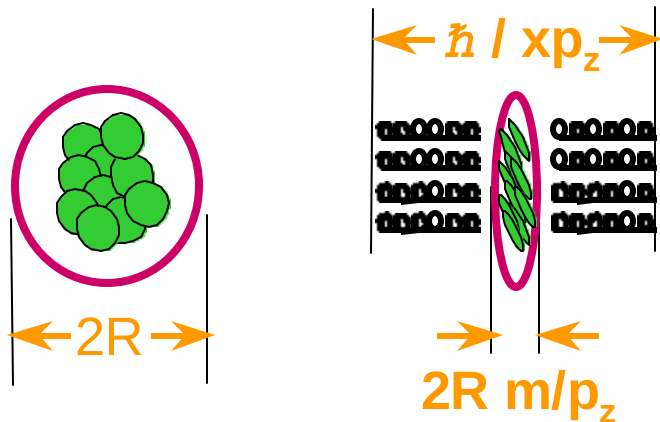
- ❑ incompatible with final-state gluon saturation model
- ❑ Good agreement with model based on *initial-state* saturation (Kharzeev and Nardi, nucl-th/0012025)



- Excellent agreement between (non-trivial) PHENIX and PHOBOS analyses of this systematic variation with nuclear overlap.



When do the gluons overlap significantly?



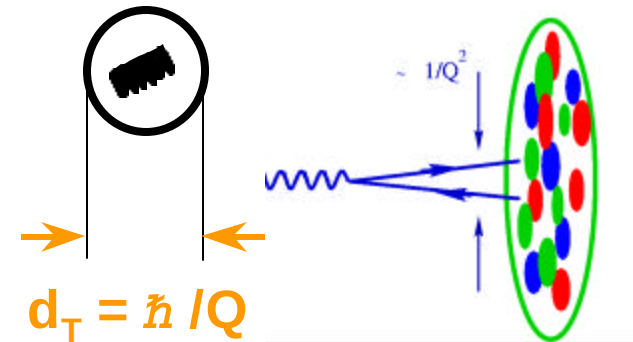
Longitudinal

So for $\hbar / mx \sim 2R$,

$\sim$ *all* constituents contribute

Parton density $r \sim A xG(x, Q^2) / pR^2$

¹ J.P Blaizot, A.H. Mueller, Nucl. Phys. B289, 847 (1987)



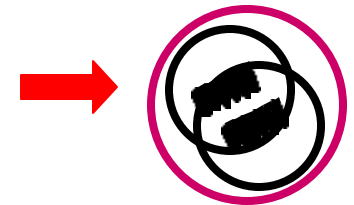
Transverse

Parton cross section

$s \sim a_S \hbar^2 / Q^2$

Saturation condition $rs \sim 1$

$Q_s^2 \sim a_S A xG(x, Q^2) / pR^2 \sim A^{1/3}$



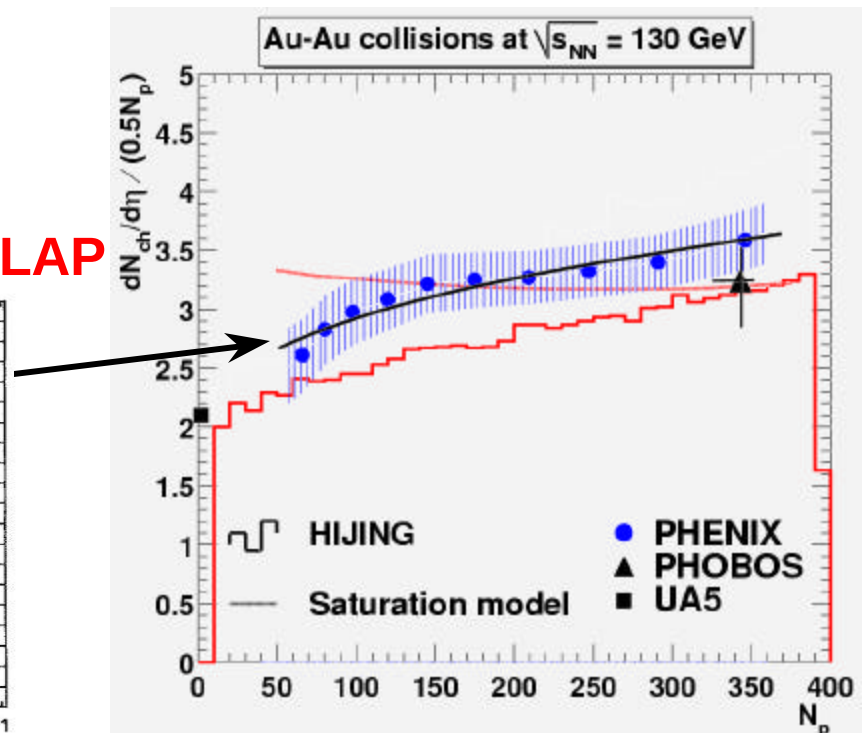
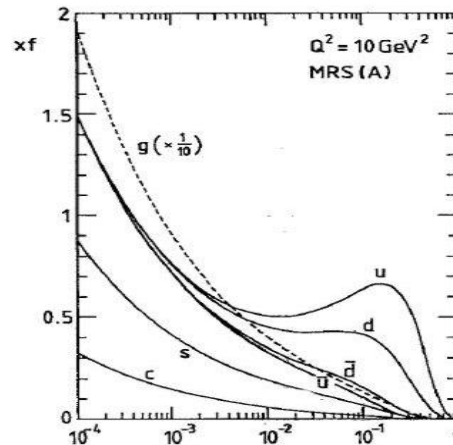
D. Kharzeev, nucl-th/0107033

● Procedure:

- Determine scale Q_S^2 as function of nuclear overlap
 - ◆ $Q_S^2 \sim A^{1/3} \sim N_{\text{part}}^{1/3}$
- Assume final-state multiplicities proportional to number of initial-state gluons in this saturated regime:
 - ◆ $dN/dh = c N_{\text{part}} xG(x, Q_S^2)$
- Note that DGLAP requires evolution of $xG(x, Q_S^2)$:
 - ◆ $xG(x, Q_S^2) \sim \ln(Q_S^2 / L_{\text{QCD}}^2)$

● Results:

- “Running” of the multiplicity yield directly from $xG(x, Q^2)$ + DGLAP





- Energy Density

- What energy densities are achieved in central collisions at RHIC?
 - How does this compare to previous experiments?

- Answered in

"Measurement of the Midrapidity Transverse Energy Distribution from $\sqrt{s} = 130$ GeV Au-Au Collisions at RHIC",

K. Adcox et al., [PRL 87 \(2001\) 052301](#),

[preprint nucl-ex/0104015](#),

or from the PHENIX site: [pdf](#), [ps](#), [ps.gz](#).



Q. How to estimate initial energy density?

A. From rapidity density

- *“Highly relativistic nucleus-nucleus collisions: The central rapidity region”, J.D. Bjorken, Phys. Rev. D27, 140 (1983).*

□ Assumes

- ◆ ~ 1-d hydrodynamic expansion

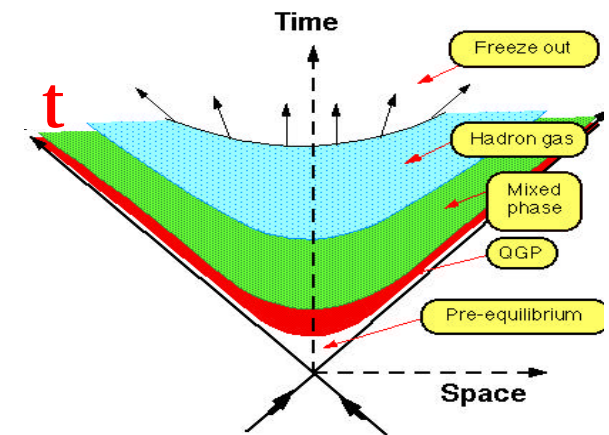
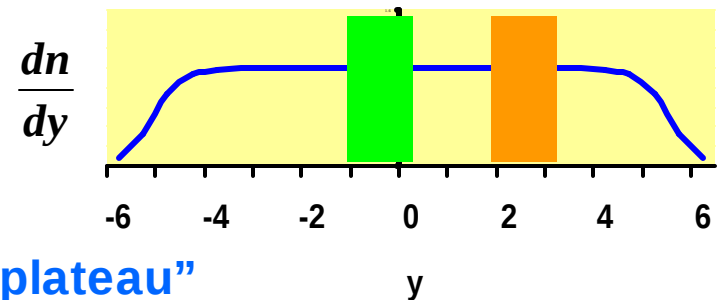
- ◆ Invariance in y along “central rapidity plateau” (i.e., flat rapidity distribution)

- ◆ Then

$$e = \frac{E}{V} \sim \frac{dE_T}{pR_T^2 dz} \sim \frac{dE_T}{pR_T^2 t dy} = \frac{1}{pR_T^2 t} \frac{dE_T}{dy}$$

since boost-invariance of matter

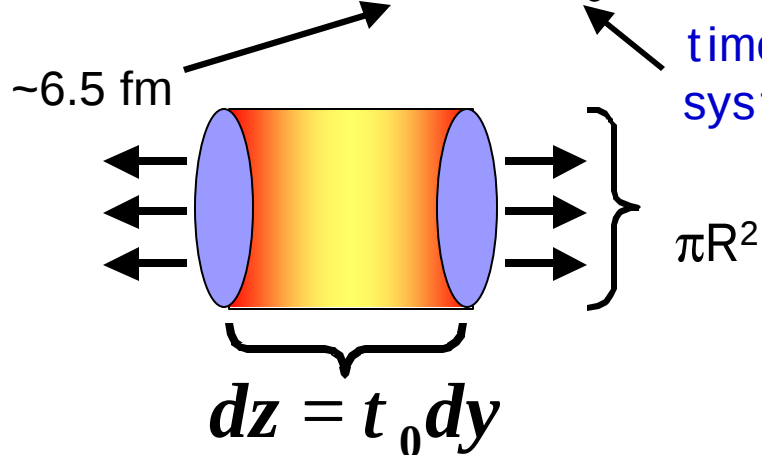
▷ $dz = t dy$ where $t \sim 1 \text{ fm}/c$



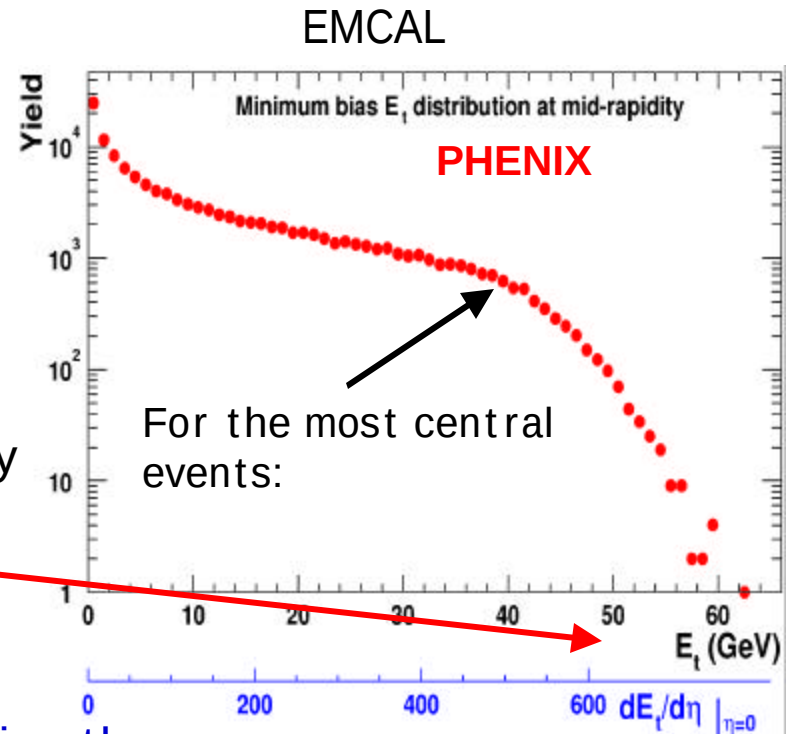
- What is the energy density achieved?
- How does it compare to the expected phase transition value ?

Bjorken formula for thermalized energy density

$$e_{Bj} = \frac{1}{p} \frac{1}{R^2} \frac{1}{t_0} \frac{dE_T}{dy}$$



time to thermalize the system ($\tau_0 \sim 1 \text{ fm}/c$)



$$\epsilon_{Bjorken} \sim 4.6 \text{ GeV}/\text{fm}^3$$

~30 times normal nuclear density
~1.5 to 2 times higher than
any previous experiments



- High Transverse Momenta Particle Production

- What is the rate for producing high p_T particles?
- How does it vary with centrality?

- Answered in

"Suppression of Hadrons with Large Transverse Momentum in Central Au+Au Collisions at $\sqrt{s} = 130$ GeV",

K. Adcox et al., Phys. Rev. Lett. 88, 022301 (2001), preprint nucl-ex/0109003, or directly from Phenix

PHENIX Energy Loss of Fast Partons



- Many approaches

- 1983: Bjorken

$$-\frac{dE}{dx} = \frac{3\sqrt{30}}{4} a_s^2 \sqrt{e} \ln \frac{4ET}{\epsilon M^2} \frac{\vec{0}}{\emptyset} \sim a_s^2 T^2 \ln \frac{4ET}{\epsilon M^2} \frac{\vec{0}}{\emptyset}$$

- 1991: Thoma and Gyulassy (1991)

$$-\frac{dE}{dx} \gg \frac{4p}{3} C_F a_s^2 T^2 \ln \frac{E}{\epsilon m_D} \frac{\vec{0}}{\emptyset}$$

- 1993: Brodsky and Hoyer (1993)

$$-\frac{dE}{dx} \gg \frac{\langle Dk_T^2 \rangle}{2}$$

- 1997: BDMPS- depends on path length(!)

$$-\frac{dE}{dx} \gg a_s \frac{C_R}{8} \frac{m_D^2}{l_g} L \ln \frac{L}{\epsilon l_g} \frac{\vec{0}}{\emptyset}$$

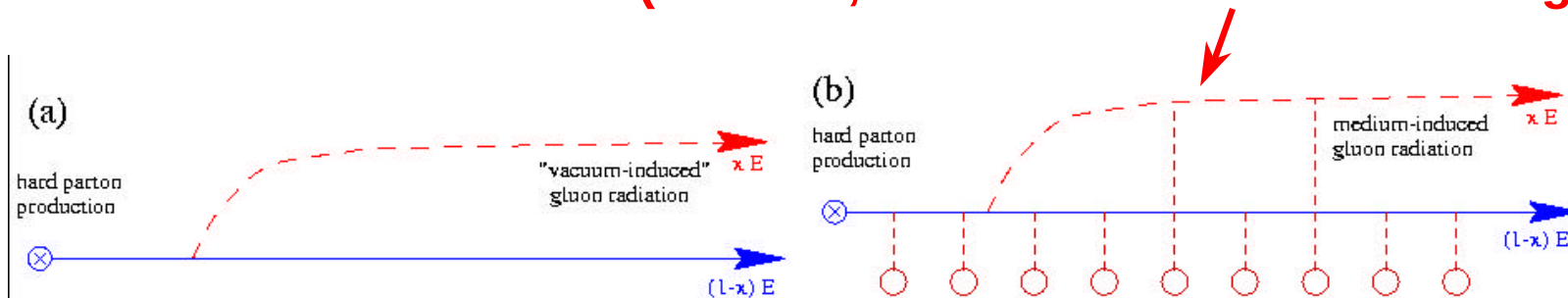
- 1998: BDMS

$$-\frac{dE}{dx} \gg a_s \frac{N_C}{4} \frac{\langle Dk_T^2 \rangle}{2}$$

- Numerical values range from

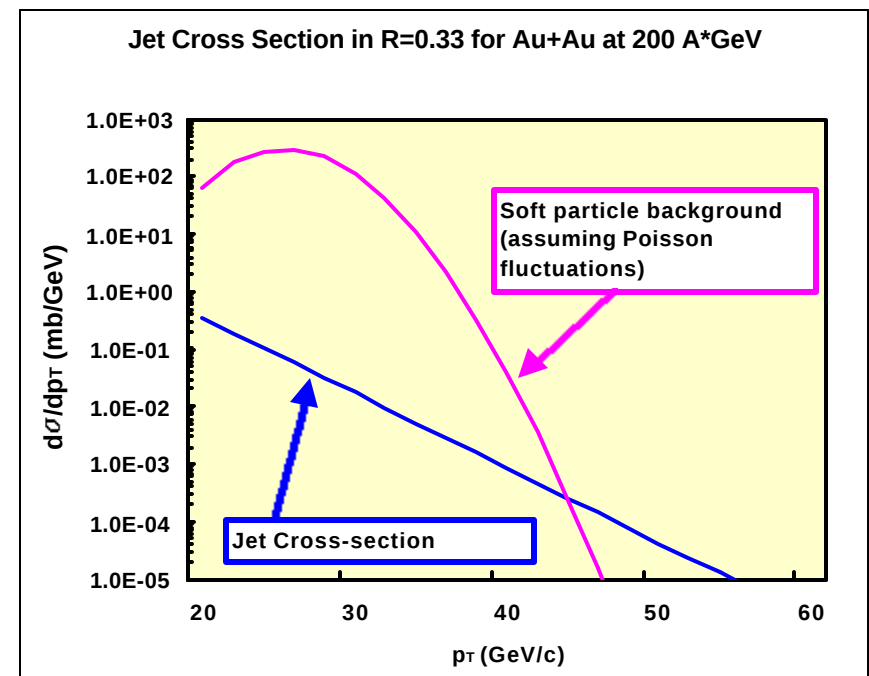
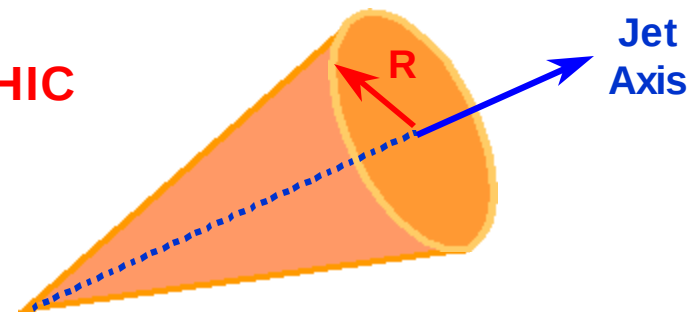
- ~ 0.1 GeV / fm (Bjorken, elastic scattering of partons)

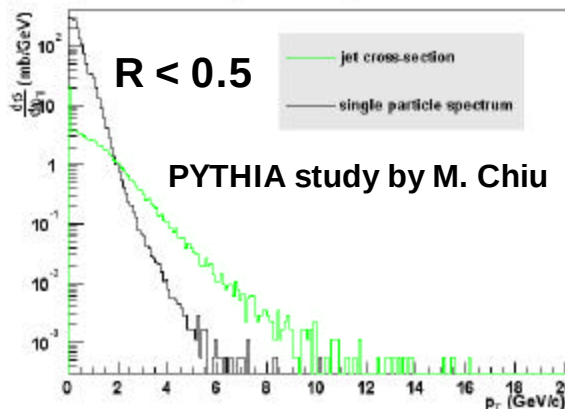
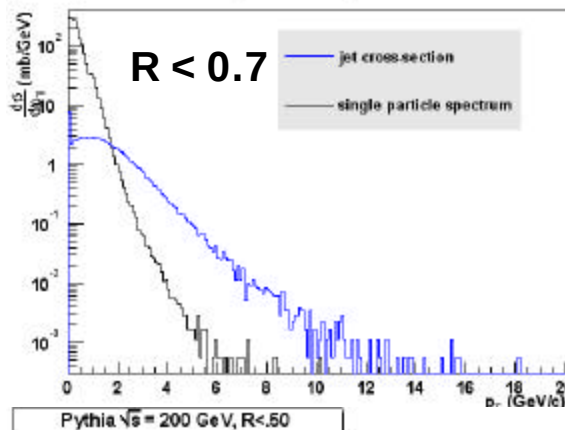
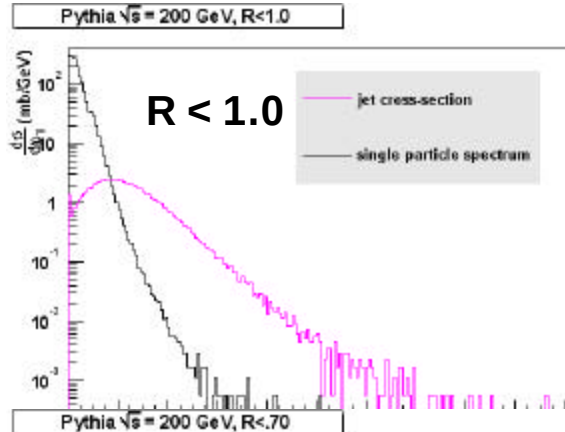
- ~several GeV / fm (BDMPS, non-linear interactions of gluons)



W.A. Zajc

- Tremendous interest in hard scattering
(and subsequent energy loss in QGP) at RHIC
 - Predictions that $dE/dx \sim$ (amount of matter to be traversed)
 - Directly due to non-Abelian nature of medium
- But:
 - “Traditional” jet methodology fails at RHIC
 - Dominated by the soft background:
 - ◆ For a typical jet cone $R = 0.33$
 $(R^2 = DF^2 + Dh^2)$
 have
 - $\langle n_{\text{SOFT}} \rangle \sim 64$
 - $\langle E_T \rangle \sim 25 \text{ GeV}$
 - ◆ Fluctuations in this soft background swamp any jet signal for $p_T < \sim 40 \text{ GeV}$:





- Why not let $R \rightarrow 0$?

- Good:
Reduces fluctuations of “underlying event”
- “Bad” :
Doesn’t contain all the energy of the jet

- But:

- For “quenching”, maximal effect expected on the “leading particle” (that is, “lost” energy probably stays in the cone, which *confuses* picture)
- High p_T single particles have a simple relation to parent parton dynamics

$$\frac{ds(\text{jet})}{dp_T} \sim \frac{A}{p_T^n}, \quad \frac{ds(\text{particle})}{dp_T} = \int_0^1 dz \frac{ds(\text{jet})}{dp_T} D(z) d(p_T - z p_T)$$

$$D(z) \sim (1-z)^m / z \quad \Rightarrow \quad \frac{ds(\text{particle})}{dp_T} = (\text{constant}) \frac{ds(\text{jet})}{dp_T} \gg \frac{1}{100} \frac{ds(\text{jet})}{dp_T}$$

Hadron Interactions, Collins and Martin



- As you like it:

$$\frac{dp^m}{dt} = \frac{e}{c} u_n F^{mn}$$

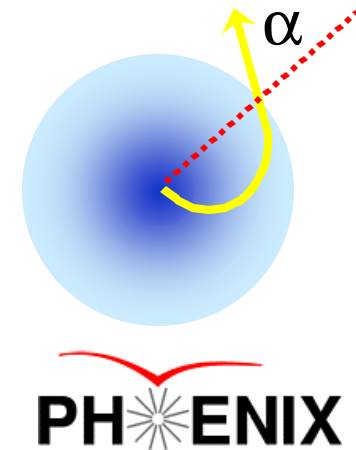
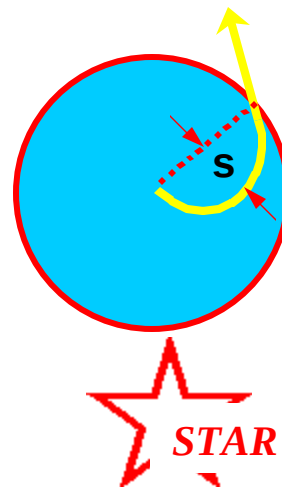
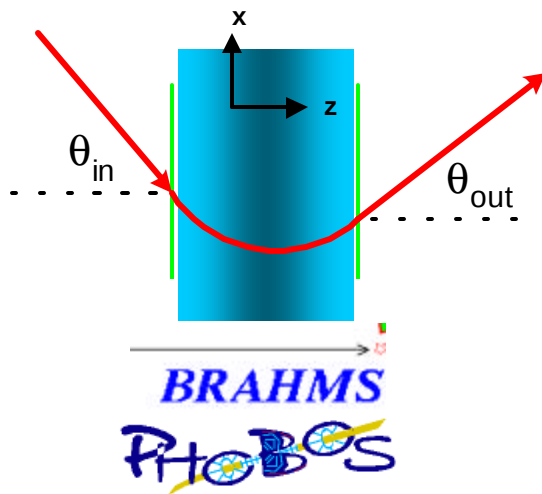
$$\frac{d\vec{p}}{dt} = \frac{e}{c} \vec{v} \times \vec{B}$$

$$\frac{d}{ds} \frac{d\vec{r}}{ds} = \frac{e}{c} \frac{d\vec{r}}{ds} \times \frac{\vec{B}}{|\vec{p}|}$$

- More usefully:

$$|\vec{p}_\perp| = \frac{e}{c} B \times R, \quad \frac{e}{c} = \frac{0.3 \text{ GeV} / c}{\text{Tesla} \cdot \text{meter}}$$

➔ 1 meter of 1 Tesla field deflects 1 GeV/c by ~17°



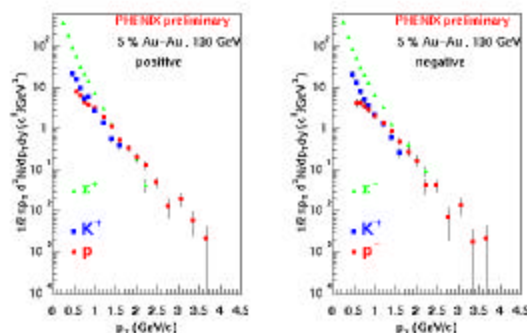
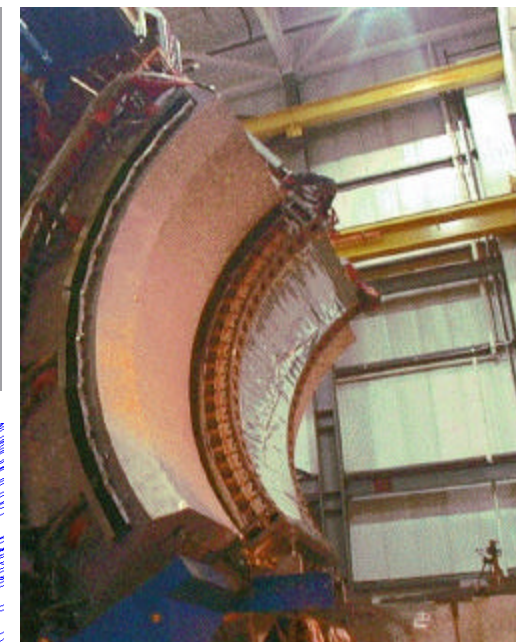
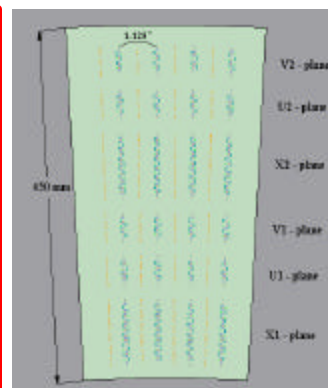
Real world: $\frac{dp}{p} = (\sim 1\%) \text{ \AA} (\sim 1\%) \times p \text{ [GeV/c]}$

~stuff in aperture

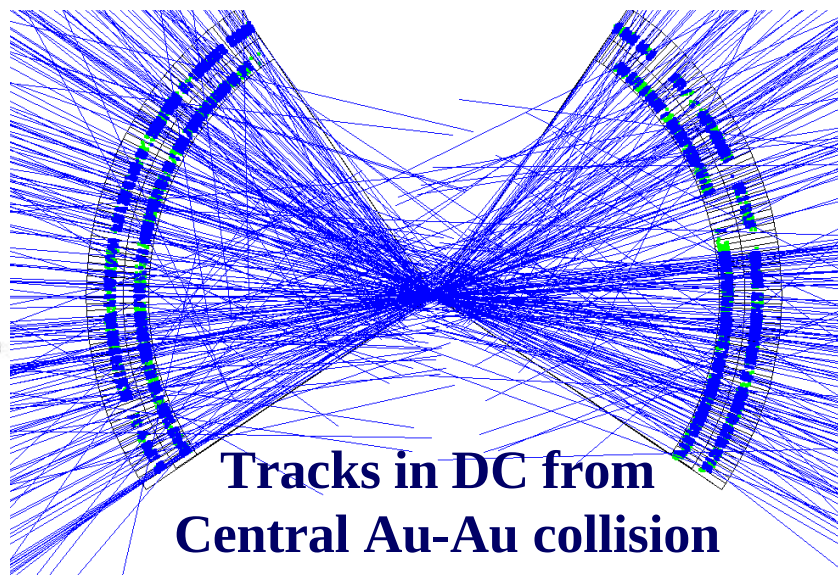
~spatial accuracy

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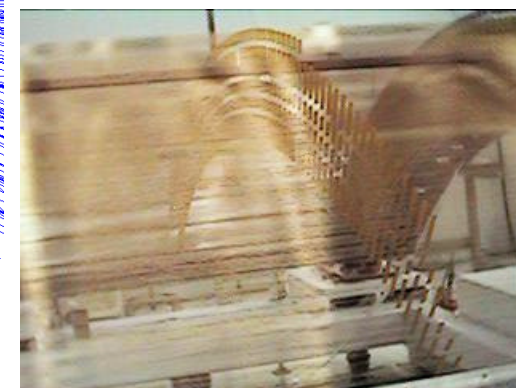
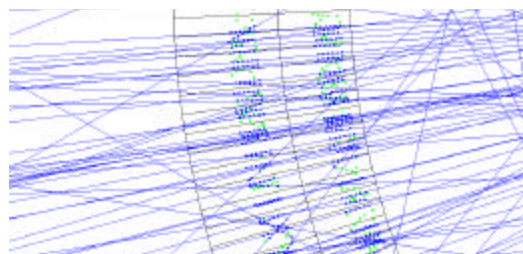
- Jet -chamber anode/cathode structure modified for HI high multiplicity
- Joint Russia/US design & construction
- All Titanium frame
- $\sigma_x = 120 \mu\text{m}$, two-track sep = 2mm



Identified particle spectra using tracking system and TOF

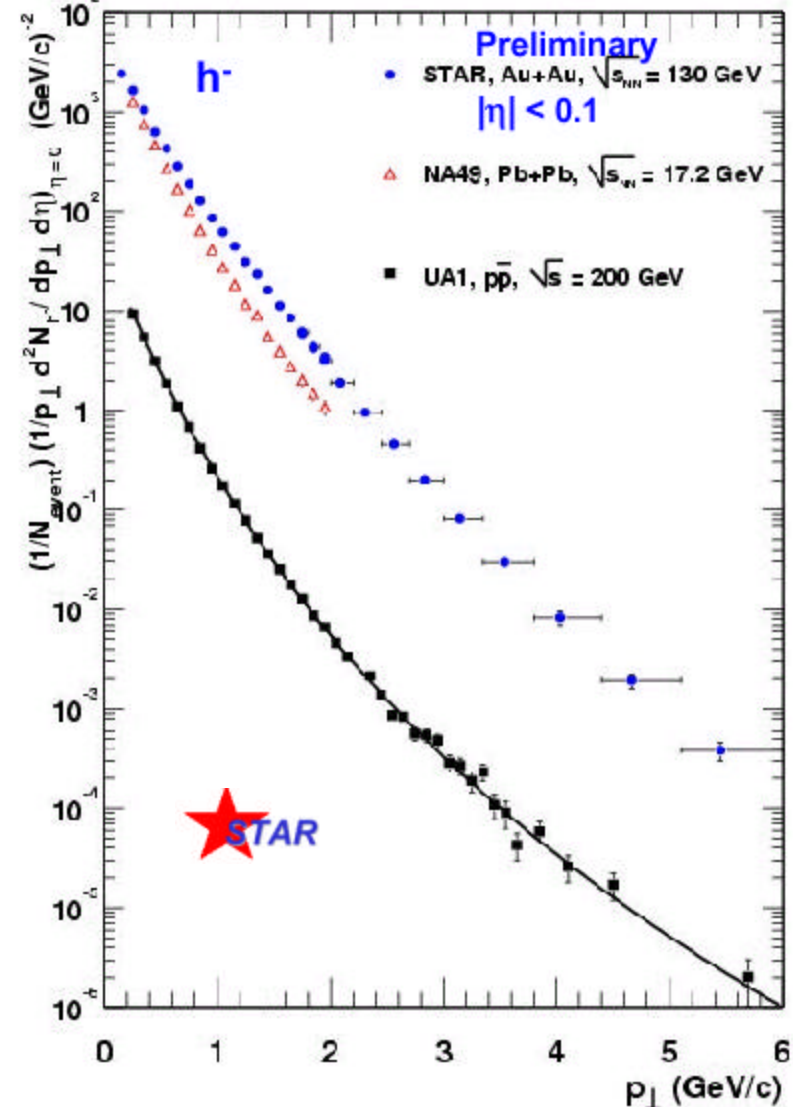
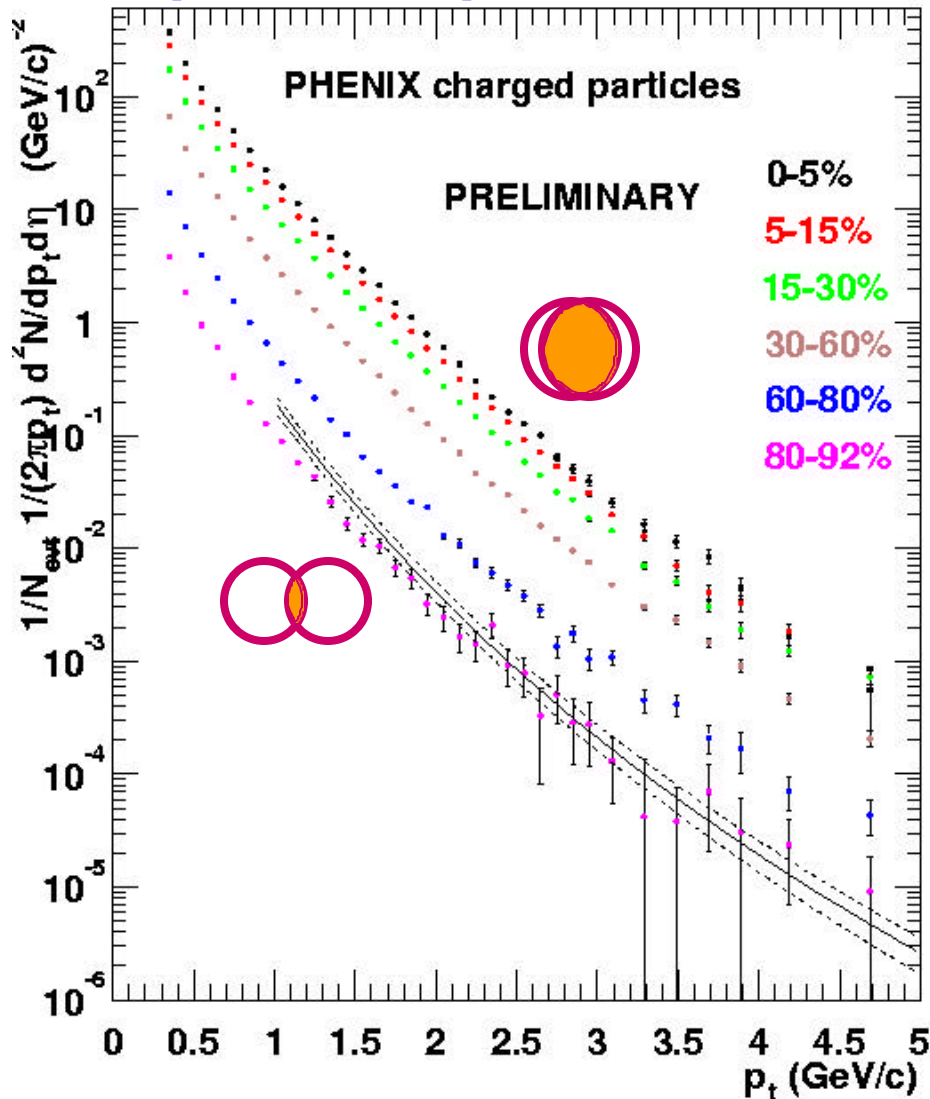


Tracks in DC from Central Au-Au collision

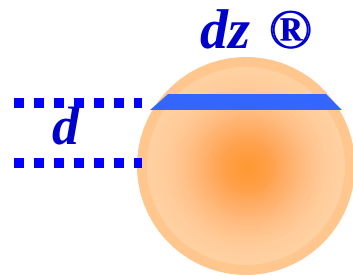


DC wires with kapton wire dividers
W.A. Zajc

Both PHENIX and STAR have measured charged particle spectra out to “small” cross sections

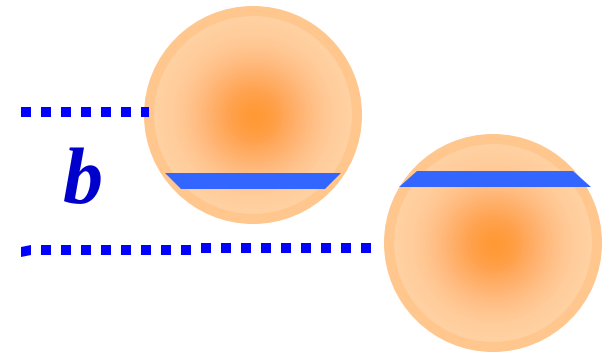


- Particle production via *rare processes* should scale with N_{coll} , the number of underlying binary nucleon-nucleon collisions



$$T_A(d) = \int_0^{\infty} r(d, z) dz$$

Thickness Function



$$T_{AB}(b) = \int_0^{\infty} T_A(\bar{s} + \frac{\bar{b}}{2}) T_B(\bar{s} - \frac{\bar{b}}{2}) d\bar{s}$$

Overlap Function

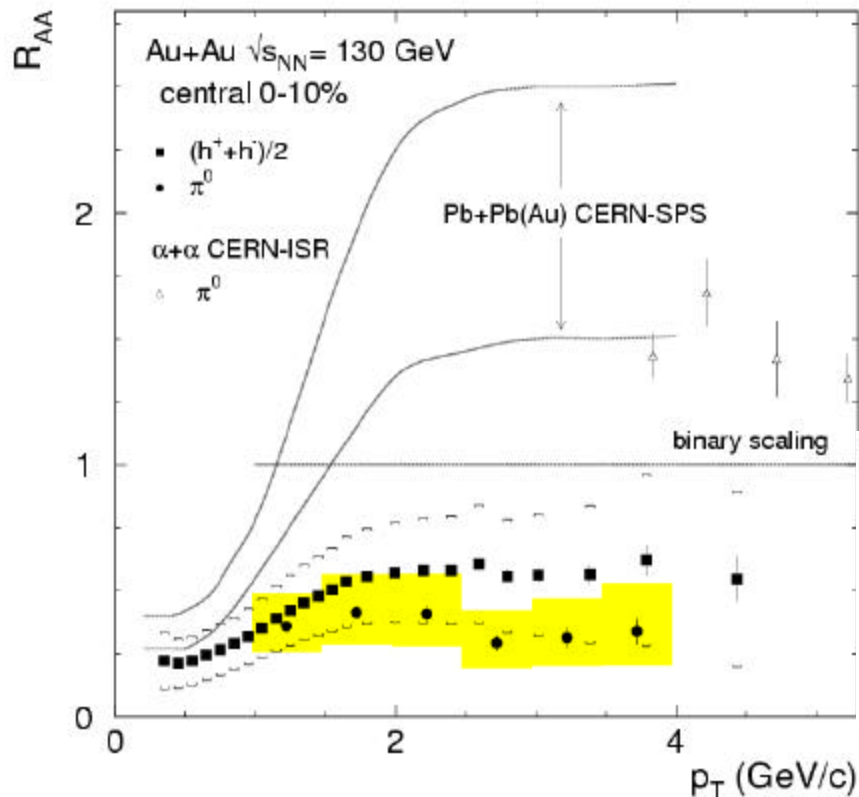
If Nucleus "A" has A constituents and Nucleus "B" has B constituents which interact with cross section s_{INT} the TOTAL cross section s_{AB} is then

$$s_{AB} = \int d^2b \left[1 - e^{-s_{\text{INT}} T_{AB}(b)} \right]$$

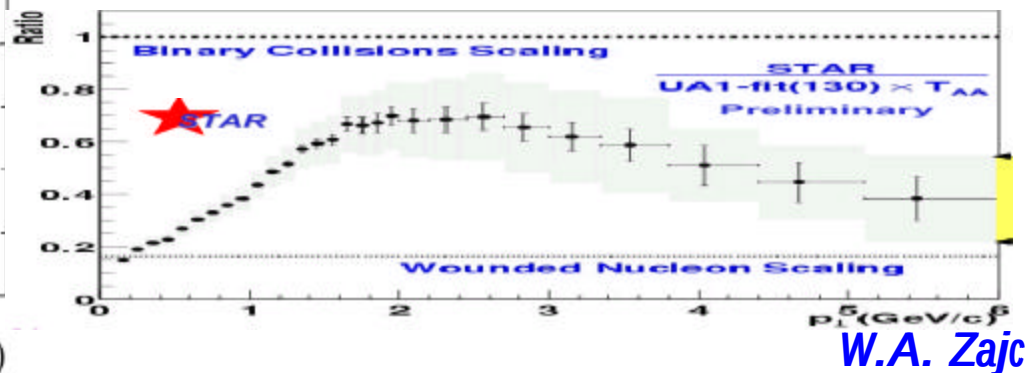
② $\gg A \times B \times s_{\text{INT}}$ for "small" s_{INT}

- ➔ Assuming no "collective" effects
- Test this on production at high transverse momentum

- NO!
- Both STAR and PHENIX data see deficit of particles at high transverse momenta, *relative to expected A*B scaling*:



Deficit opposite sign of enhancements previously seen in nuclear collisions (Cronin effect)





- A good example of a “combinatoric” background
- Reconstruction is *not* done particle-by-particle
- Recall: $p^0 \otimes gg$ and there are $\sim 200 p^0$ s per unit rapidity

□ So:

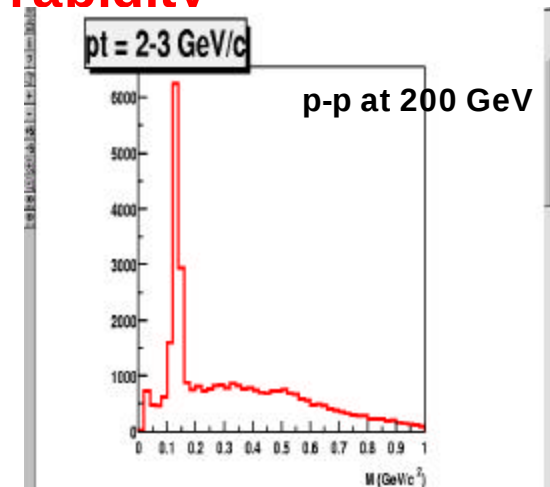
$$p^0_1 \otimes g_{1A} + g_{1B}$$

$$p^0_2 \otimes g_{2A} + g_{2B}$$

$$p^0_3 \otimes g_{3A} + g_{3B}$$

$$p^0_N \otimes g_{NA} + g_{NB}$$

PHENIX
 π^0 reconstruction
 $p_T > 2 \text{ GeV/c}$
 Asymmetry < 0.8



□ Unfortunately, nature doesn't use subscripts on photons

➔ N correct combinations: $(g_{1A} g_{1B}), (g_{2A} g_{2B}), \dots (g_{NA} g_{NB}),$

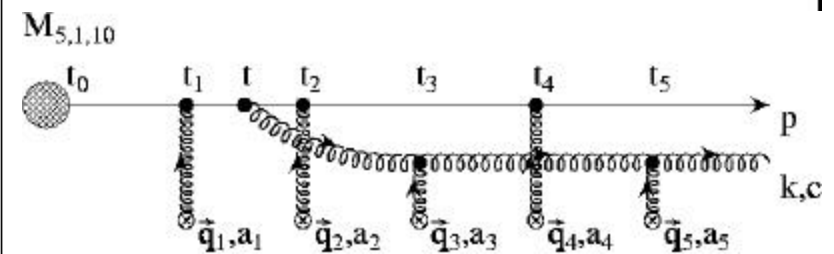
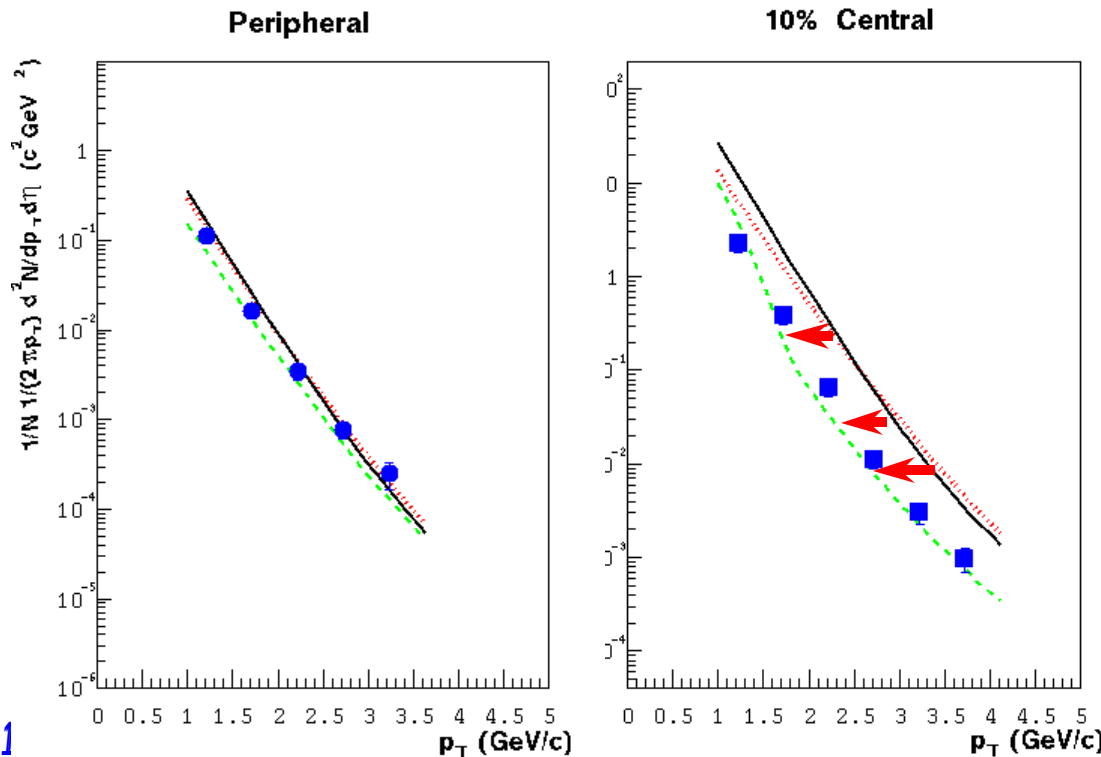
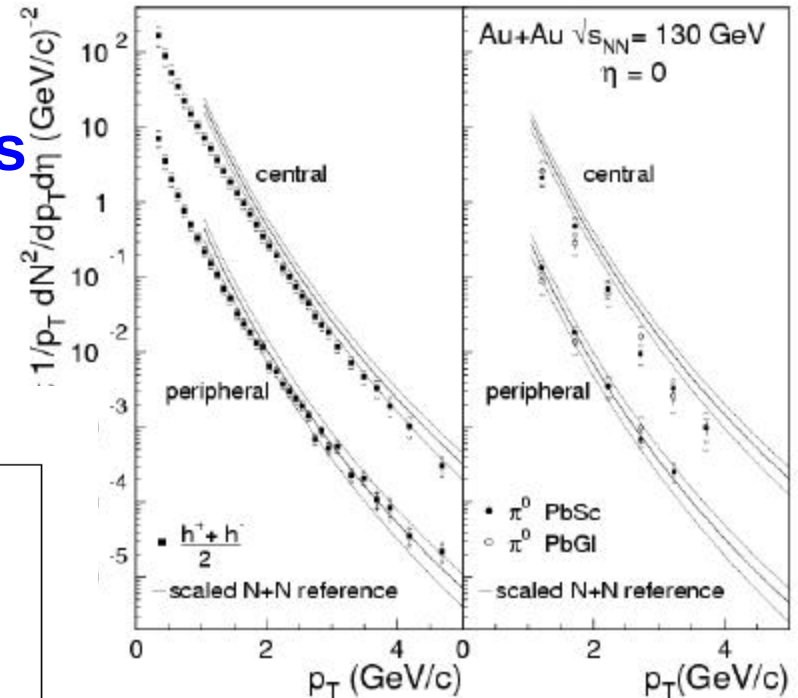
➔ $N(N-1)/2 - N$ *incorrect* combinations $(g_{1A} g_{2A}), (g_{1A} g_{2B}), \dots$

✗ Incorrect combinations $\sim N^2$ (!)

✓ Solution: Restrict N by p_T cuts

use high granularity, high resolution detector

- PHENIX sees a *larger* effect in the p^0 spectrum:
- Consistent with crude estimates of additional energy loss in a *deconfined medium*





- Some would (have) claimed just that
- To be done before this can be experimentally established:

- ❑ Improve p_T range (Run-2)
- ❑ Measure p-p spectrum (Run-2)

arXiv:nucl-th/0104035 10 Apr 2001

Discovery of Jet Quenching at RHIC and the Opacity of the Produced Gluon Plasma

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The predicted quenching of STAR and PHENIX in preliminary theory [1] of QCD radiative energy loss produced in $Au + Au$ collision is two Cronin enhancement of π^0 of $p_T \sim 2 - 4$ GeV π^0 report static plasma opacity $L/\lambda_g \approx$

1. Energy Loss of High E

The moderately high $p_T < 6$ π^0 production [2], reported at

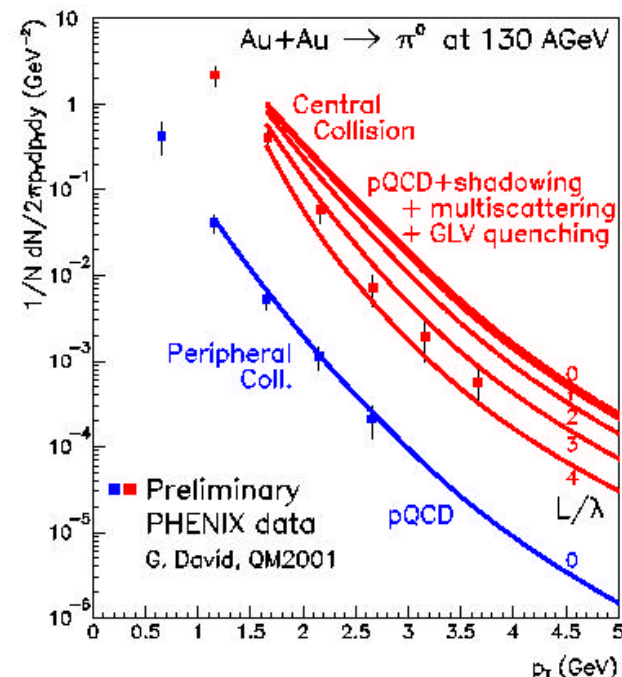


Figure 3. Pion production in Au+Au collision including jet quenching with opacity $L/\lambda = 1, 2, 3, 4$. Data are from Ref. [2].

➡ Study in “cold nuclear matter” with p-A collisions (Run-2?)

❑ Vary system size (Run-2???)



- Particle Composition

- What are the relative yields of pions, kaons and protons?
- How do their yields depend on transverse momentum?
- How do their yields depend on centrality?

- Answered in

"Centrality dependence of π^{+-} , K^{+-} , p and $\bar{p}$ production from $\sqrt{s}=130$ GeV Au + Au collisions at RHIC",

K. Adcox et al., Submitted to PRL,
[preprint nucl-ex/0112006](#),

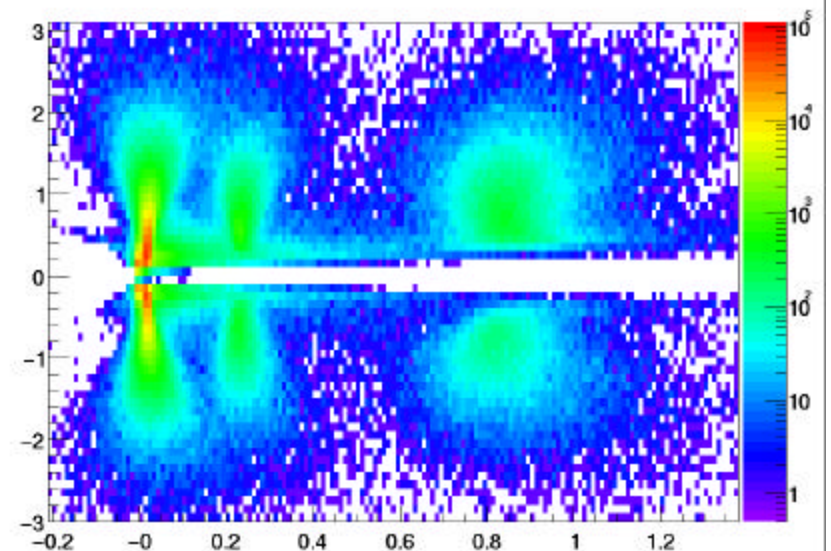
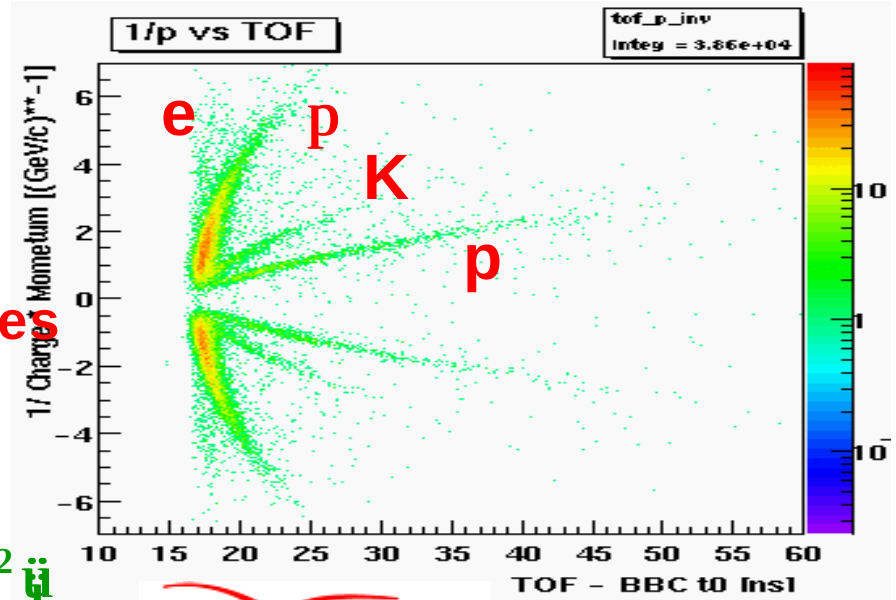
- The most direct way

- Measure β by distance/time
- Typically done via scintillators read-out with photomultiplier tubes
- Time resolutions ~ 100 ps
- Exercise: Show

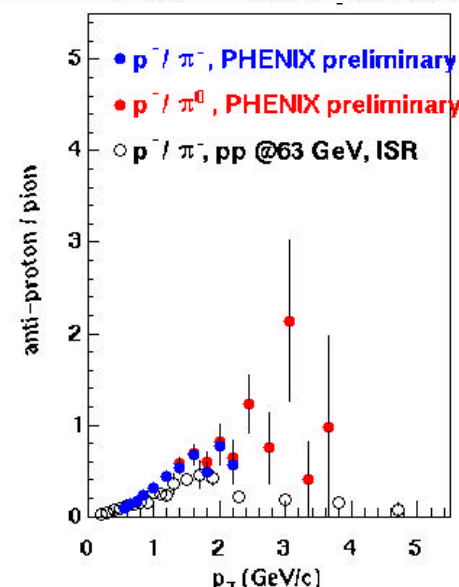
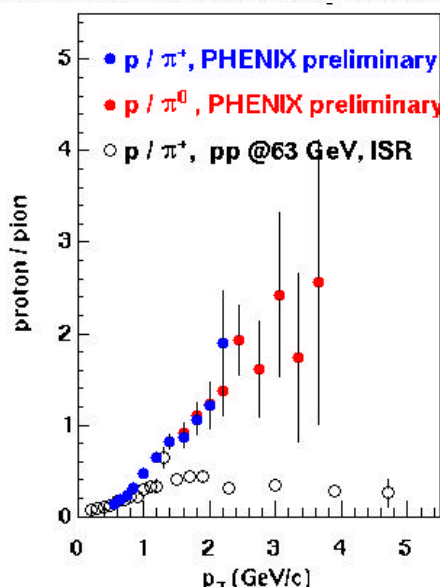
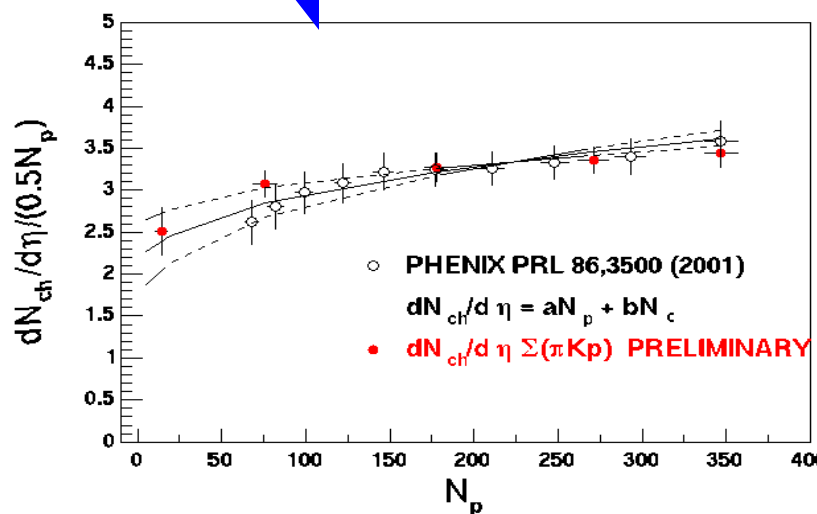
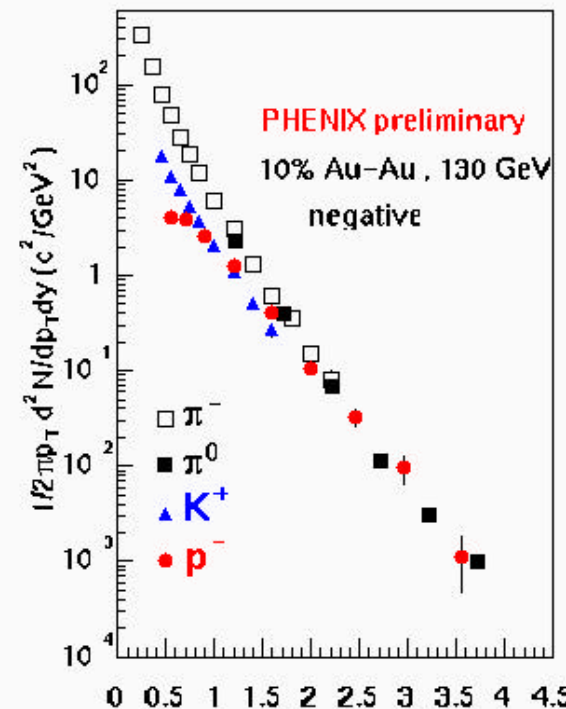
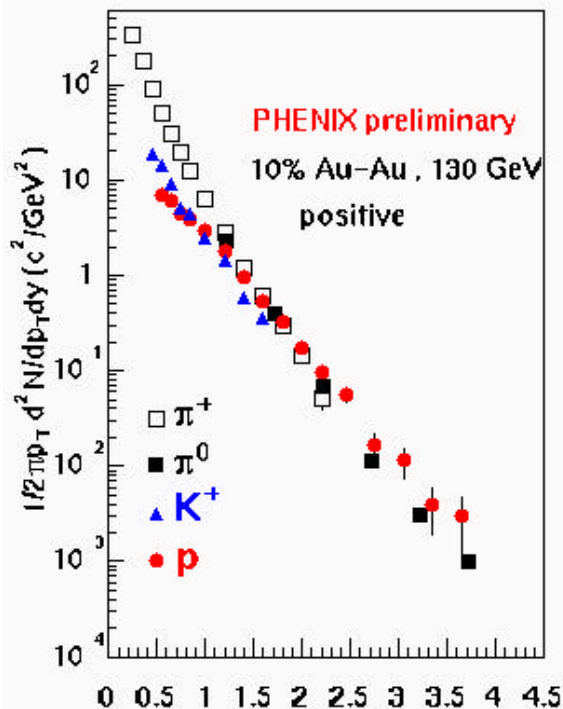
$$\frac{d}{dt} \left(\frac{dm}{m} \right)^2 = \frac{d}{dt} \left(\frac{dp}{p} \right)^2 + g^4 \left(\frac{d}{dt} \left(\frac{dt}{t} \right)^2 + \frac{d}{dt} \left(\frac{ds}{s} \right)^2 \right)$$

- Performance:

- $dt \sim 100$ ps on 5 m flight path
- ➔ P/K separation to ~ 2 GeV/c
- ➔ K/p separation to at least 4 GeV/c

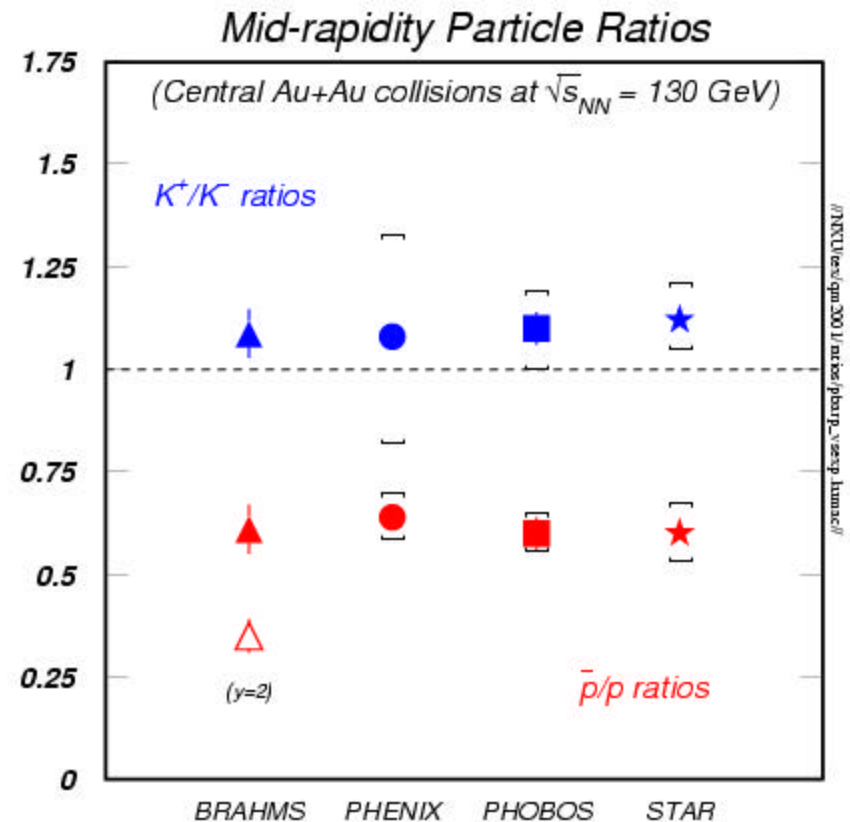
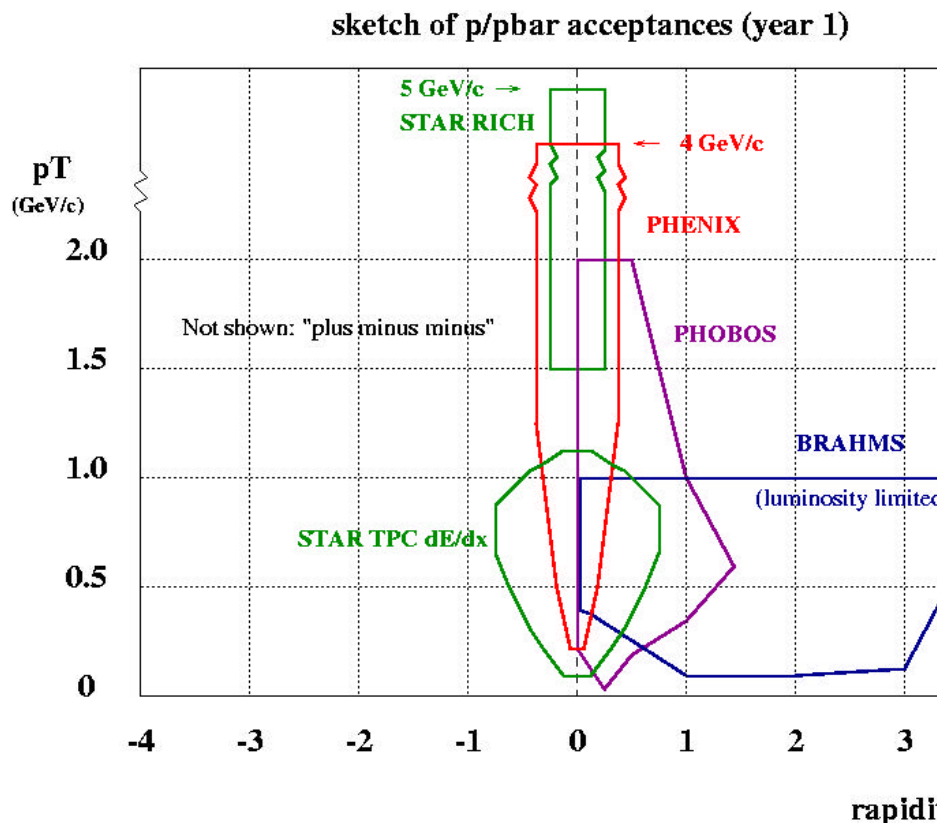


- Charged pions match nicely with p^0 's
 - Supports conclusion of preferential suppression of pions
 - "High" p_T dominated by protons and kaons??
- Summed hadrons consistent with our first result on multiplicity



- Overlaps in detector capabilities a feature of RHIC program

- Excellent agreement in common observables provides confidence

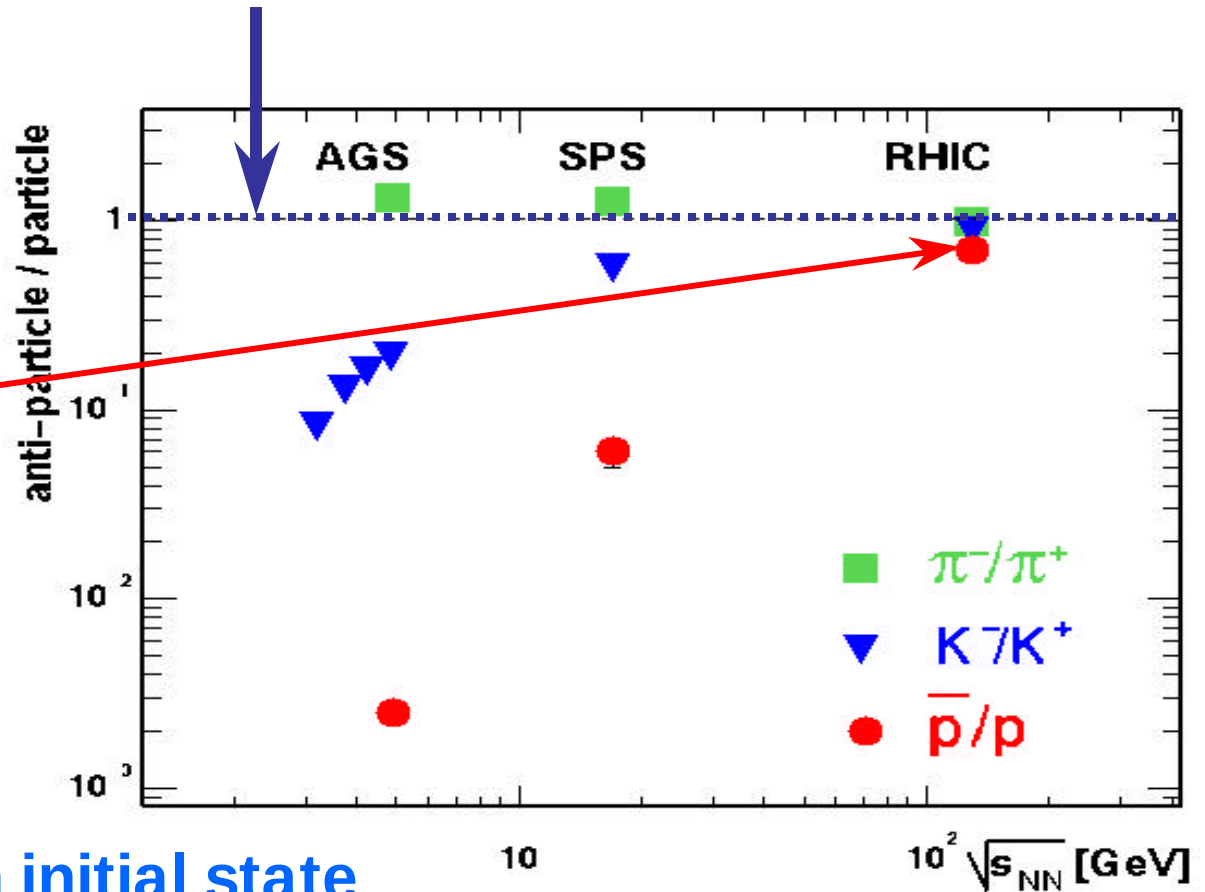




- Early Universe:

- Anti-proton/proton = 0.9999999999

- We've created "pure" matter approaching this value
- For the first time in heavy ion collisions, more baryons are pair-produced than brought in from initial state





- This is a sample of the many physics topics addressed by PHENIX in Run-1.

- Additional results available on

□ Flow

□ Fluctuations in

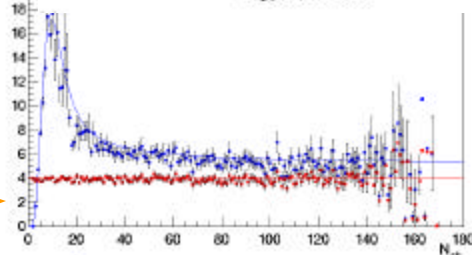
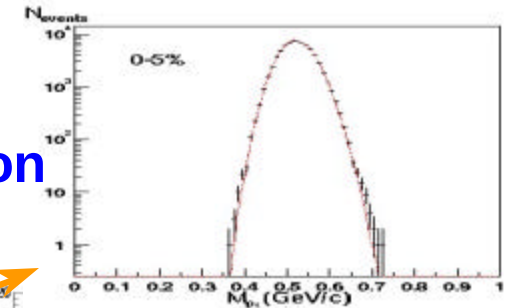
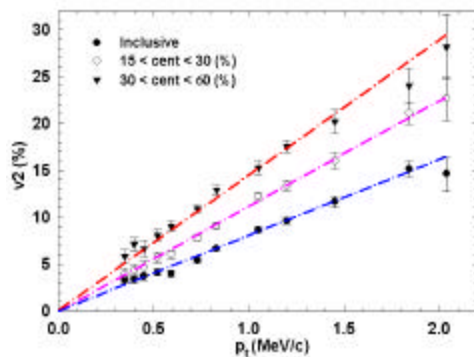
◆ $\langle p_T \rangle$

◆ Charge

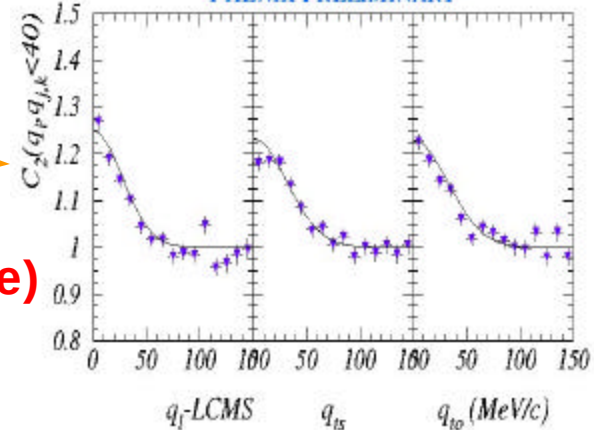
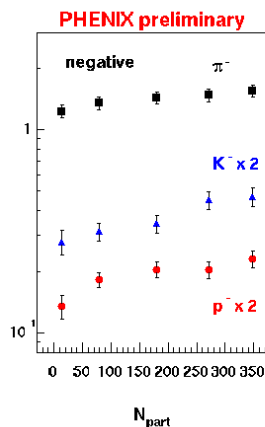
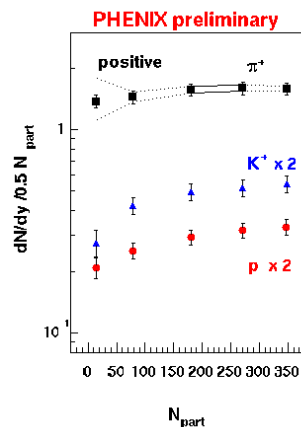
□ Centrality dependence of particle yields

□ HBT

□ Charm (next lecture)



$$\begin{aligned} \langle R^2 \rangle - \langle R \rangle^2 & \propto N_{ch} & (R = N_+/N_-) \\ 4(\langle Q^2 \rangle - \langle Q \rangle^2) & / N_{ch} & (Q = N_+ - N_-) \end{aligned}$$





- ✓✓✓✓● How many particles are produced?
More than ever produced previously
- ✓✓✓✓● Have we made contact with the early universe?
Closer than all previous heavy ion experiments
- ✓✓✓ ● Have we made sufficiently high energy densities for the phase transition?
Yes
- ✓✓✓ ● Is there evidence that the dense matter behaves collectively?
Yes
- ✓✓ ● Are there results consistent with the formation of a new state of matter?
Perhaps ??????!

(1 ✓ = 1 experiment)