

# Geometry Rules:

## Heavy Ion Physics of Heavy Quarkonia and Harmonic Flow



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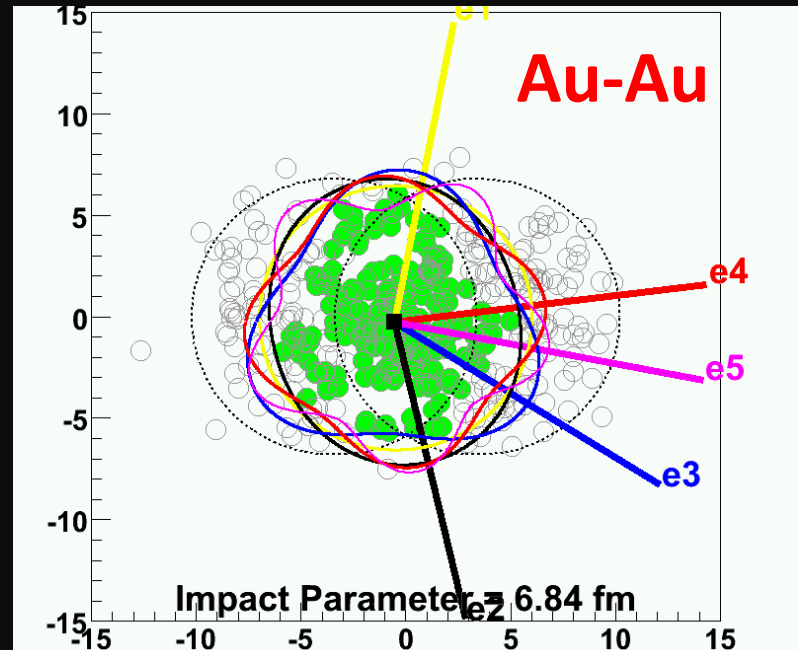
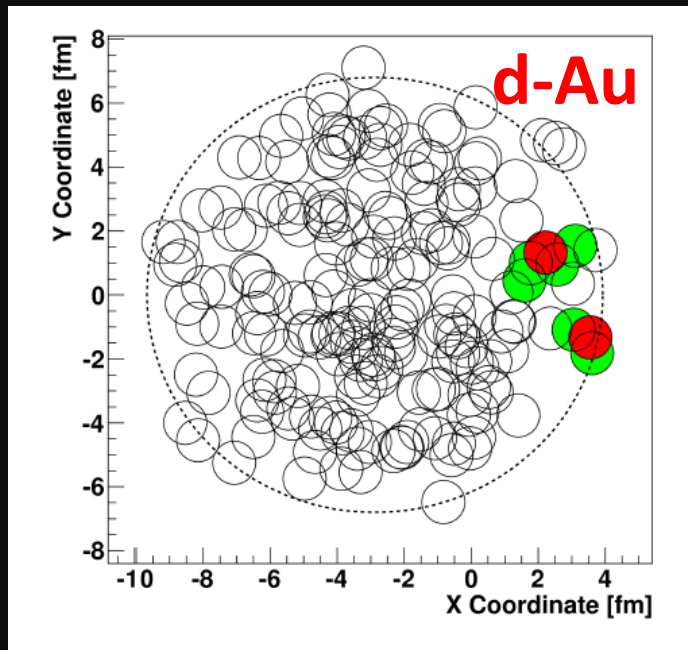


Focus on two topics in today's talk...

1. Heavy Quarkonia ( $J/\psi$ ) in p-p, d-Au, Au-Au

2. Harmonic Flow of the bulk QGP

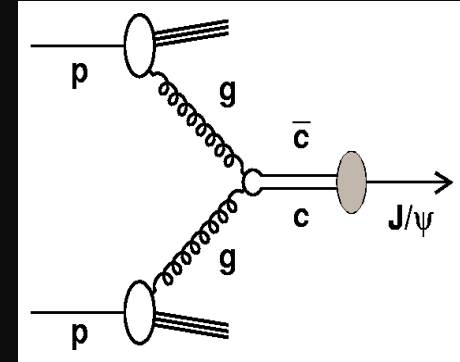
Highlight how geometric dependencies provide key insights...



# Heavy Quarkonia

## 1. $J/\psi$ Production Mechanisms

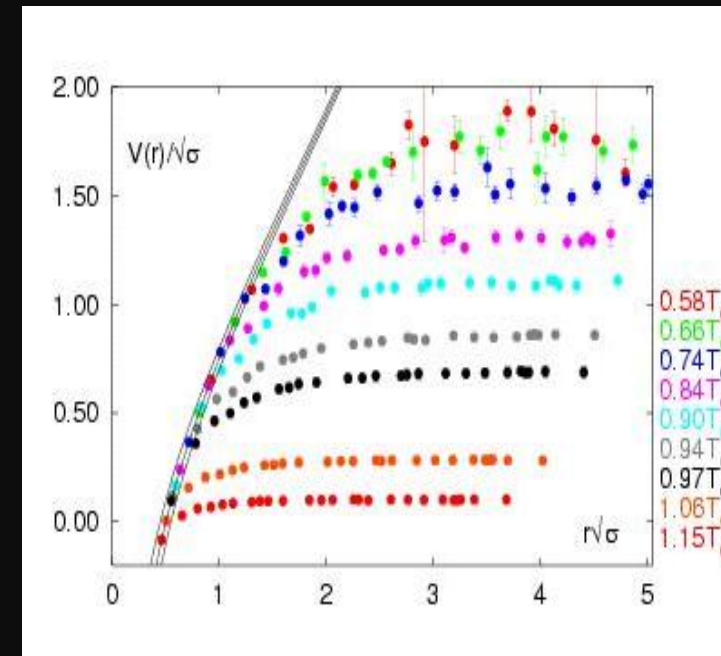
Not understood and not addressed in this talk  
see excellent review in arXiv:1010.5827



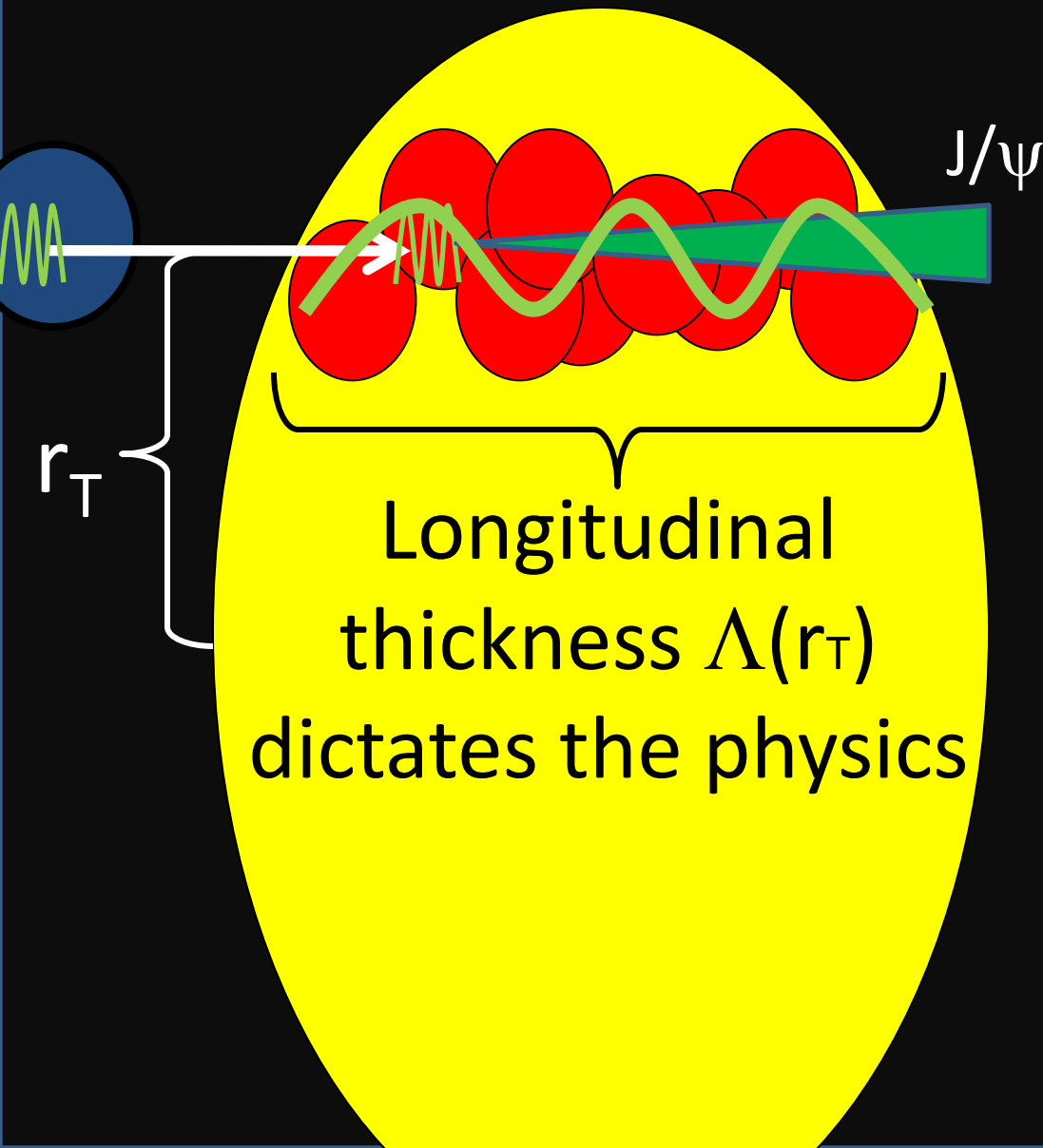
## 2. Time evolution into color neutral $J/\psi$ state (d-Au reactions are an excellent test)

## 3. Color screening in QGP (Au-Au reactions are the test)

Highlight results from PHENIX Experiment  
d-Au (Alex Linden-Levy) and  
Au-Au (Matt Wysocki Ph.D. thesis) and  
theory paper (JN, Frawley, Linden-Levy, Wysocki,  
arXiv:1011.4534)



# J/ψ Modifications in p-A Reactions



gluon-gluon  $\rightarrow c \bar{c}$

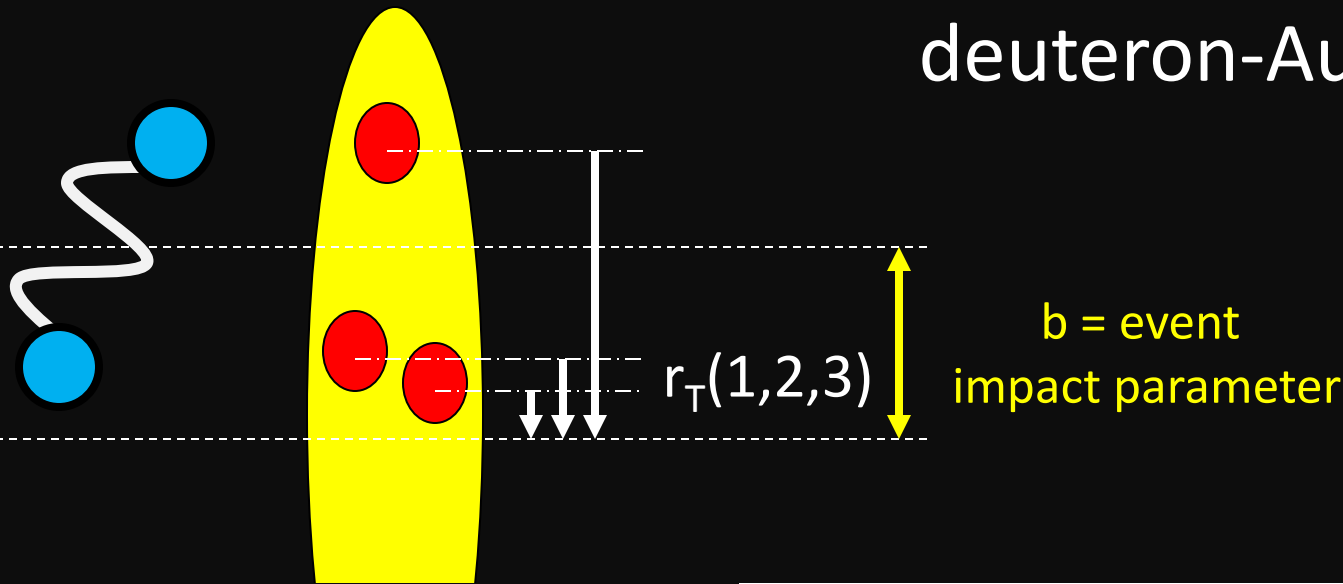
$c \bar{c}$  evolves to J/ψ on a timescale of 0.3 fm/c

At RHIC energies, the J/ψ forms outside the Lorentz contracted nucleus

Simplest modifications

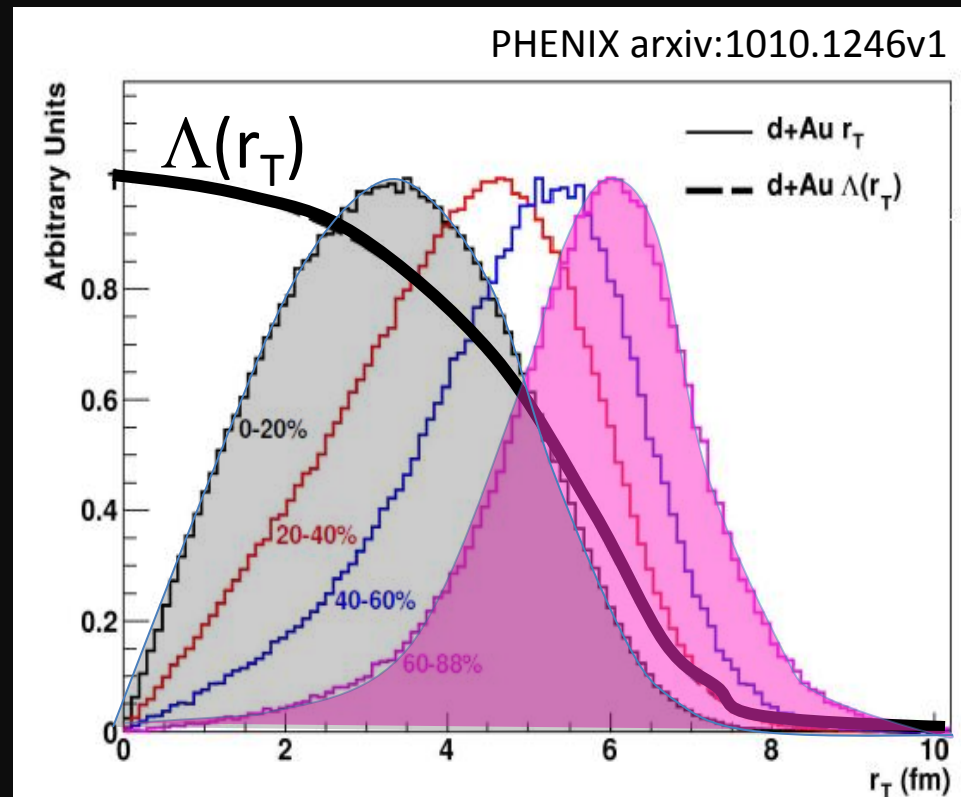
1. Shadowing nPDF changes initial charm pair production.
2. Breakup of charm pair in nucleus

# deuteron-Au collision

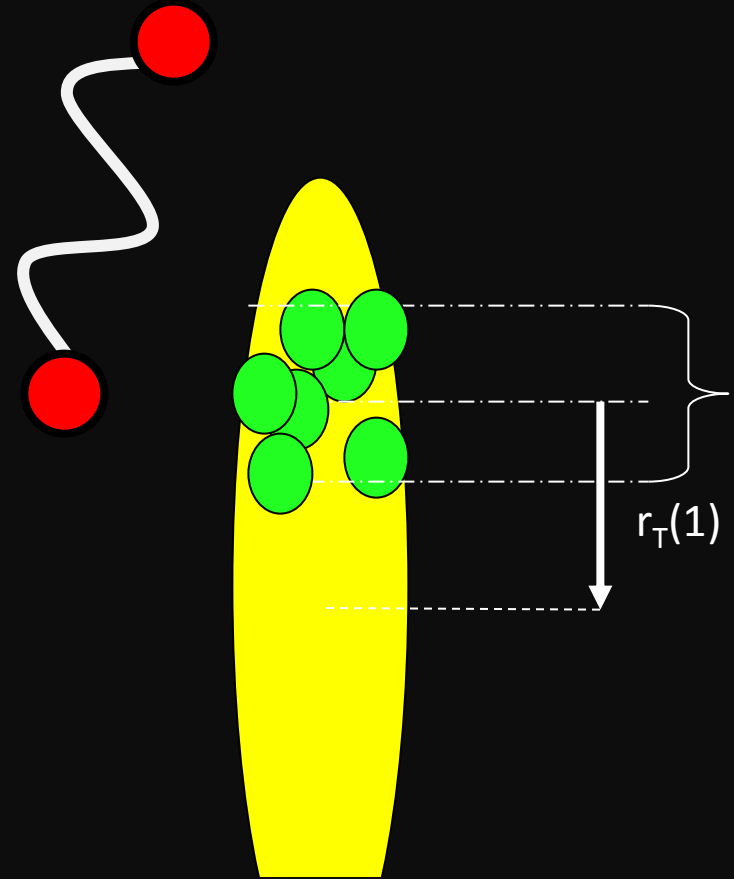
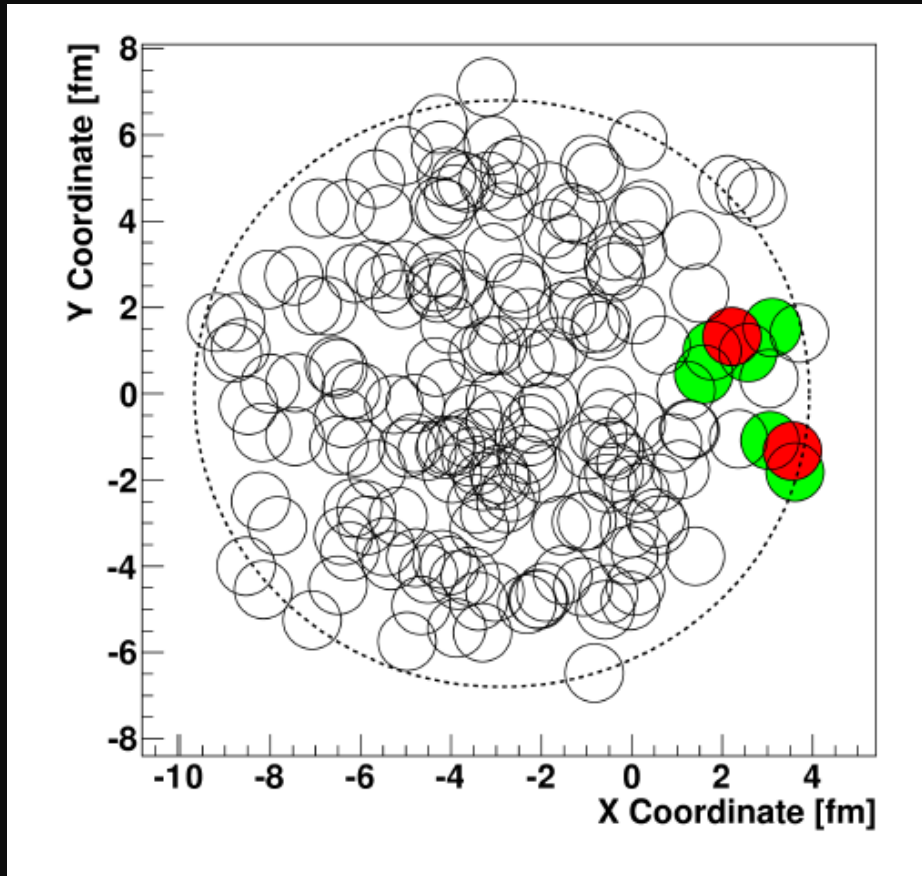


PHENIX collected a high statistics d-Au @ 200 GeV data set in 2008.

Using multiplicity in Beam-Beam Counter ( $3.1 < \eta < 3.9$ ) we select on d-Au centrality, then use Glauber MC to determine  $r_T$  distributions.

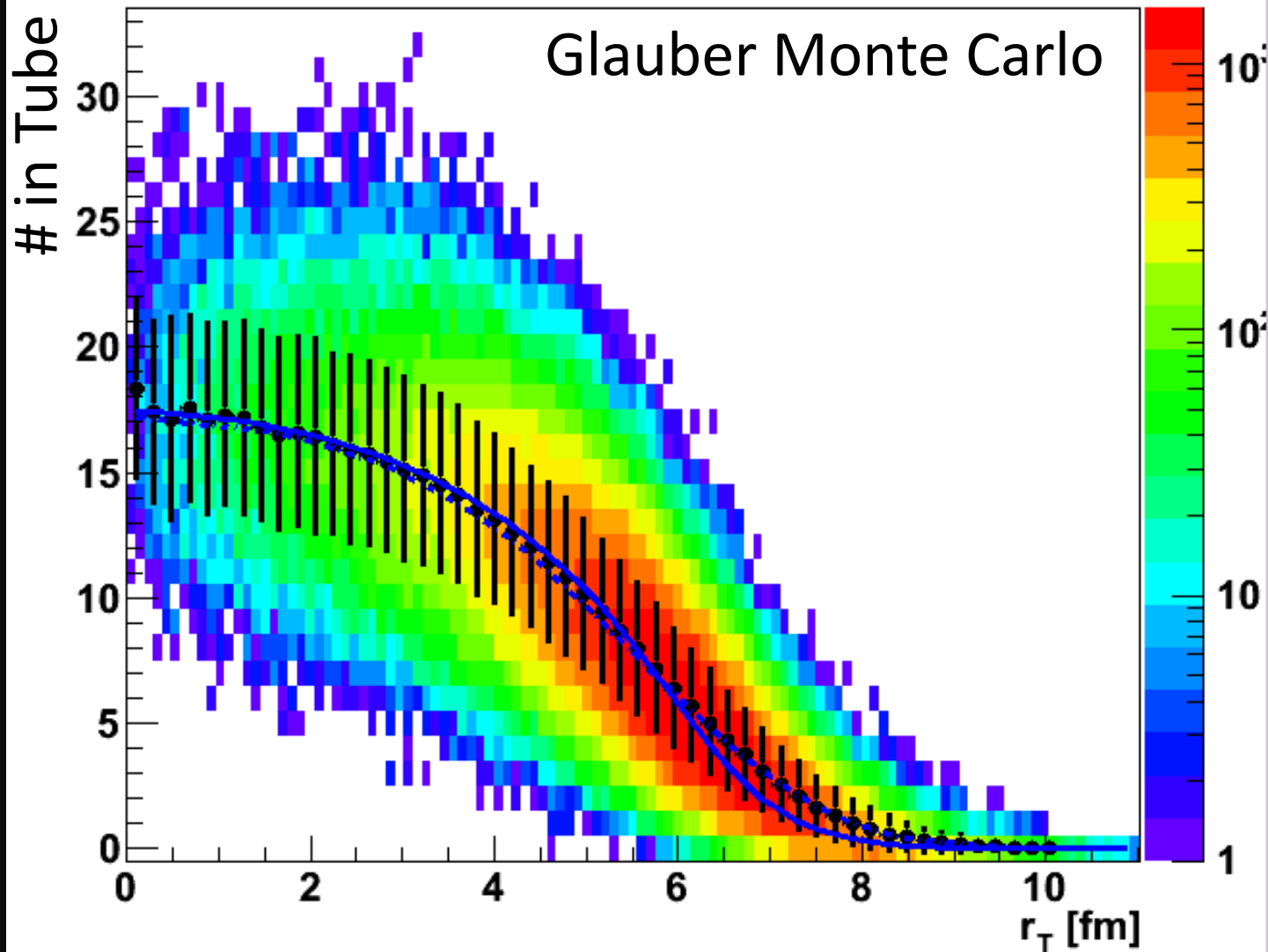


The nucleus is quite lumpy.  
Longitudinal thickness varies event-to-event.



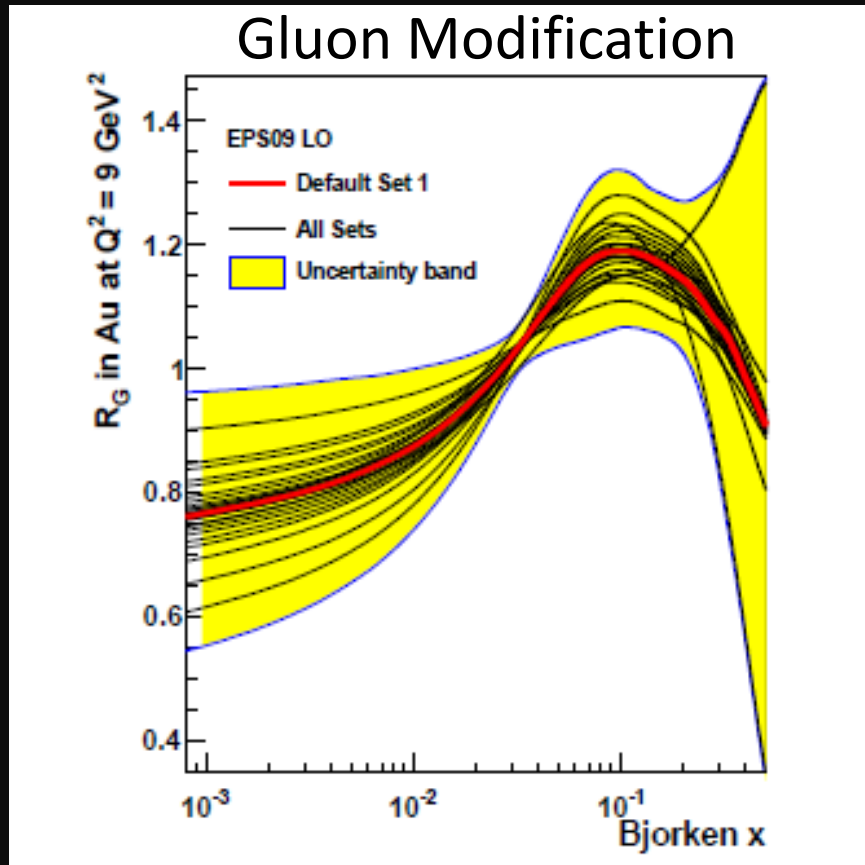
Thus, calculate the thickness  $\Lambda(r_T)$  for each sampled charm pair production point by counting the number of nucleons in a tube ( $N_{\text{tube}}$ ).





At the center of the nucleus,  $\langle N_{\text{tube}} \rangle \approx 18$  and the rms  $\approx 5$ .  
*Never believe any calculation with a smooth nucleus.*

# Nuclear Modified Parton Distribution Functions



EPS09 provides constrained nPDFs and with quantified uncertainties.

However, there is no  $\Lambda(r_T)$  dependence (only average modification)

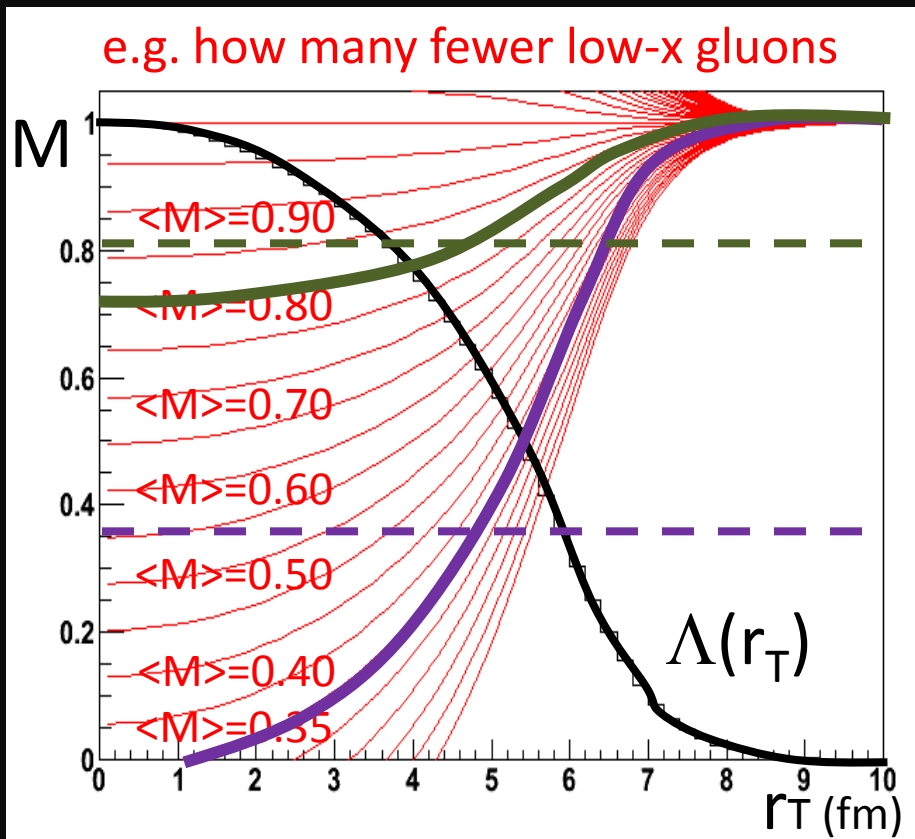
Many calculations in the literature assume that the modification is linear in the thickness

What does that look like?



# Geometric Mapping

Modification (M) linear in  $\Lambda(r_T)$  for different average  $\langle M \rangle$



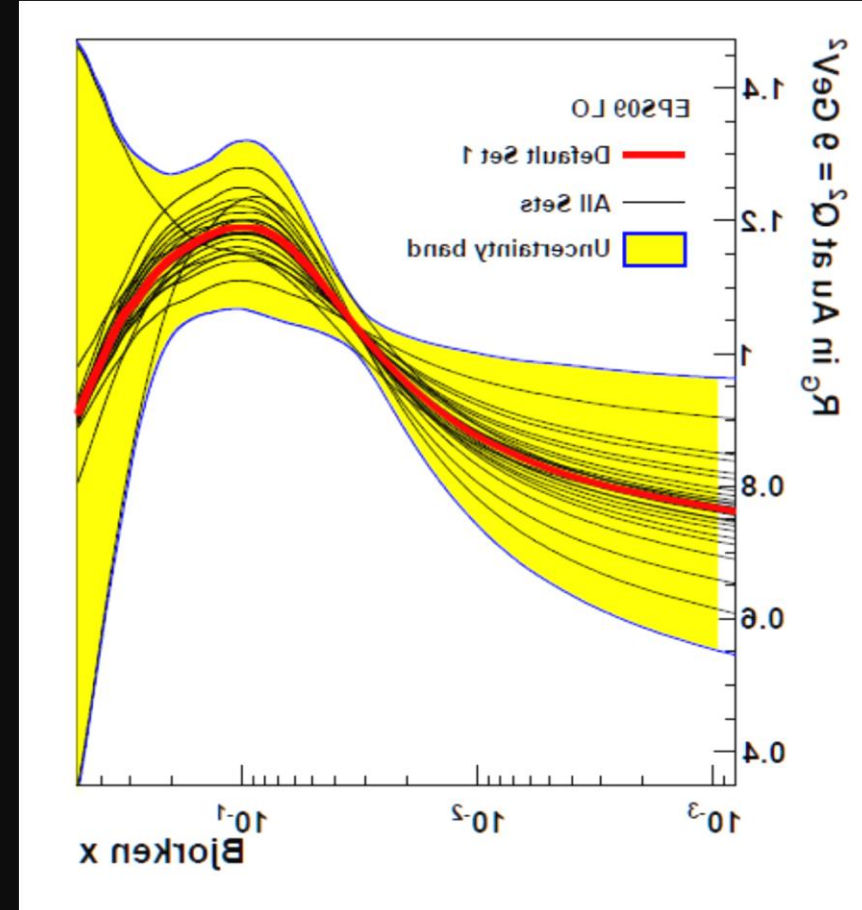
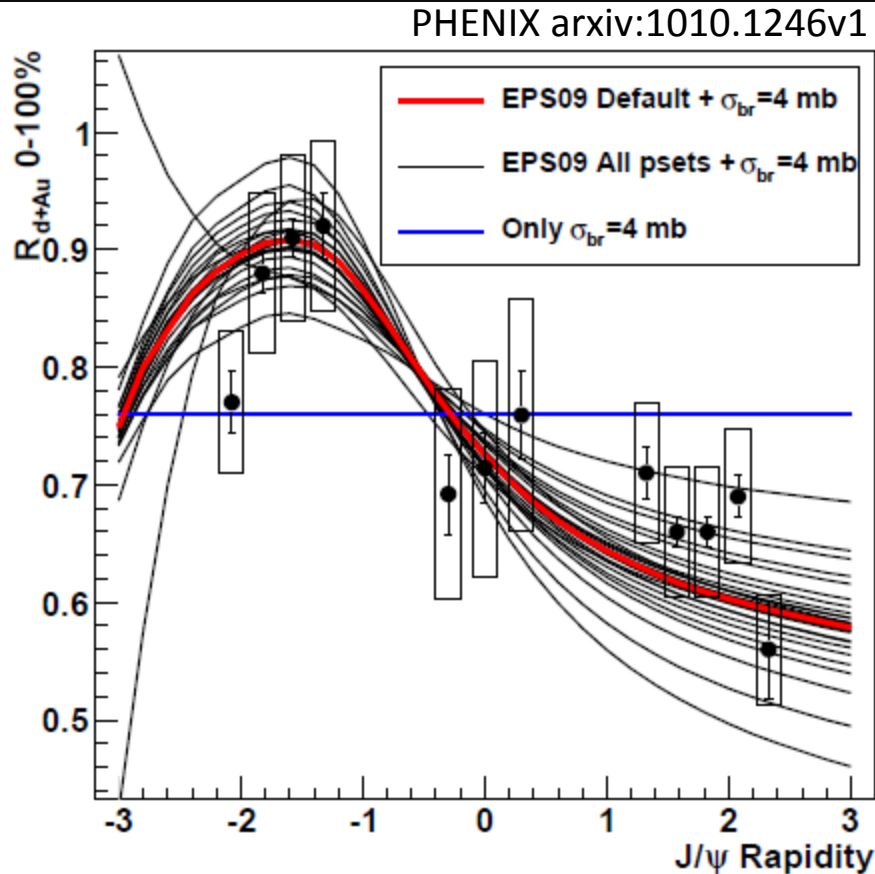
$$\text{Linear : } M(r_T) = 1.0 - a\Lambda(r_T)$$

For  $\langle M \rangle \approx 0.35$ , one is forced to have a negative value at small  $r_T$ .

This is unphysical.

Also try other forms...

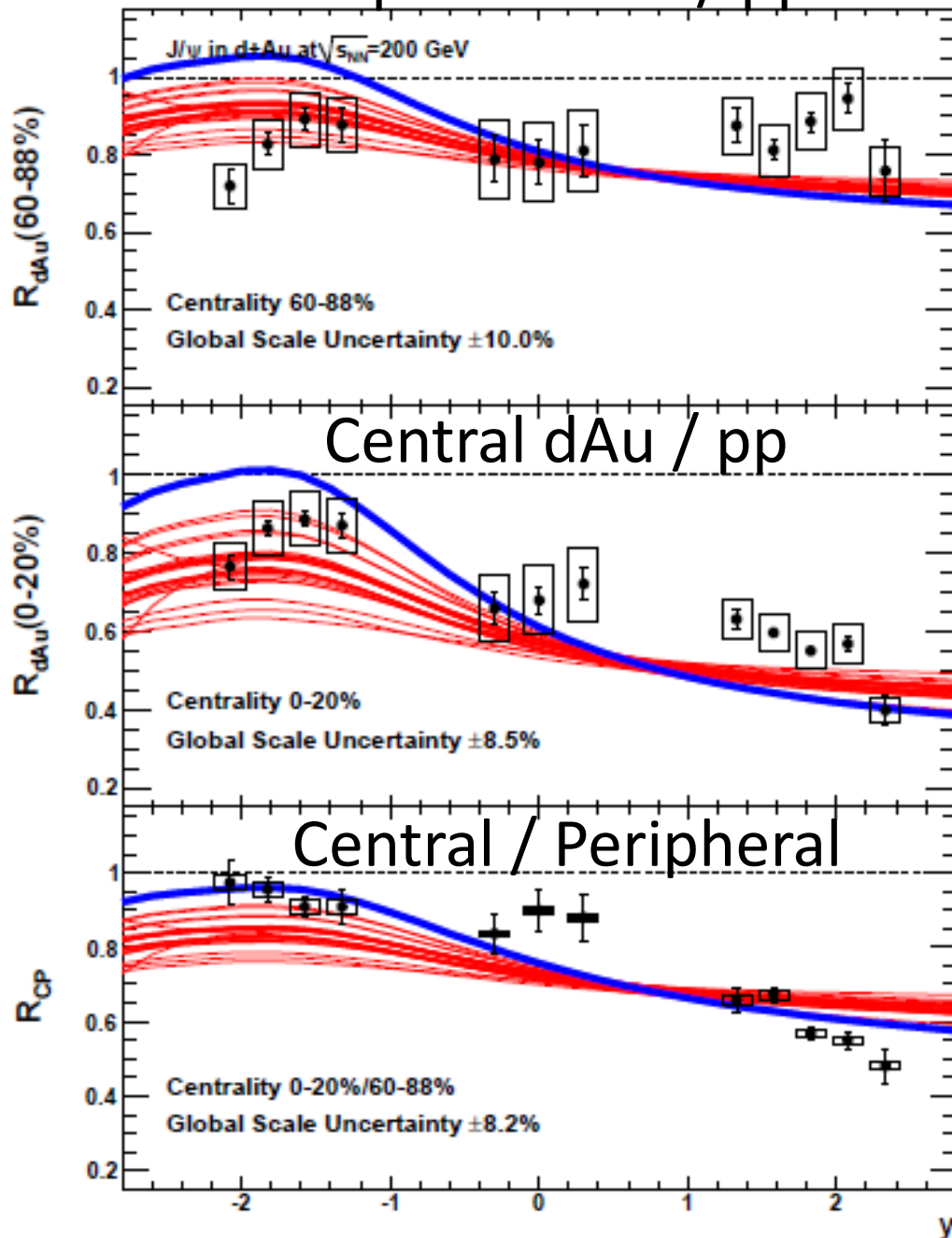
# J/ψ Data Averaged Over All Geometries



We test all EPS09 nPDF variations and find reasonable agreement within uncertainties by including  $\sigma_{br} = 4 \text{ mb}$ .

Rapidity dependence comes entirely from nPDF.

# Peripheral dAu / pp



For each EPS09 nPDF set,  
fit the data including  
statistical and systematic  
uncertainties to find the  
best  $\sigma_{br}$ .

Best fit EPS09 pset = 17  
and  $\sigma_{br} = 3.2$  mb.  
However, the fit is very  
poor (pvalue  $< 0.0001$ ).

Data cannot be  
described by any  
combination of EPS09  
nPDF and  $\sigma_{br}$ !

# Geometry Test

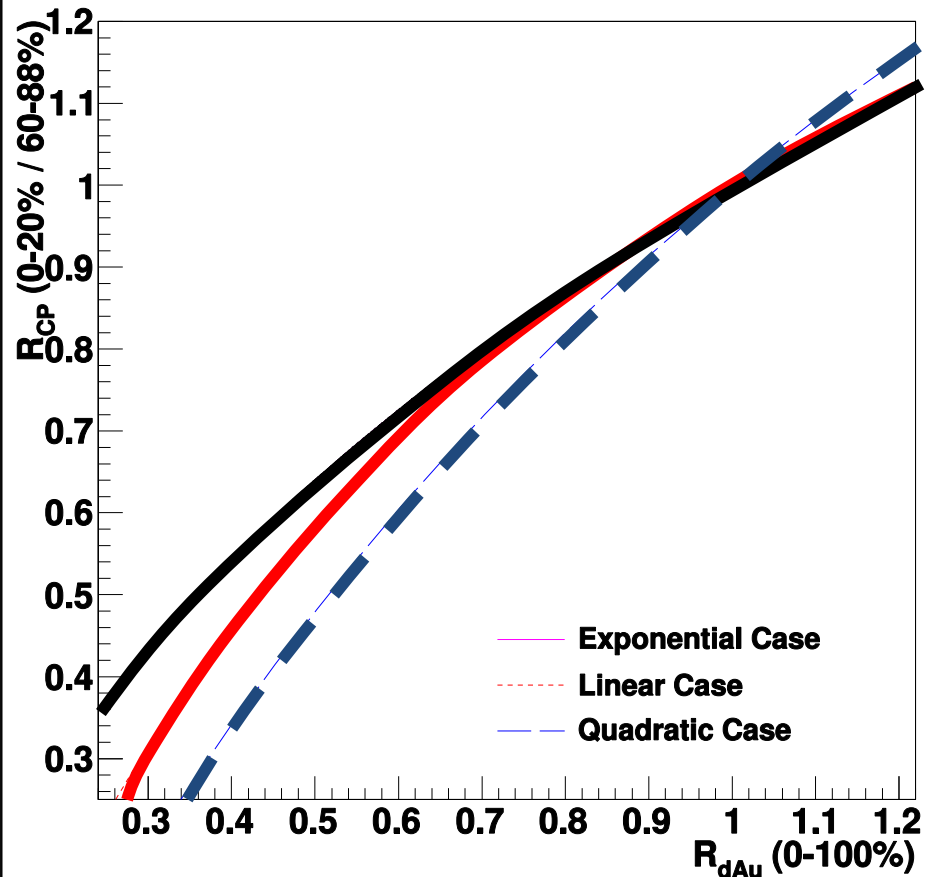
Linear :  $M(r_T) = 1.0 - a\Lambda(r_T)$

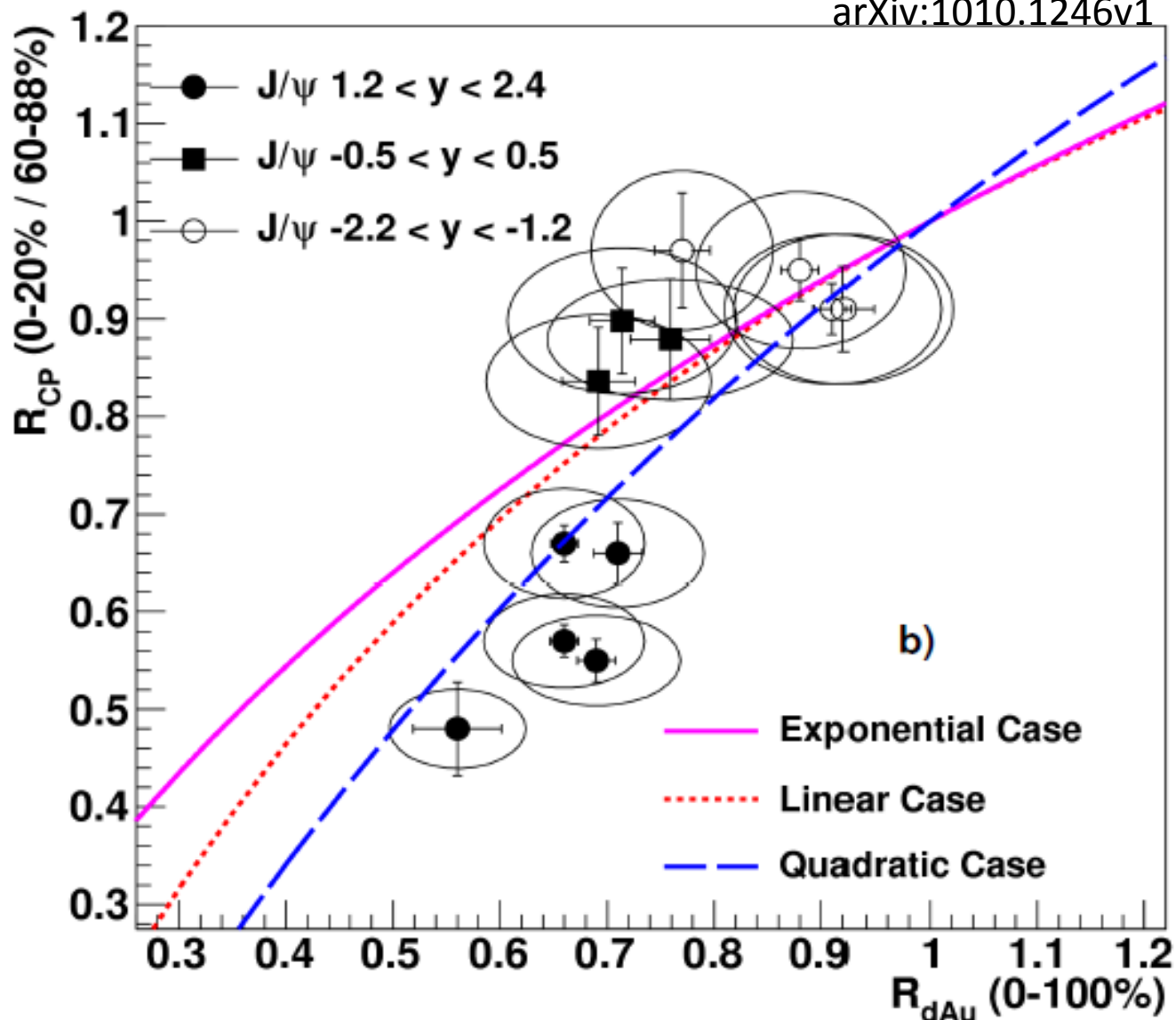
Quadratic :  $M(r_T) = 1.0 - a\Lambda(r_T)^2$

Exponential :  $M(r_T) = e^{-a\Lambda(r_T)}$

For each value of  $\underline{a}$  and the distribution for  $\Lambda(r_T)$  one can calculate  $R_{dAu}$  for every centrality case.

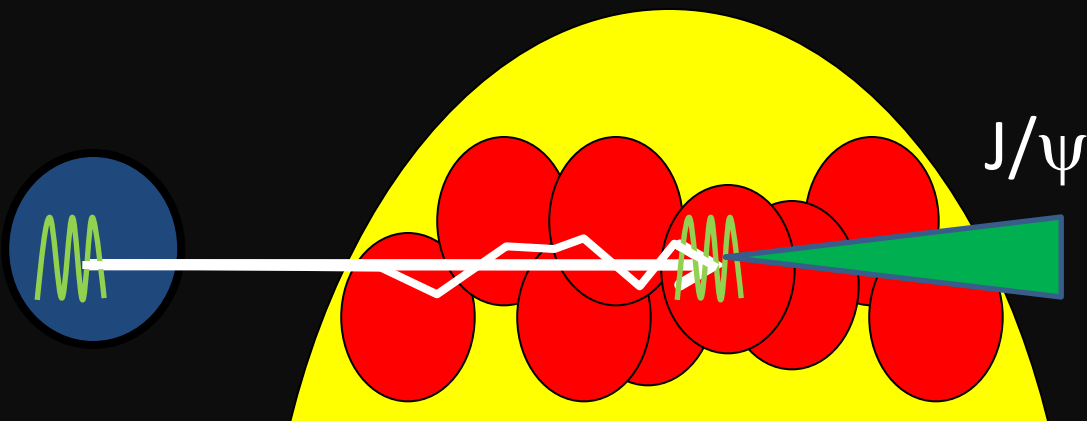
For the Linear Case, for any values of  $\underline{a}$ , there is a unique relationship between the average suppression ( $R_{dAu}$ ) and the steepness versus centrality ( $R_{CP}$ ).



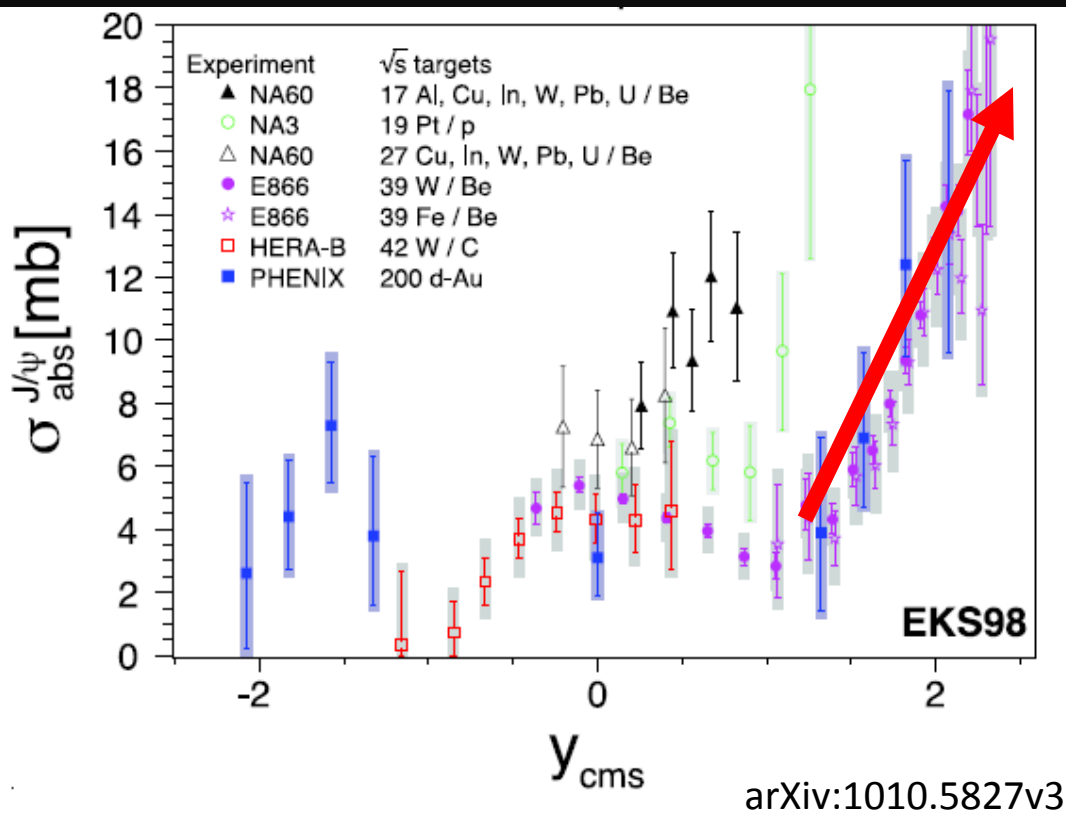


Data driven proof that non-linear (in the thickness) physics effects must contribute to  $J/\psi$  modification.

# Initial State Parton Energy Loss



Energy loss of incoming gluon prior to hard scatter creating charm pair



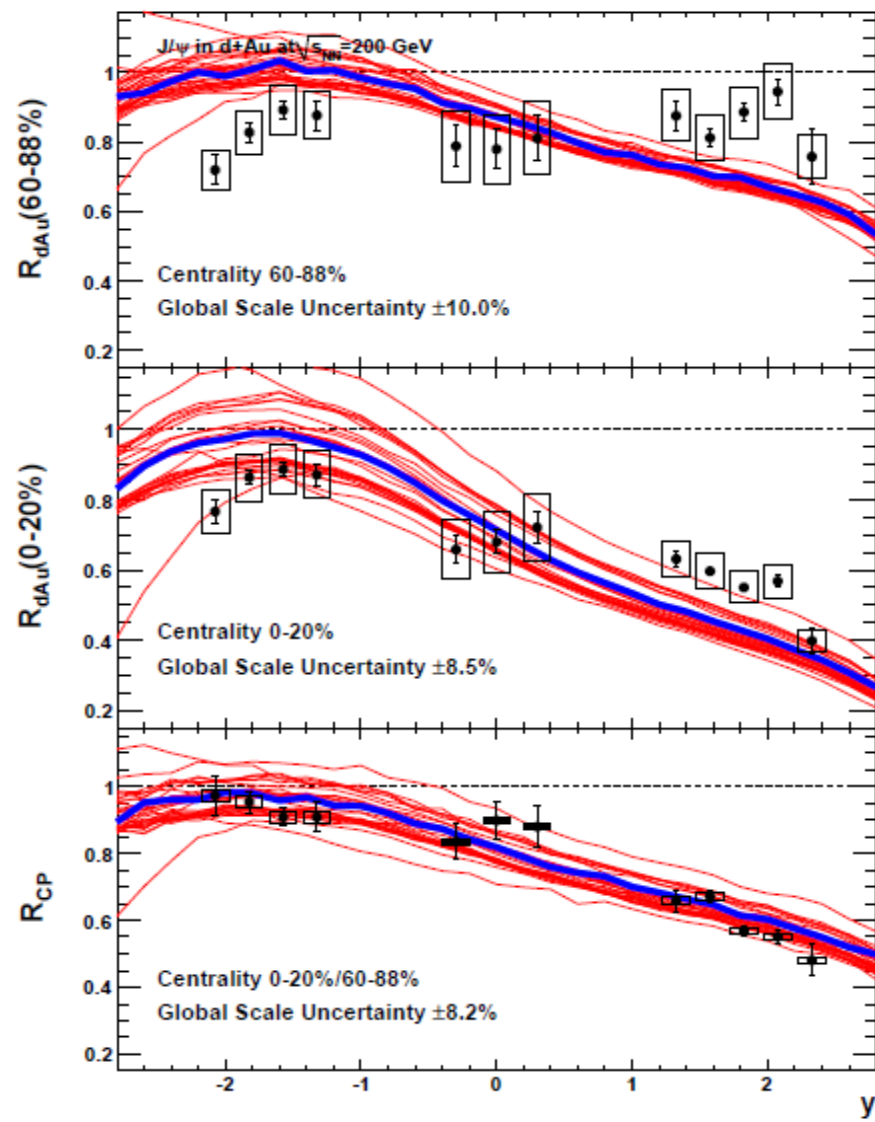
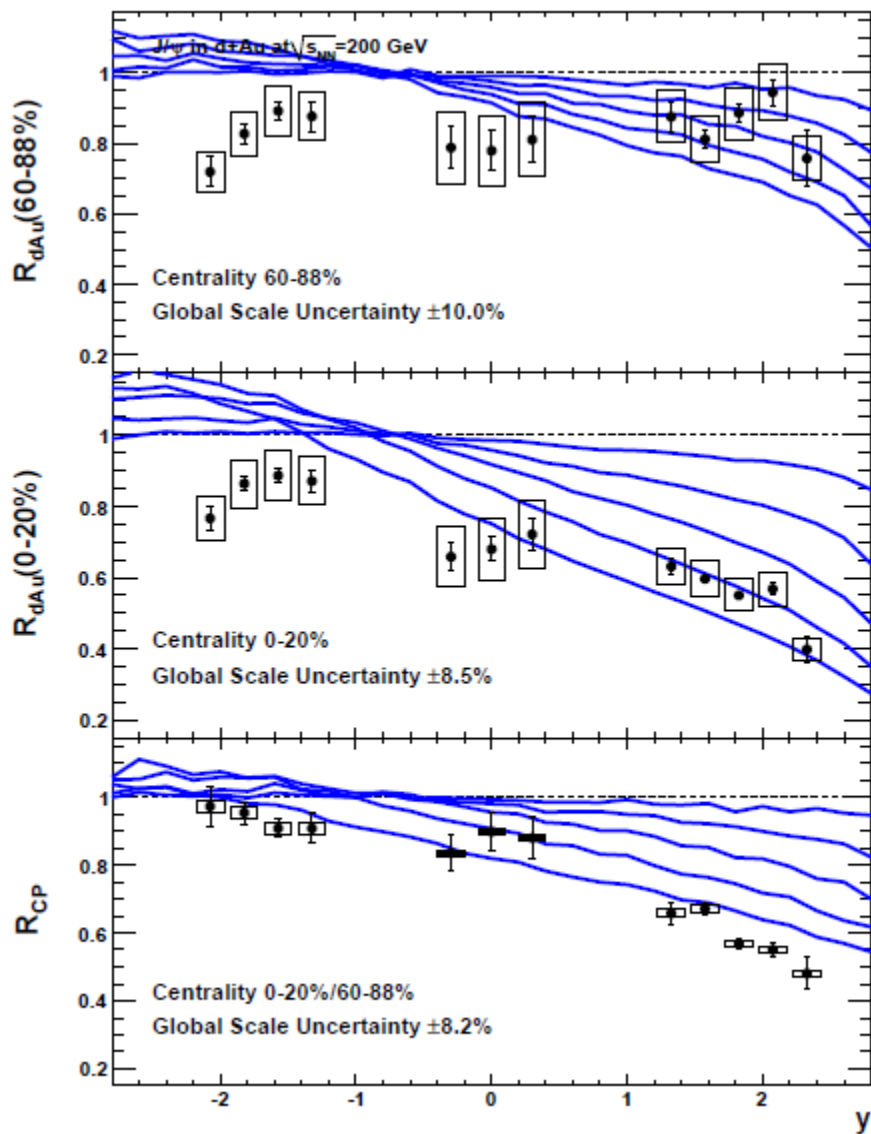
Fit each rapidity point to a separate  $\sigma_{br}$ , yields rapid increase of  $\sigma_{br}$  at forward rapidity.

Initial state energy loss of high x partons?

# Initial parton energy loss in our calculation ( $\Delta E \propto L$ or $L^2$ )

Only e-loss ( $\Delta E \propto L$ )

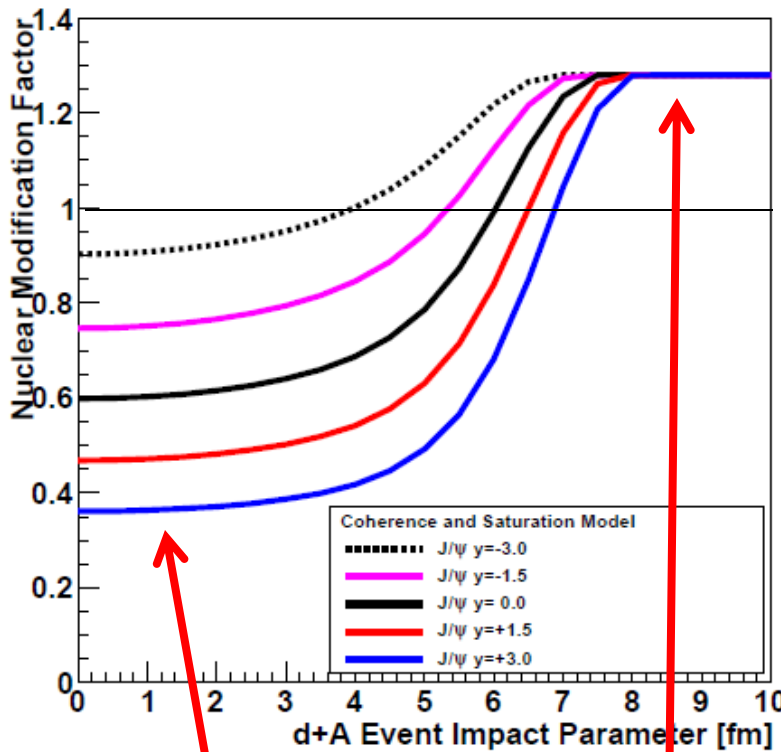
Best fit e-loss, nPDF,  $\sigma_{br}$





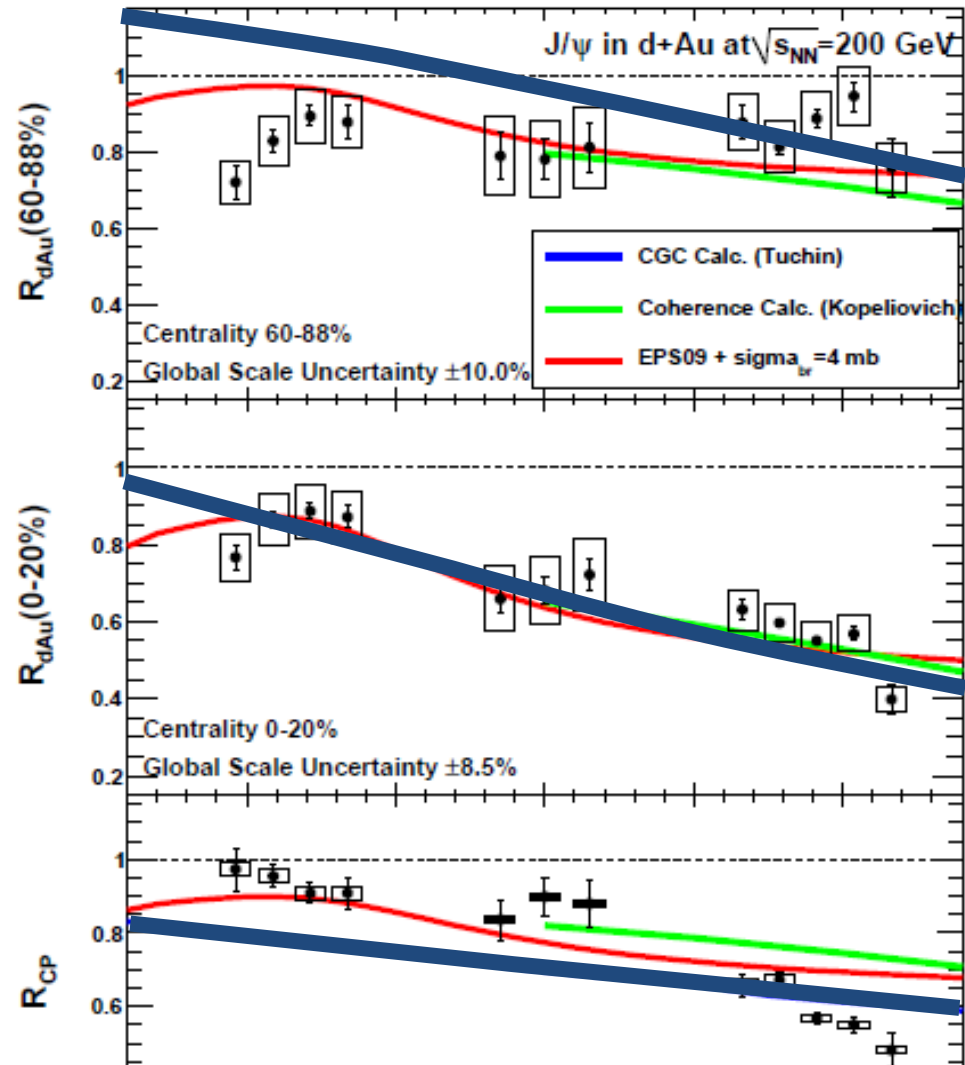
# Coherence Calculations

Color Glass Condensate  
K. Tuchin



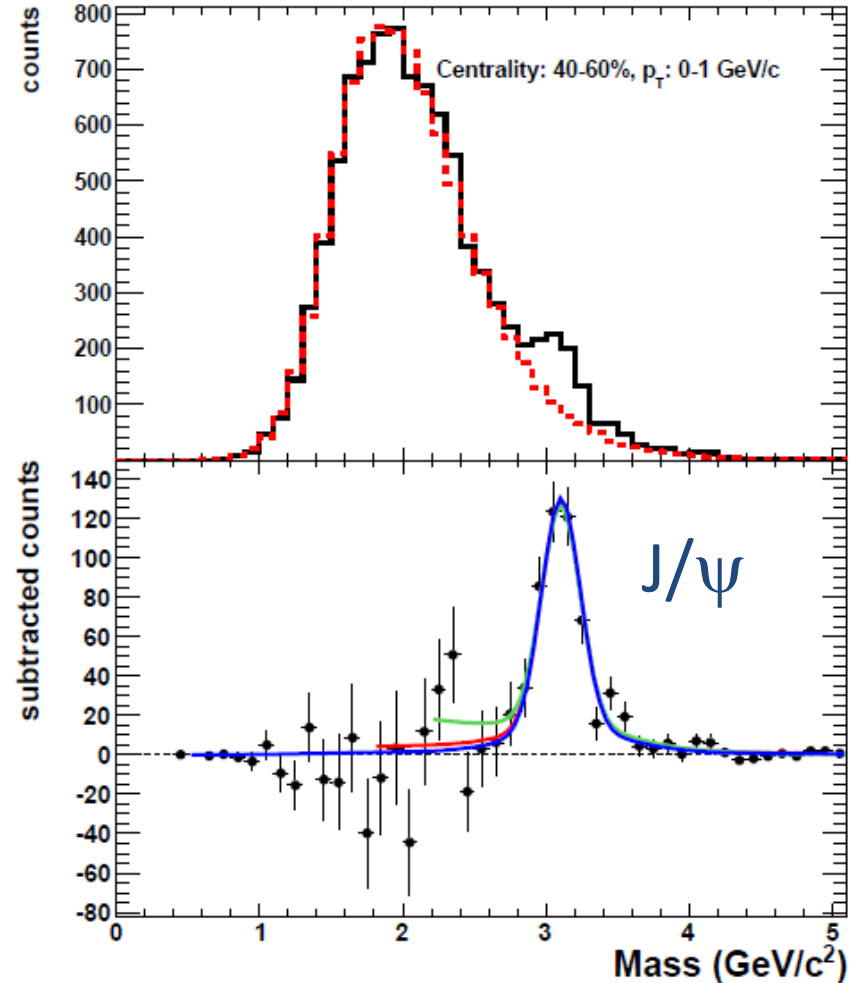
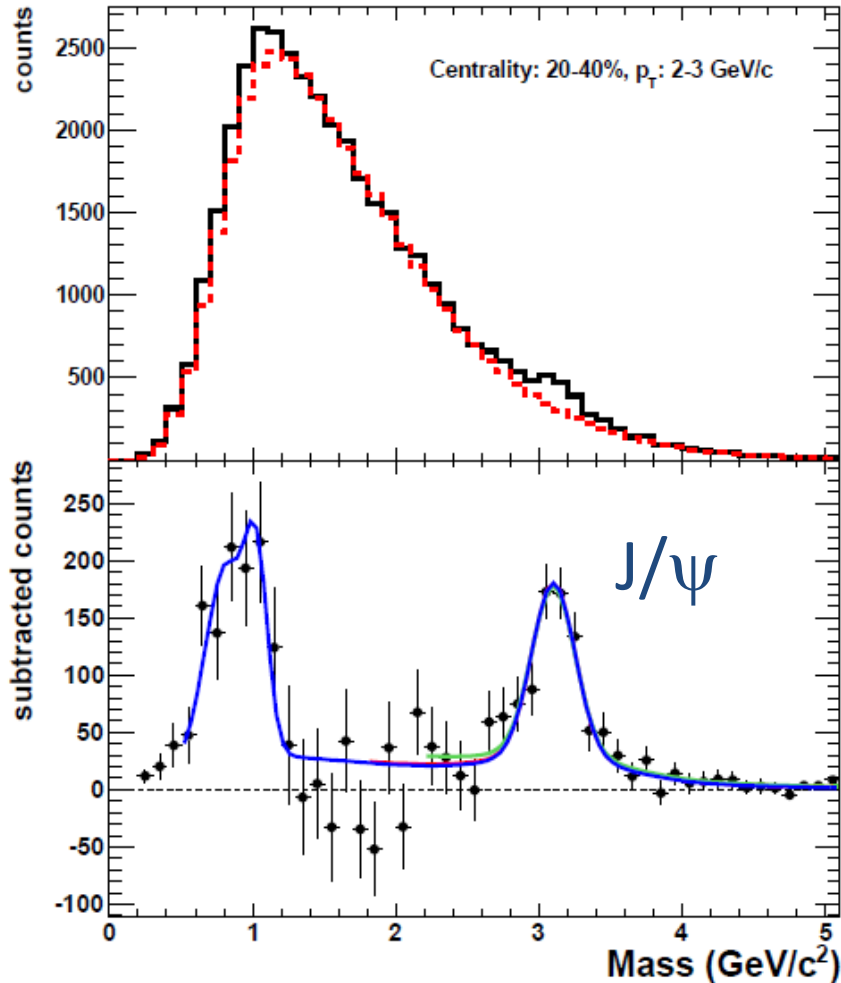
Gluon Saturation

Double gluon exchange



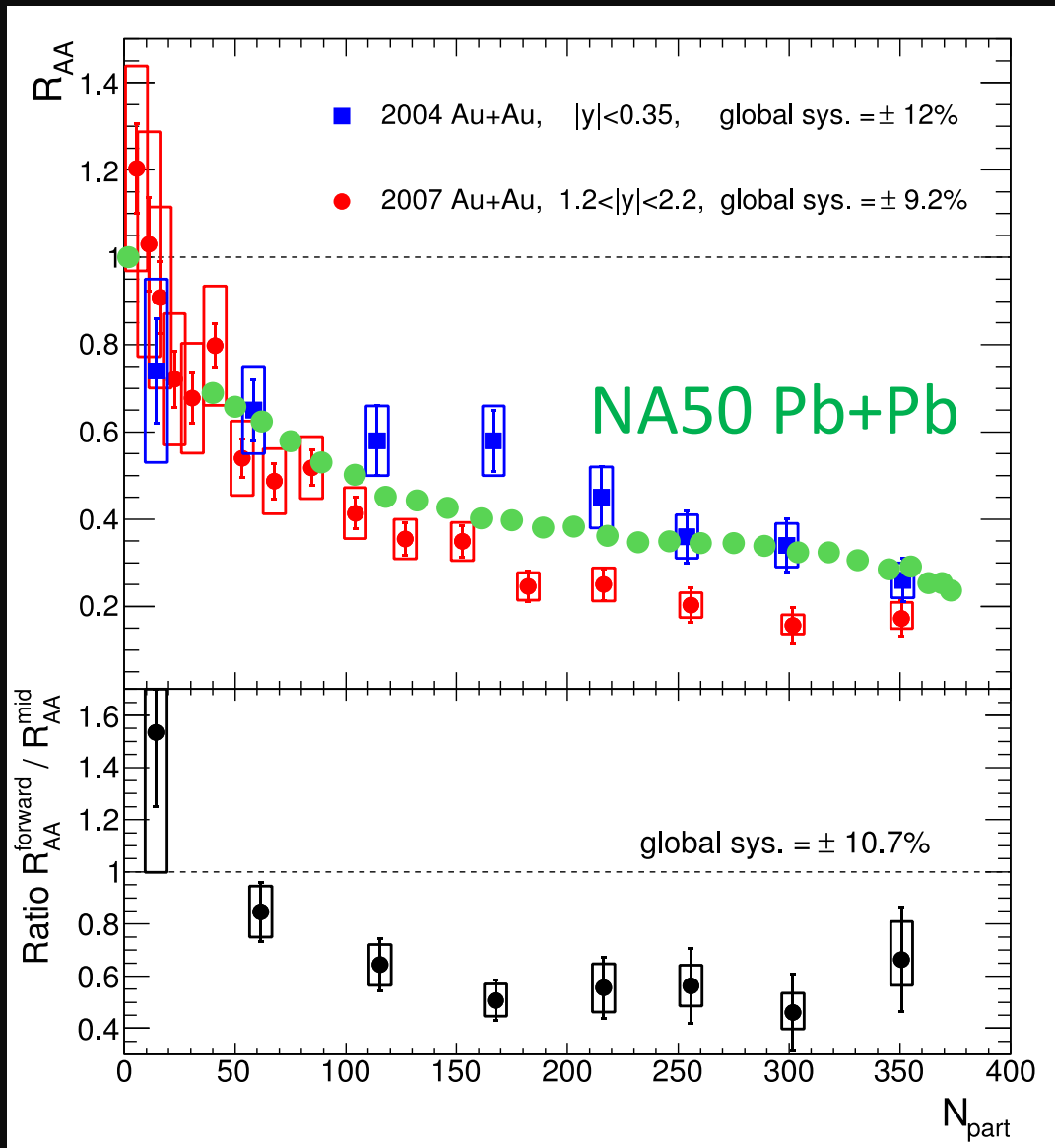
But peripheral makes no physics sense.

# Gold-Gold Data (arXiv:1103.6269 and Wysocki Ph.D.)



Challenging analysis with low signal to background.  
Over 3 billion Au+Au events recorded for analysis.

# PHENIX Data



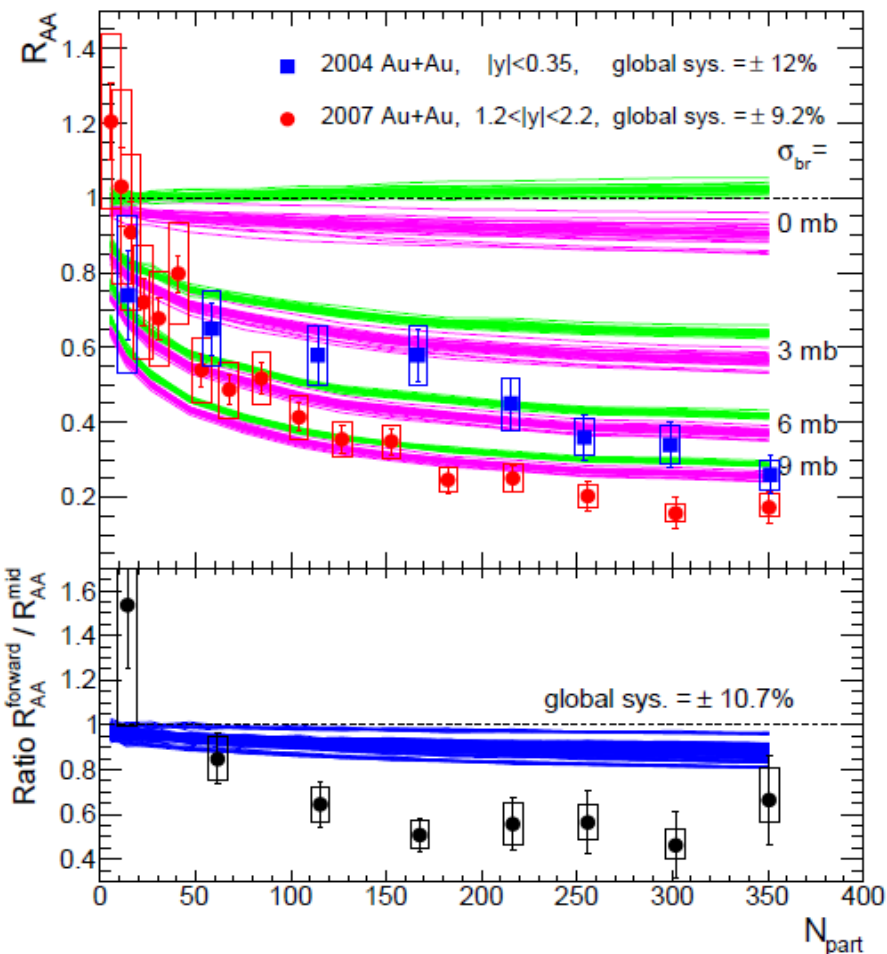
Very similar suppression at CERN-SPS and RHIC at mid-rapidity.

New data confirms with precision the larger suppression at forward rapidity.

Is this just a cold nuclear matter effect?

# Extrapolate Cold Nuclear Matter Effects

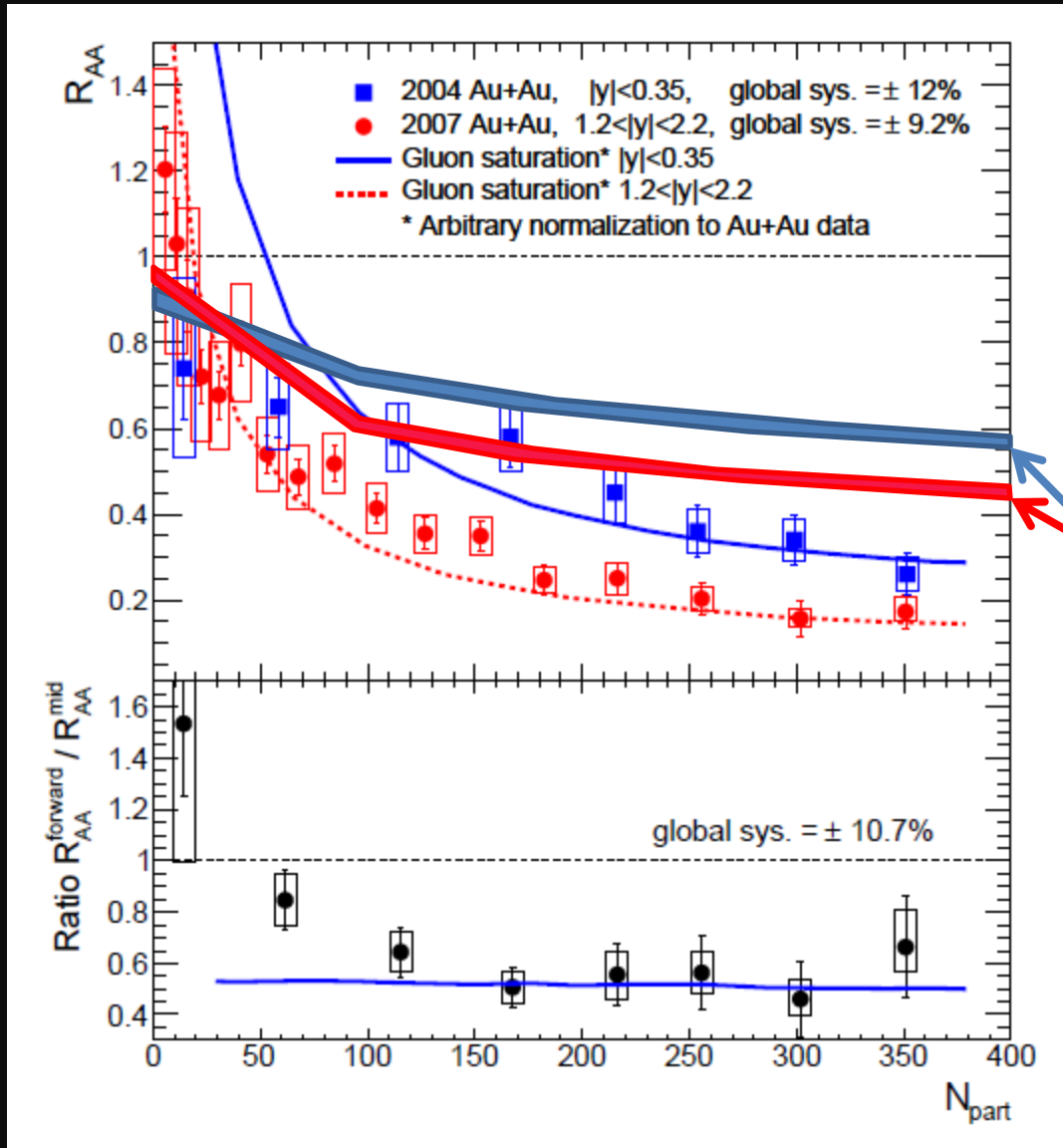
We calculate with all EPS09 parameterizations and  $\sigma_{br}$ .  
Almost no additional forward suppression – cancellation  
of low-x gluon ( $\downarrow$ ) and high-x gluon ( $\uparrow$ ).



Including energy loss  
matched to d-Au data does  
very little to change the  
picture.

**Hot nuclear matters  
(perhaps screening)  
effects are definitely  
important.**

# Gluon Saturation (Color Glass Condensate)



Appears to describe data and forward suppression.

However, arbitrary normalization to data.

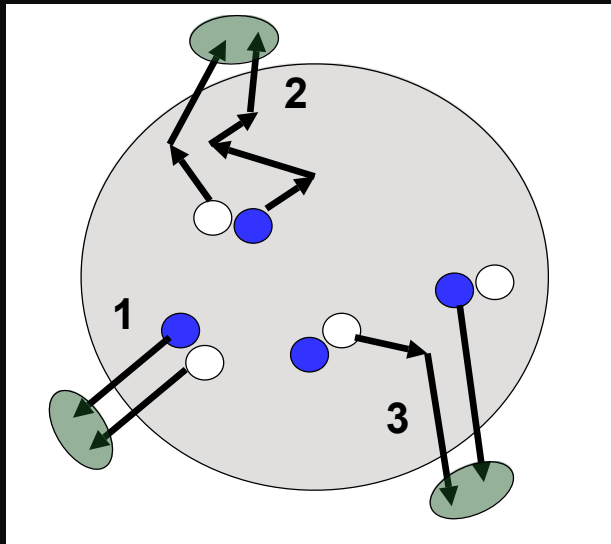
New calculation from M. Nardi et al. indicate much smaller effect.

At Hard Probes 2010, she declared that there are definitely final state effects.

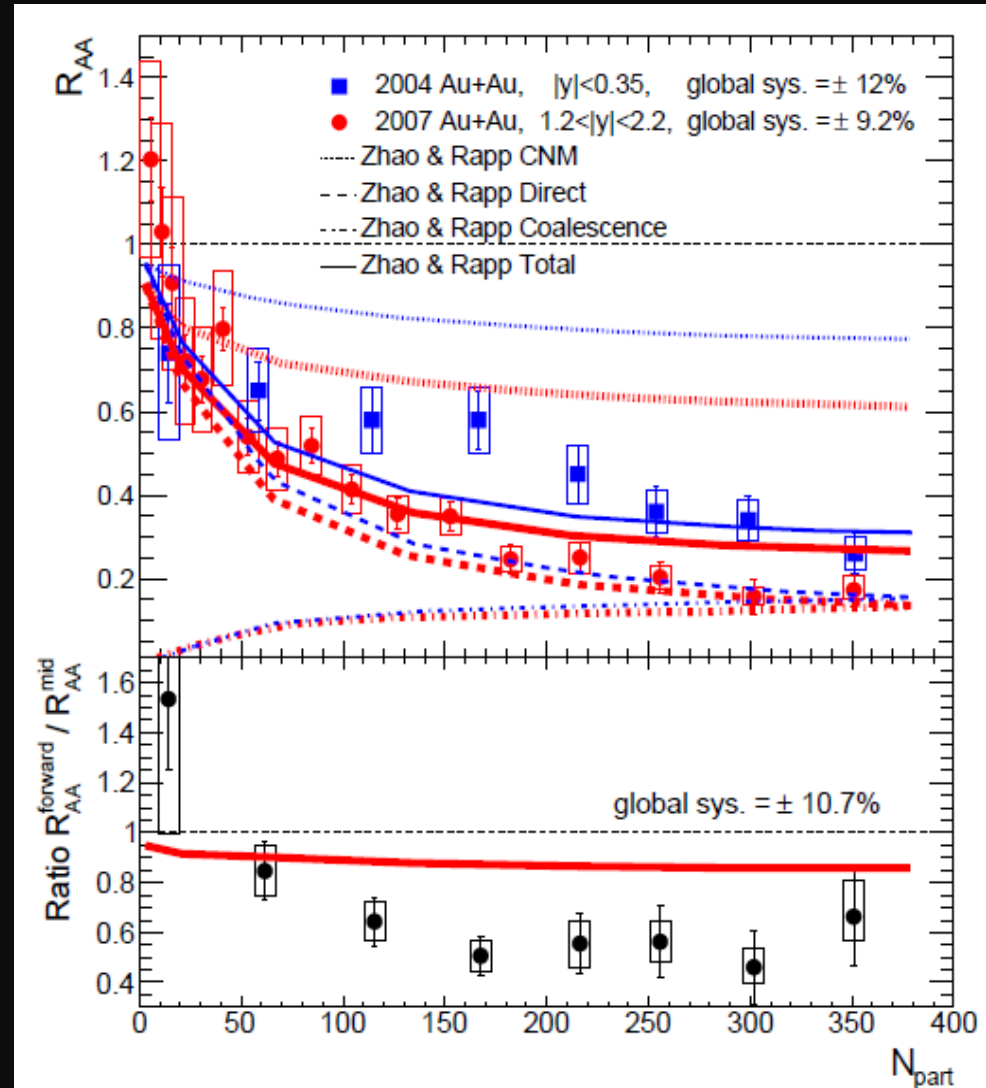
# Charm Recombination

Similar  $R_{AA}$  as lower energies due to larger screening loss compensated by later recombination (?)

Less recombination at forward rapidity since lower  $c\bar{c}$  density.



No effect, recombination of “diagonal pairs” dominates



# Quarkonia Summary

Precision d-Au and Au-Au data sets

Non-linear physics is needed, perhaps e-loss or saturation

No complete description of d-Au data set

Au-Au indicates hot nuclear matter effect, and exact quantification remains elusive

206 citations for PHENIX PRL (2007), but no breakthrough

Near future data from lower energies 39 GeV, 62 GeV  
from PHENIX and LHC data may shed light...

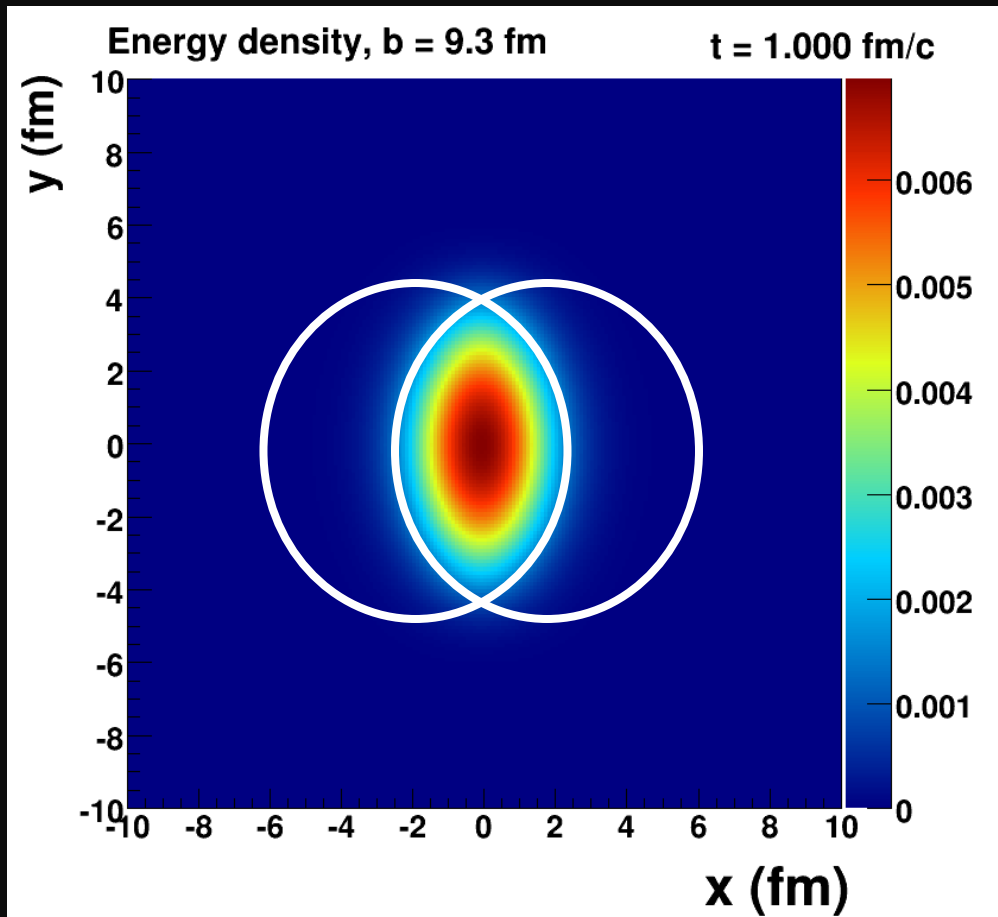
Upgrades to push for even more forward measurements  
including Drell-Yan to disentangle effects.



# Harmonic Flow and Constraining QGP Properties

Our mental picture of heavy ion collisions has evolved...

Smooth initial conditions

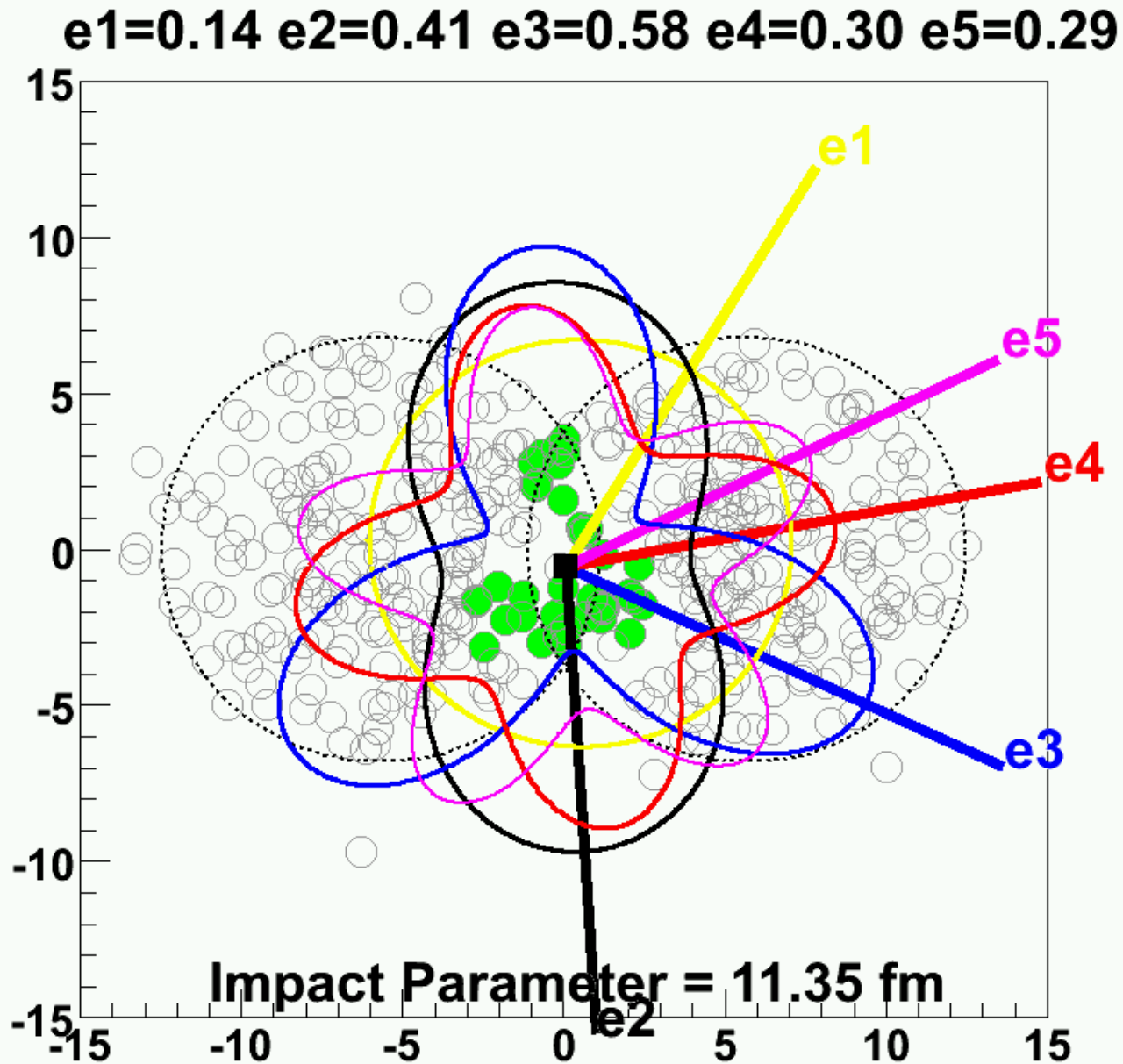


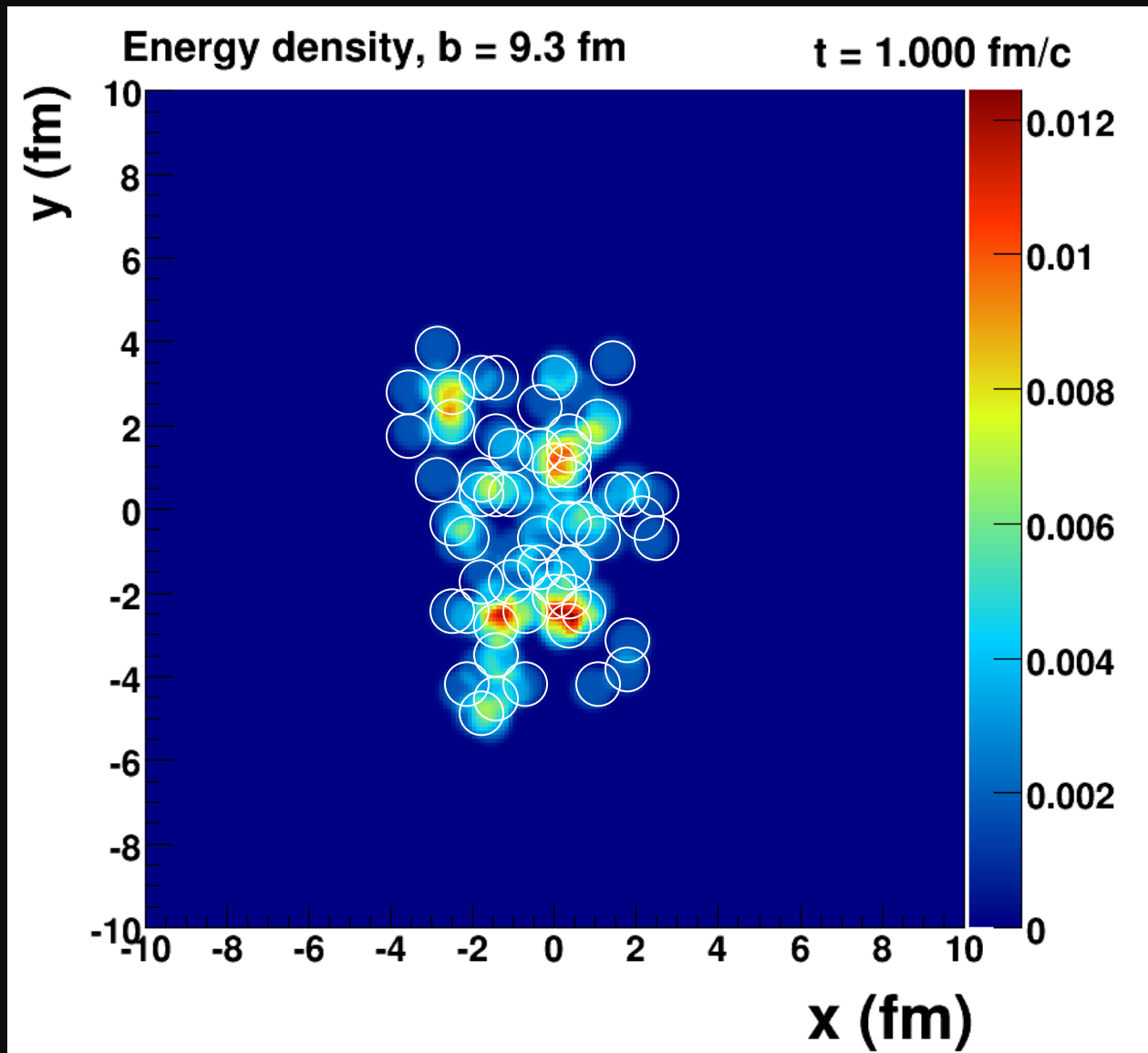
Need to know initial geometric shape.

Initial Eccentricity  $\varepsilon_2$  results in  $\cos(2\phi)$  flow called  $v_2$ .

Measured  $v_2$  is sensitive to key QGP properties like the shear viscosity  $\eta/s$ .

# Reality = Lumpy Initial Conditions

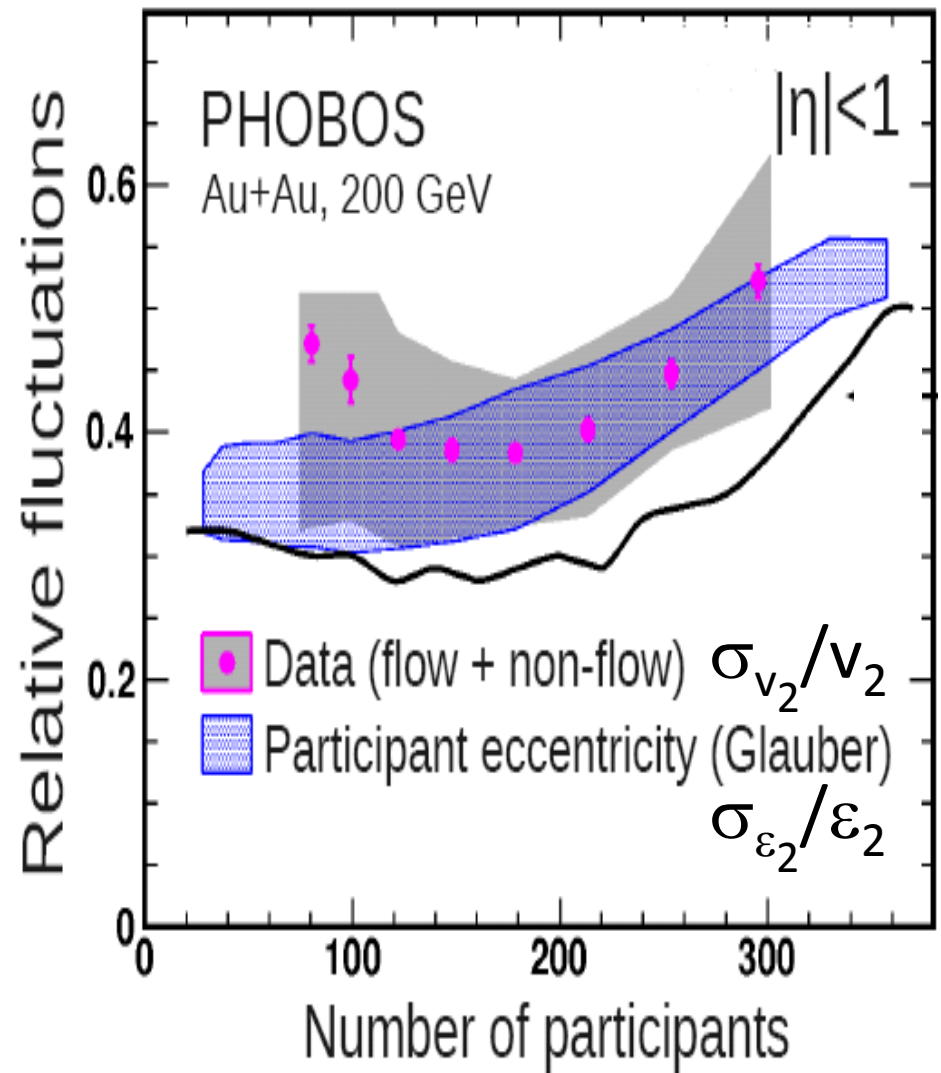
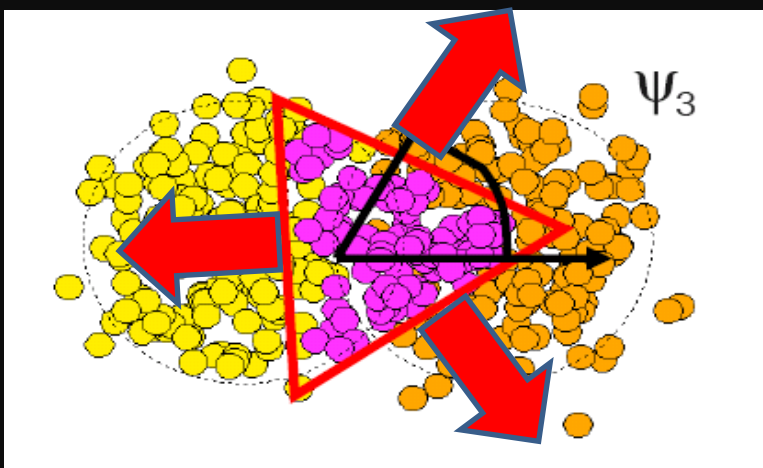
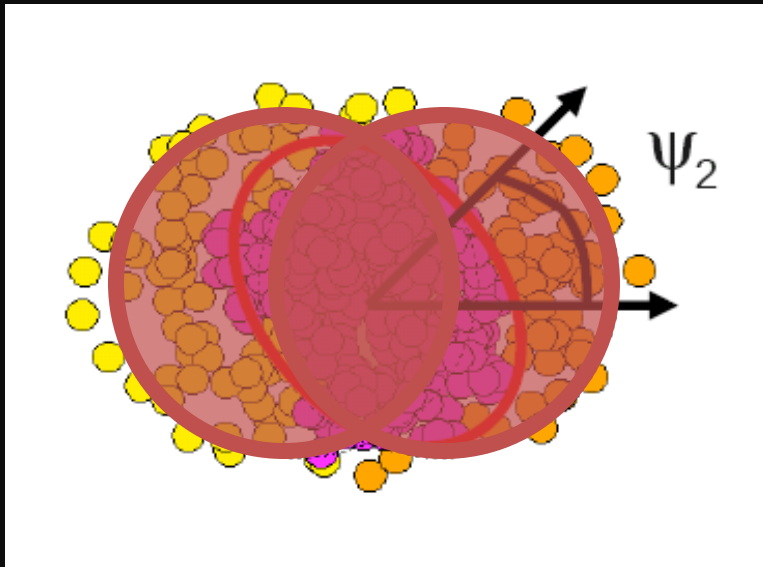




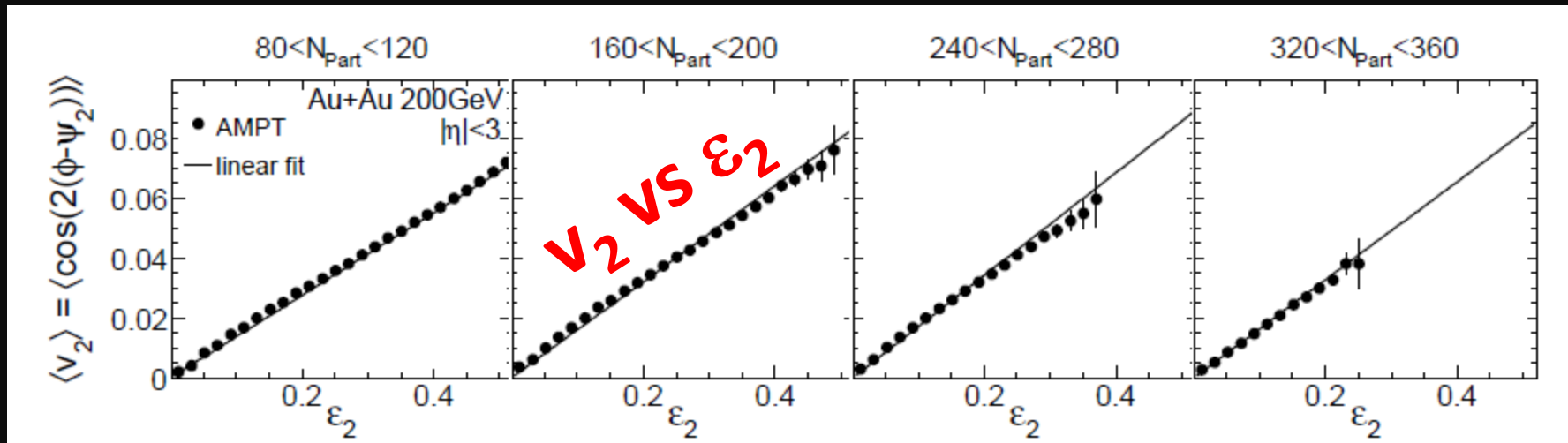
Romatschke=viscous hydrodynamics, McCumber=lumpy conditions + animation

# Fluctuations Dominate

Flow dictated by participating nucleon geometry



# Spatial moments translate into momentum anisotropy moments

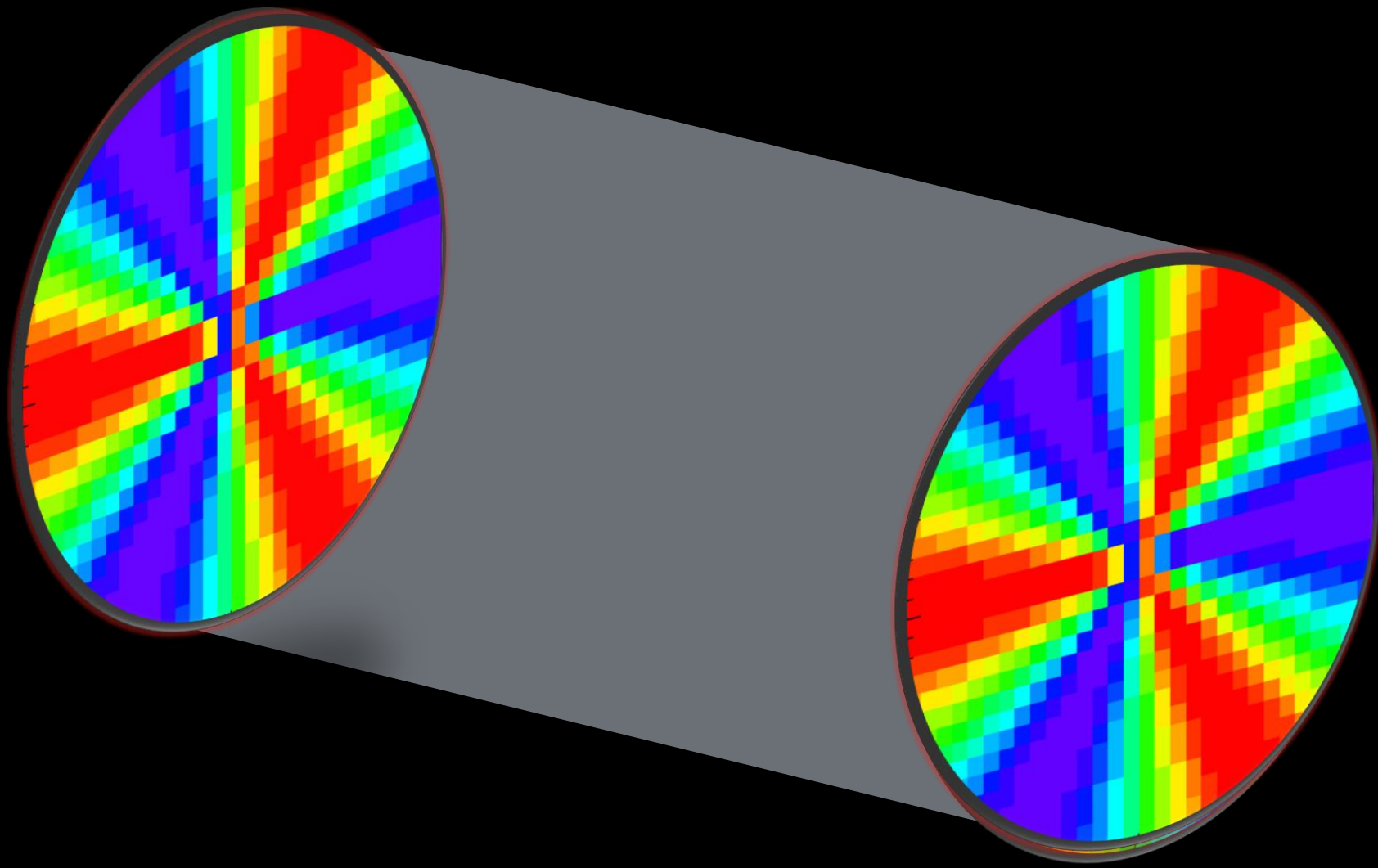


$v_3$  is as irrefutable as  $v_2$  from  $\epsilon_2$ .

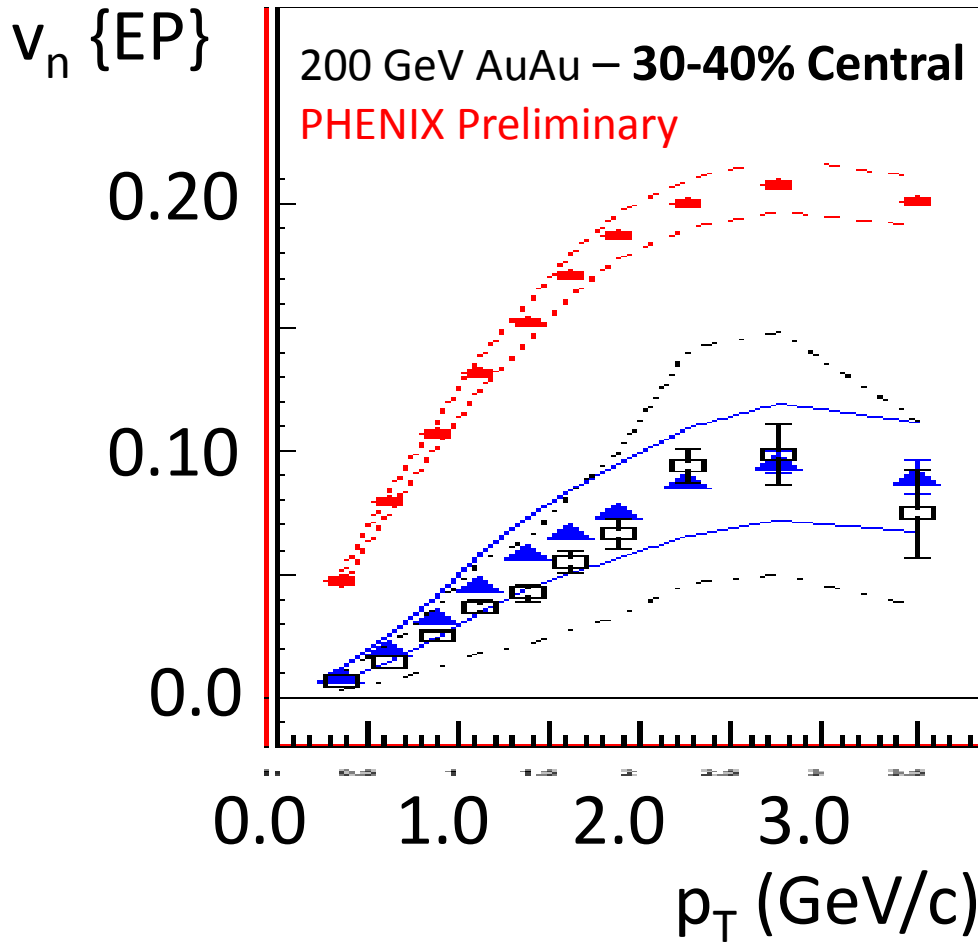
Now we must quantify the effects and implications.

# Measuring $v_3$

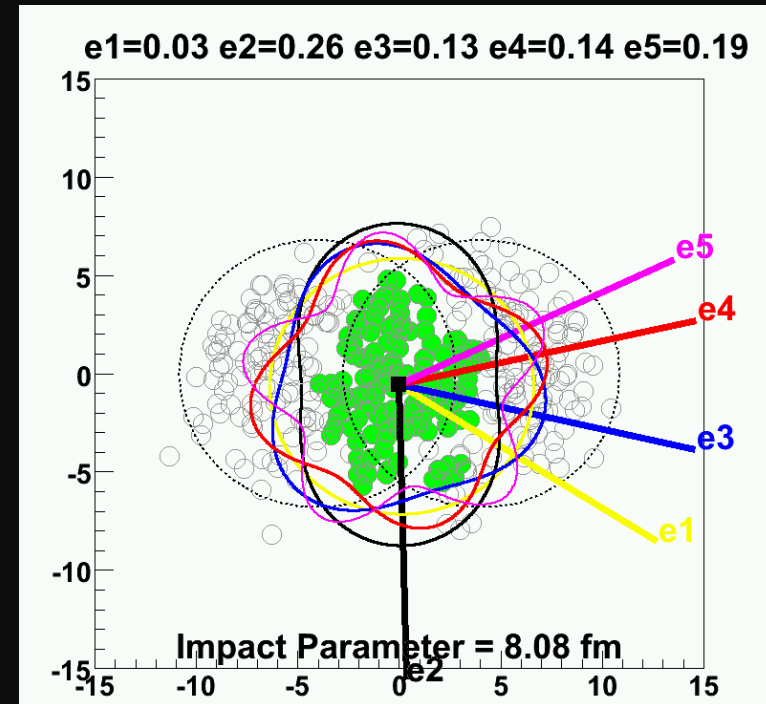
First, can detectors separated by  $\Delta\eta = 2$  or even  $\Delta\eta = 6$  measure event-by-event the 3<sup>rd</sup> order participant plane?



- $v_2 \{\Phi_2 \text{ forw.}\eta\}$
- ▲—  $v_3 \{\Phi_3 \text{ forw.}\eta\}$
- $v_4 \{\Phi_4 \text{ forw.}\eta\}$



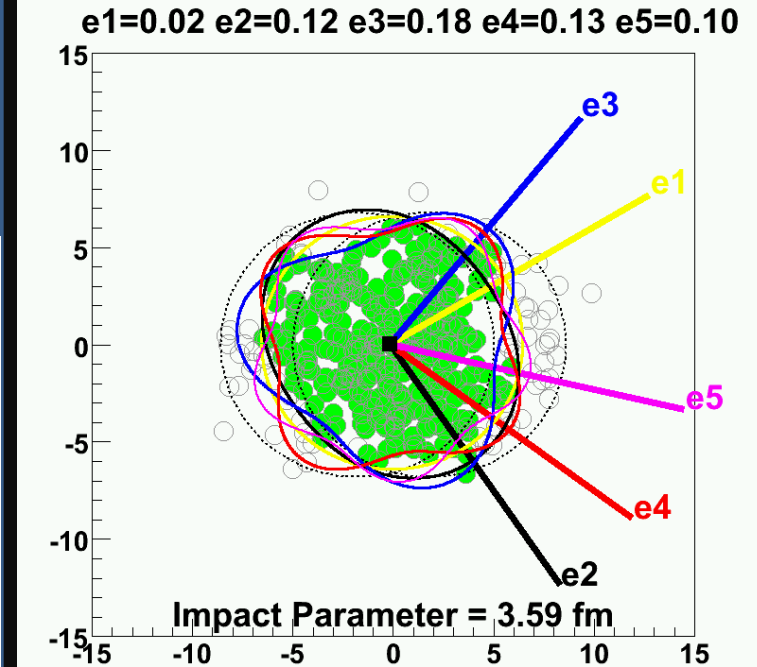
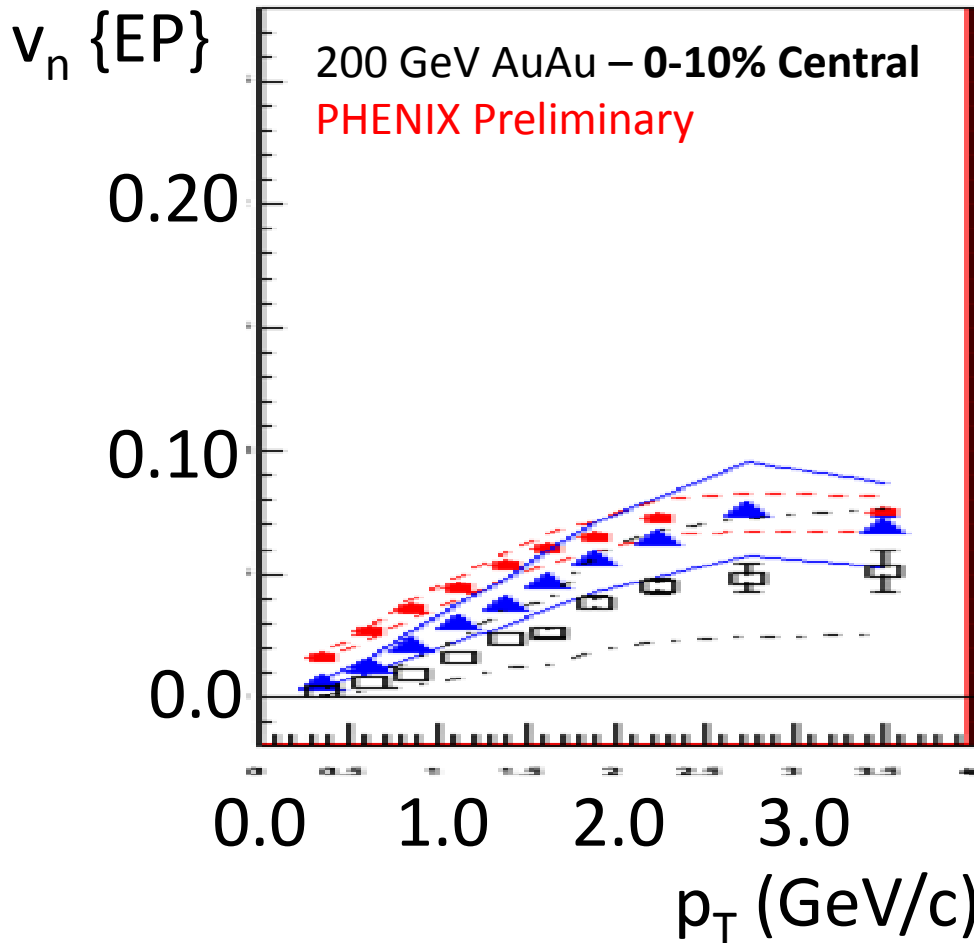
Systematic uncertainties defined by the variations with  $\Phi_n$  from different  $\Delta\eta$  and from different methods.



$$\varepsilon_2 \approx 2 \times \varepsilon_3 \approx 2 \times \varepsilon_4$$



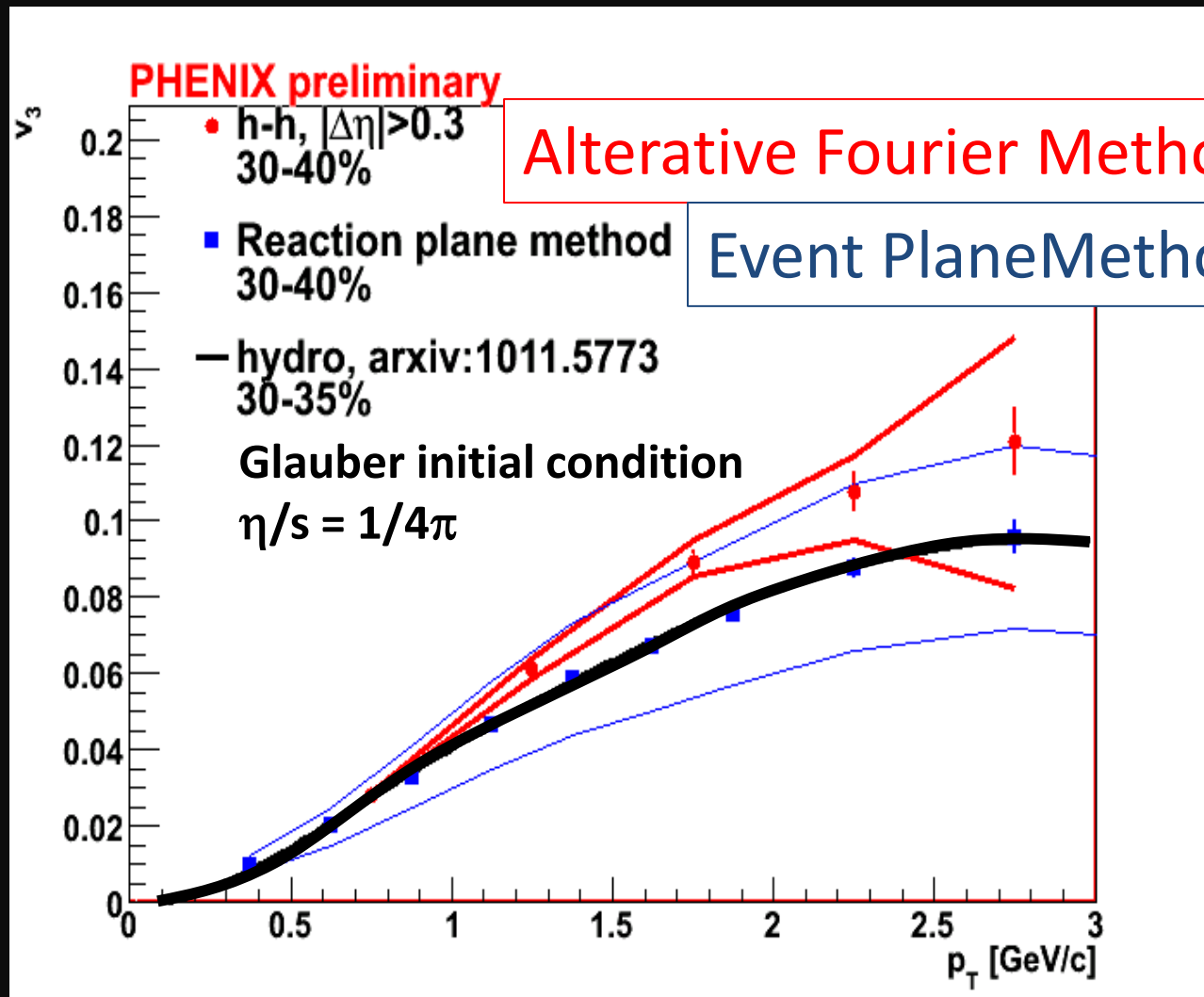
- $v_2 \{\Phi_2 \text{ forw.}\eta\}$
- ▲—  $v_3 \{\Phi_3 \text{ forw.}\eta\}$
- $v_4 \{\Phi_4 \text{ forw.}\eta\}$



$$\varepsilon_2 \approx \varepsilon_3 \approx \varepsilon_4$$

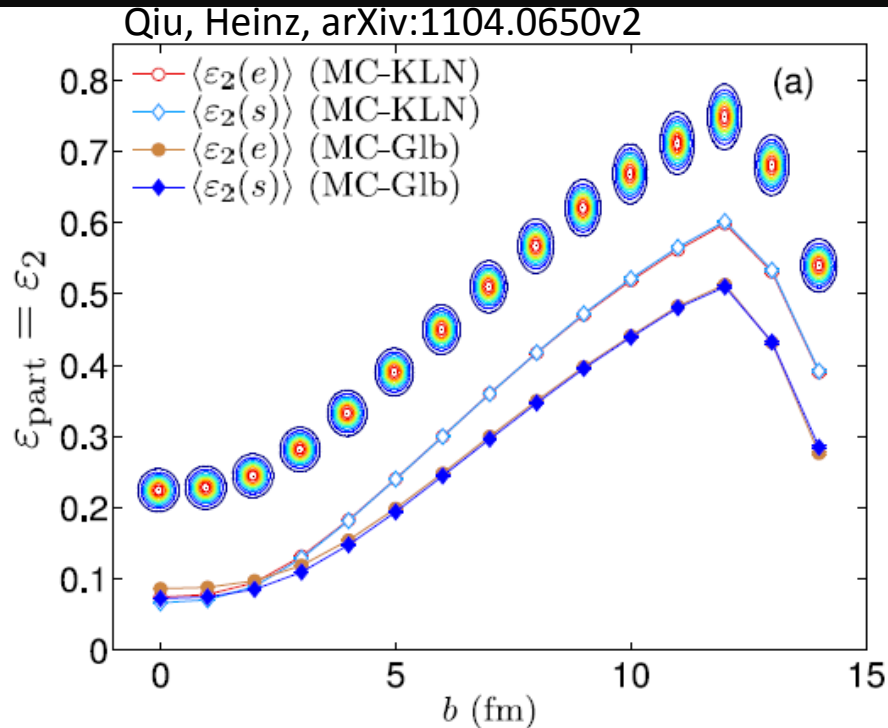
Forces the question of  
how many moments  
are important? 5,6,7<sup>th</sup>  
Like white noise?  
Also 1<sup>st</sup> moment!

# What do we Learn?



Good agreement with viscous hydrodynamics calculations  
(ask Matt Luzum for calculation details)

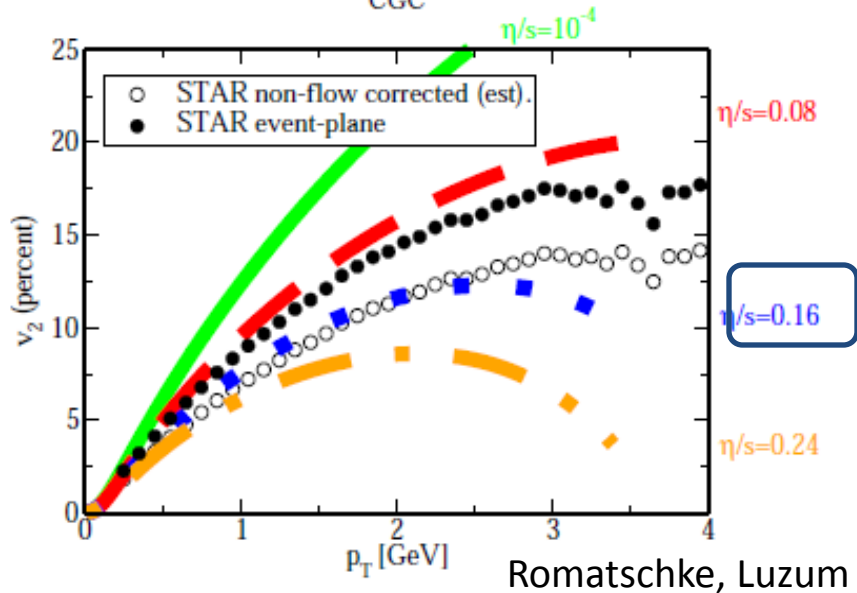
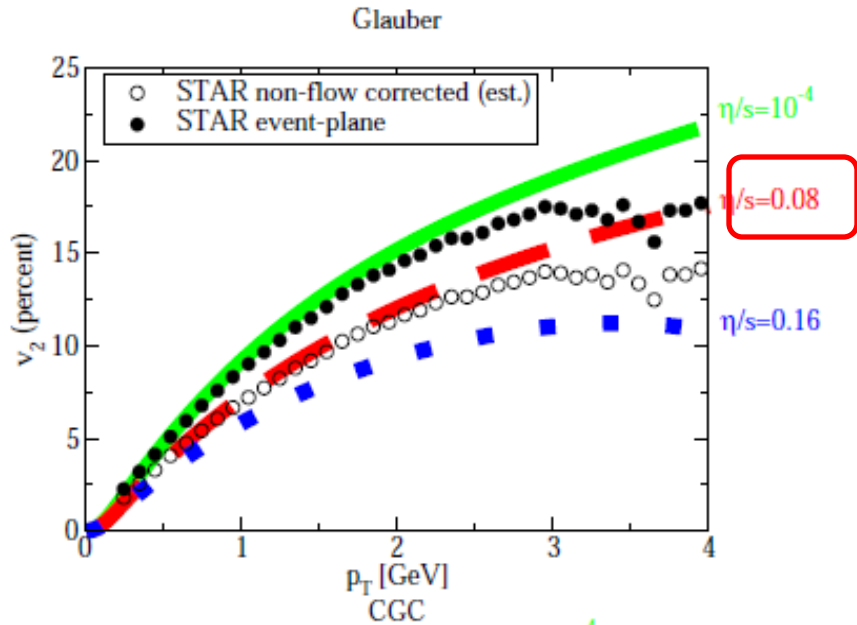
# Key Point



Color Glass Condensate versus Glauber results in different  $\varepsilon_2$  by  $\sim 20\%$ . This is because gluon saturation reweights local entropy density – and thus different initial shape.

**However, for  $\varepsilon_3$ , CGC and Glauber are nearly identical!**

# Why is this so important?



Glauber initial  $\varepsilon_2$  requires  
 $\eta/s \approx 1/4\pi$  to match  $v_2$  data

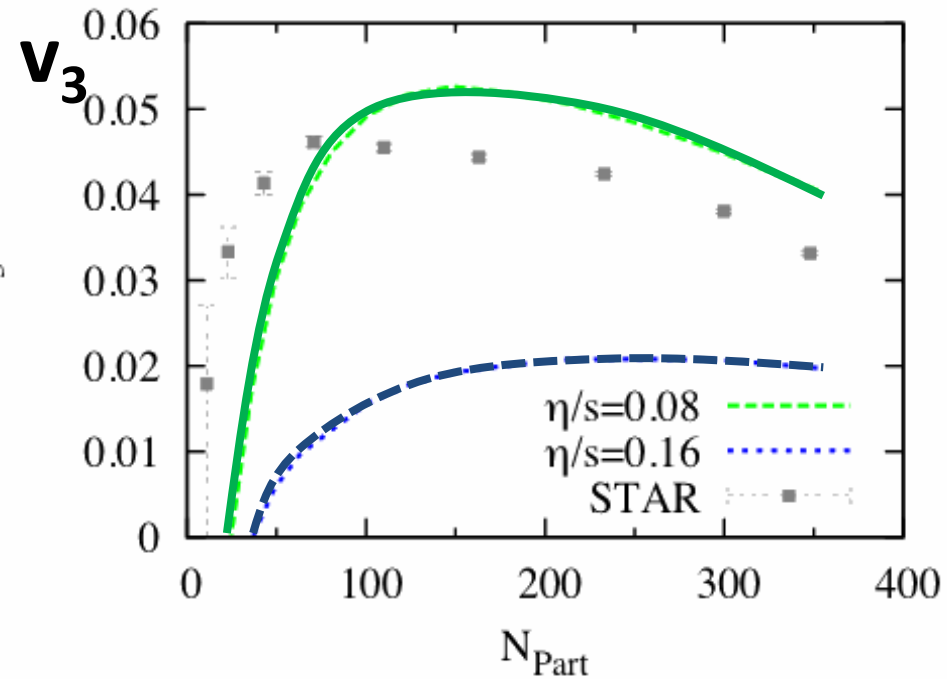
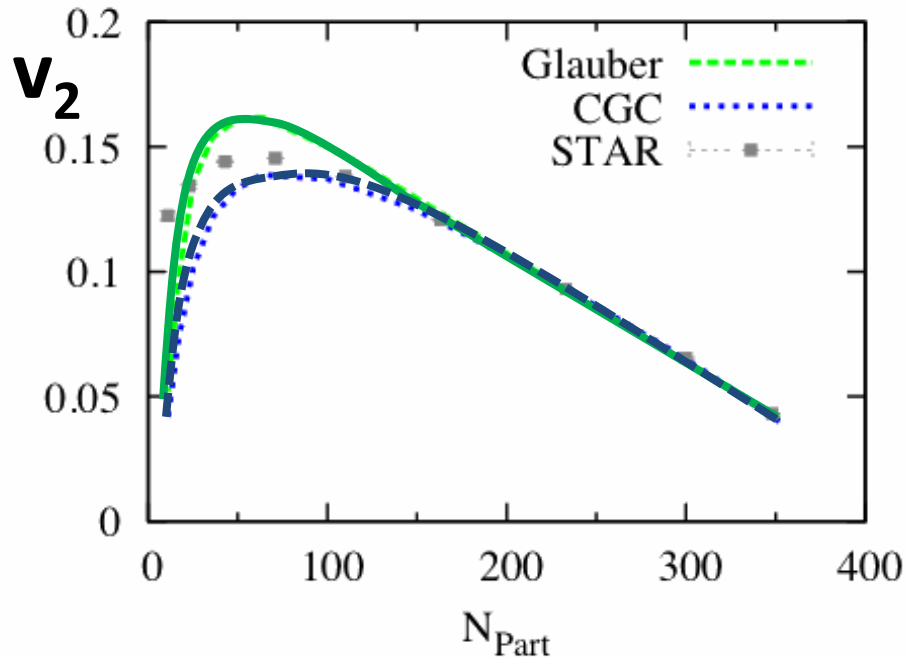
Color Glass Condensate initial  
 $\varepsilon_2$  is larger by 20% and thus  
requires more dissipation to  
get the same flow and thus  
 $\eta/s \approx 2/4\pi$

This ambiguity has troubled  
the field for 3 years.

If both have the same  $\varepsilon_3$   
then CGC with  $\eta/s = 2/4\pi$   
should give much lower  $v_3$

# Further Proof

Luzum, arXiv: 1007.5469

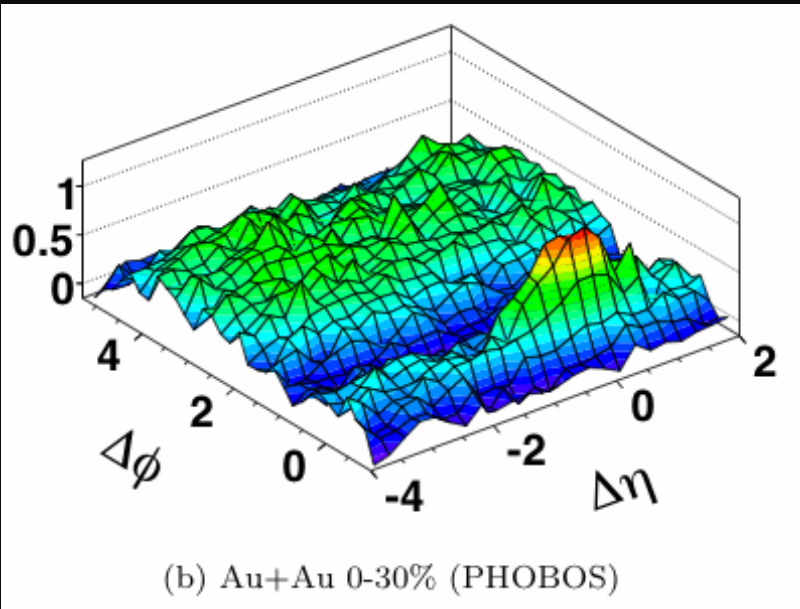


STAR states in their paper that  $v_3$  flow is not the explanation.

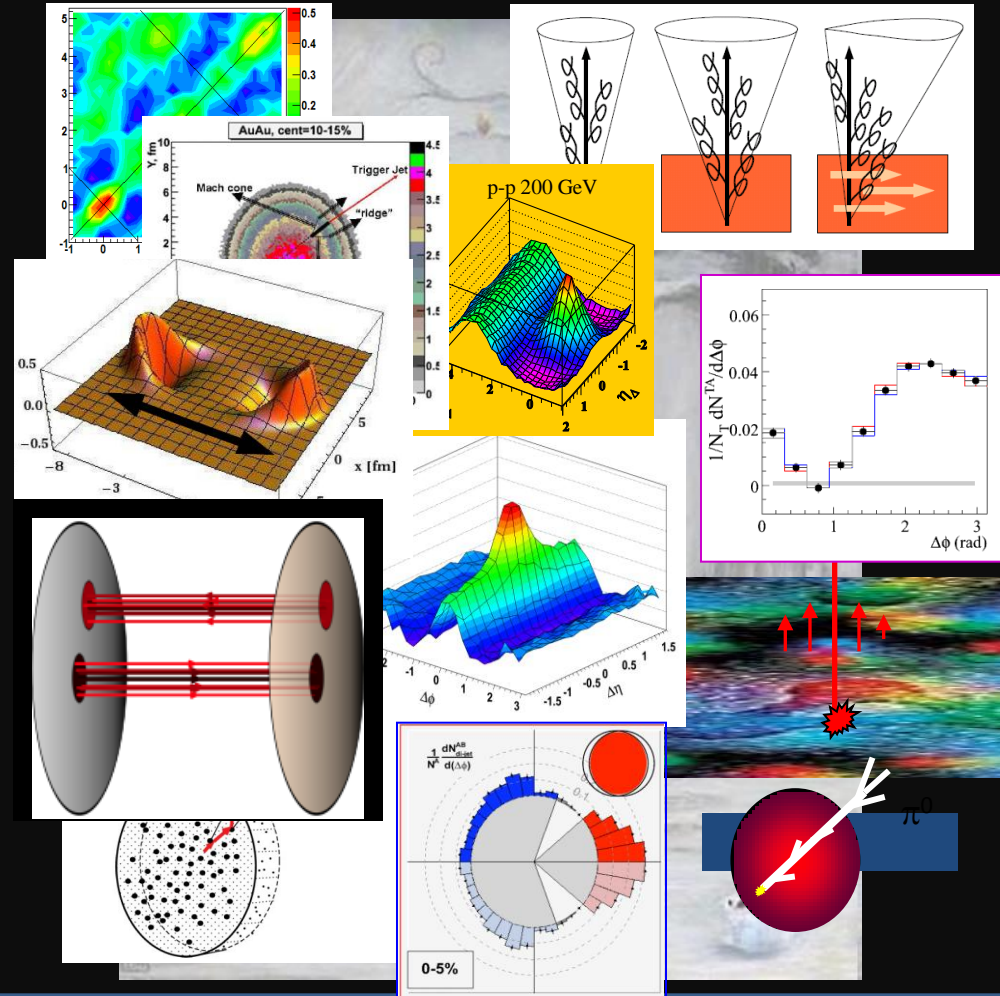
“Thus, we conclude that the away-side double-peak structure is not an artifact of the ZYAM flow subtraction procedure used in this analysis.”

STAR: arXiv:1010.0690v1

# “Ridge” and “Shoulders”



Features in two-particle correlations that have generated a lot of excitement.

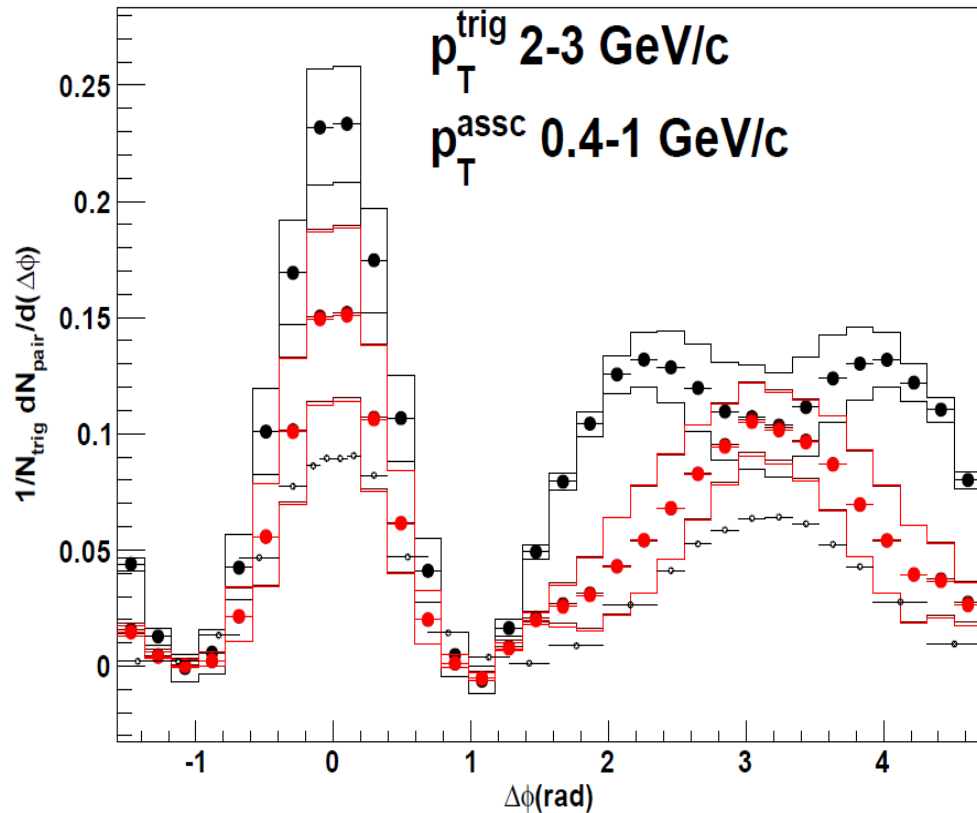


Two years ago (QM09)  
I gave a talk with my prediction.

The “death” of the ridge  
and shoulders.

# So what about the “Ridge” and “Shoulders”?

PHENIX Au+Au 0-20% Central



## Black points

PHENIX published with  $v_2$  background modulation and ZYAM.

## Red points

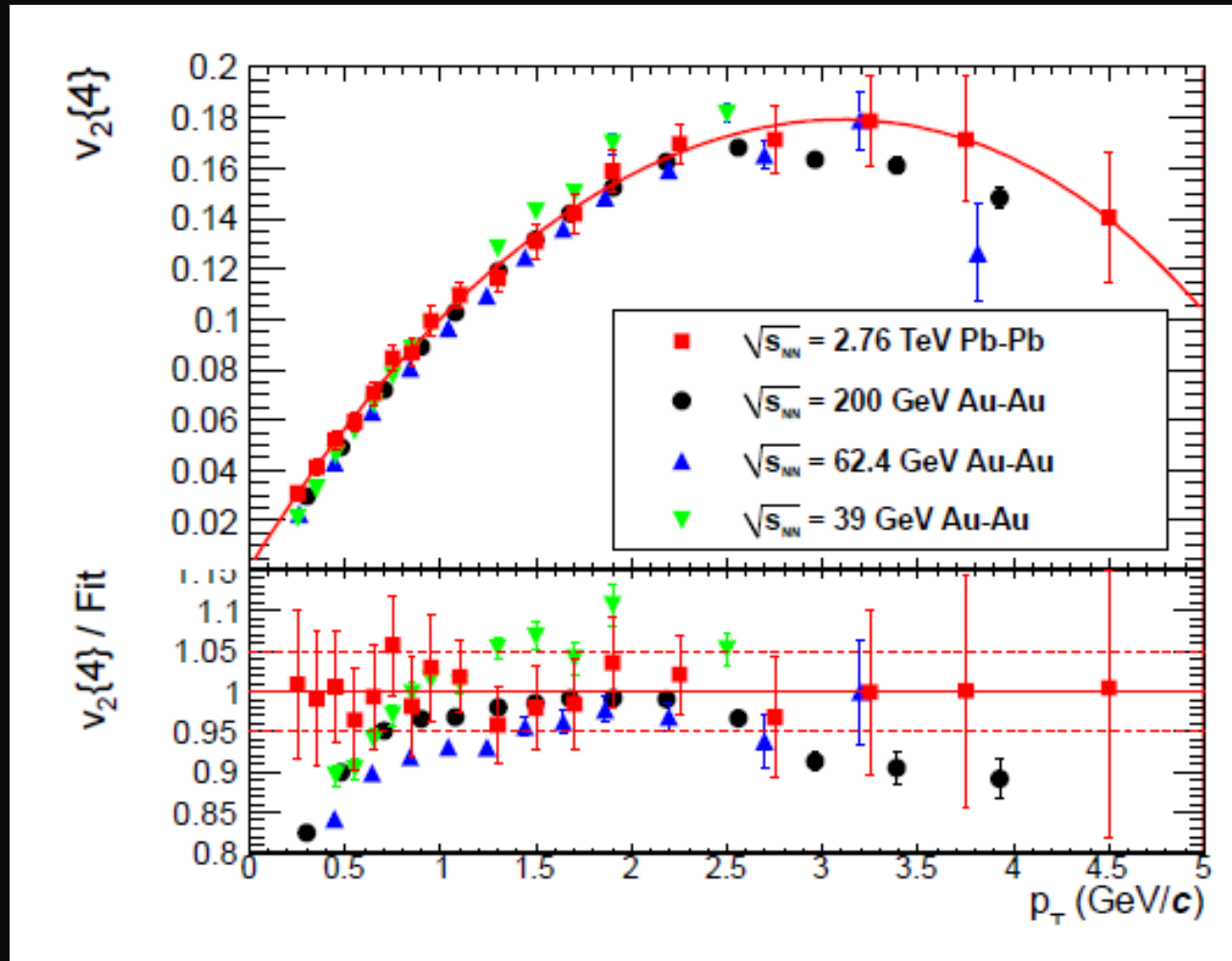
(NOT PHENIX OFFICIAL)  
Use AMPT  $v_3/v_2$  ratio and include  $v_2$  and  $v_3$  background modulation.  
Calculation by A. Adare

Dominant ridge and shoulders will be gone.

Detailed careful analysis needs  $v_1 \dots v_5$  and method checks.



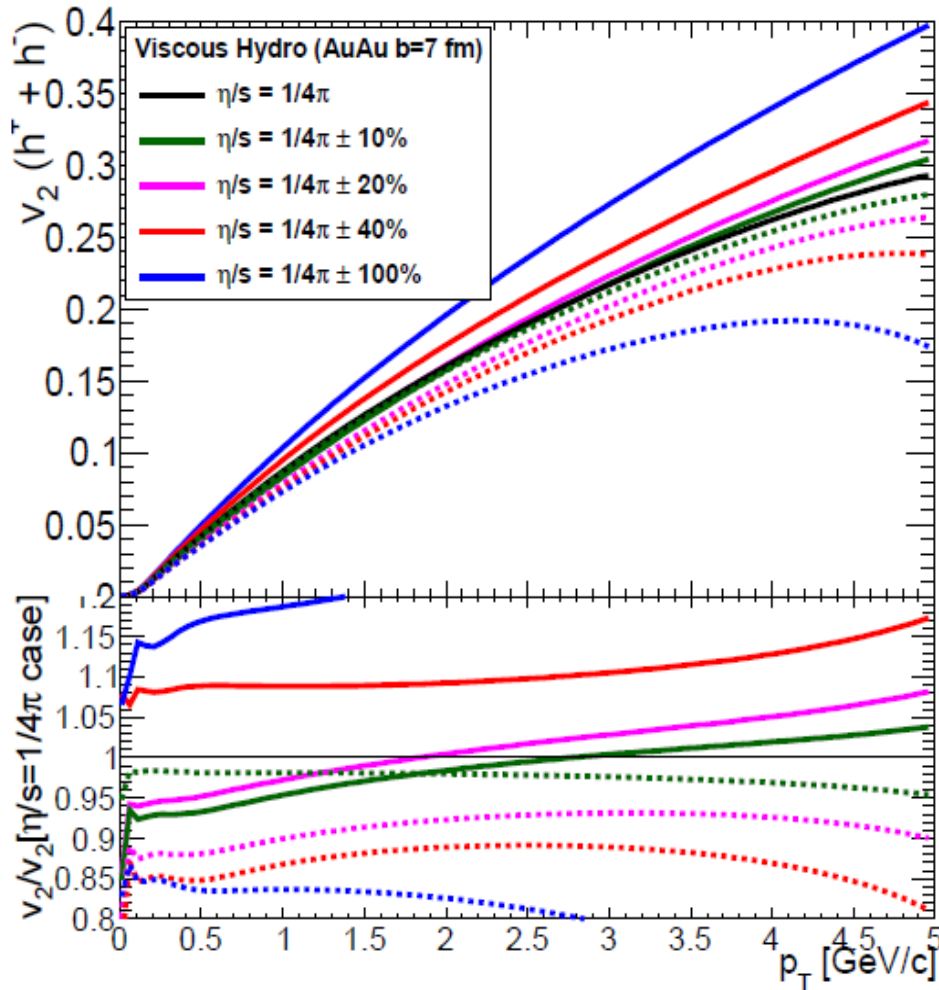
# What if we crank up the energy?



Bill Zajc, Ian Bearden and I wondered what is required to get this remarkable agreement (arXiv:1102.0680).

Use Romatschke's publicly available viscous hydro code...

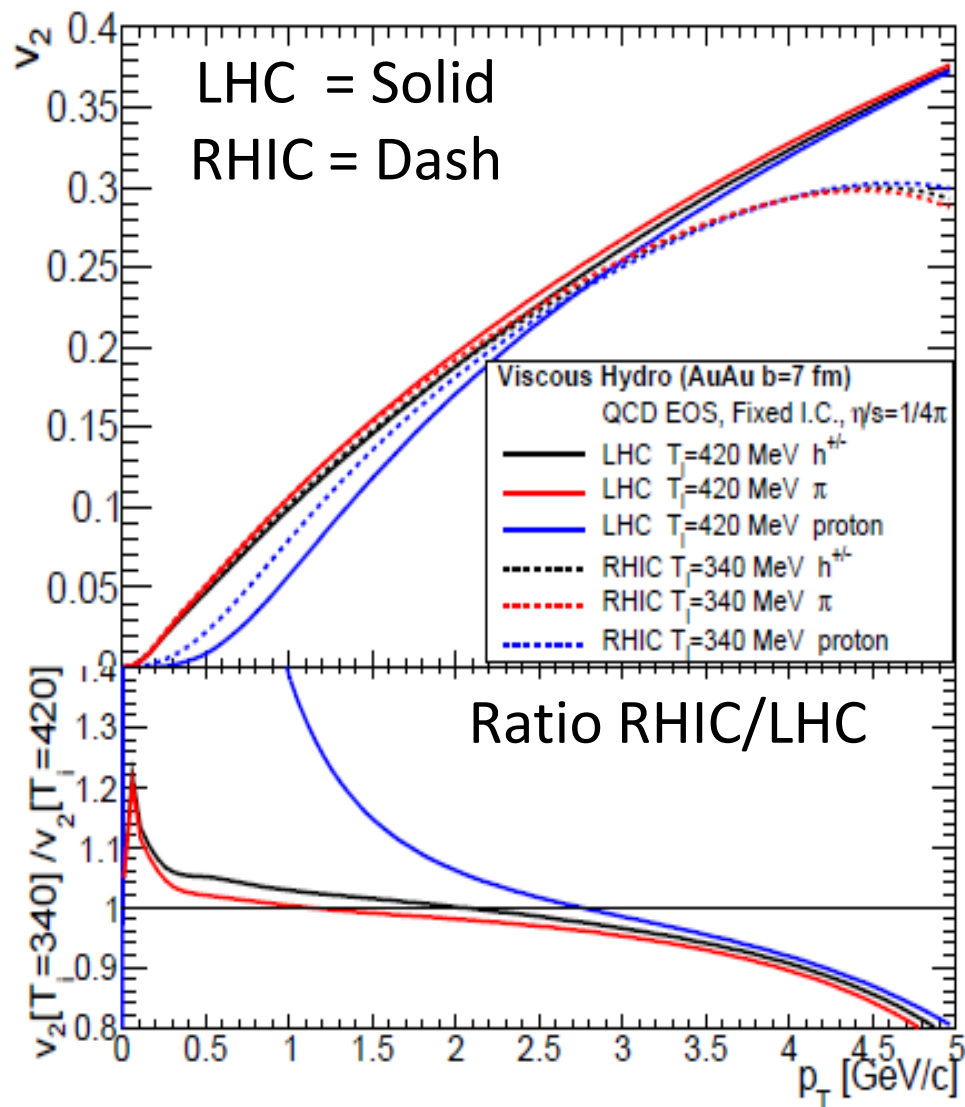
Q1: How much does a constant  $\eta/s$  need to change to get a 5% change in  $v_2$ ?



$\Delta(\eta/s) \sim 20\% \rightarrow \Delta v_2 \sim 5\%$   
 $\Delta(\eta/s) \sim 40\% \rightarrow \Delta v_2 \sim 10\%$   
 $\Delta(\eta/s) \sim 100\% \rightarrow \Delta v_2 \sim 25\%$

Always take ratios.  
 What looks like little difference at low  $p_T$ , turns out to be very sensitive (and where experiments have the smallest uncertainties).

Q2: How does  $v_2$  change for initial  $T = 420$  MeV (LHC) and  $T = 340$  MeV (RHIC)?



Black = all hadrons

Very similar  $v_2(p_T)$  except larger viscous effects for  $p_T > 3$  GeV/c

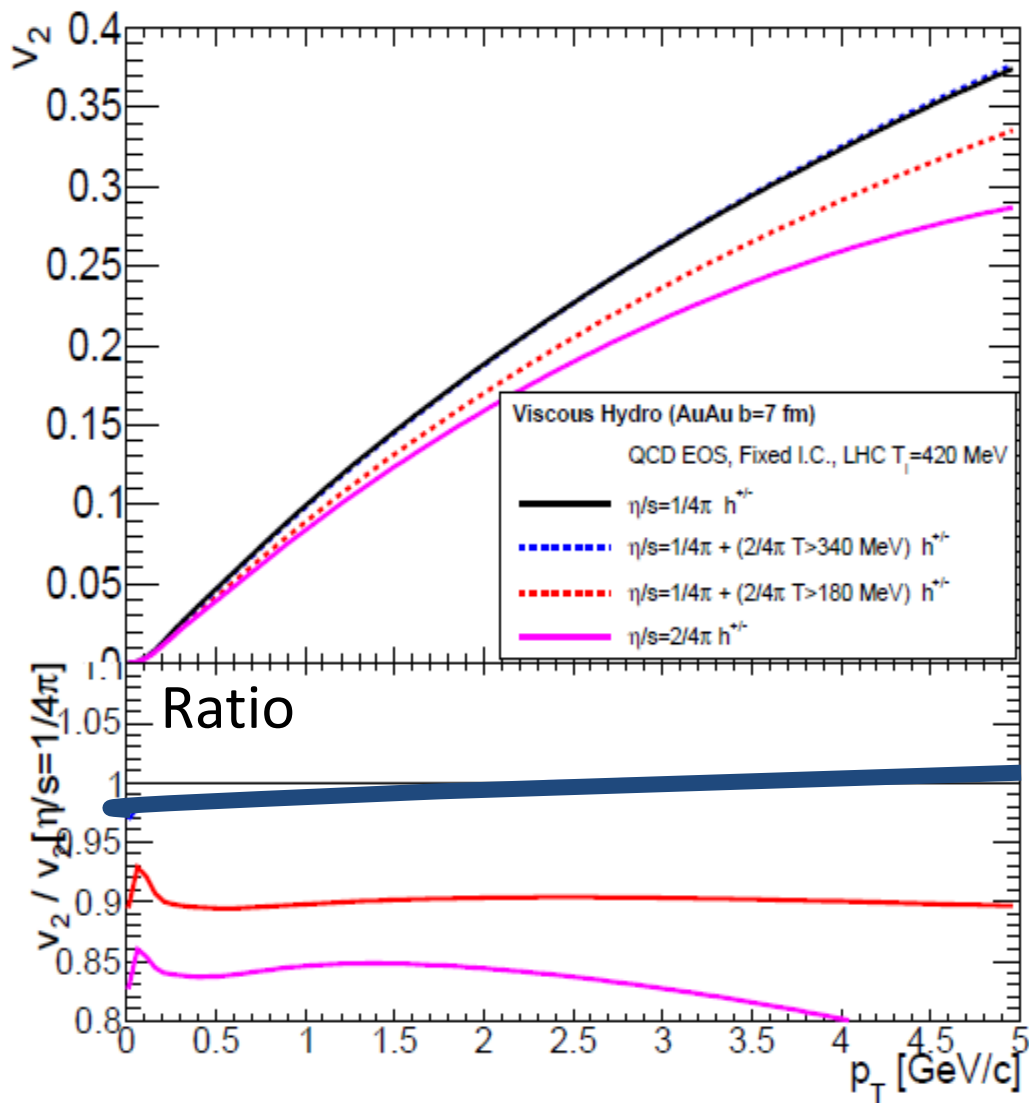
Blue = protons

Large difference for  $v_2(p_T)$  due to larger radial boost at LHC temperatures.

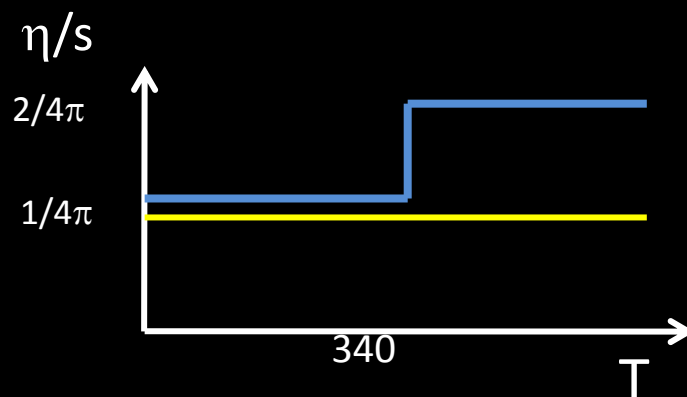
Solid prediction.

Previously noted with ideal hydrodynamics by Kestin, Heinz

Q3: What if  $\eta/s$  is larger for  $T > 340$  MeV  
(just the range sampled at the LHC in early times)?



Consider case I:



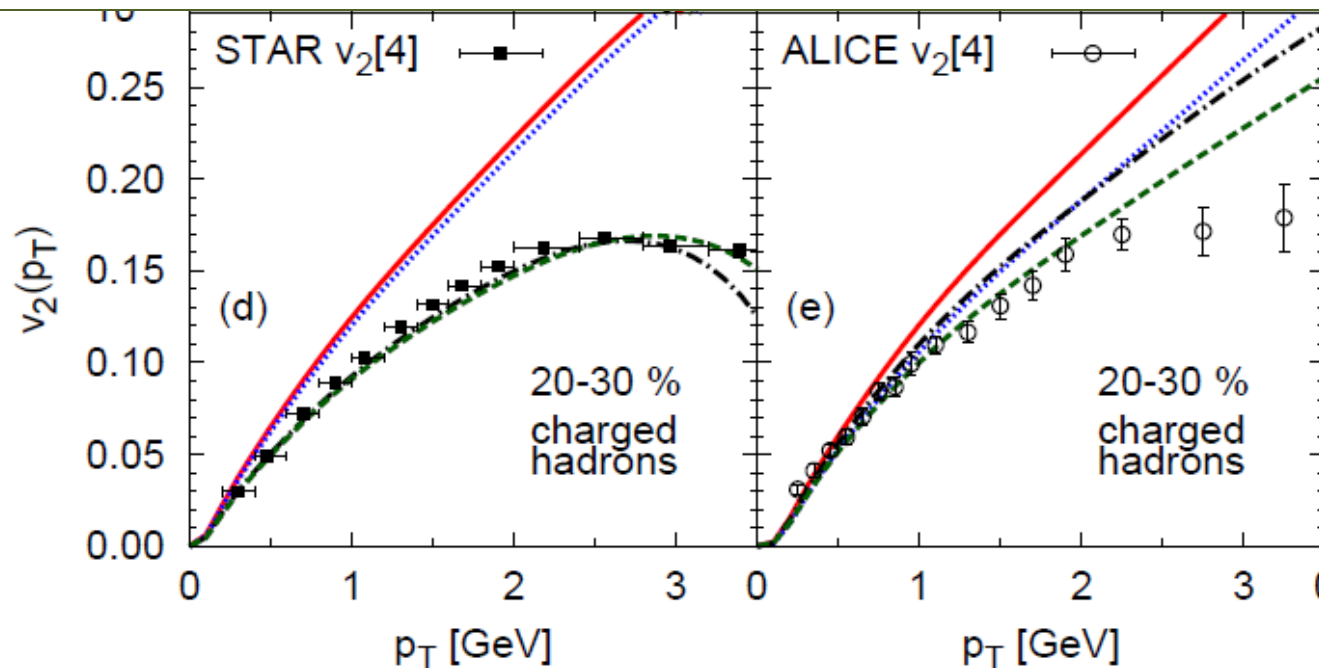
No change in  $v_2(p_T)$ !

Earliest LHC time not  
effected by factor 2  
change in  $\eta/s$ .

# Recent study of $\eta/s$ (T) – arXiv:1101.2442

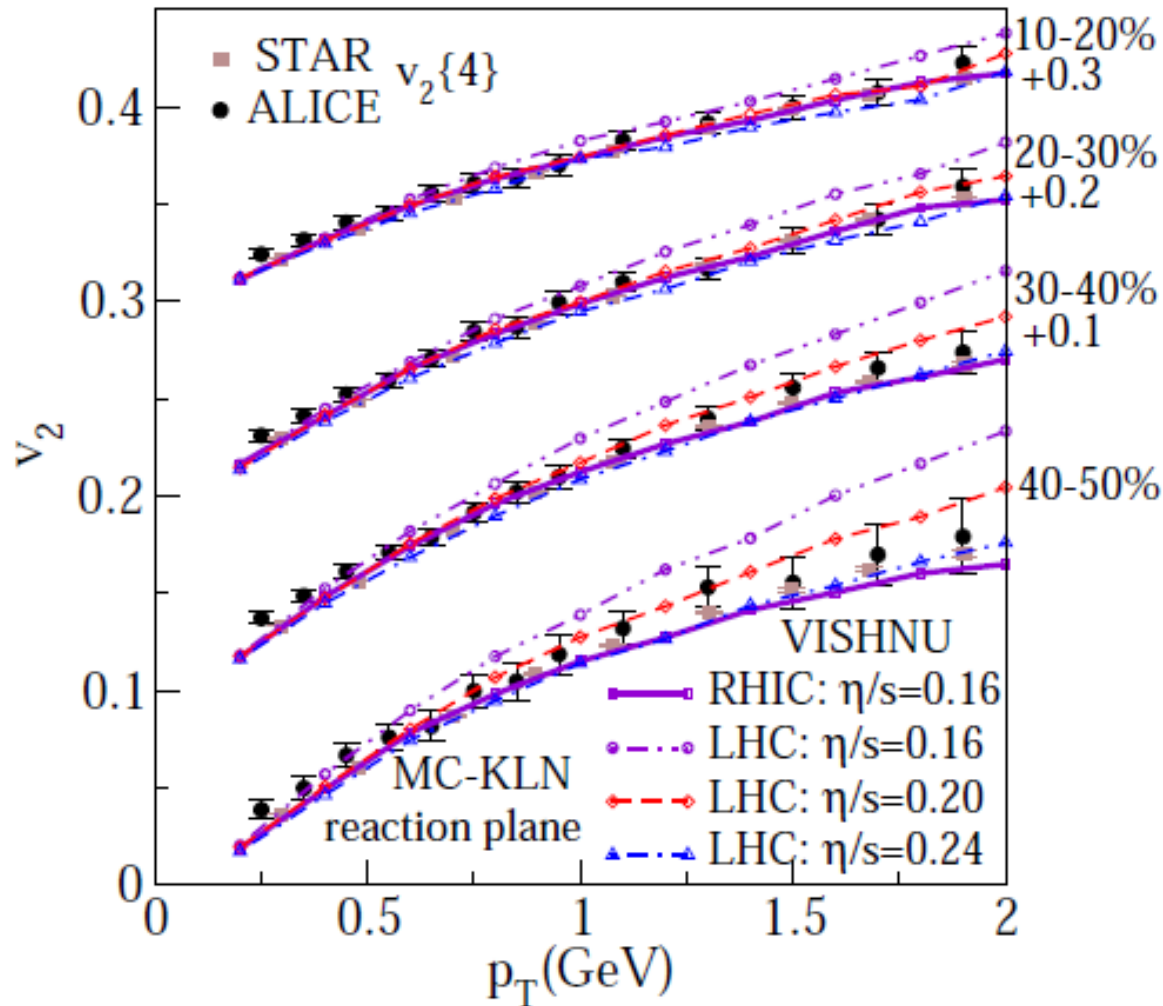
Conclusion that RHIC  $v_2$  is dominated by  $\eta/s$  below  $T_c$  and LHC  $v_2$  by  $\eta/s$  above  $T_c$  seems an unlikely fine tuning problem.

In fact, if one takes the ratios, the green-dash curve fails to describe the differences.



# Viscous Hydrodynamics + Hadron Cascade

## Fixed $\eta/s$ case for QGP



LHC needs larger QGP  $\eta/s$  to get the same  $v_2(p_T)$ .

However, the  $\eta/s$  in the same T range at RHIC cannot be different.

Fine tuning problem again.



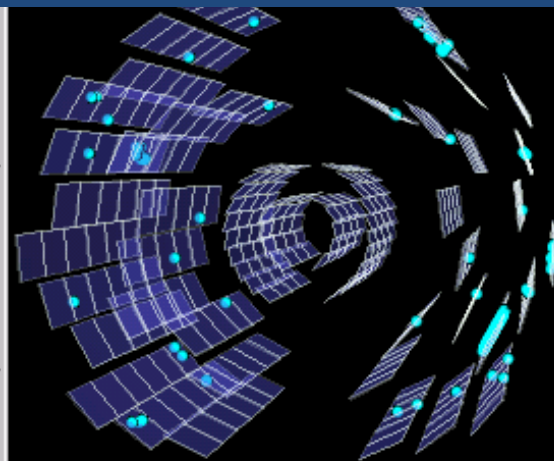
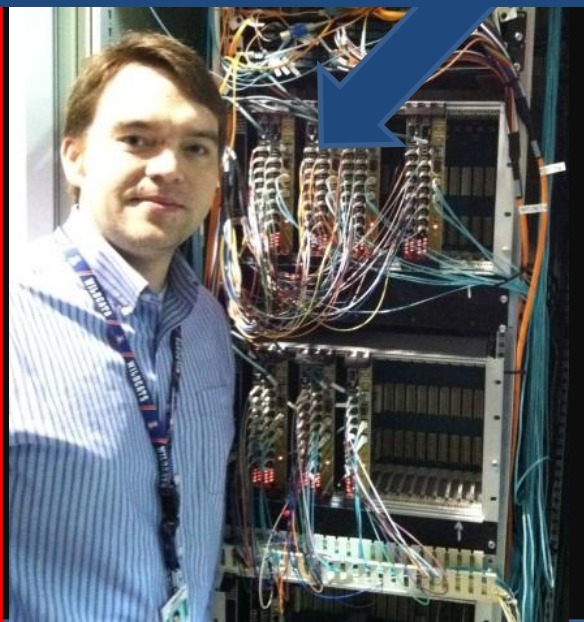
# Another Class of $\eta/s$ Constraints

## “Does the Charm Flow at RHIC?”

S. Batsouli, S. Kelly, M. Gyulassy (Columbia U.) , J.L. Nagle (Colorado U.) . Dec 2002. 11pp.  
Published in **Phys.Lett.B557:26-32,2003.**

Flow of charm and beauty quarks may yet provide the best constraints. PHENIX Silicon Detector now operational.

10 Gigabit link Data Collection system  
Colorado effort led by Mike McCumber



# PHENIX Decadal Plan

## Major Upgrade Proposed Extending into EIC Era

### The PHENIX Experiment at RHIC

*Decadal Plan 2011–2020*

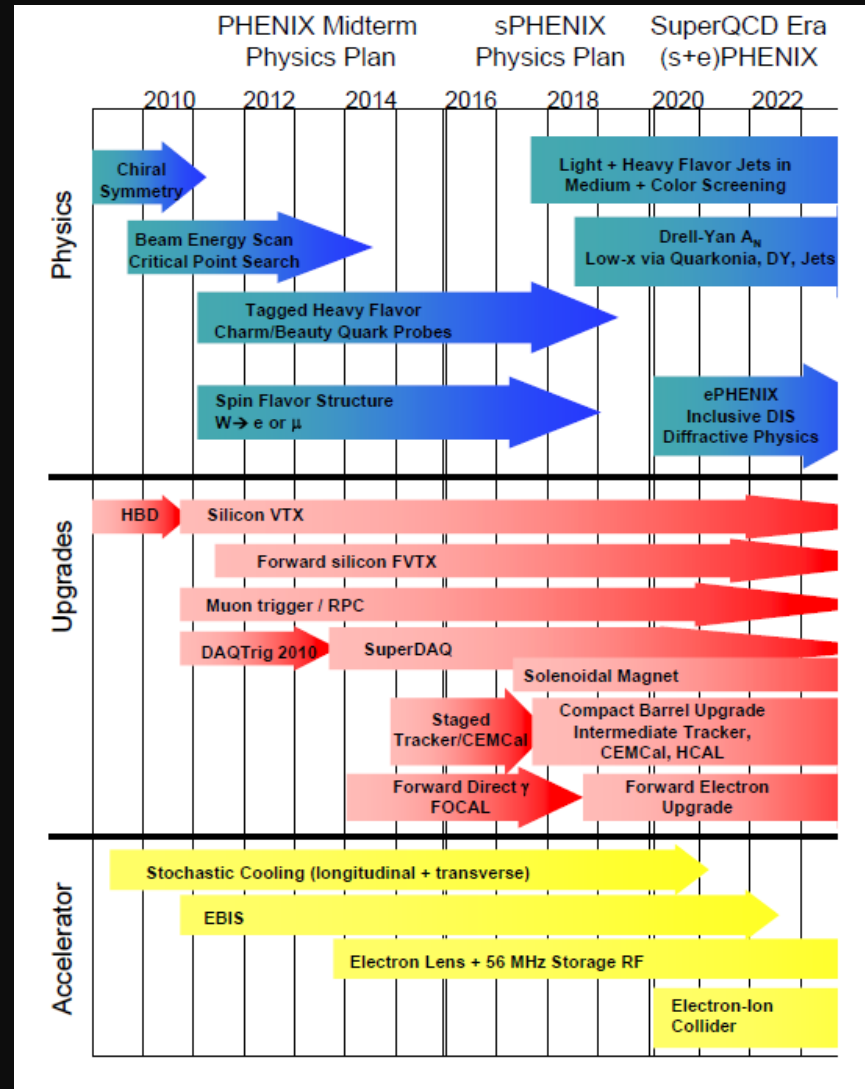
Brookhaven National Laboratory

Relativistic Heavy Ion Collider

October, 2010



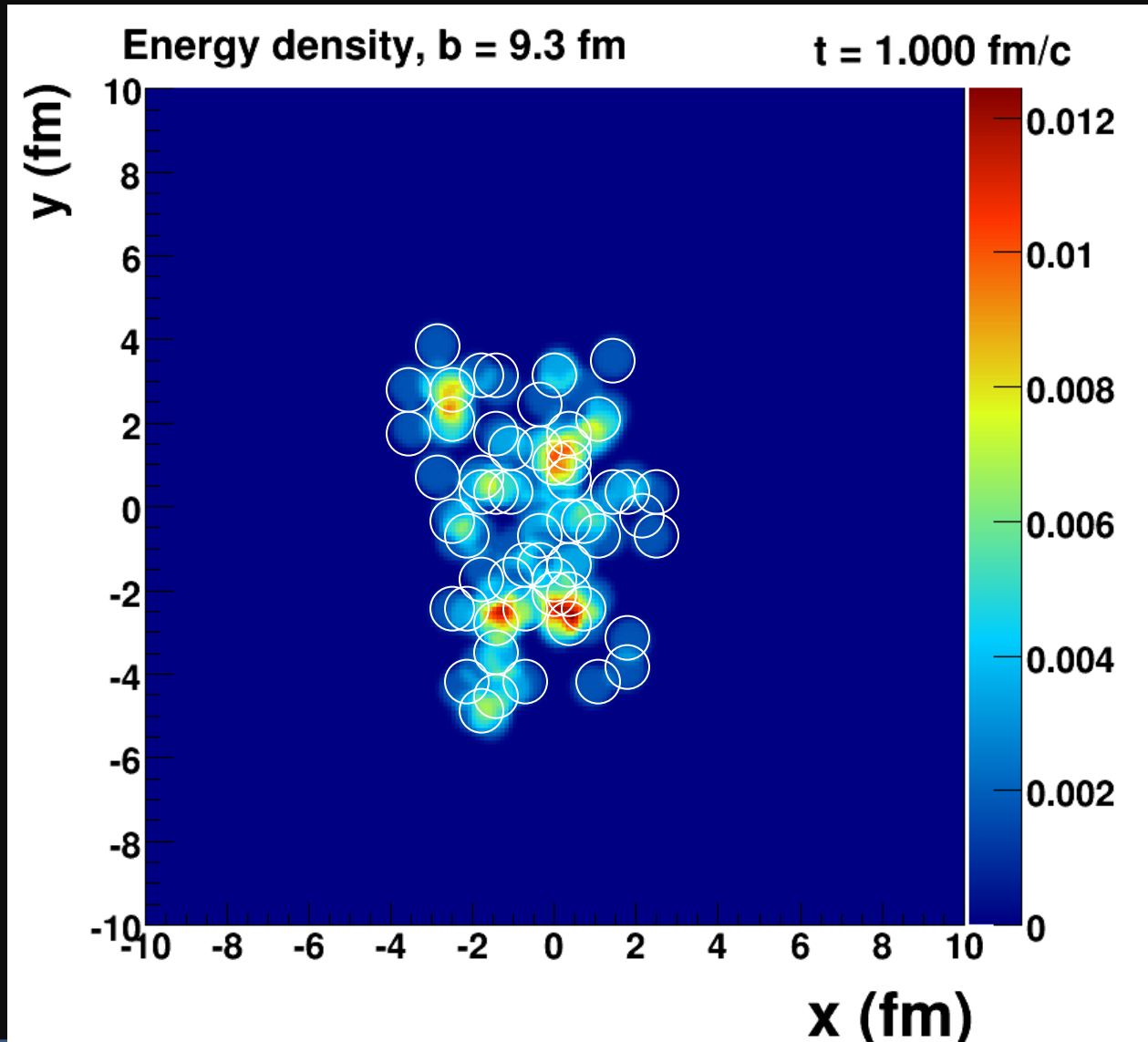
We are interested in  
feedback, suggestions,  
involvement.





# Summary:

Deviations from smooth geometry have major impact in all areas of heavy ion physics



EXTRAS

# Heavy quarks and charmonium at RHIC and LHC within a partonic transport model

J. Uphoff, K. Zhou, O. Fochler, Z. Xu, C. Greiner

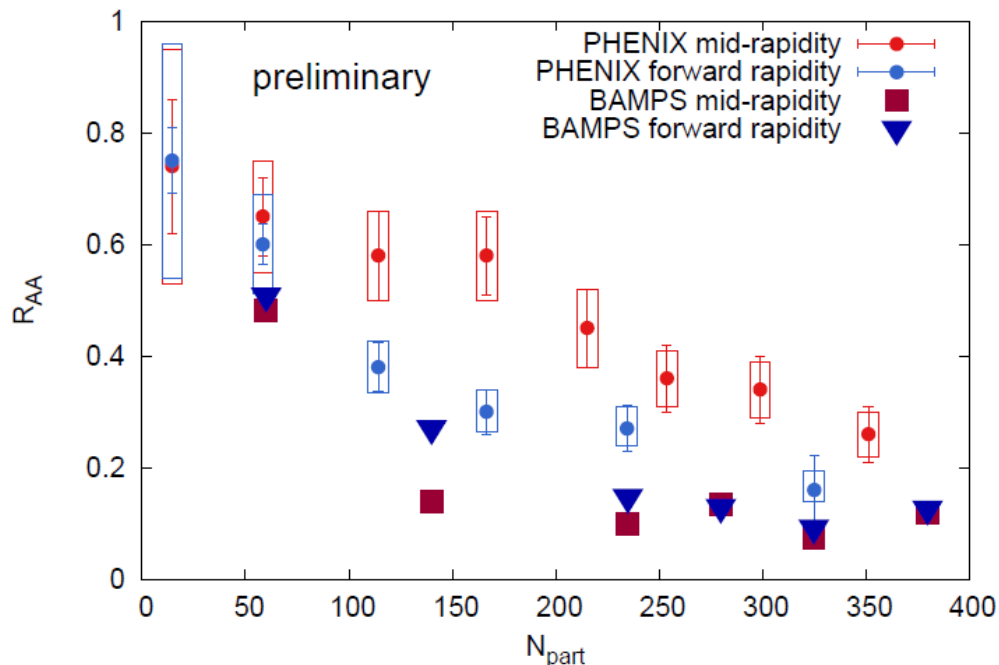
Cold nuclear matter effects linear in thickness...

Should confront dAu data...

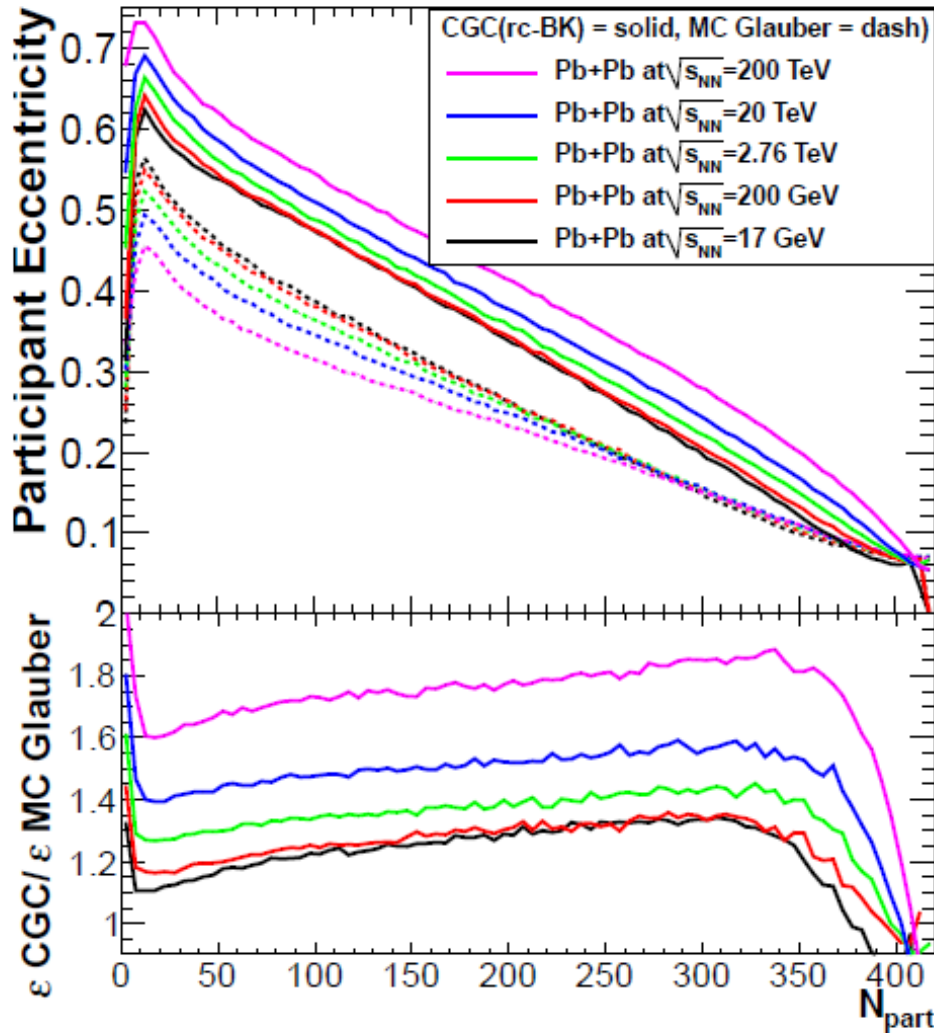
$$\mathcal{R}_i^A(\mathbf{x}_T, x, \mu_F) = 1 + N_{A,\rho} [R_i^A(x, \mu_F) - 1] \frac{T_A(\mathbf{x}_T)}{T_A(0)}.$$

Far too much J/ $\psi$   
suppression.

Less suppression at  
forward rapidity.



Q4: If the initial  $\varepsilon_2$  is Glauber, does that predict a change in  $v_2(p_T)$  as a function of colliding energy?



Glauber predicts a 5% change from RHIC (200 GeV) to LHC and a 10% change for RHIC (39 GeV).

Color Glass Condensate has an even smaller energy dependence.

That was puzzling to me since saturation effects should be larger for the lower-x partons at LHC.