

Decrease \sqrt{s} : Saturation turns off ($Q_s(x) \sim \Lambda_{QCD}$), hydro phase shorter ($T_{initial}$ decreases, τ_0 increases), maybe jumps with ($c_s, \eta/s$ at deconfinement), Knudsen number $\sim (TR)^{-1}$ higher, chemical composition different ($\tau, Kn \sim \mu_B?$)

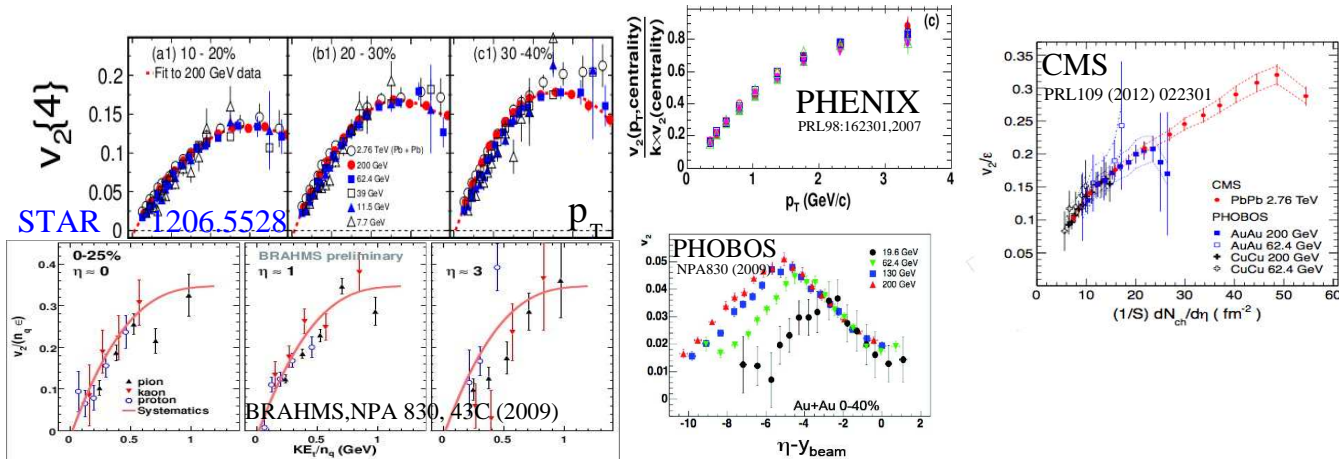
Increase rapidity: Saturation effects larger, hydro phase shorter (as above)

Decrease system size Slightly less saturation, $T_0(\sim N_{part}^{1/3})$, bigger R_0^{-1} Knudsen number
(pA : Higher y and smaller size!)

Increase p_T first role of flow is increased, then tomographic regime

Abundance of experimental data makes “toy models” useful

Low p_T harmonics the experimental situation



Here is what we know experimentally

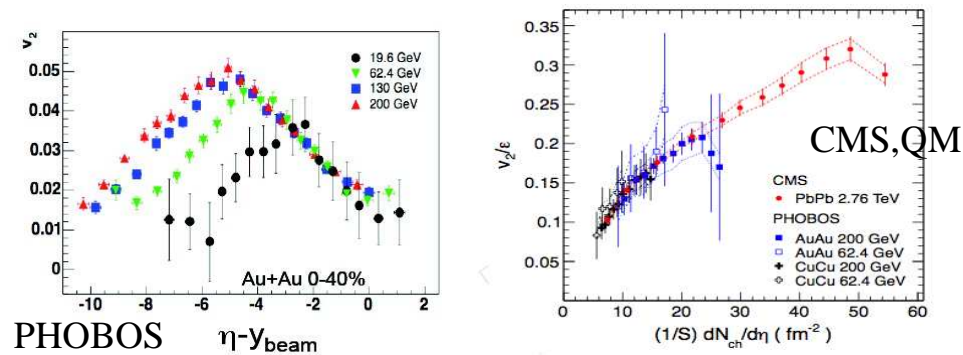
$$v_2 \simeq \epsilon(b, A) F(p_T) \quad , \quad \langle v_2 \rangle \simeq \int dp_T F(p_T) f(p_T, \langle p_T \rangle_{y,A,b,\sqrt{s}})$$

$F(p_T)$ universal for all energies , $f(p_T)$ tracks mean momentum, $\sim \frac{1}{S} \frac{dN}{dy}$

This is an experimental statement, as good as the error bars

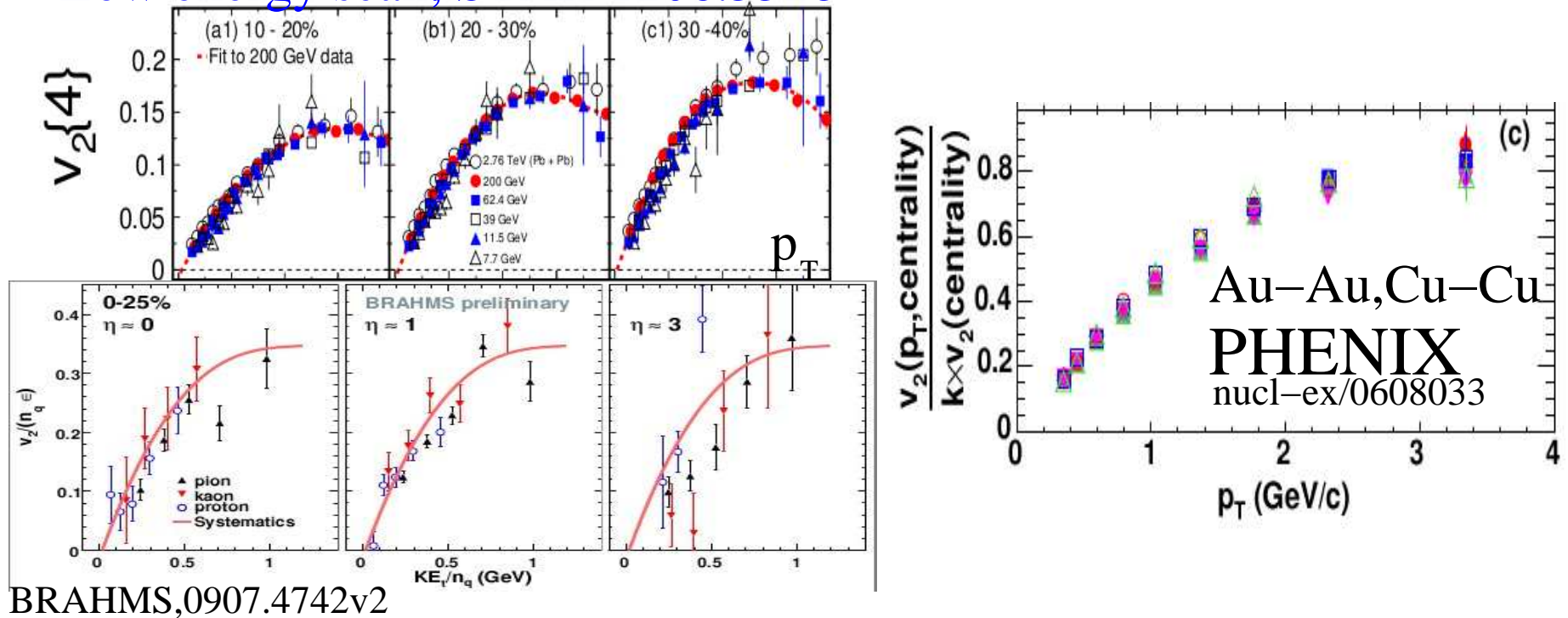
$$\frac{v_n}{\epsilon_n} \sim c_s f \underbrace{\left(\frac{1}{T_f^3 \tau_0 R^2} \frac{dN}{dy} \right)}_{\sim \tanh(\dots)} \left(1 - \mathcal{O}(1) \frac{\eta}{s} \underbrace{\frac{1}{TR}}_{\text{no scaling}} \right)$$

No change in **any of these terms** observed, from SPS to RHIC!



No sign of Knudsen number. v_3 of pA,AA **amazing**, as $\epsilon_3^{pA} \simeq \epsilon_3^{AA}$, but Kn $\mathcal{O}(10)$ bigger. Also no $\tau_0(\sqrt{s})$ seen.

Low energy scan, STAR 1206.5528



Puzzle2: Why is $v_2(p_T)$ constant (at least at **high** p_T)?

NB: this means $v_2(p_T) / \langle v_2 \rangle$ independent of N_{part}, y, \sqrt{s}

Cooper-Frye

$$v_2(p_T) = \int d\phi \cos(2\phi) \left(E - p_T \left(\frac{dt}{dr} + \Delta \frac{dt}{dr}(\phi) \right) \right) e^{\left(-\frac{\gamma(E - p_T(u_T + \delta u_T(\phi)))}{T} \right)}$$

$$\simeq \int d\phi \cos^2(2\phi) \left[\underbrace{e^{-\frac{\gamma(E - p_T u_T)}{T}}}_{=0} - \underbrace{p_T \Delta \frac{dt}{dr}}_{\epsilon p_T} + \underbrace{\frac{\gamma \delta u_T(\phi) p_T}{T}}_{\sim \frac{\delta v_T}{T} p_T \sim \epsilon p_T / T} + \mathcal{O}(\epsilon^2) + \mathcal{O}(Kn) \right]$$

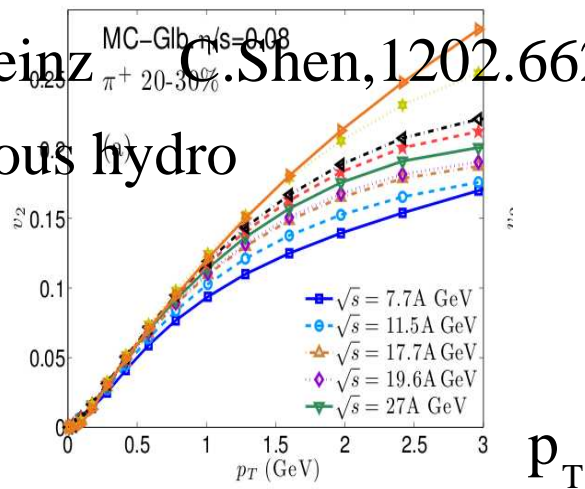
As long as $\frac{\delta v_T}{T} \sim \epsilon s^0$ (close to saturation of ϵ_p), $v_2(p_T)$ independent of \sqrt{s}

In ideal and long-lived limit $\frac{\eta}{sT} \ll R$ in $v_2(p_T) \sim \epsilon \tanh\left(\epsilon \frac{p_T}{T_f}\right) \xrightarrow{p_T \gg T} p_T^0$

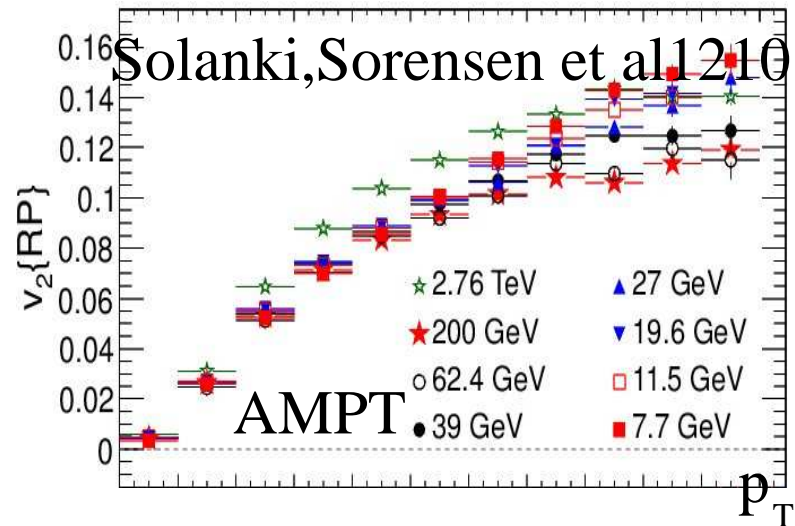
NB: deviations $\sim p_T$, more prominent at @high p_T

U.Heinz C.Shen, 1202.6620

viscous hydro



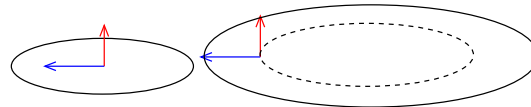
Solanki,Sorensen et al 1210.0512



This is why **all hydro and transport models** generally give a systematic shift of $v_2(p_T)$ **going up with p_T** , as long as $p_T \sim \mathcal{O}(1) \langle p_T \rangle$. Thus this data-model disagreement **not likely to go away!**

Go to high p_T : At $Kn(pT - T) \geq 1$ tomographic regime

Tomographic v_n , unlike hydrodynamic v_n , depends on size as well as density



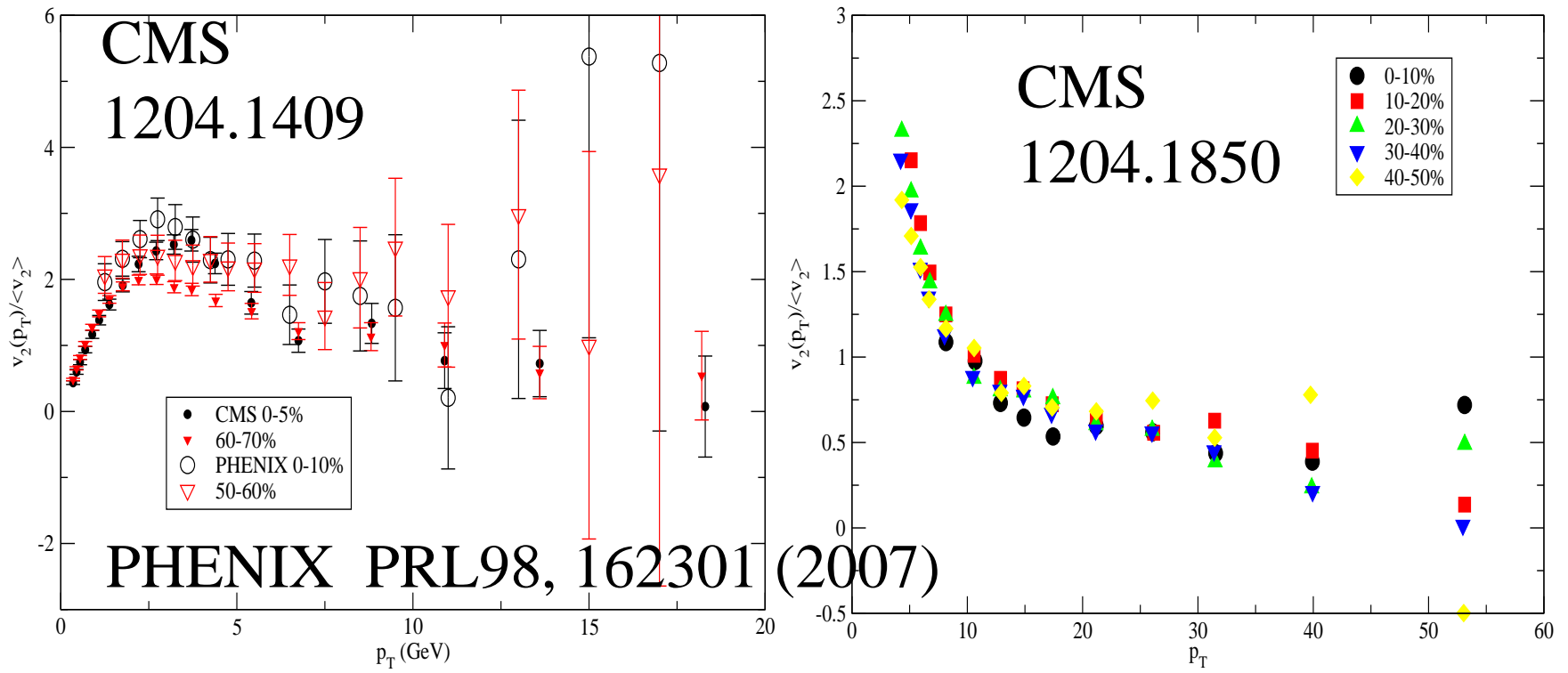
Take, as an initial condition, an elliptical distribution of matter at a given ϵ_n , run jets through it and calculate v_n . Now increase R with constant ϵ_n

$$\frac{v_n}{\epsilon_n} \Big|_{\text{tomo}} \rightarrow \frac{\text{Surface}}{\text{Volume}} \rightarrow 0, \quad \frac{v_n}{\epsilon_n} \Big|_{\text{hydro}} \rightarrow \text{constant}$$

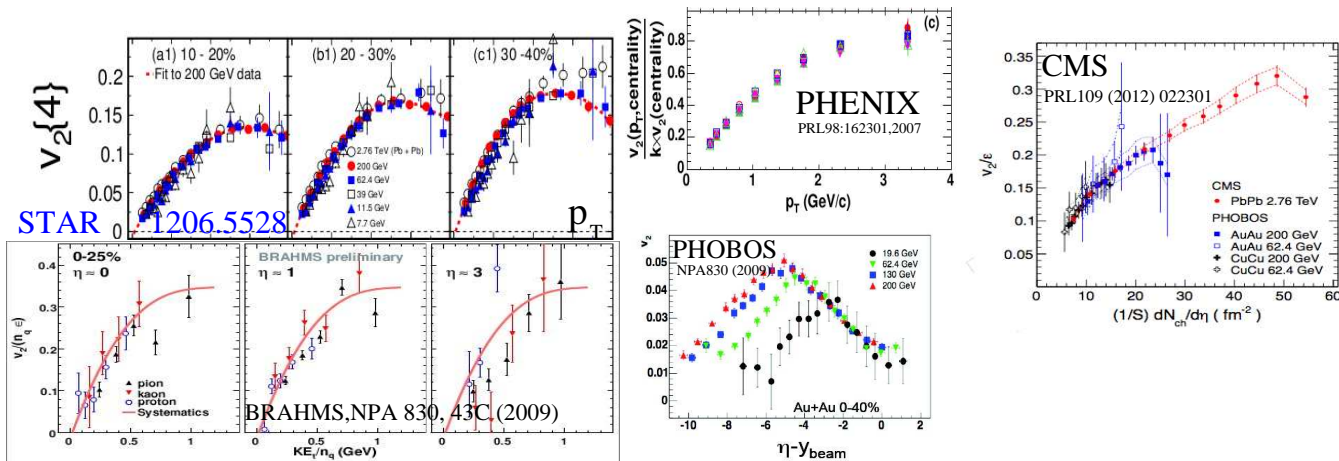
Role of “size” different in tomo vs hydro regime (M.Gyulassy, B.Betz, GT).

$$\frac{dE}{d\tau} \sim \kappa E^a T^b \tau^c \Rightarrow \Delta E \sim \langle E^x \rangle \langle T^y \rangle \langle R \rangle^z \neq f\left(\frac{1}{S} \frac{dN}{dy}\right)$$

Volcanoes and waterfalls can only make this worse, as $\epsilon_n(T \sim T_c) \neq \epsilon_n^{\text{total}}$



But , it seems no scaling break up to $p_T \sim 3.5$ GeV, maybe til 20 GeV. Split@intermediate p_T , reunification (?) higher (but big errors!). **If** no scaling violation for all momenta, puzzle for tomography



Anti-conclusion: Scalings puzzling for all popular models

Unless scalings break (**Experimentalists?**) model constrains on scalings more useful than "my model with 10 parameters fits this data at this \sqrt{s} ". Remember Bjorken scaling/partons. QCD a descendant of a scaling non-trivial with bootstrap/Smatrix/...

More info: [arXiv:1208.5996](https://arxiv.org/abs/1208.5996), [arXiv:0911.4775](https://arxiv.org/abs/0911.4775) (PRC), [nucl-th/0702013](https://arxiv.org/abs/nucl-th/0702013) (PRC), extensive calculation with ABC model using TITAN supercomputer in progress!