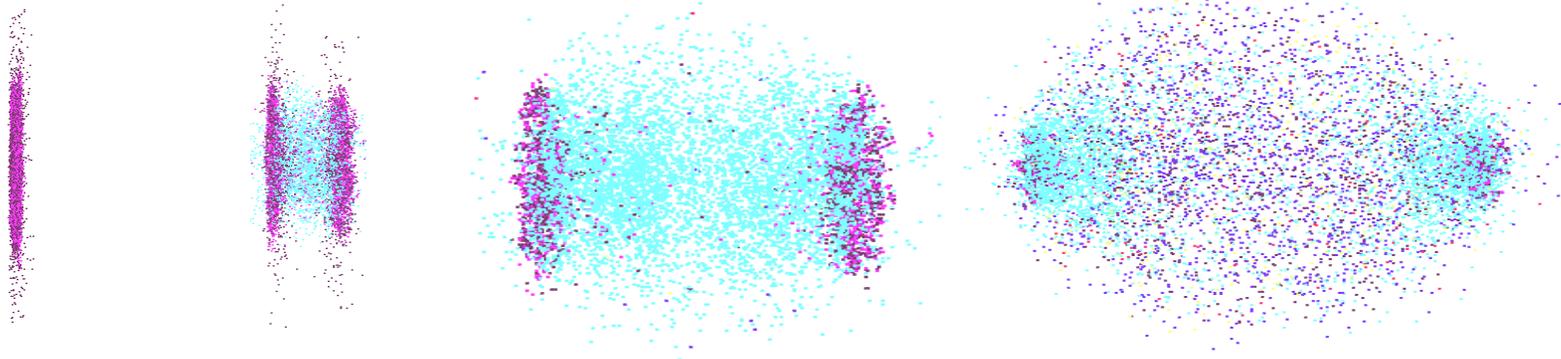


---

# What do we learn from PHENIX high pt results



Jiangyong Jia

State University of New York at Stony Brook



For the **PHENIX** Collaboration

---

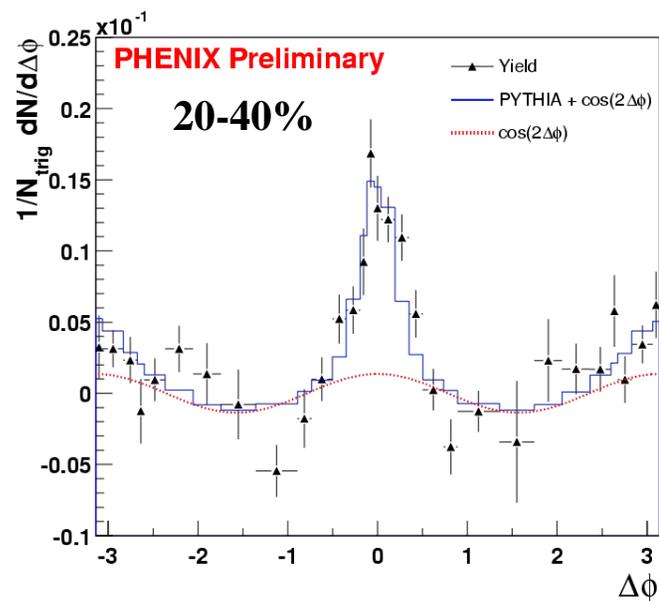
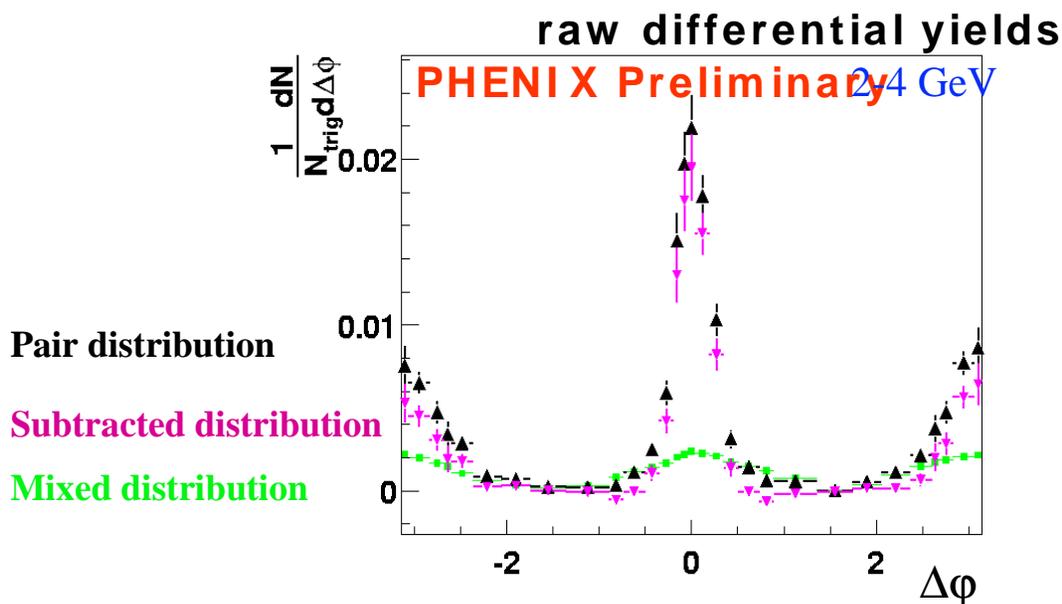
# Outline

---

- **Quick overview of high  $p_T$  results from 200GeV**
  - Jet-like angular correlation identified
  - High  $p_T$  suppression continues to highest  $p_T$  measured
  - Large  $v_2$  at high  $p_T$
- **High  $p_T$  proton and antiproton from PHENIX**
  - Analysis
  - Yield and scaling behavior
- **What have we learned**
- **Comparison with simple geometrical model calculation.**
  - Consistent with high  $p_T$  yield and di-jet suppression
  - Insufficient to describe  $v_2$
- **Conclusion**

# Direct observation of jets

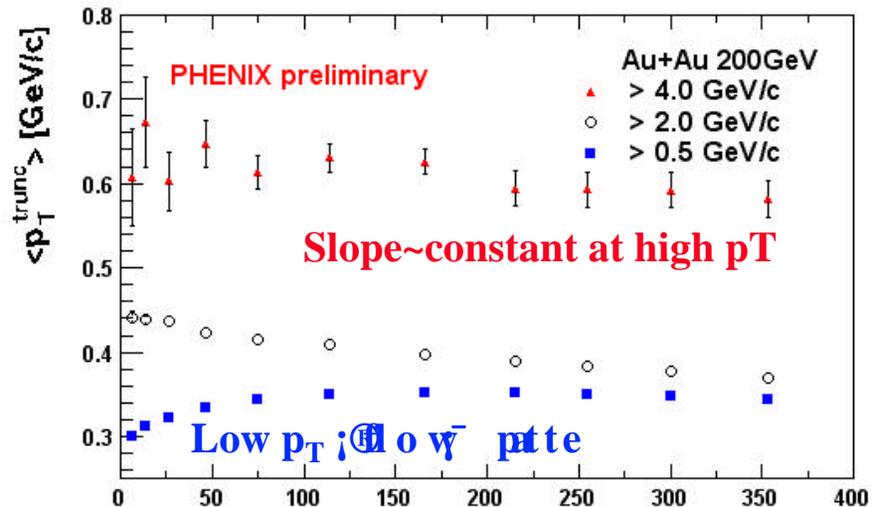
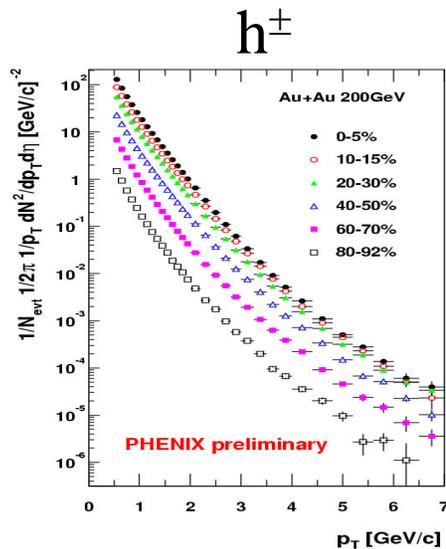
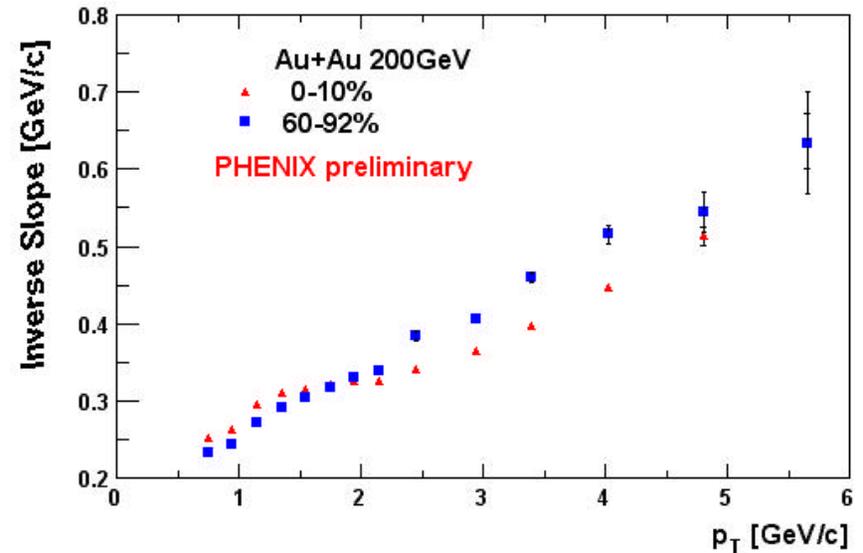
- **Leading particle + angular correlation**
  - Trigger photon  $E_\gamma > 2.5$  GeV
  - Correlate with partner charged particle
- **Jet like near angle correlation for p+p and Au+Au**
  - Similar width
  - Jet fragmentation not modified much?



STAR measure the same and away side jet yields 4 D. Hardtke

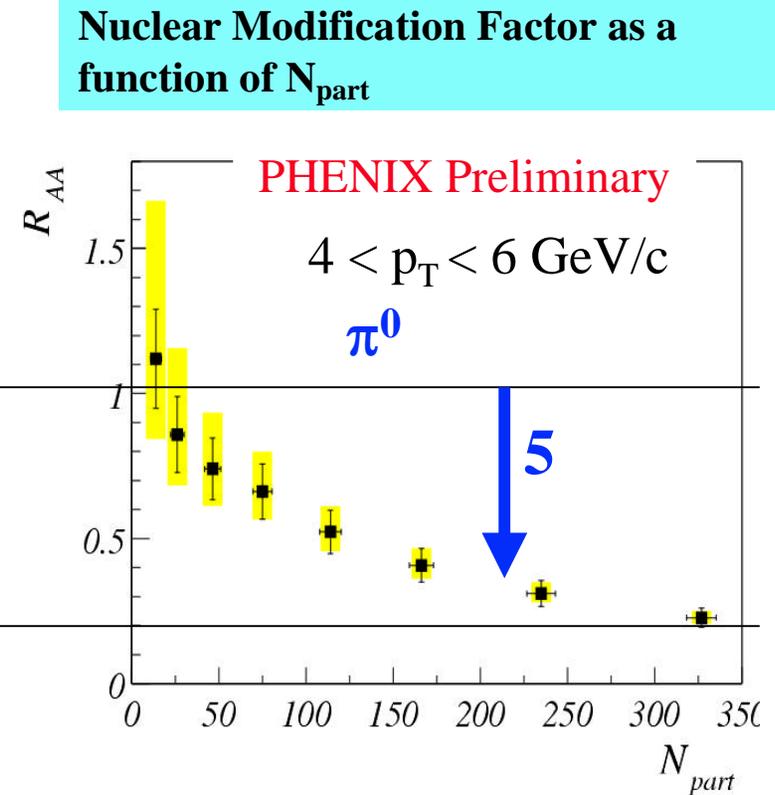
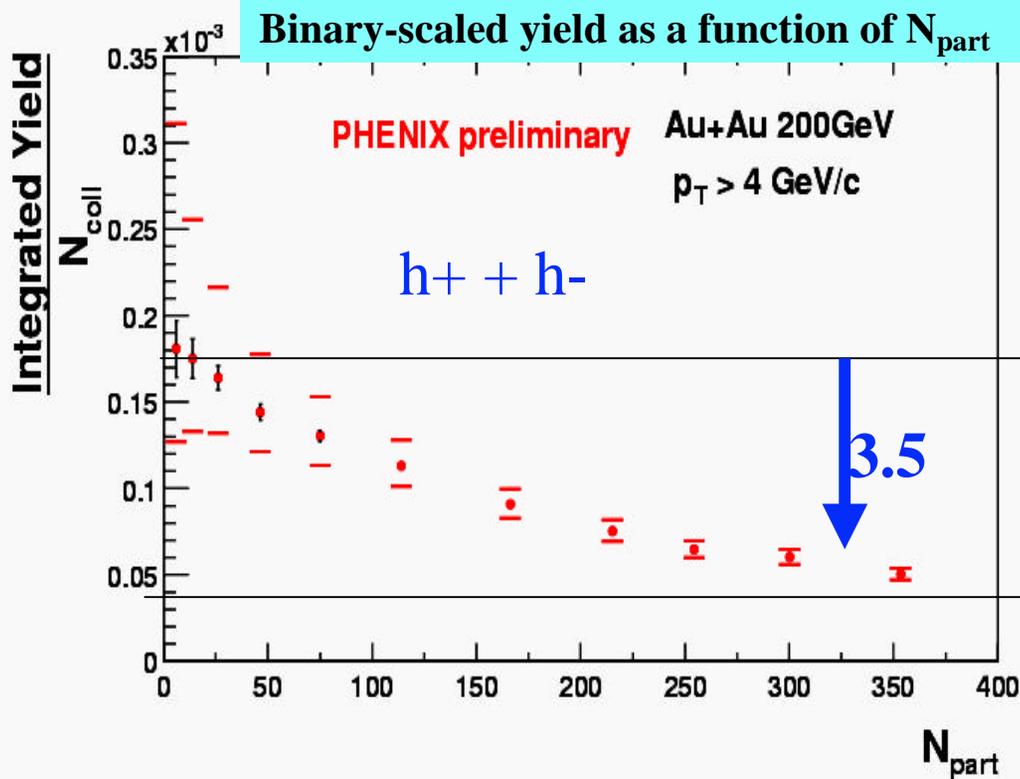
# Spectra shape at high pT

- Inverse slope increase with pT
  - Spectra is more power law
- $\langle p_T^{\text{trunc}} \rangle = \langle p_T \rangle - p_T^{\text{min}}$  as function of centrality
  - Nearly flat at high pT
- $p_T > 4 \text{ GeV/c}$  are jet like



# Centrality dependence of high $p_T$ suppression

- Continues decrease as function of centrality
- Charged : factor  $\sim 3.5$  from peripheral to central
- $\pi^0$  : factor  $\sim 5$  from peripheral to central

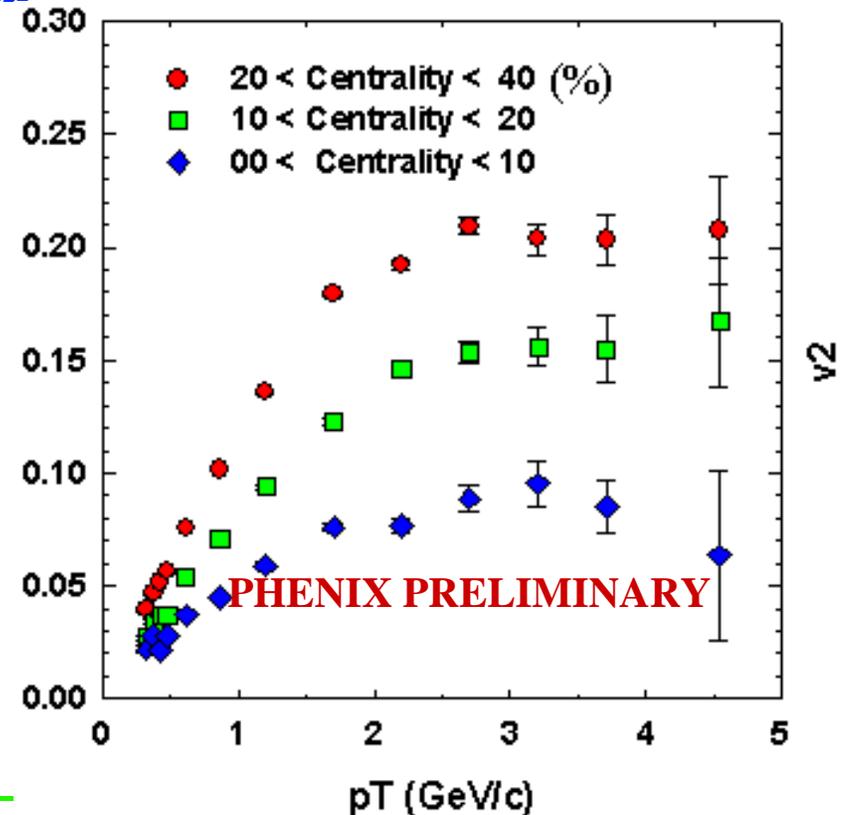
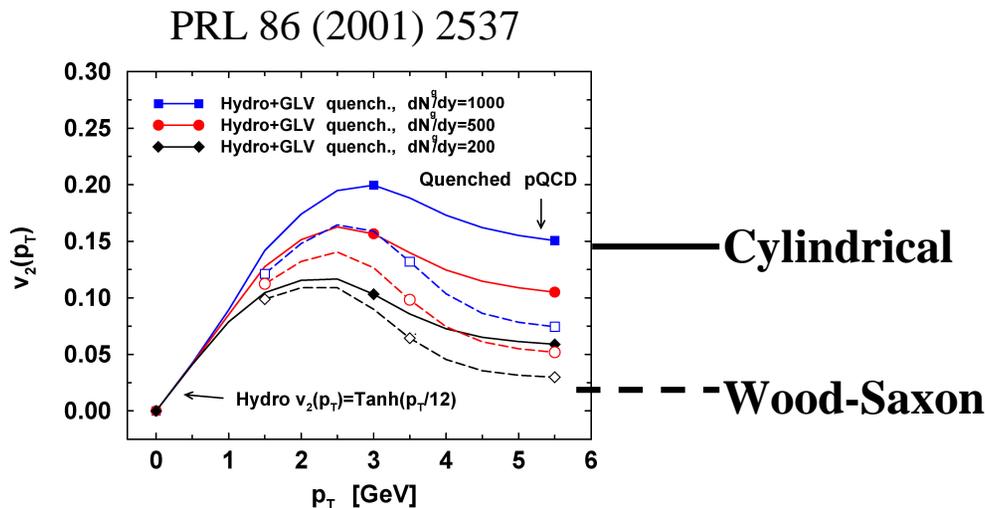


# Anisotropy in momentum space = $v_2$

- **$v_2$  saturates at high  $p_T$**
- **Can be described by pure energy loss?**
- **hydro pic:  $v_2$  created by pressure.**
- **Jet pic:**
  - **Different energy loss along different path**
  - **Sensitive to geometrical profile**
  - **Decrease with  $p_T$ ?**

$$\sqrt{s_{NN}} = 200 \text{ GeV}$$

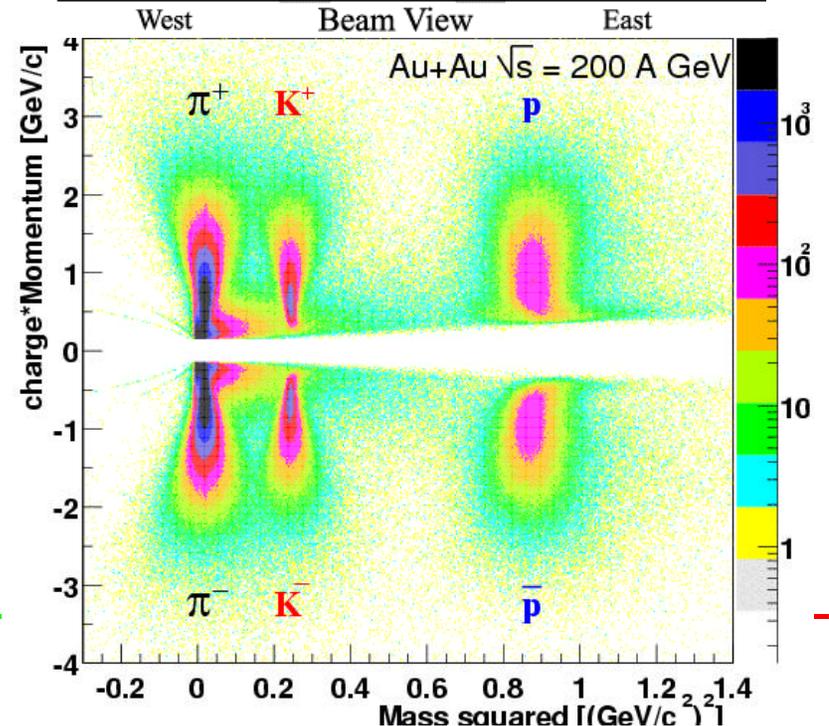
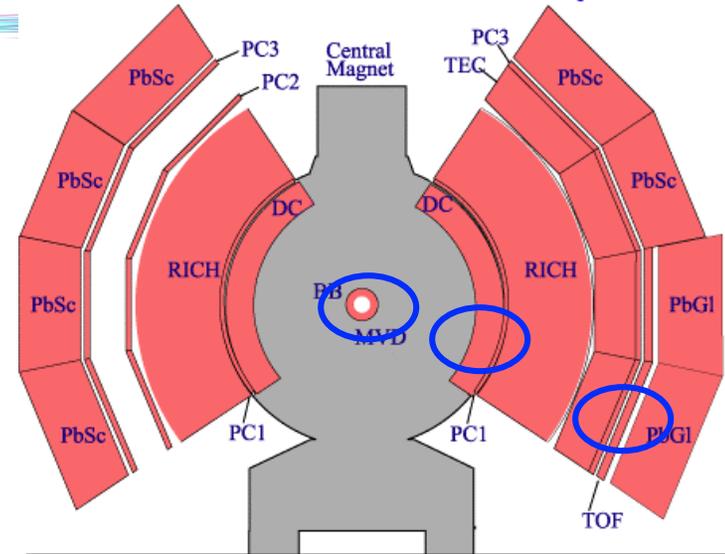
M. Gyulassy, I. Vitev and X.N. Wang



# Analysis of high $p_T$ proton and antiproton

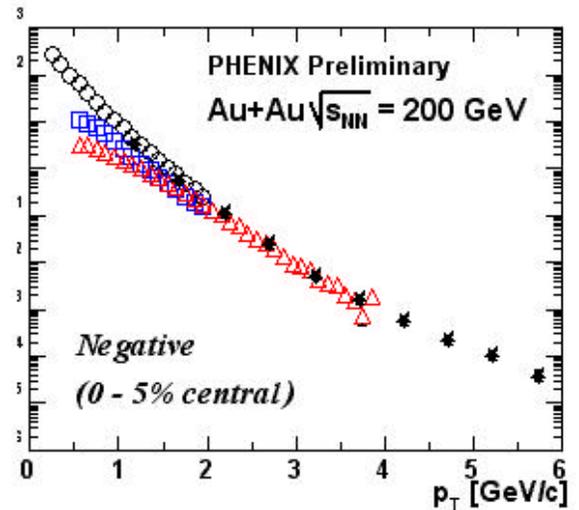
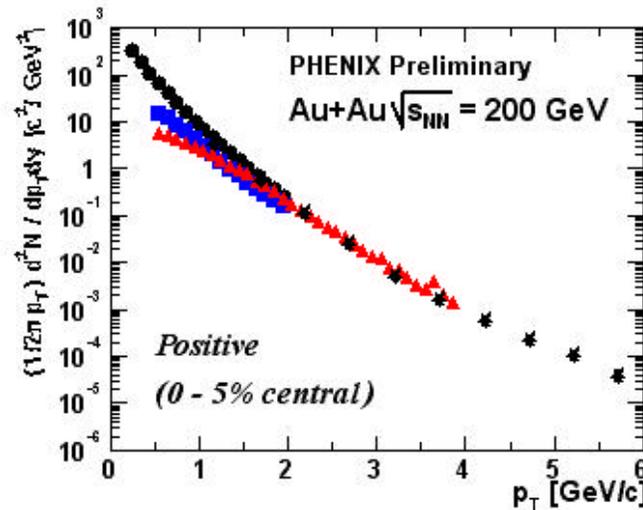
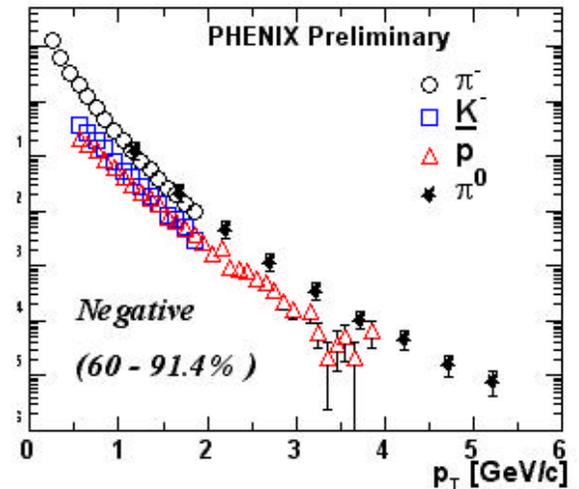
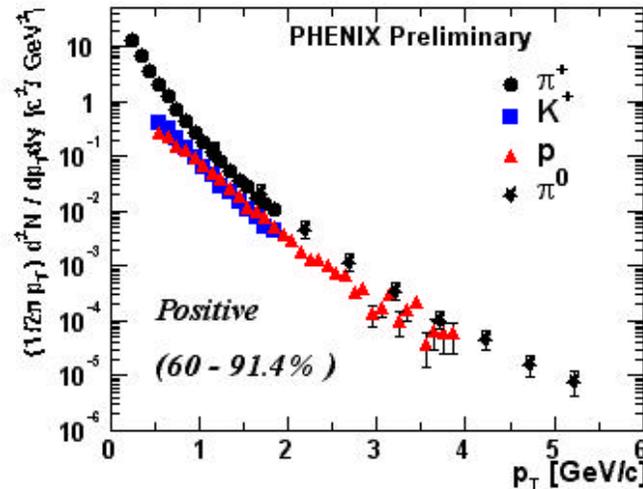
PHENIX Detector - Second Year Physics Run

- Using High Resolution TOF PID device and Drift Chamber.
- Making  $p_T$  dependent  $2\sigma$  cut in squared mass
- Range and Systematic Error
  - proton, pbar: up to 4GeV/c
  - $p_T$  dependent: 11%
  - Overall normalization: Central 18%, Peripheral 16.4%



## Identified hadron spectra at $\sqrt{s} = 200$ GeV

- Excellent agreement between charged and neutral pions
- anti(proton) yields increase with centrality relative to the pion yield at high  $p_T$

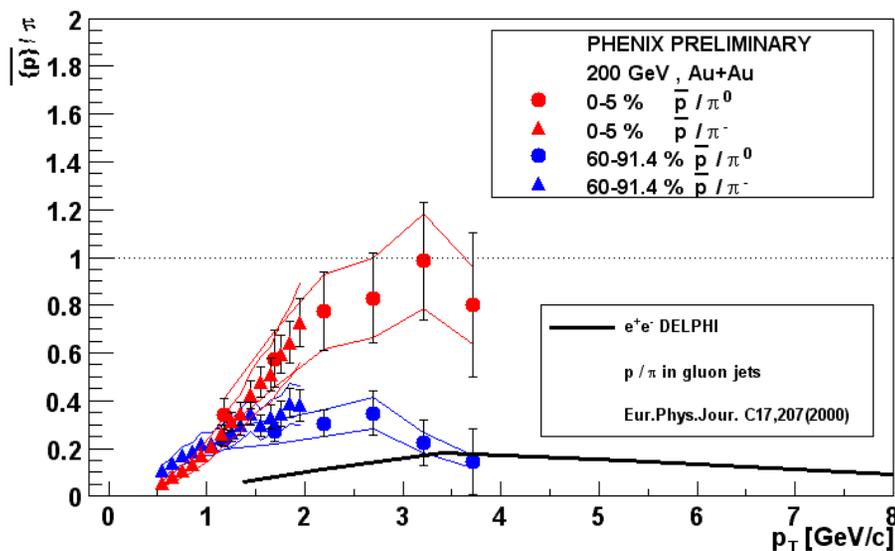
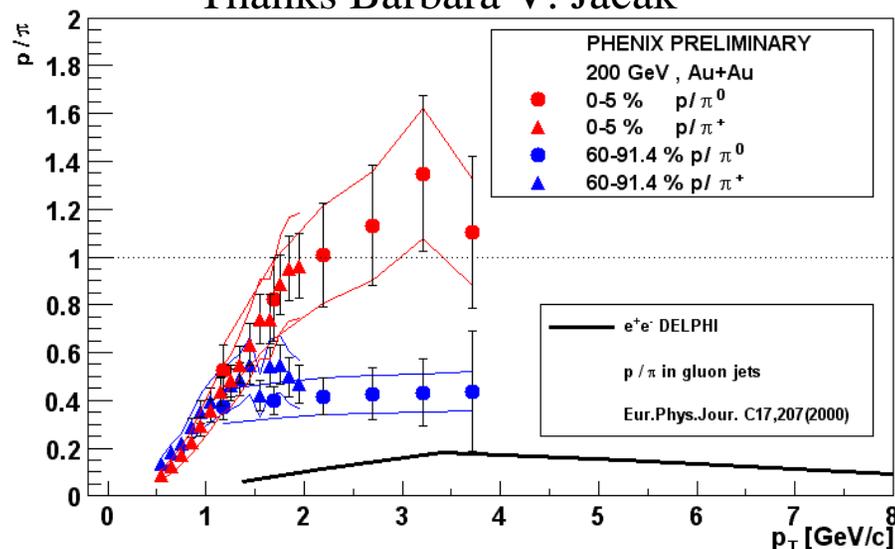


# Centrality and $p_T$ dependence of $p/\pi$

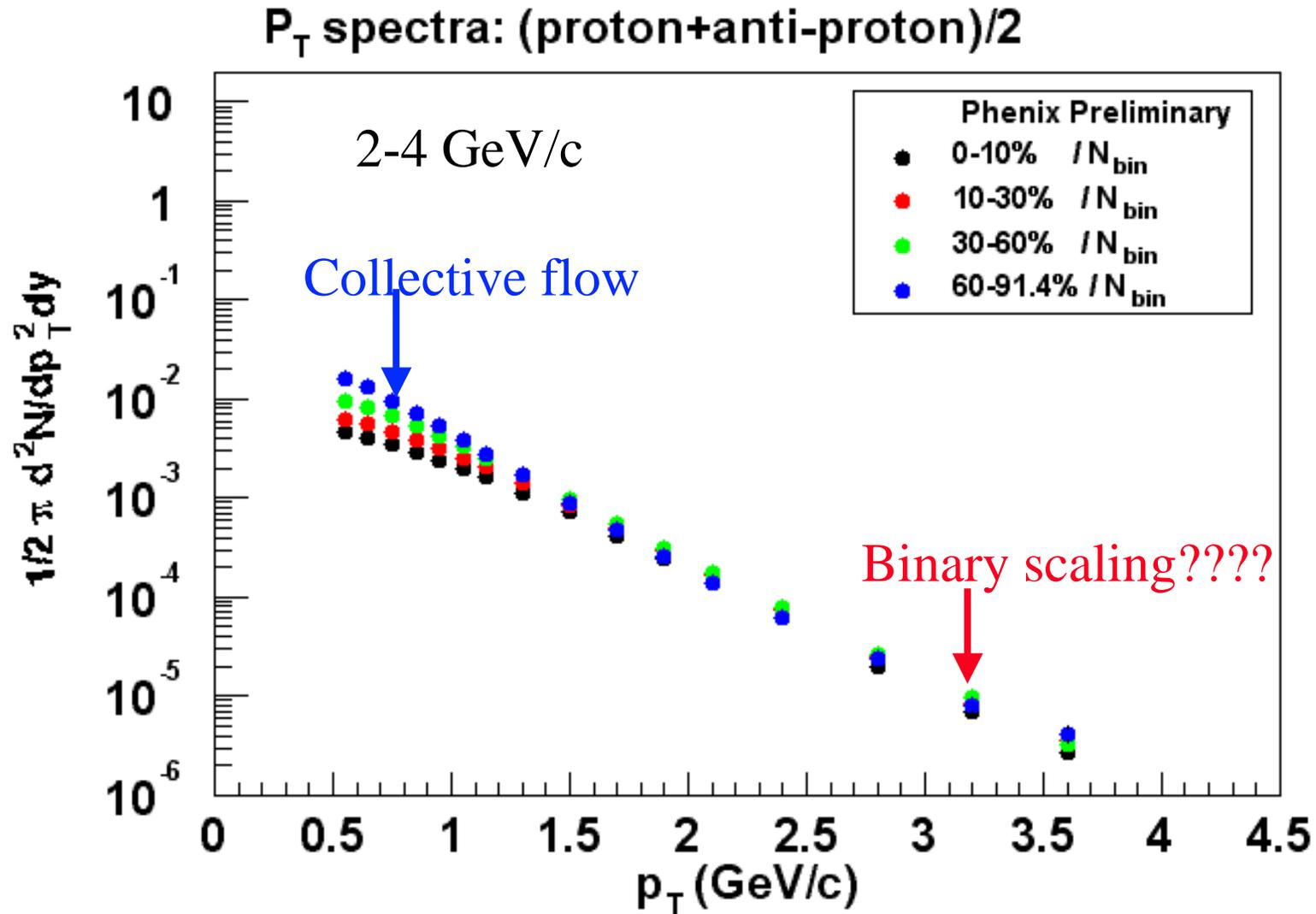
- $p/\pi$  ratio for AuAu at  $p_T > 2 \text{ GeV}$ 
  - $\sim 1$  in central collisions
  - $\sim 0.4$  in peripheral collisions
  - $\sim 20\%$  lower for anti-proton
- Above gluon jet value

If the observed anti(proton) are from fragmentation, then it is not standard fragmentation

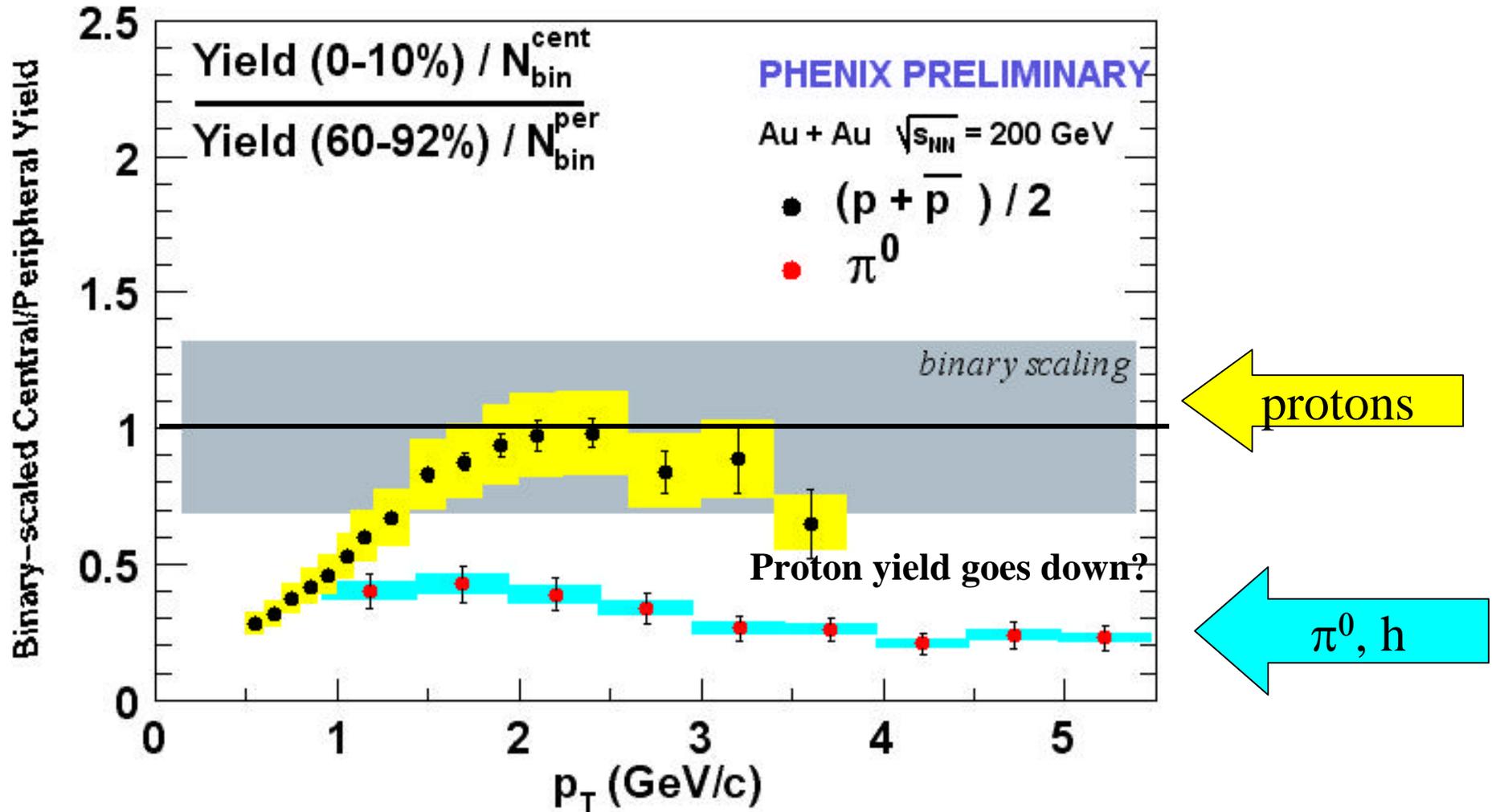
Thanks Barbara V. Jacak



# Scaling of the $\bar{p}$ p spectra

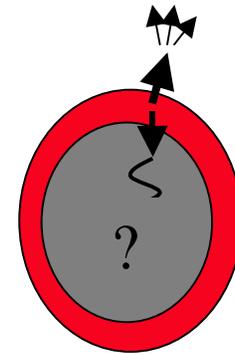


# Are proton and antiproton suppressed?



# What have we learned?

- **Strong centrality dependence of suppression,  $v_2$ , away side jets at high  $p_T$** 
  - Suggestive of surface emission of jets or mono-jet production?
  - Saturation(initial state), Jet quenching(final state)
- **Binary scaling of protons yield between 2-4 GeV/c**
  - Soft or hard or in between?
  - Baryon transport



Inspired by E. Shuryak Phys.Rev. C66 (2002) 027902

- **Which effect can be explained by jet absorption and collision geometry?**
  - Consistent with high  $p_T$  yield and di-jet suppression.
  - Insufficient to describe  $v_2$
  - Can't describe proton yield since there is no absorption

# Simple Toy Geometrical Model for Jet absorption

- Jets are absorbed in overlap region according to

$$f = e^{-k \int \rho dl}$$

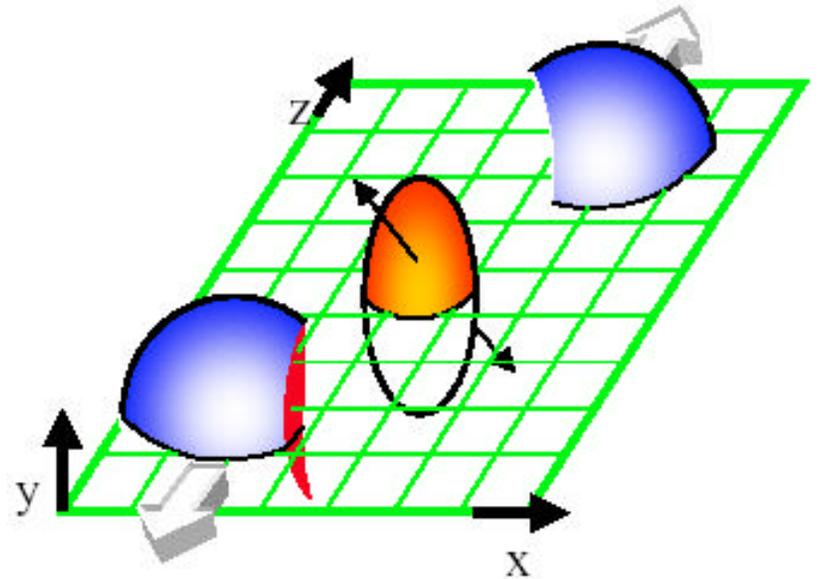
- Density of matter in transverse plane determined by participant density

$$\rho \sim \rho_{Npart}$$

- Azimuth isotropic di-jets production according to binary collision profile

$$\rho_{Ncoll}$$

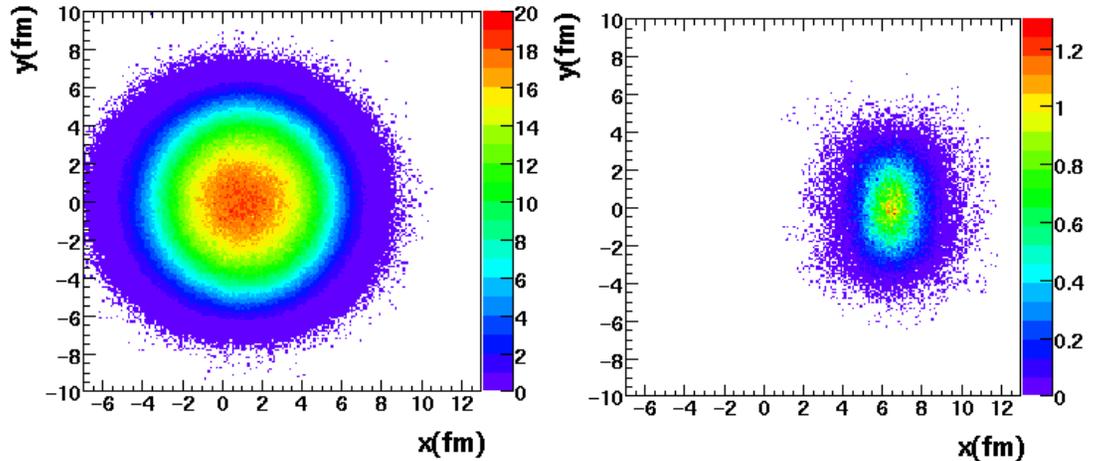
- Hadron spectra are related to jet distribution via parton-hadron duality



# Density distribution calculated from glauber model

## Wood-saxon nuclear density profile

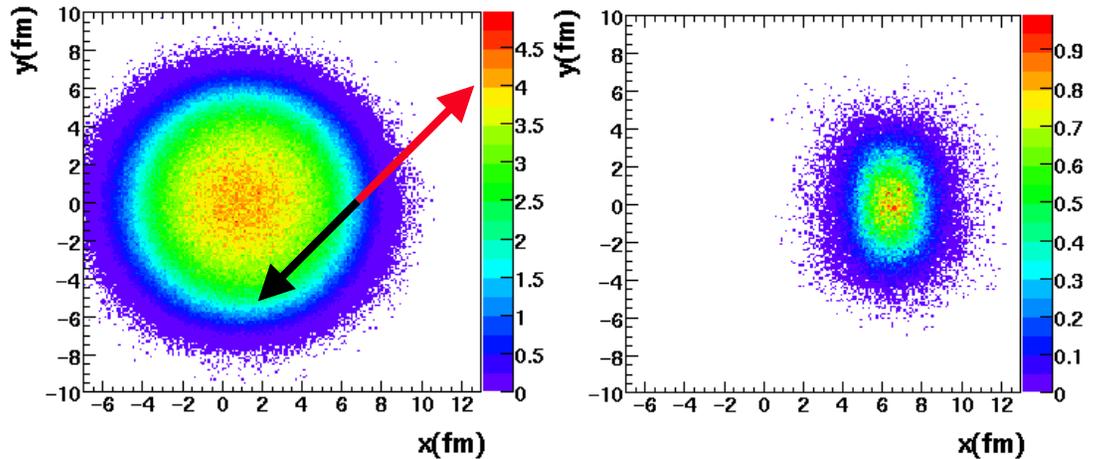
$\rho_{Npart}$



0-5%

75-80%

$\rho_{Ncoll}$



# Model Assumptions to Compare with Data

- **Model is simplistic and can only describes main features of geometry**

- Assume all particles above 4 GeV/c from jets
- Assume jets are back-to-back
- Ignore the fragmentation
- Contains no pT and  $\sqrt{s}$  and flavor dependence
- Try different kind of absorption :  $f = e^{-\kappa I_i}$
- $\kappa$  is the only free parameter

**Normal nuclear absorption**

$$I_1 = \int_0^{\infty} dl \rho(x, y)$$

**Energy loss style absorption**

$$I_2 = \int_0^{\infty} dl l \rho(x, y)$$

**Absorption + expansion,  $l_0 \sim 0.2$  fm**

$$I_3 = \int_0^{\infty} dl \frac{l_0}{l_0 + l} \rho(x, y)$$

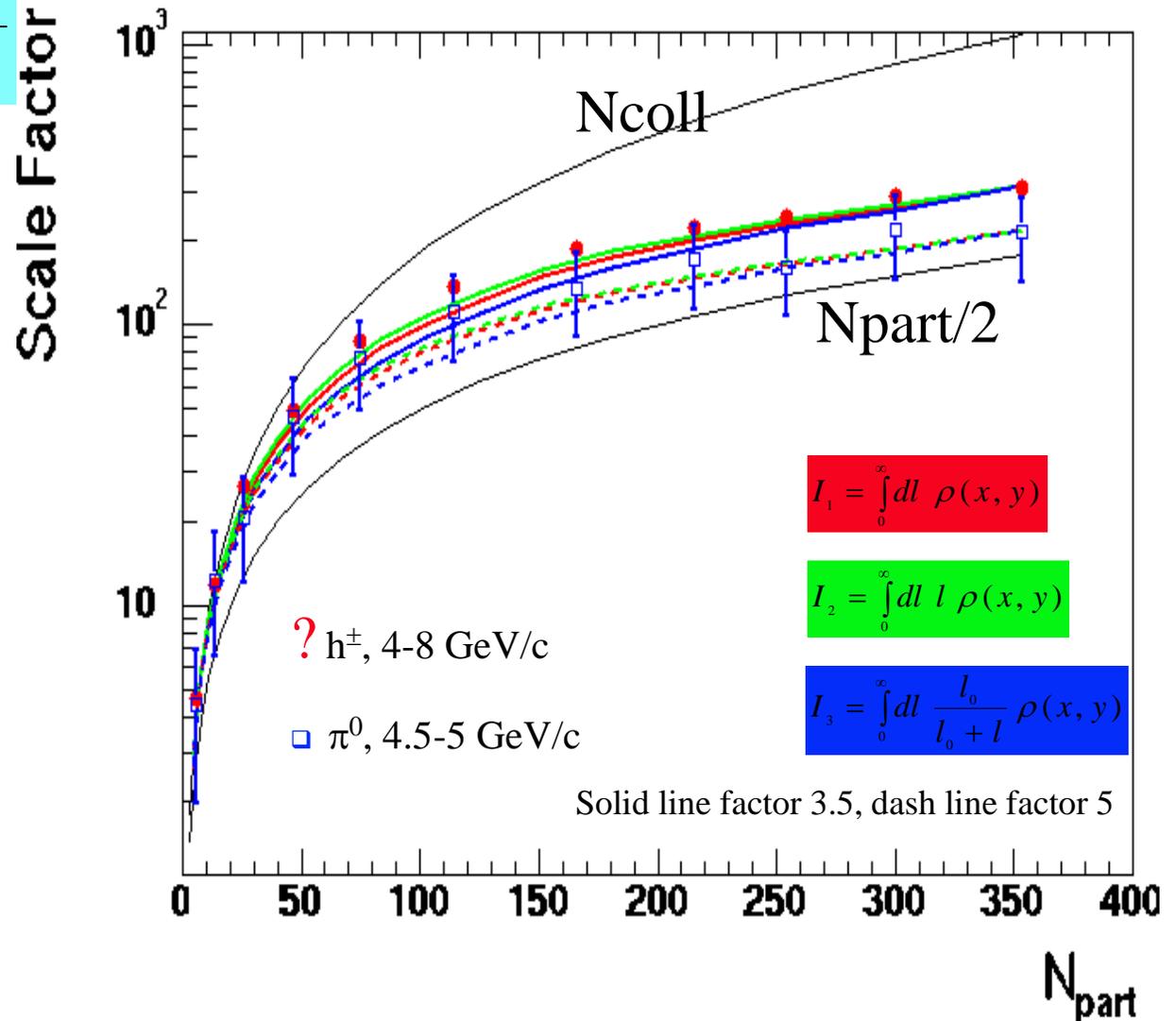
- **Adjust absorption strength  $\kappa$  to reproduce suppression factor of 3.5(charged) and 5( $\pi^0$ ) in most central collisions.**

- Predict centrality dependence of yield suppression
- Predict centrality dependence of di-jet suppression
- Predict centrality dependence of  $v_2$

## Scaling behavior of the yields

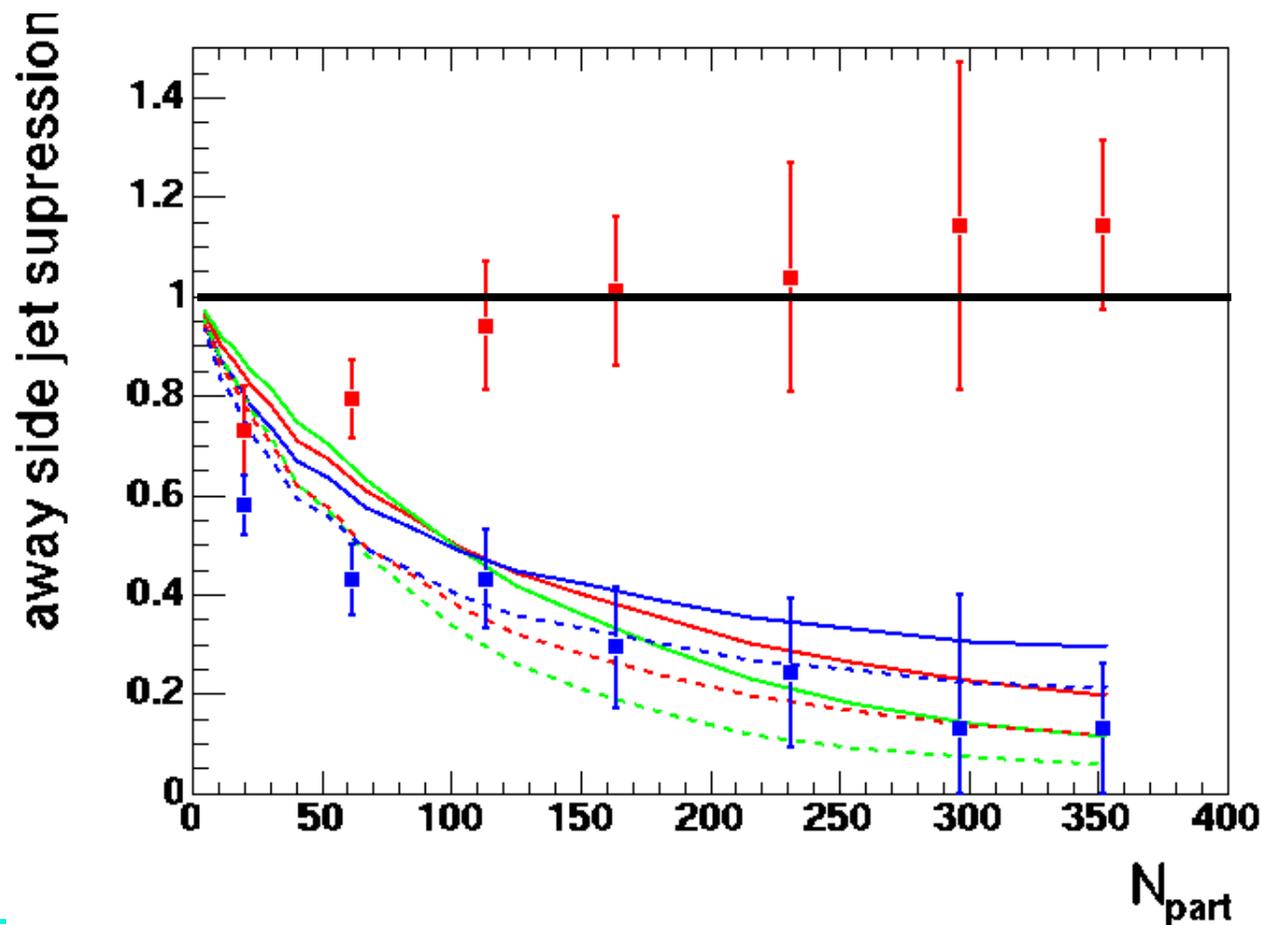
$$S_y(b) = \frac{y_{AuAu}(N_{part})}{y_{AuAu}(peri) \text{ or } y_{pp}}$$

- Qualitatively describe the centrality dependence of the yield
- Not sensitive to the absorption pattern



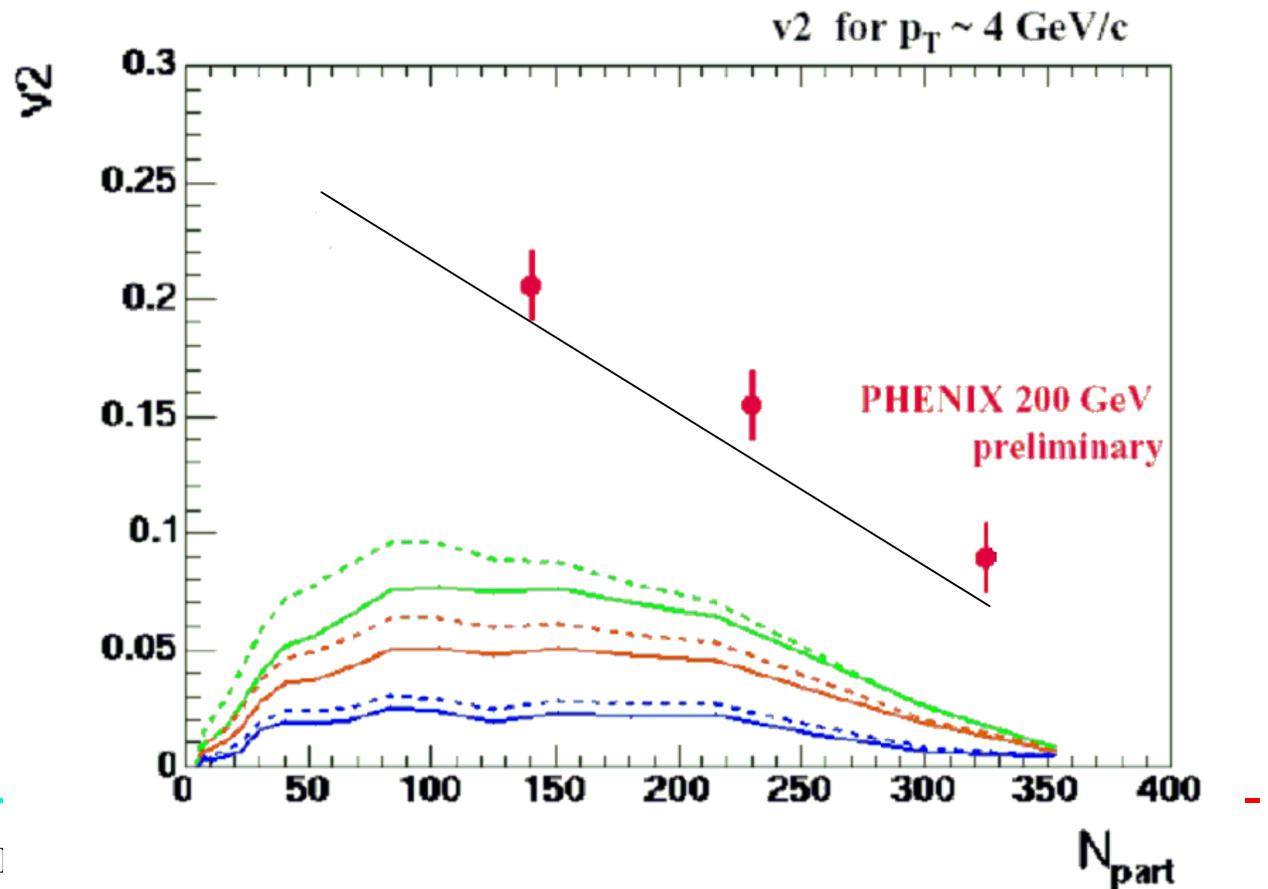
## Centrality dependence of back-back jets

- Compare STAR back-back jet measurement for  $4 < p_T < 6$  GeV/c
  - By construction, same side jet will always be 1
  - Qualitatively describe the centrality dependence of the away side



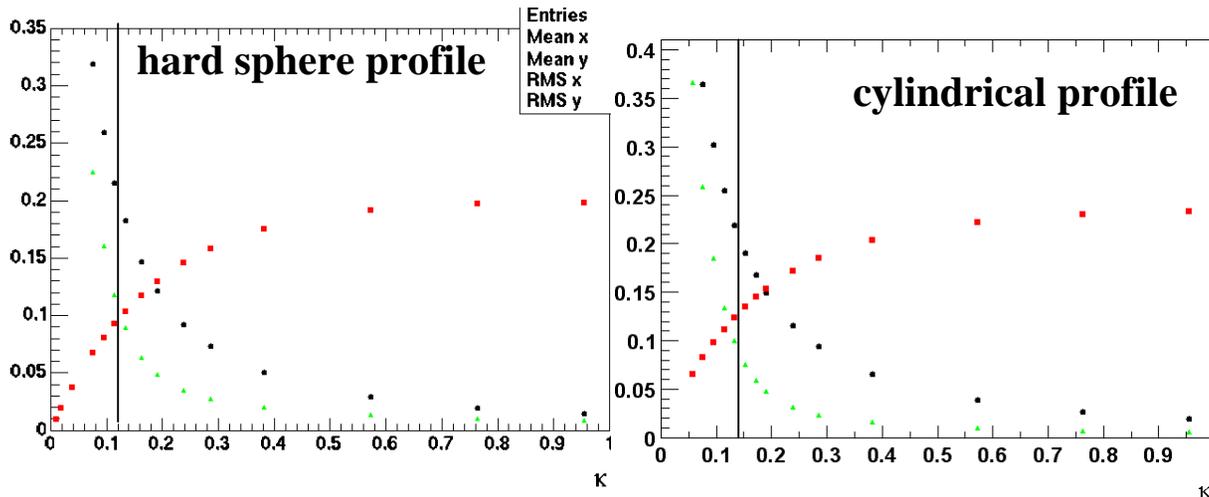
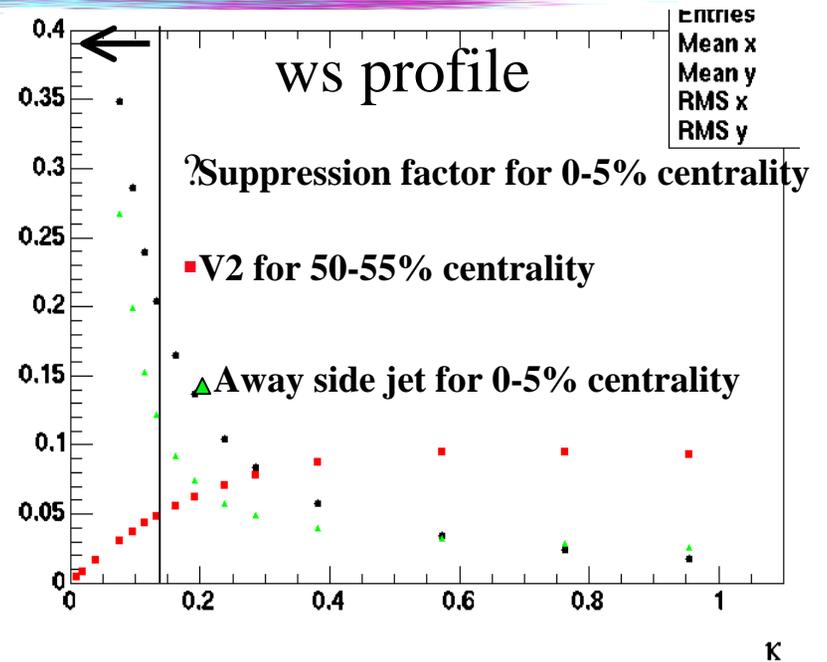
## Centrality dependence of $v_2$

- NB: in these models,  $v_2$  are purely geometrical origin, does not consider detailed modification of momentum distribution
- Geometrical anisotropy is insufficient to produce observed  $v_2$  by jet absorption
- $V_2$  is sensitive to absorption model, smallest for longitudinal expanding source
- $V_2$  will increase using hard sphere or cylindrical nuclear profile, but unphysical



## Varying $\kappa$

- Maximum  $v_2$  produce by surface.
- Different geometrical profile have different geometrical limit.
- $v_2$  is below geometrical limit at experimental observed suppression value



For sphere, surface volume is

$$\frac{4\pi R^2 d}{4/3\pi R^3} = \frac{3d}{R} \sim 26\%$$

S. Voloshin

cyl profile :  $v_{\max} = \sin(2\alpha) / 6\alpha$

# Conclusion

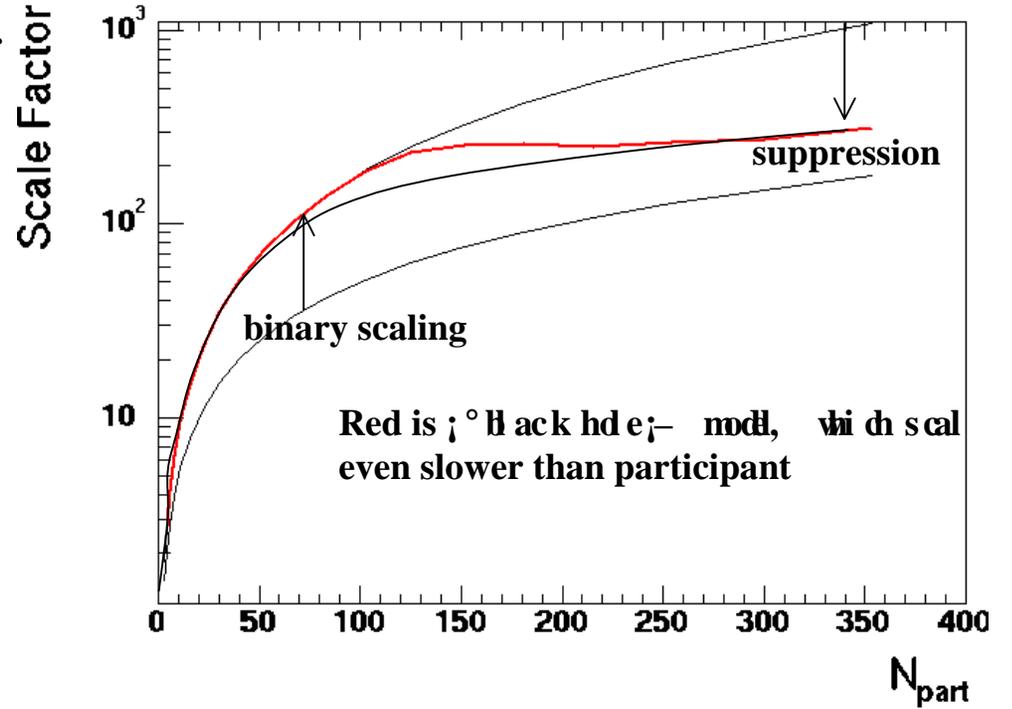
---

- **Exciting high pT results from 200GeV AuAu**
  - Direct observation of jets
  - Strong high pT suppression
  - **Puzzle1 : Large  $v_2$  at high pT**
  - **Puzzle2: High pT proton and antiproton**
    - Binary scaling between 2-4GeV/c
    - Mysterious particle composition at high pT.
- **Jet absorption gives characteristic centrality dependence of observables**
  - Consistent with high pT yield and di-jets suppression
  - Insufficient to describe  $v_2$  and proton yield

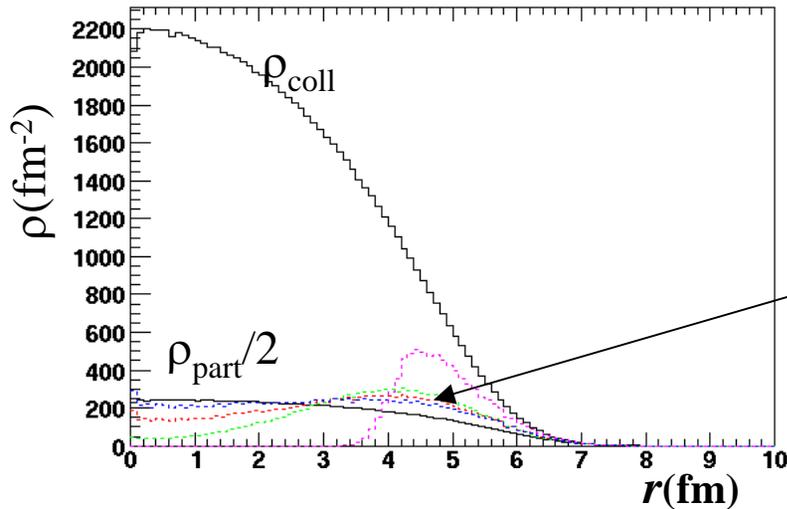
# Discussion on Scaling

- Absorption model naturally lead to stronger suppression in large  $\rho_{N_{coll}}$  region
  - Reduce binary scaling

$$\frac{\rho_{N_{coll}}}{\rho_{N_{part}}} \sim \sqrt{R^2 - r^2}$$



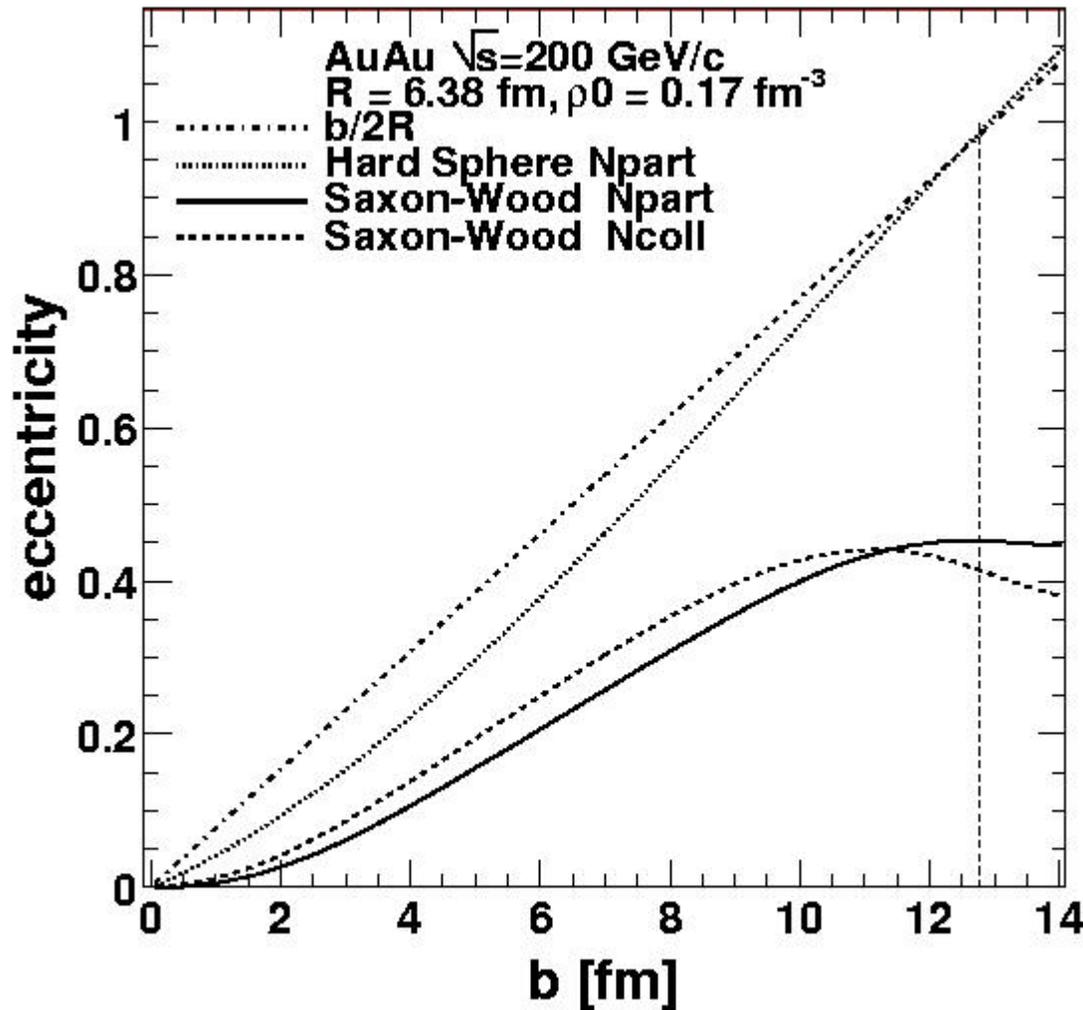
Density distribution for surviving jets ( $b < 3\text{fm}$ )



Suppression can lead to Npart scaling

# Hard sphere vs Saxon-Woods eccentricity

Made by Jan. Rak



$$\epsilon = \frac{\langle y^2 - x^2 \rangle}{\langle y^2 + x^2 \rangle}$$

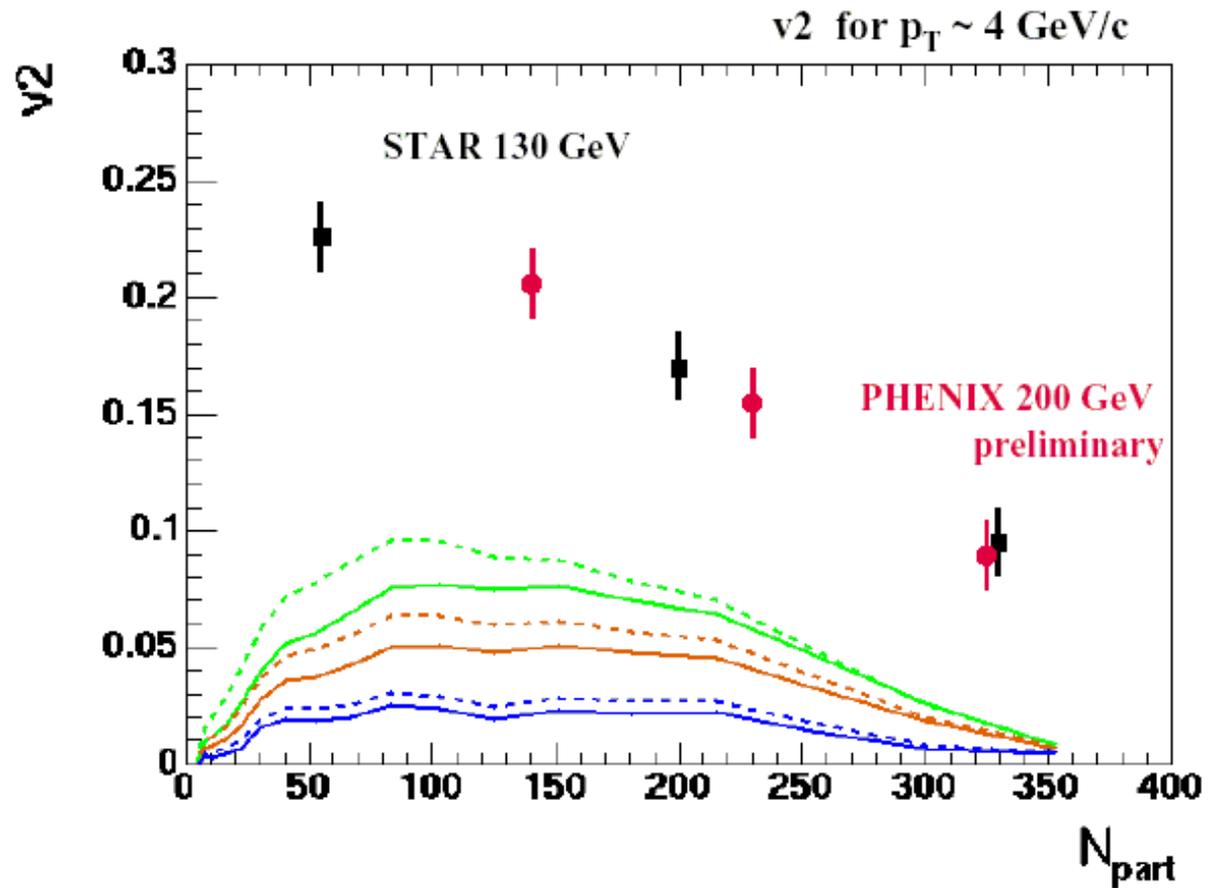
$$v_2(p_{\perp}) = \frac{\langle p_y^2 - p_x^2 \rangle}{\langle p_y^2 + p_x^2 \rangle}$$

- sheer almond  $\epsilon = b/2R$
- HS  $\epsilon \approx b/2R$
- Npart  $\epsilon \approx \sqrt{\text{cent}}/2$

HS over predict eccentricity by factor of 2.

Any meaningful geometrical estimate **has to be inferred from SW**

$$\begin{aligned}
P_{\text{esc}} &= \int_0^\infty dh \exp(-\rho\sigma h / \cos\theta) \propto \cos\theta|_{\sigma \rightarrow \infty} \\
v_2 &= \int d\phi P_{\text{esc}} \cdot \cos(2\phi) = \int_{-\alpha_m}^{\alpha_m} d\alpha \int_{-\pi/2}^{\pi/2} d\theta \cos\theta \cos 2(\theta + \alpha) / (2 \cdot 2\alpha) \\
&= \frac{\sin 2\alpha}{4\alpha} \int_{-\pi/2}^{\pi/2} d\theta \cos\theta (1 - 2\sin^2\theta) = \frac{\sin 2\alpha}{4\alpha} \left(2 - \frac{4}{2}\right) \\
&= \frac{\sin 2\alpha}{6\alpha}
\end{aligned}$$



# Some Simple Mathematics

- We look at particle yield above 4 GeV/c

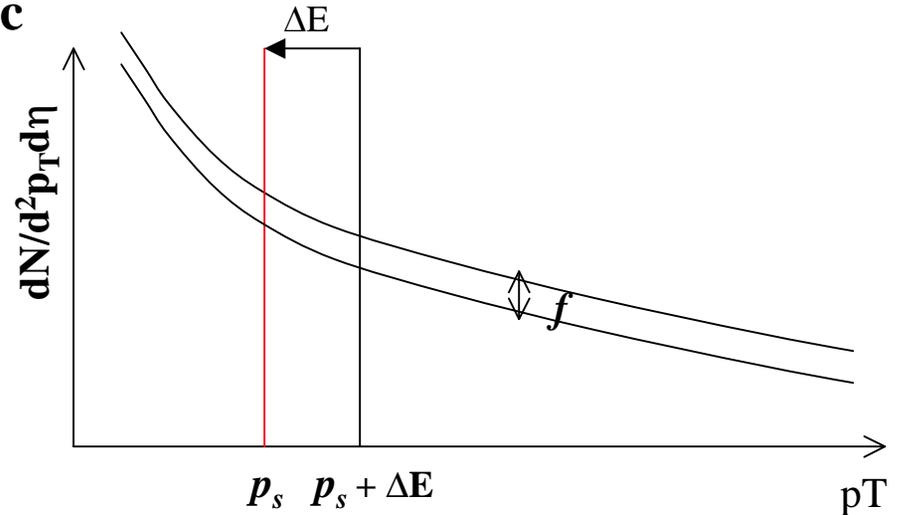
- Spectra have similar shape as pp

Assume spectra is:

$$\frac{dN}{d^2 p_T d\eta} \sim e^{-p_T/T}$$

Integrated yield above  $p_s$  is:

$$I(p_s) = \frac{e^{-p_s/T}}{T}$$



Assume energy loss  $\Delta E(p_T)$  is independent of  $p_T$ :

$$I_1(p_s) = \frac{e^{-(p_s+\Delta E)/T}}{T} = f I(p_s), \text{ where } f^{-1} = e^{\Delta E/T} \text{ is suppression factor}$$

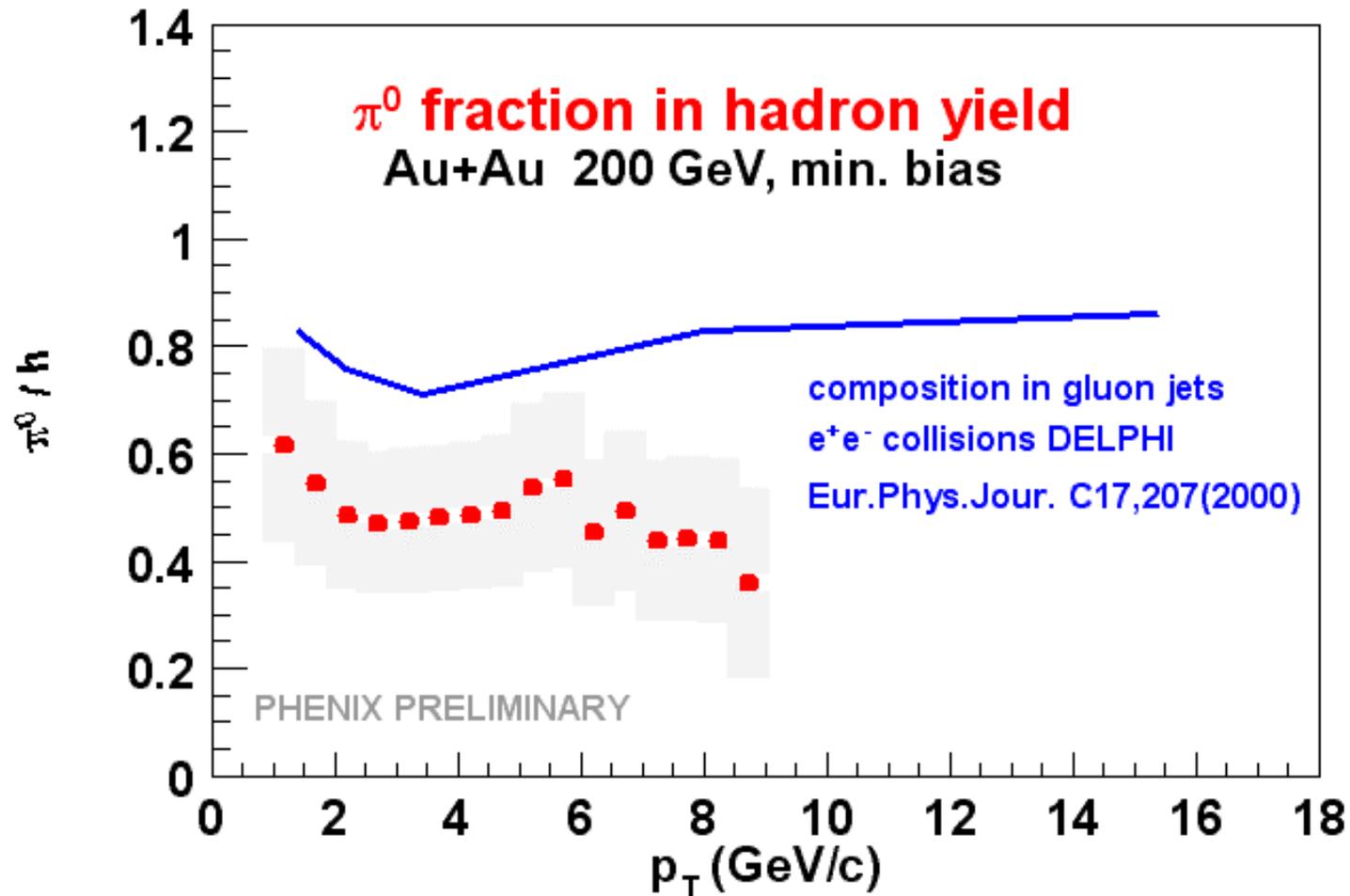
Left shift(energy loss) is effectively a down shift(absorption), one can obtain the similar result

$$f = e^{-k \int \rho dl} \quad \mathbf{E. Shuryak} \quad \kappa \rightarrow 1/T$$

For power law shape, assumption is good for small  $\Delta E$ ,  $\Delta E \leq 1/T \sim 1 \text{ GeV}$  for  $p_T=10 \text{ GeV/c}$

Hadron spectra and jet spectra are related via parton-hadron duality

## $\pi/h$ at higher $p_T$



# STAR yield per trigger

