

Some Trigger Oldies but Goldies
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♡ For Matthias' trigger question, I want an E_T trigger which doesn't suffer distortion from 'fragmentation' or hadronization, compared to a cluster counting trigger with a low threshold which I think is better done in the mvd.

♡ As for EMCAL thresholds, normally I would use 4:

- 1) highest for single cluster = γ or π^0
- 2) lower for back-to-back π^0 or γ pairs
- 3) lowest EMCluster seed in each arm for e^\pm pair trigger, but we know this threshold, it must be 1 GeV or less to see the J/Ψ
- 4) 1.2 to 1.5 GeV EMCluster threshold as a seed for the single e^\pm trigger for charm. Again, we know this threshold.

Maybe the trigger 2 can have the same threshold as trigger 4, but it is unlikely. Alternatively, if we are lucky we can use trigger 3) threshold for singles also.

♡ I had discussed an electron trigger algorithm in [PN85 6/10/93](#). It matches clusters in the EMCAL with an energy $E > E_{\text{thresh}}$ with rings in the RICH. The only complication is that the EMcal measures η , ϕ while the RICH measures η and $\lambda = \delta + \phi$, where δ is the local tangent angle of the trajectory relative to the radial direction, also called α , β (???) in PHENIX. The good news is that since the EMCAL gives you the E of the cluster in question you know $\delta = \lambda_{\text{RICH}} - \phi_{\text{EMCAL}}$.

1.1 Track from the origin

For a particle emitted from the origin $\tau = 0$ and $\lambda_N = \phi_0$, the azimuthal angle of emission:

$$\sin \delta = \frac{R}{2\rho} \quad (5)$$

$$\sin(\phi - \phi_0) = \frac{R}{2\rho} = \sin \delta \quad (6)$$

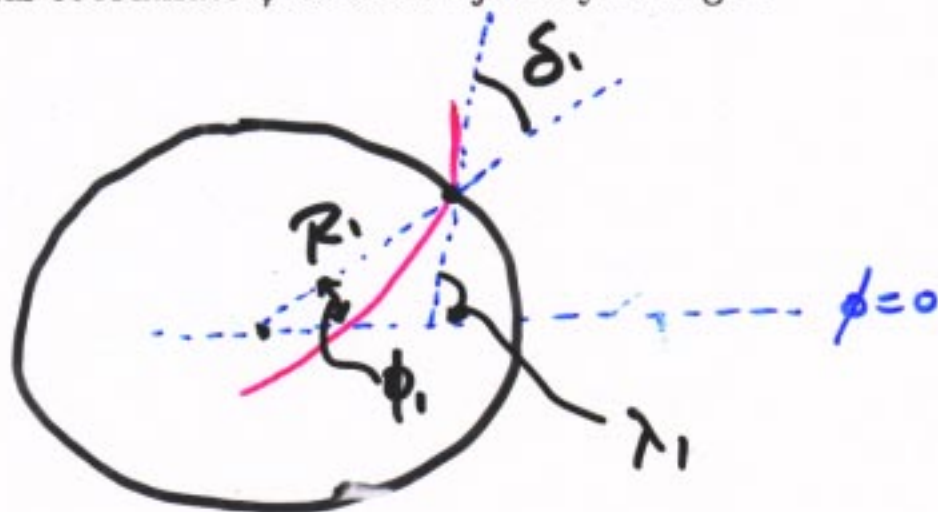
1.2 Field-free region

These equations work just as well for a field free region, inner radius R_i , outer radius R_o in which case $\rho_{i0} = \infty$, with the result

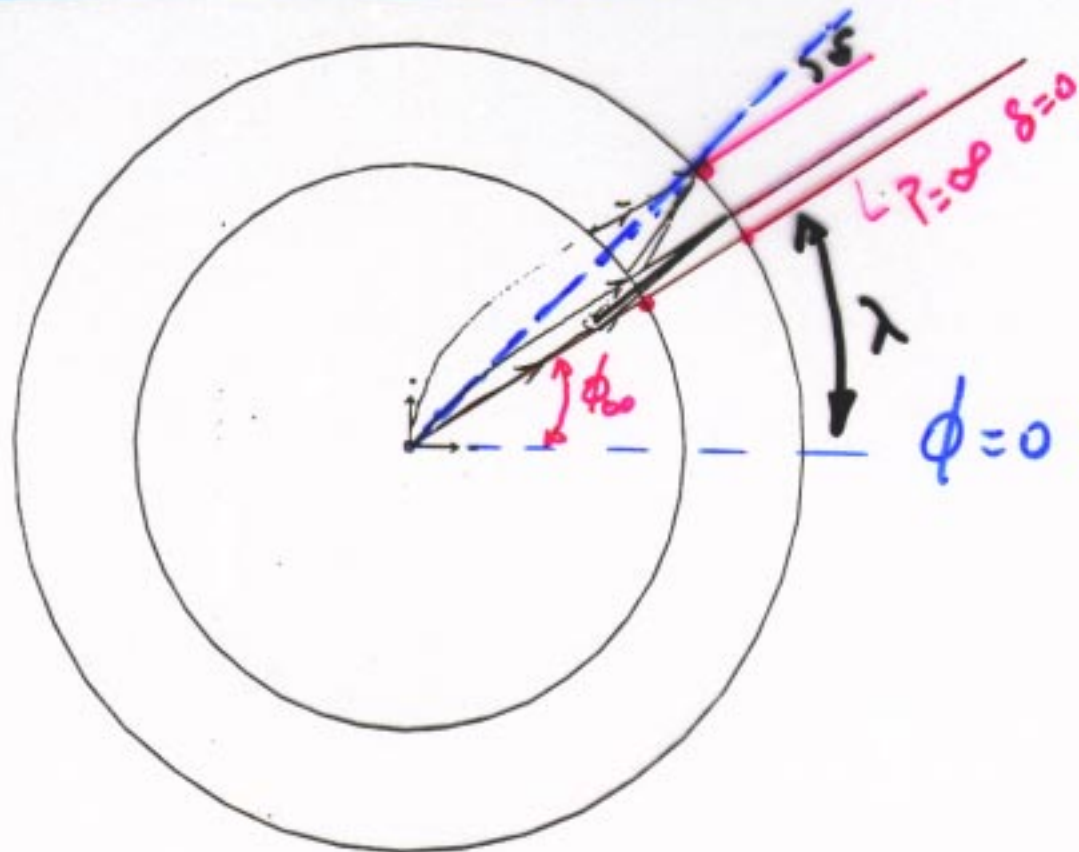
$$\sin \delta_o = \frac{R_i}{R_o} \sin \delta_i \quad (7)$$

$$\phi_o - \phi_i = \delta_i - \sin^{-1}\left(\frac{R_i \sin \delta_i}{R_o}\right) \quad (8)$$

Note that the trajectory vector does not change in the field free region, but the local tangent angle δ to the radius changes because the azimuthal coordinate ϕ of the trajectory changes.

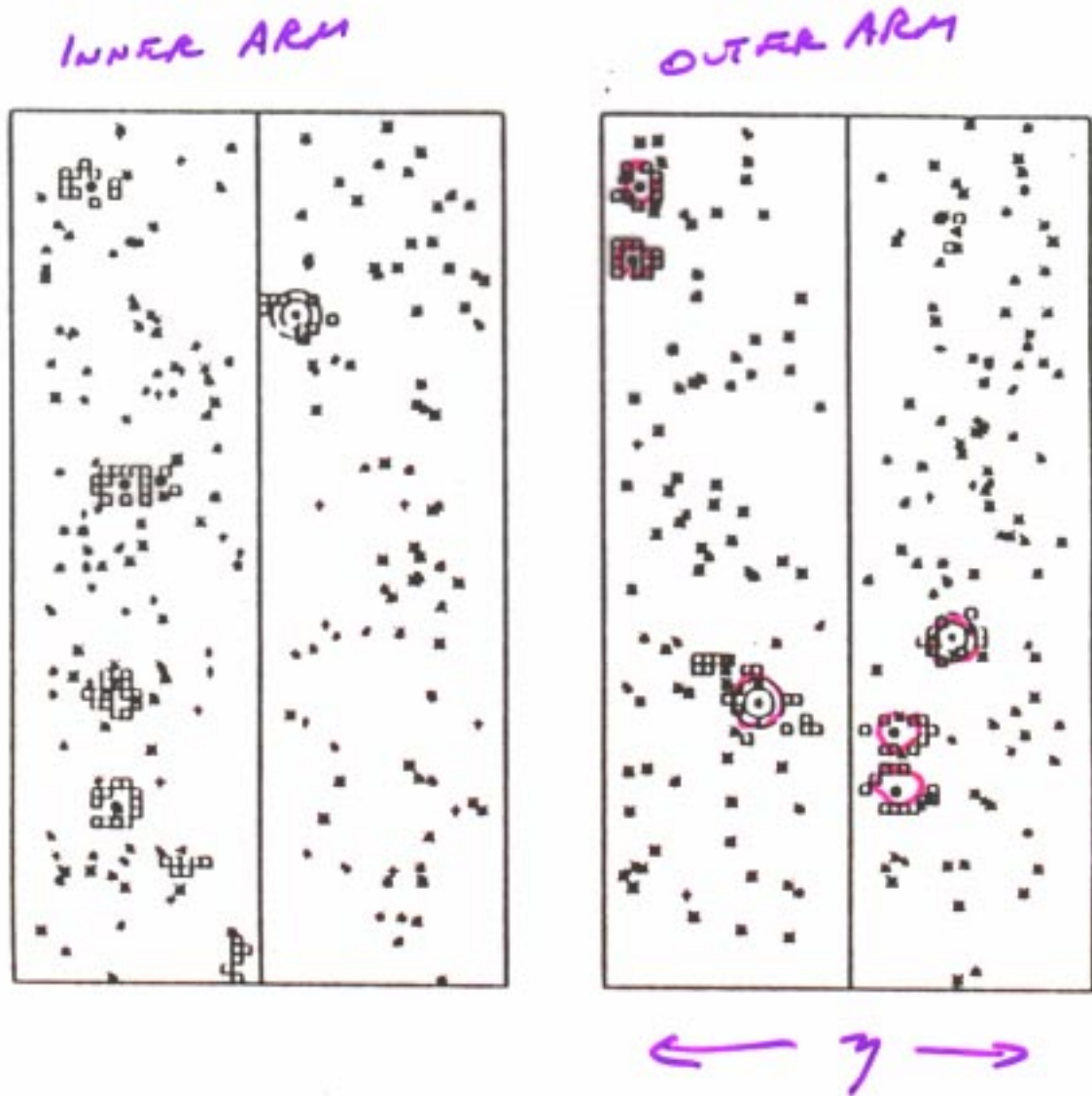


RICH IMAGES $\lambda \neq \theta$ NOT $\phi \neq \theta$



- ALL 3 TRAJECTORIES IMAGE AT THE SAME POINT ϕ_{∞} = AZIMUTHAL COORDINATE AT RICH FOR INFINITE MOMENTUM TRACK
- 3 TRAJECTORIES DRAWN ALL HAVE THE SAME $\lambda = \phi_{\infty}$
- FOR TRACKS WITH THE SAME $\lambda = \phi_{\infty}$
 $\phi_{\infty} - \phi = \delta = \sin^{-1} R/2p$

Figure 2.43: A typical hit pattern on the Cherenkov detector for a simulated Au + Au central collision event for a charged particle multiplicity of $dN_c/dy = 1000$. The hit pixel in the



$$\lambda = \phi_{\infty}$$

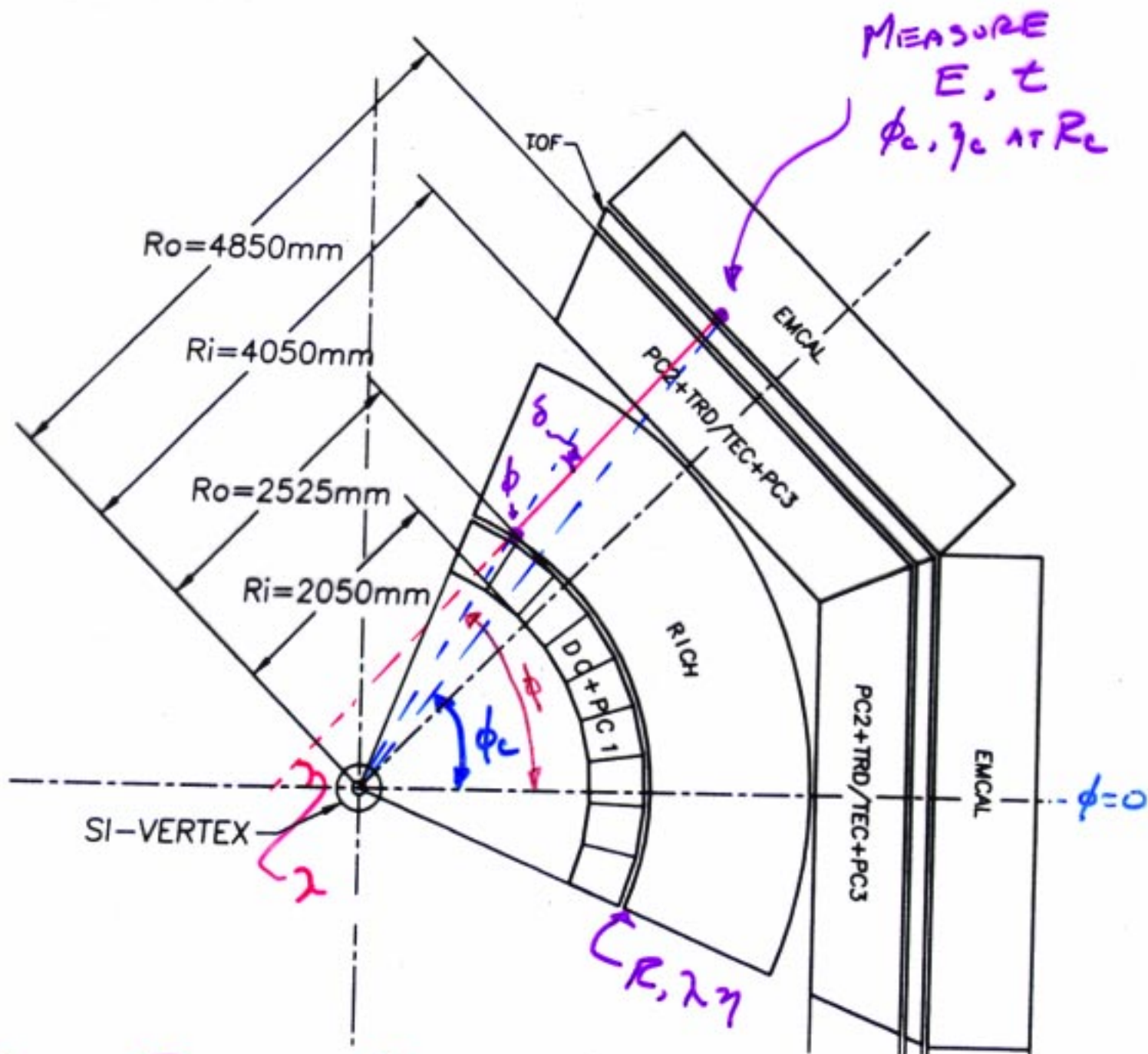
$$\gamma_{RICH} \approx \gamma_{CAL}$$

$$\phi_{CAL} = \lambda_{RICH} - \lambda_{cut}^{-1} \left(\frac{P_{cut} \Delta S}{R_c} \right)$$

IDEALLY

$$\Delta S = \frac{R}{z_p} = \frac{P_{CUTOFF}}{P_T}$$

$$P_{CUTOFF} = \left(2 \cdot \frac{R}{z} \right) \times \dots$$



ASSUME R IS AT THE BOUNDARY OF B OR IN FIELD-FREE REGION
 Figure 6.1: Schematic of PHENIX detector including the tracking subsystem.

$$\phi_c - \phi = \delta - \sin^{-1}\left(\frac{R \sin \delta}{R_c}\right) \quad \lambda = \phi + \delta$$

$$\boxed{\phi_c - \lambda = \sin^{-1}\left(\frac{R \sin \delta}{R_c}\right)} \quad \sin \delta \sim \frac{p_T^{\text{cutoff}}(R)}{p_T}$$

$\rightarrow 0$
 AS p_T INCREASES

MJT E-trigger strategy PN85 6/10/93

- Match λ , η hits on RICH
- To η , ϕ_c hits on EMCal
 - within 1 ns of $\beta = 1$
 - charged only
 - for different slices of E
 - or for $E > E_{\text{thresh}}$
- Then Match p and E when tracks are available
- Then Reject Dalitz & Conversions
 - when HBD is available
 - or with low magnetic field setting
- Also Select Dalitz and Conversions!
 - see 'Charm in PHENIX—a signal or a background?'
 - <http://www.phenix.bnl.gov/~sapi/charm96.ps.gz>
 - see many experts largely concentrated in light/heavy PWG