

Jet-associated deuteron production in p-p collisions at 13 TeV with ALICE at the LHC



Brennan Schaefer Oak Ridge National Laboratory LANL Seminar - 16 Feb 2018

contents

- motivation why look for deuterons in jets?
- the ALICE Detector
- analysis procedure: PID, correlations, corrections
- data and findings



the nature of deuteron production is a long standing mystery in HE physics





The Coalescence Model

deuteron anisotropy measurements are consistent with coalescence picture



v₂ is approximately additive from constituents to composite particles



deuteron coalescence model measurements

$$B_2 = \frac{\pi}{2} \frac{\frac{dN_{\rm d}}{dy}}{\left(\frac{dN_{\rm p}}{dy}\right)^2} \frac{1}{\int_0^\infty \frac{f^2(p_{\rm T})}{p_{\rm T}} dp_{\rm T}}$$



Centre-of-mass energy (GeV)



are deuterons made in jets?

Ν

Nucleons from jet and bulk can in principle combine; but what is their spatial relationship? <u>Where</u> do jet fragments form?

Coherence volume

In order to coalesce, the p and n must be emitted (last scattering) at a separation similar to the deuteron size.

How do jets produce baryons? Are they correlated or anti-correlated with other baryons at close momentum?

Jet



Rulk

is coalescence more or less likely inside jets?





key questions include:

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are deuterons made in jets?
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if so, what is B₂ inside jets, background?

plan of attack:

analyze 2-particle correlation withdeuterons, protons in yr: 2015-1713 TeV p-p dataset, 1B minbias events.

(p-p events minimize contribution from flow)















The <u>Inner Tracking System</u> has six layers of silicon.

The <u>Time Projection Chamber is the</u> world's largest (88m³).

The <u>Time Of Flight</u> has 157k channels.



From a Google search, this is how you'd think we enter ALICE.





How we normally get inside the detector.









parametrization comes from Bethe-Bloch







 2σ TOF-PID cut and sidebands

1.0 GeV/c lower limit is used to avoid secondary deuterons



deuteron identification cuts are made using raw m² data







cuts use mean, width functional fittings



DGł



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deuteron candidates with $1.0 < p_T < 4.4 \text{ GeV/c}$ are in about 1/4500 p-p 13TeV events







deuteron to hadron(3.0+ GeV) correlation





is the correlation due to impurities?

 $Y_{per-trigger}^{deuteron}$

CorrelatedYield TriggerCount

> ____back _____signal + back

3. Find impurity ratio with mass-squared fitting.

1. Subtract uncorrelated pairs "underlying event" with ZYAM method (zero yield at minimum).

2. Correlated yield is divided by the total trigger hadron count.

4. Select side-band regions above and below candidate region.

5. Subtract purity weighted (*side-band correlated yield / trigger*).





∆**∳ (rad)**



 $Y = Y_{per-trigger}^{deuteron}$







B/(B+S) of separate charges is combined using statistically weights



B/(*B*+*S*) of separate charges is combined using statistically weights

4.5

4.5

5

10⁵

10⁴

10³

10²

10

Brennan Schaefer

pair reconstruction efficiency is accounted for using event mixing

mixing tracks are selected from the same p_T interval,

and event multiplicity

Extraction of Correlated Jet Pair Signals in Relativistic Heavy Ion Collisions

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$$C\left(\Delta\phi\right) = \frac{\frac{d\langle n_{same}^{AB} \rangle}{d\Delta\phi}}{\frac{d\langle n_{mix}^{AB} \rangle}{d\Delta\phi}} \frac{\int \frac{d\langle n_{mix}^{AB} \rangle}{d\Delta\phi} d\Delta\phi}{\int \frac{d\langle n_{same}^{AB} \rangle}{d\Delta\phi} d\Delta\phi} \tag{1}$$

$$c = \frac{real_bin}{mix_bin} \cdot \frac{mix_sum}{real_sum}$$

Per-trigger deuteron yields

!! remaining work: duplicate per-trigger yields for protons

 $B_2 \sim \frac{deuteron_corrected_yields}{(proton_corrected_yields)^2}$

the production of light (anti)-nuclei is simulated using PYTHIA afterburners

³He measurements!

the production of light (anti)-nuclei is simulated using PYTHIA afterburners

model comparisons

also anti-triton!

basic cross check

selecting pure deuteron associates

Charge Inclusive Mass-Squared vs p_

basic cross check 5.0+ GeV/c trigger hadron

backup

$$Y_{corrected} = \left(\frac{c_{deuteron}}{N_{trig}} - \frac{back}{signal + back} \frac{N_{deut_candidate}}{N_{side_band}} \frac{c_{side_band}}{N_{trig}}\right) \frac{1}{eff. \cdot accept.}$$

 $^{1}\sigma = \sigma_{c}$

$$^{2}\sigma = \frac{back}{signal + back} \frac{N_{deut_candidate}}{N_{side_band}} \sigma_{c_{side_band}}$$

$$^{3}\sigma = c_{side-band} \frac{back}{signal + back} \frac{1}{N_{side-band}} \sigma_{N_{deut_candidate}}$$

$${}^{4}\sigma = c_{side-band} \frac{back}{signal + back} \frac{N_{deut_candidate}}{\left(N_{side-band}\right)^{2}} \sigma_{N_{side-band}}$$

$${}^{5}\sigma = c_{side-band} \frac{signal}{(signal + back)^{2}} \frac{N_{deut_candidate}}{N_{side-band}} \sigma_{back}$$

$${}^{6}\sigma = c_{side-band} \frac{back}{(signal + back)^{2}} \frac{N_{deut_candidate}}{N_{side-band}} \sigma_{signal}$$

$$\sigma_{c} = \sqrt{\sigma_{area_|\phi|<0.7}^{2} - \sigma_{area_ZYAM}^{2}}$$

statistical uncertainties are taken as Poissonian

$$\sigma_{Y} = \frac{1}{N_{trig}} \frac{1}{eff. \cdot accept.} \sqrt{\sum^{i} \sigma^{2}}$$

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PHYSICAL REVIEW LETTERS

DEUTERON PRODUCTION IN HIGH-ENERGY COLLISIONS

R. Hagedorn CERN, Geneva, Switzerland (Received August 29, 1960)

finding it as small as Ω (or smaller) is given by the integral over the deuteron wave function:

$$\int_{\Omega} |\psi_d|^2 dV.$$

Therefore, whereas each pion and nucleon is represented in S by a factor Ω/V , the deuteron gives rise to a factor

$$\frac{\Omega}{V} = \frac{\Omega}{V} \int_{\Omega} |\psi_d|^2 dV.$$
(3)

Taking a Hulthén-type wave function with hard core (the latter has little influence) gives⁵ $\int_{\Omega} |\psi_d|^2 dV \approx \frac{1}{5} - \frac{1}{10}$. With this value for Ω_d one finds the results given in Tables I and II. They show, though they do not apply directly to the experiments (no *pp* collisions have so far been analyzed with respect to deuterons), that this "elementary production" yields the correct orders of magnitude. The differences may be due to the presence of nuclear matter, to an anisotropy of nucleons in the center-of-mass frame (the transformation c.m. to lab involves the simplifying assumption of isotropy), and to "peripheral" collisions.

It is satisfying that the above picture, which treats the deuteron as a quasi-elementary particle (at least makes no assumptions about how it is formed), can be supported by a kinematical consideration. It can be shown that the above formula can be made plausible also by asking the condition that a neutron and a proton leave the interaction region with a relative momentum, which Table II. The (deuteron/proton) ratio in the lab system at 15.9° in *pp* collisions with 25-Gev primary energy.

Momentum of d and p in Gev/ c	$\Omega_d = \frac{1}{5}$	$\Omega_d = \frac{1}{10}$	Experiment <i>p</i> -Al ^a
2	0.002%	0.001%	approximately
4	0.6 %	0.3 %	constant, 2%, b
6	2.8 %	1.4 %	between 2.6 and
8	7.8 %	3.9 %	5.5 Gev/c
10	8.2 %	4.1 %	

^aSee reference 1.

^bNote that the experiment refers to p-Al collisions and the theory to pure p-p collisions.

will be explained elsewhere. Even a detailed final-state interaction treatment⁶ leads back to essentially the statistical formula with Ω_d as in (3). The phase-space integrals are computed rigorously (apart from the statistical errors of a Monte Carlo method⁷). All computing work was done on Ferranti-Mercury computers⁸ (partly by the author but) mainly by Dr. W. Laskar, University College, London. The author is very grateful to him for not only carrying out the machine runs but also writing up all data tapes.

Many thanks are due to the University College, London, for offering computing time on their Mercury, to Dr. J. von Behr, CERN, for transforming the c.m. spectra to the lab system, and to many colleagues for discussions, in particular to J. von Behr, G. Bernardini, F. Cerulus, G.

Z -> d-d Phys. Lett. B 639, 192 (2006).

alternatives to the coalescence model....?

