

# $J/\psi$ production in Cu+Cu and Au+Au collisions at RHIC-PHENIX

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**Abstract.** The  $J/\psi$  is considered to be among the most important probes for the deconfined quark gluon plasma (QGP) created by relativistic heavy ion collisions. While the  $J/\psi$  is thought to dissociate in the QGP by Debye color screening, there are competing effects from cold nuclear matter (CNM), feed-downs from excited charmonia and bottom quarks, and regeneration from uncorrelated charm quarks. Measurements that can provide information to disentangle these effects are presented in this paper.

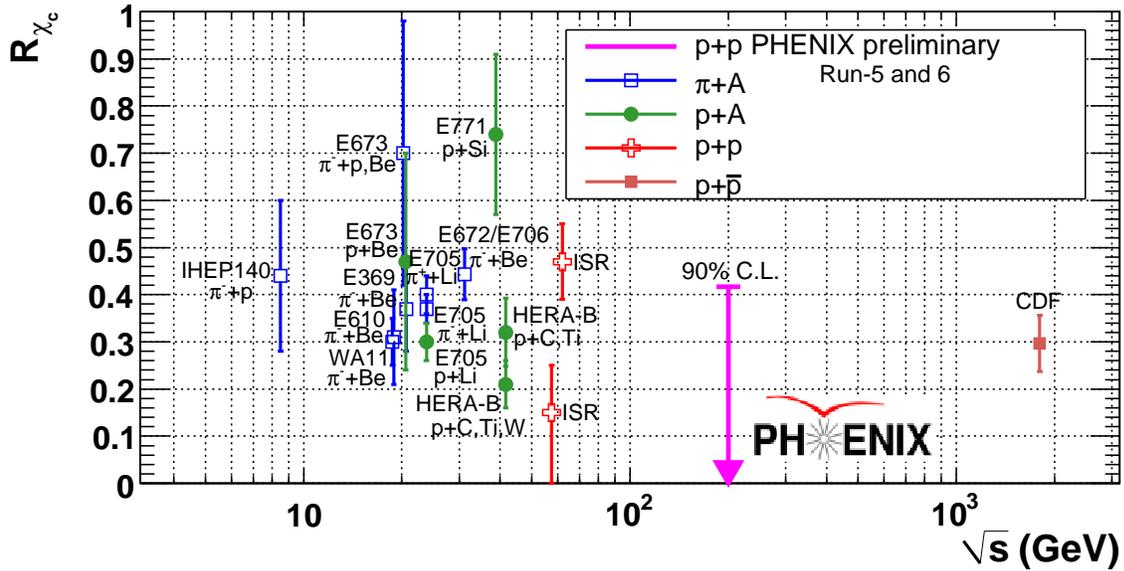
## 1. Introduction

The yield of heavy quarkonia is expected to be suppressed in the QGP due to the Debye screening of the color charge [1]. The  $J/\psi$  is especially promising because of its large production cross section and di-lepton decay channels, which make it easily detected. The PHENIX experiment at RHIC is able to detect the  $J/\psi$  at midrapidity ( $|y| < 0.35$ ) via its decay to  $e^+e^-$  and at forward rapidity ( $1.2 < |y| < 2.2$ ) via its  $\mu^+\mu^-$  decay. Models of  $J/\psi$  production in heavy ion collisions at RHIC energy contain a number of important competing effects, including modification of the  $J/\psi$  yield by the CNM effects, destruction of  $J/\psi$  due to interactions with thermal gluons in the QGP, reduced feed-down from excited charmonium states that melt just above the QGP transition temperature, bottom quark decay and enhancement of the yield due to coalescence of uncorrelated charm pairs [2]. The PHENIX Au+Au data at  $\sqrt{s_{NN}} = 200$  GeV showed that  $J/\psi$  suppression at forward rapidity is larger than that at midrapidity and the suppression at midrapidity is similar to that observed by NA50 at SPS in Pb+Pb collisions at  $\sqrt{s_{NN}} = 17.3$  GeV [3]. However, these results are not well understood theoretically. Systematic study of  $J/\psi$  production in heavy ion collisions across the entire range of  $N_{part}$  is needed to disentangle the competing effects. PHENIX recorded Cu+Cu collisions at  $\sqrt{s_{NN}} = 200$  GeV to obtain precise data in the range  $N_{part} \leq 126$ , where Au+Au data is limited by statistics and systematic uncertainty [4]. Feed-downs into  $J/\psi$  in  $p + p$  collisions were measured at the same energy and are important in the picture of the sequential dissociation of quarkonia. The elliptic flow of  $J/\psi$  can set a constraint on coalescence of charm quarks and the measurement is reported in [5].

## 2. Measurement of feed-downs into $J/\psi$ in $p + p$ collisions

An inclusive  $J/\psi$  measurement in  $p + p$  collisions at  $\sqrt{s}=200$  GeV has been reported by PHENIX [6]. Recently, PHENIX separately measured the feed-down contributions of  $\chi_c$ ,  $\psi'$  and bottom quarks to  $J/\psi$  in  $p + p$  collisions at midrapidity ( $|y| < 0.35$ ) at the same energy. The preliminary results are described in this section.

The fraction of  $J/\psi$  from the  $\chi_c$  decay,  $R_{\chi_c}$ , was measured via the decay chain of  $\chi_c \rightarrow J/\psi + \gamma \rightarrow e^+e^- \gamma$  with  $4.1 \times 10^3$  reconstructed  $J/\psi$ . Branching ratios of  $\chi_{cJ}$  to  $J/\psi + \gamma$  are 1.3% ( $\chi_{c0}$ ), 36% ( $\chi_{c1}$ ) and 20% ( $\chi_{c2}$ ) [7], and the contribution of  $\chi_{c0}$  was neglected in the analysis. The 90% confidence level upper limit obtained for  $R_{\chi_c}$  is 0.42. Figure 1 shows the upper limit of  $R_{\chi_c}$  with results from other experiments.

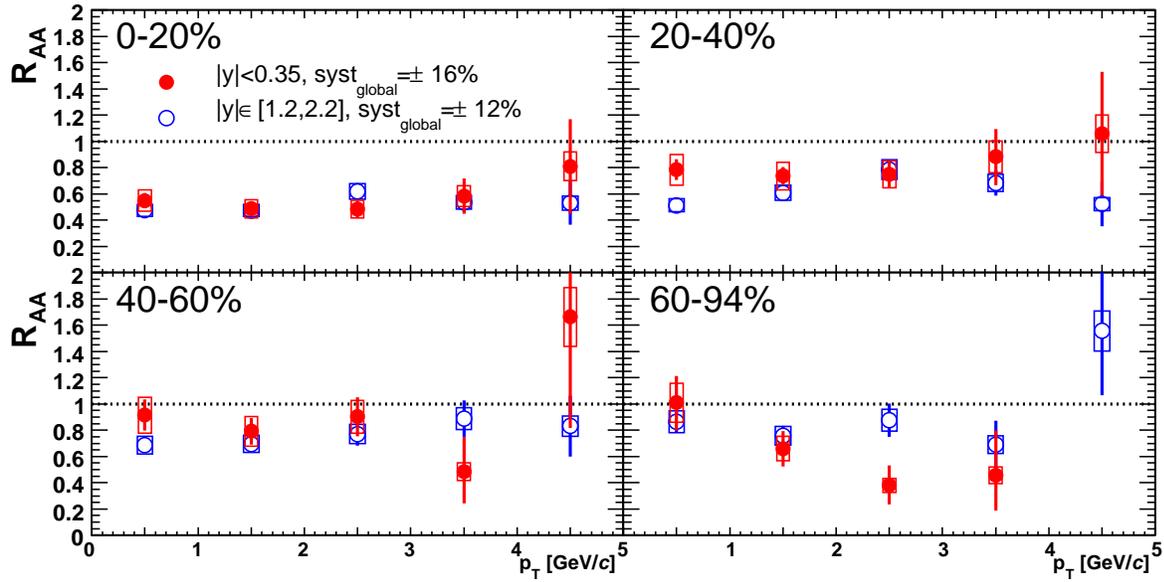


**Figure 1.** The fraction of  $J/\psi$  from the  $\chi_c$  decay,  $R_{\chi_c}$ , as a function of the center of mass energy,  $\sqrt{s}$ . The 90% confidence level upper limit of  $R_{\chi_c}$  in  $p + p$  collisions at  $\sqrt{s}=200$  GeV which is the PHENIX preliminary result is represented by the arrow.

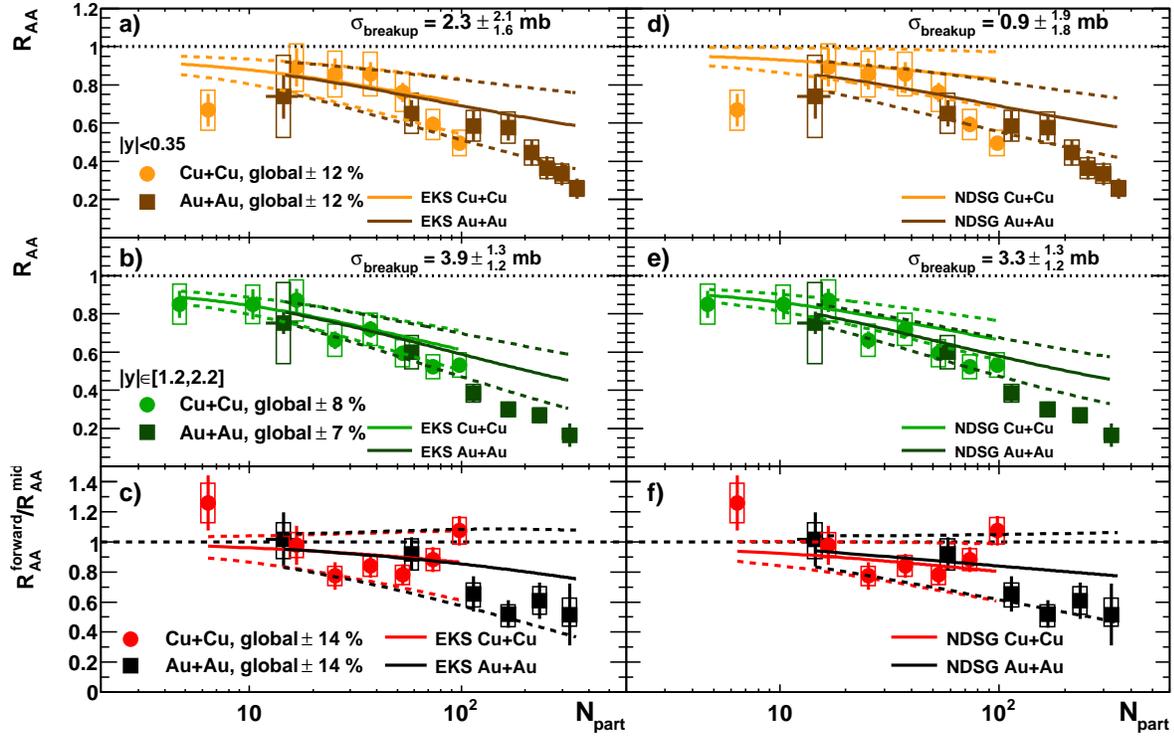
The ratio of cross sections of  $\psi'$  to  $J/\psi$  was measured in the  $e^+e^-$  decay mode and the feed-down fraction of  $J/\psi$  from  $\psi'$  decay is  $0.086 \pm 0.025$ . The cross section of bottom quarks was measured by electron-hadron correlation [8] and the feed-down fraction of  $J/\psi$  from decay of bottom and anti-bottom quarks is  $0.036^{+0.025}_{-0.023}$  for  $p_{T,J/\psi} > 0$  GeV/c. Branching ratios of [7] were used to obtain these feed-down fractions. A set of theoretically predicted feed-down fractions [9] is almost consistent with the PHENIX preliminary results.

## 3. Measurement of $J/\psi$ in $Cu+Cu$ and $Au+Au$ collisions

The final results of the  $J/\psi$  measurement in  $Cu+Cu$  collisions are reported here and in [4] with the published results in  $Au+Au$  collisions [3]. The numbers of observed



**Figure 2.**  $R_{AA}$  vs transverse momentum  $p_T$  for  $J/\psi$  production in Cu+Cu collisions via  $e^+e^-$  decay at midrapidity (closed circles) and  $\mu^+\mu^-$  decay at forward rapidity (open circles). Centrality bins are written in figures.



**Figure 3.** a)  $R_{AA}$  vs  $N_{part}$  for  $J/\psi$  production in Cu+Cu and Au+Au collisions at midrapidity with the prediction curves [13] with the EKS shadowing model [14]. b) The same figure at forward rapidity. c) Forward/mid rapidity  $R_{AA}$  ratio. d, e, f) The same figures with the prediction curves with the NDSG shadowing model [15]. The breakup cross sections for the prediction curves are obtained from fits to the  $d$ +Au data for each rapidity bin.

$J/\psi$  in Cu+Cu collisions at mid and forward rapidity are about 2050 and 9000, respectively. Compared to the Au+Au data, the Cu+Cu data provides a more precise  $R_{AA}$  measurement for  $N_{part} < 100$ , where the CNM effects might be dominant. Figure 2 shows  $R_{AA}$  as a function of  $p_T$  for each centrality bin in Cu+Cu collisions at mid and forward rapidity. Suppression by a factor of two is seen at both mid and forward rapidity in the most central collisions. No strong  $p_T$  dependence of  $R_{AA}$  is observed. Analysis is ongoing to extend  $R_{AA}$  beyond  $p_T=5$  GeV/c, where the STAR experiment has a moderate acceptance [10].

Figure 3 shows  $R_{AA}$  at mid and forward rapidity and the forward/mid rapidity  $R_{AA}$  ratio in Cu+Cu and Au+Au collisions. The Cu+Cu and Au+Au data agree well in the overlap region. The results of the  $J/\psi$  measurement in  $d$ +Au collisions, which is the primary tool to investigate the CNM effects, were updated and are reported in [11, 12]. PHENIX has interpreted the  $d$ +Au data and extrapolated prediction of the CNM effects to heavy ion collisions based on calculations of R. Vogt including nuclear shadowing and nuclear breakup [13]. While the nuclear breakup cross sections of  $J/\psi$ ,  $\sigma_{breakup}$ , does not vary with rapidity in the framework of the calculations, PHENIX treated  $\sigma_{breakup}$  as a rapidity dependent “effective” parameter and determined it by fitting to the  $d$ +Au data [4]. The theoretical prediction curves with  $\pm 1\sigma$  bands obtained for the rapidity dependent “effective”  $\sigma_{breakup}$  with two nuclear shadowing models (EKS [14] and NDSG [15]) are also shown in figure 3. The expected CNM effects from the PHENIX “effective” model at forward rapidity are larger than those at midrapidity. The  $J/\psi$  suppression below CNM effects seems to start at  $N_{part} \sim 200$  (100) at mid (forward) rapidity. Rapidity narrowing in central Au+Au collisions is consistent with the expected CNM effects.

While the current uncertainty of the CNM effects is large, it is expected to be reduced by the  $d$ +Au collision data collected in 2007–2008 whose statistics are 30 times larger than the currently available  $d$ +Au statistics.

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