



# $k_T$ Asymmetry in Longitudinally Polarized $p+p$ Collisions at PHENIX

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# Origin of Jet $k_T$

Net transverse momentum of a di-hadron pair in a factorization ansatz

$$\langle k_T^2 \rangle_{di-hadron} = \langle k_T^2 \rangle_{initial} \oplus \langle k_T^2 \rangle_{hard} \oplus \langle k_T^2 \rangle_{frag}$$



- *initial* – e.g., fermi motion of the confined quarks or gluons, initial-state gluon radiation...
- *hard* – final-state radiation (three-jet)...
- *frag* – fragmentation transverse momentum  $j_T$

## Another Possibility:

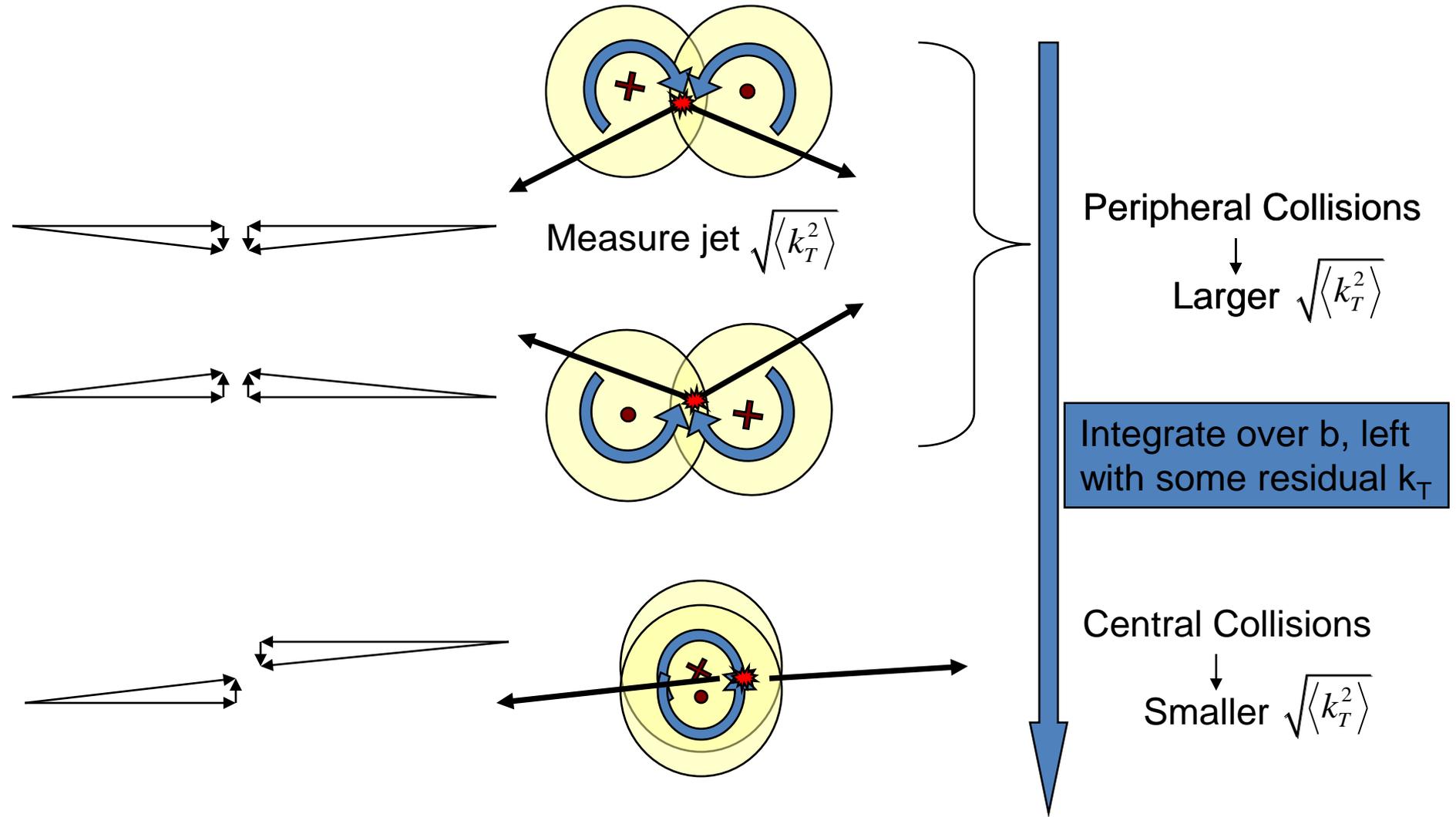
- Spin-correlated transverse momentum - partonic orbital angular momentum
- Can examine the helicity dependence of each term:

$$\Delta \langle k_T^2 \rangle_{initial} \neq 0$$

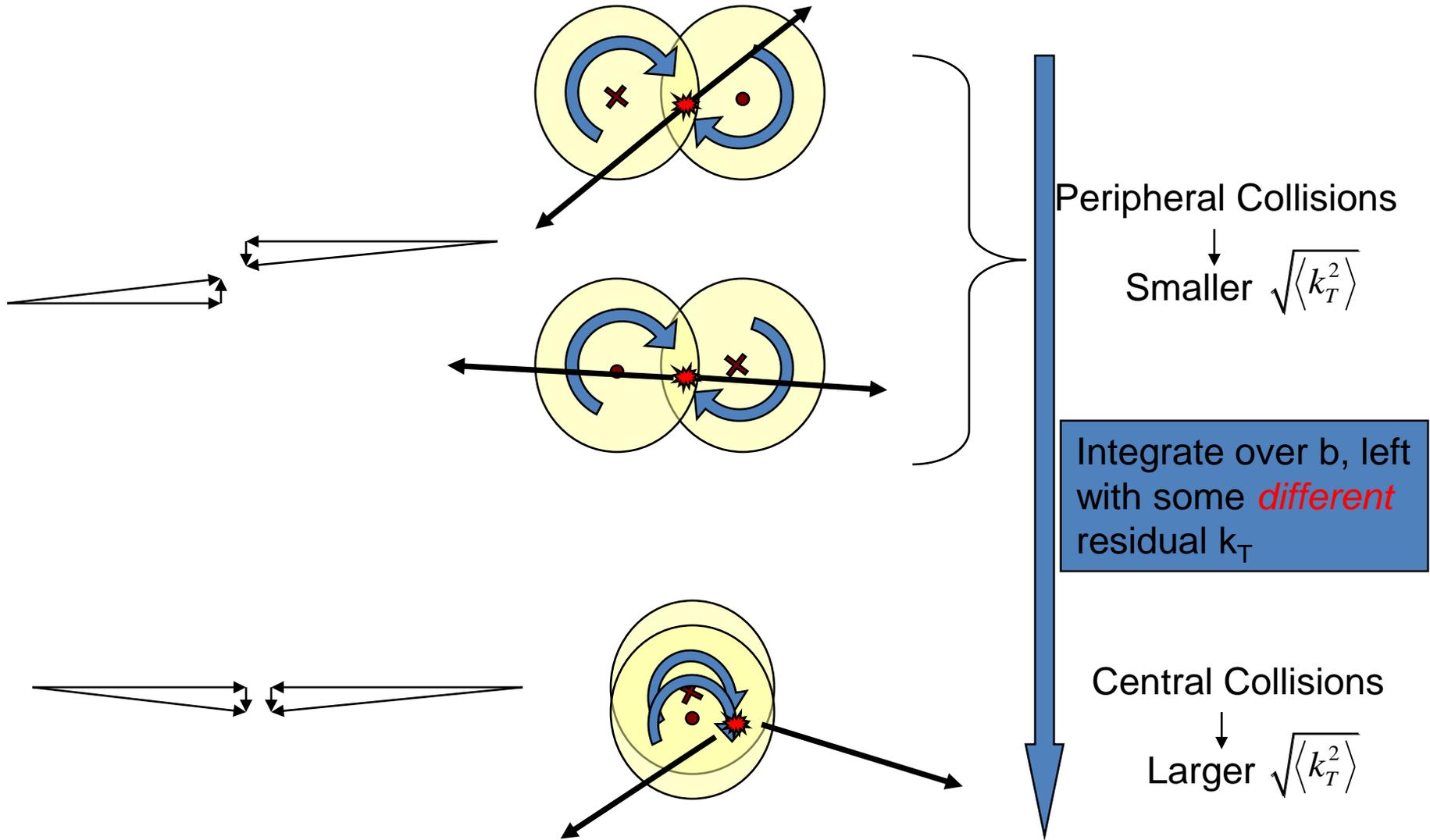
$$\Delta \langle k_T^2 \rangle_{hard} \neq 0$$

$$\Delta \langle k_T^2 \rangle_{frag} \approx \mathcal{G} \left( \frac{1}{Q} \right)$$

# Mechanism for $\Delta \langle k_T^2 \rangle_{initial} \neq 0$



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# Integrating over Impact Parameter

Drell-Yan pair production in polarized  $p+p$  collisions.

(Meng Ta-Chung et al. Phys. Rev. D 1989)

$$p_t^2(b, \theta_P, \theta_T; \pm) = \left| \vec{k}_{PR} \right|^2 + \left| \vec{k}_{TR} \right|^2 \pm \left| \vec{k}_{PR} \right| \left| \vec{k}_{TR} \right| \cos(\theta_P + \theta_T)$$

Averaging over  $D(b)$ :

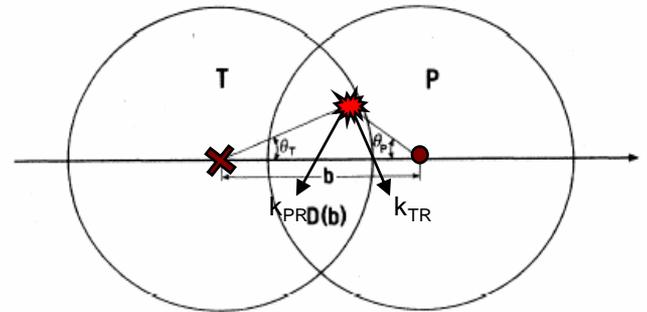
$$\Delta \langle p_t^2(b) \rangle = \langle p_t^2(b, +) \rangle - \langle p_t^2(b, -) \rangle$$

Integrating over impact parameter  $b$ :

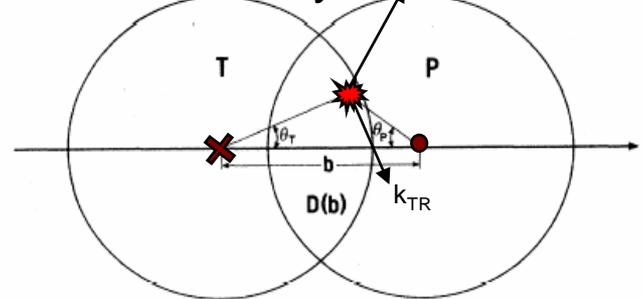
$$\Delta \langle p_t^2 \rangle \approx 1.9 \langle k_R \rangle^2$$

$$\langle k_R \rangle \equiv \langle \left| \vec{k}_{PR} \right| \rangle = \langle \left| \vec{k}_{TR} \right| \rangle$$

Like Helicity



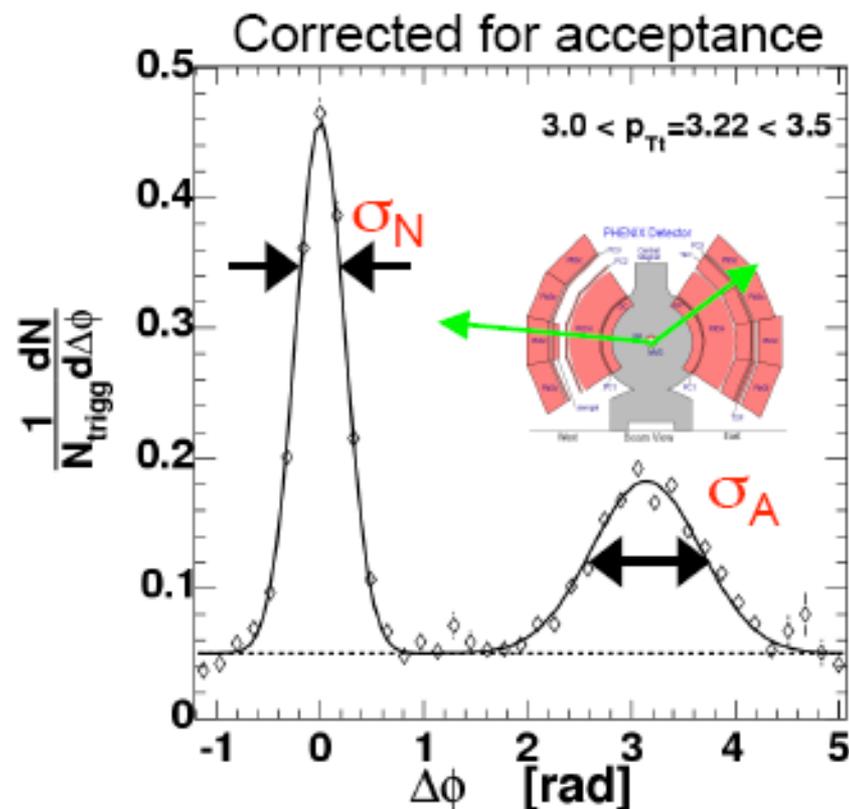
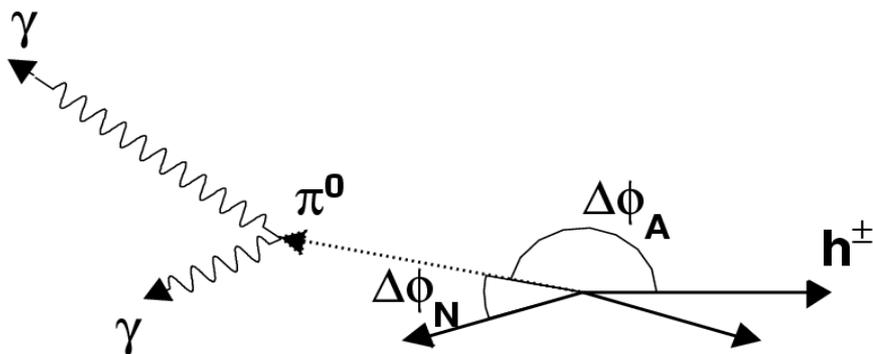
Un-like Helicity



# Measuring jet $k_T$ - di-hadron correlation

## $\pi^0$ - $h^\pm$ azimuthal correlation

Trigger on a particle, e.g.  $\pi^0$ , with transverse momentum  $p_{Tt}$ . Measure azimuthal angular distribution *w.r.t.* the azimuth of associated (charged) particle with transverse momentum  $p_{T\alpha}$ . The strong **same** and **away side peaks** in p-p collisions indicate di-jet origin from hard scattering partons.

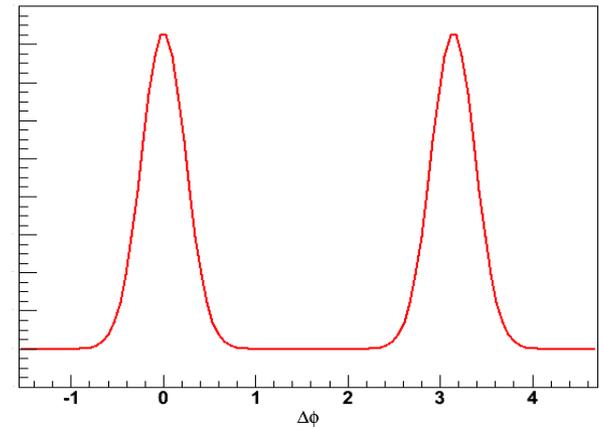
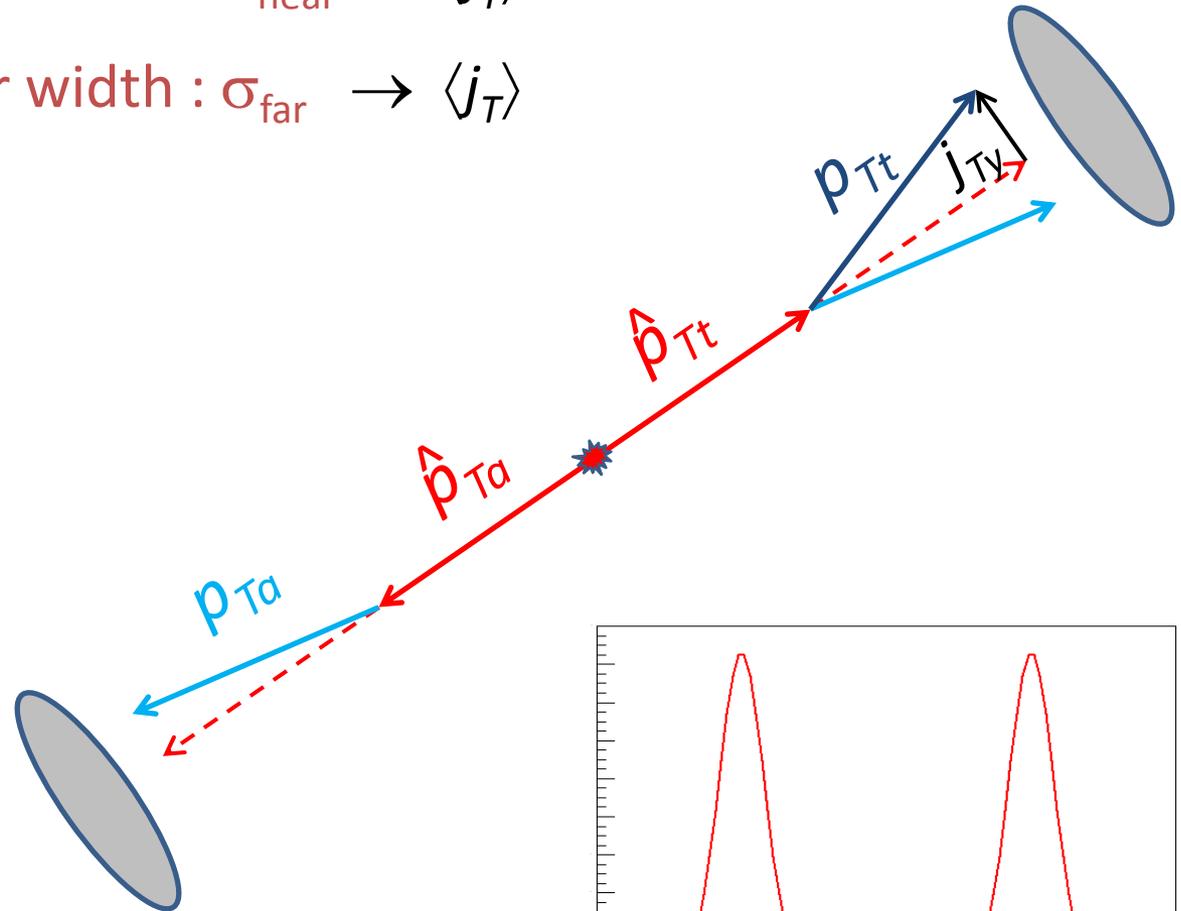


# Measuring jet $k_T$ - di-hadron correlation

Intra-jet pairs angular width :  $\sigma_{\text{near}} \rightarrow \langle j_T \rangle$

Inter-jet pairs angular width :  $\sigma_{\text{far}} \rightarrow \langle j_T \rangle$

- Zero  $k_T$

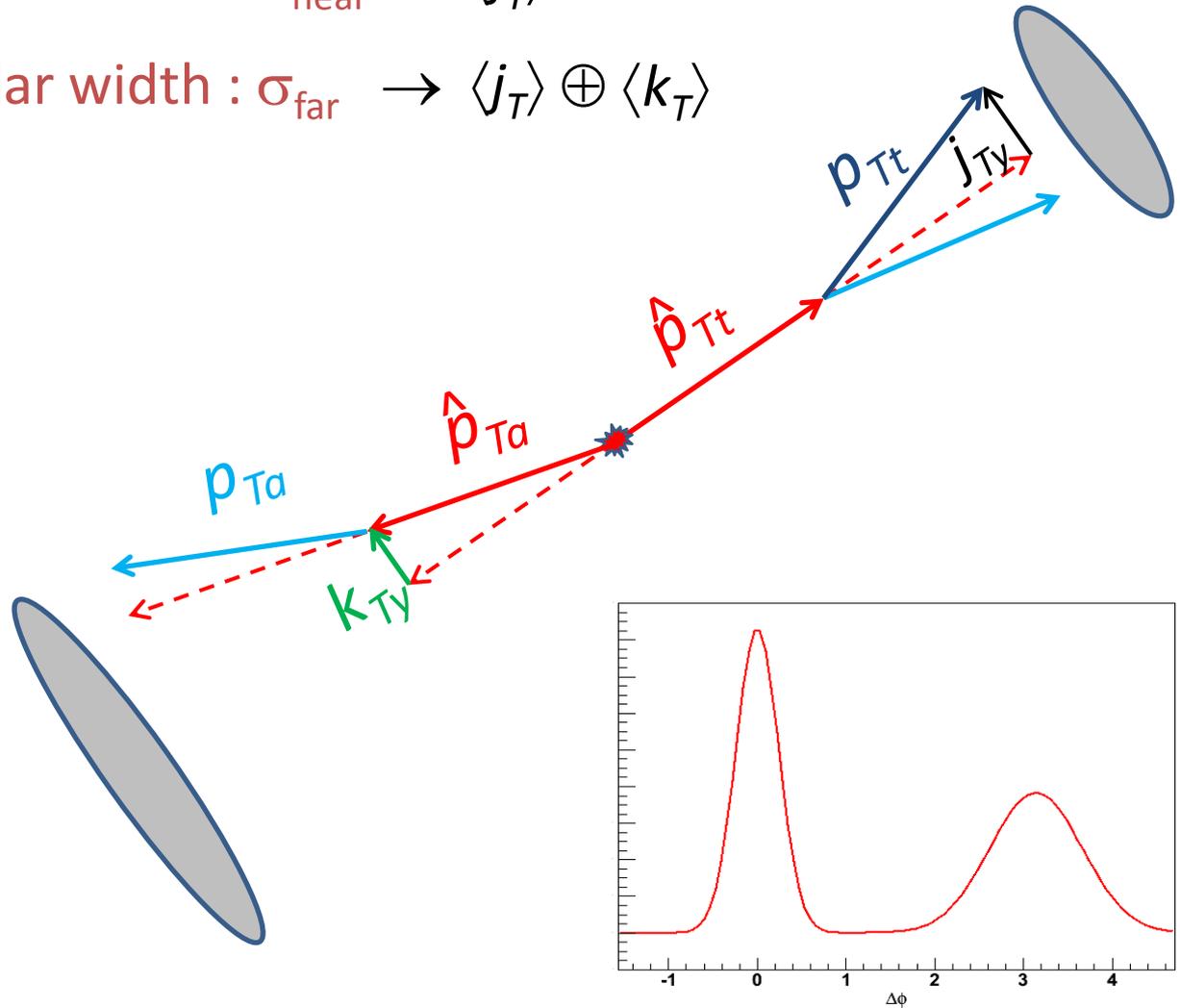


# Measuring jet $k_T$ - di-hadron correlation

Intra-jet pairs angular width :  $\sigma_{\text{near}} \rightarrow \langle j_T \rangle$

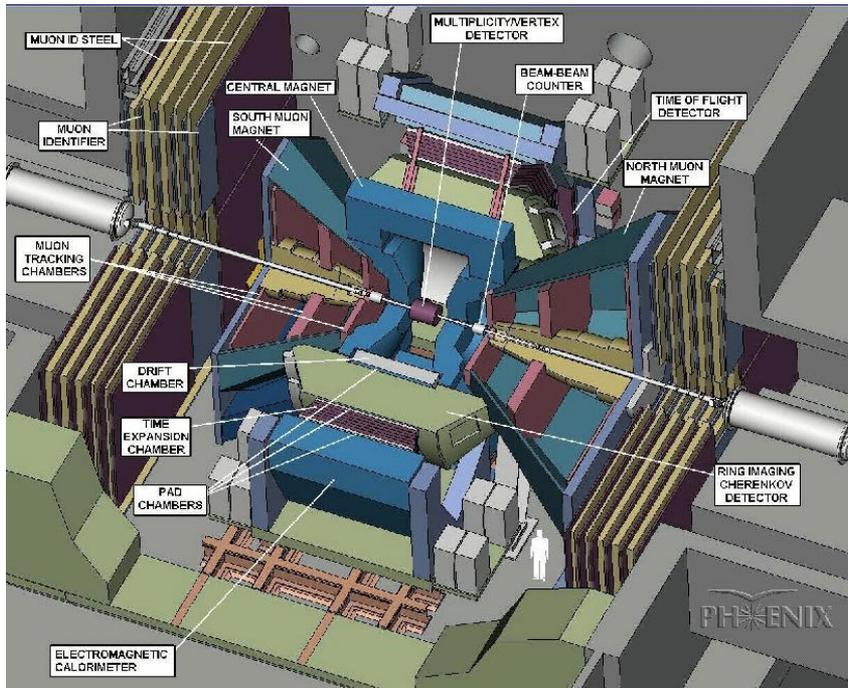
Inter-jet pairs angular width :  $\sigma_{\text{far}} \rightarrow \langle j_T \rangle \oplus \langle k_T \rangle$

- Non-Zero  $k_T$



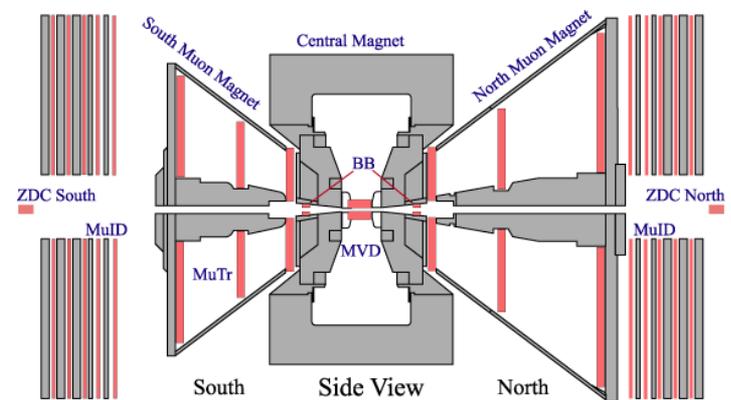
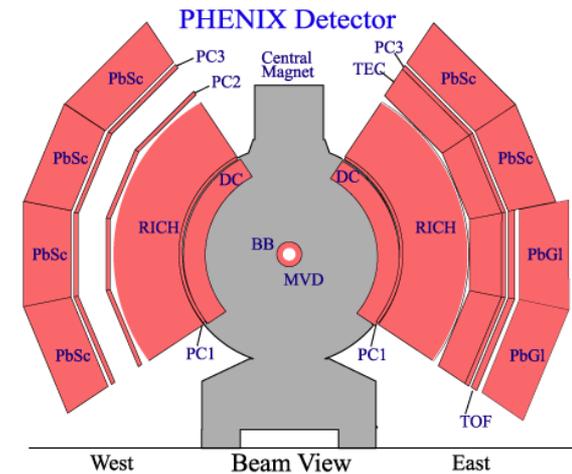


# PHENIX Detector



## •Philosophy:

- High resolution at the cost of acceptance
- High rate capable DAQ
- Excellent **trigger capability** for rare events



## •Central Arms:

$$|\eta| < 0.35, \Delta\phi = 2 \times 90^\circ$$

Charged particle ID and tracking;  
photon ID. EM Calorimetry.

## Muon Arm:

$$1.2 < |\eta| < 2.4$$

Muon ID and tracking.

## •Global Detectors

Collision trigger  
Collision vertex  
characterization  
Relative luminosity  
Local Polarimetry

# Data Selection

Analyzed data sample from 2005 & 2006 RHIC Runs.

- Integrated Luminosity:  $2.5 \text{ pb}^{-1}$  (Run 5) +  $6.4 \text{ pb}^{-1}$  (Run 6)
- Average beam polarization: 50% (Run 5), 57% (Run 6)
- $\pi^0$  transverse momentum:  $2 < p_T < 10 \text{ GeV}$
- $h^\pm$  transverse momentum:  $2 < p_T < 5 \text{ GeV}$
- $|\eta| < 0.35$

# $\pi^0$ Identification

$$\pi^0 \rightarrow \gamma + \gamma$$

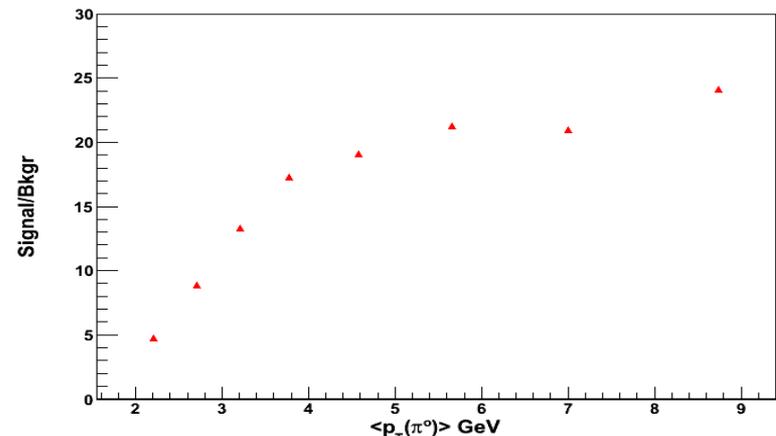
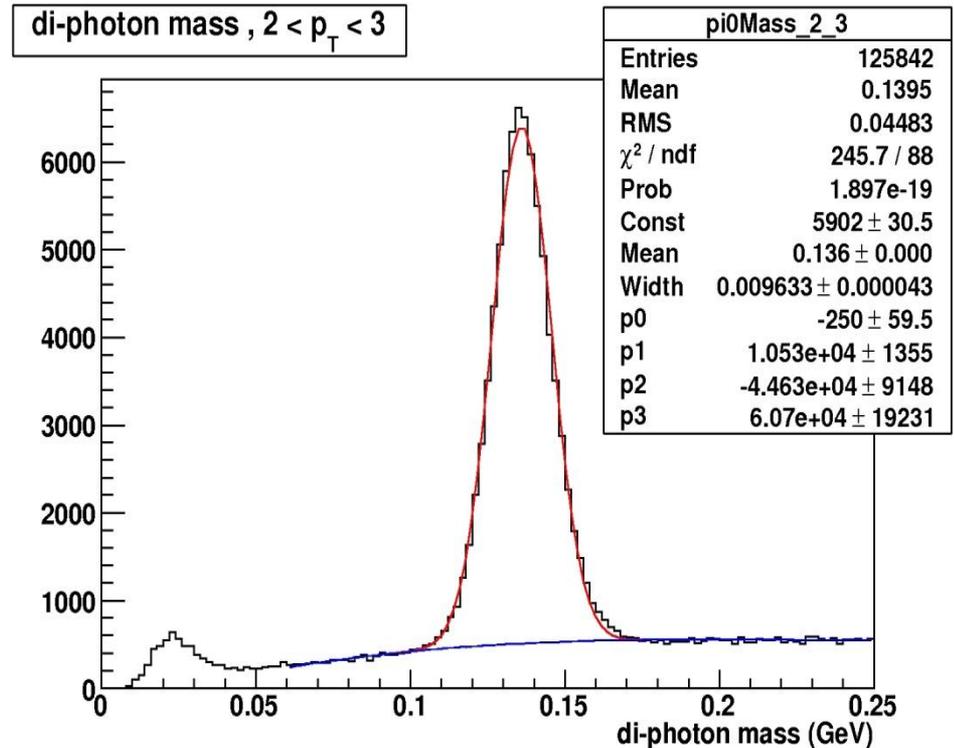
The pions are selected within

$$M_{\pi^0} \pm 2.0\sigma_{\pi^0}$$

Photon trigger:

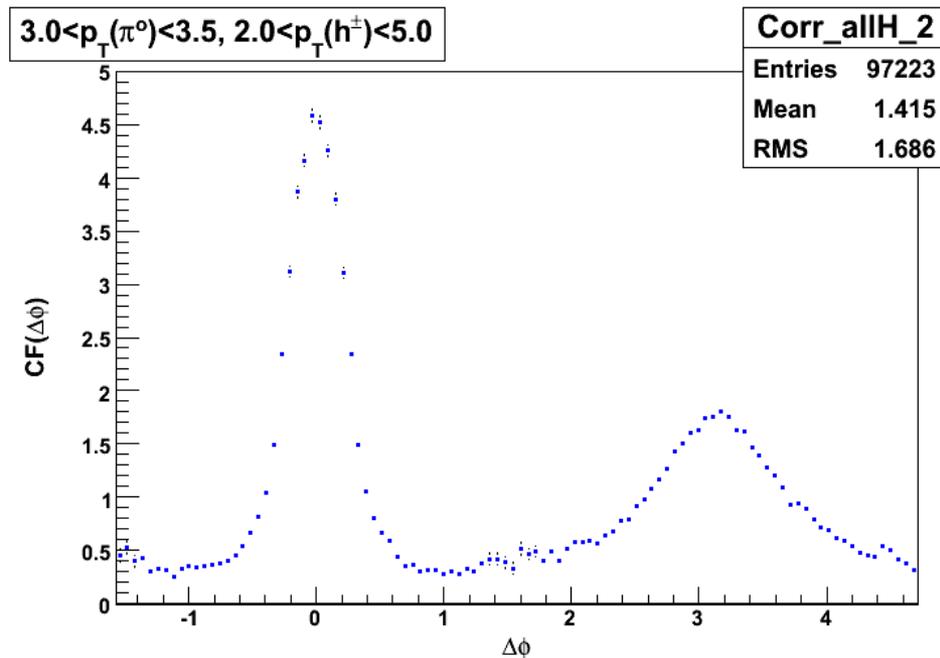
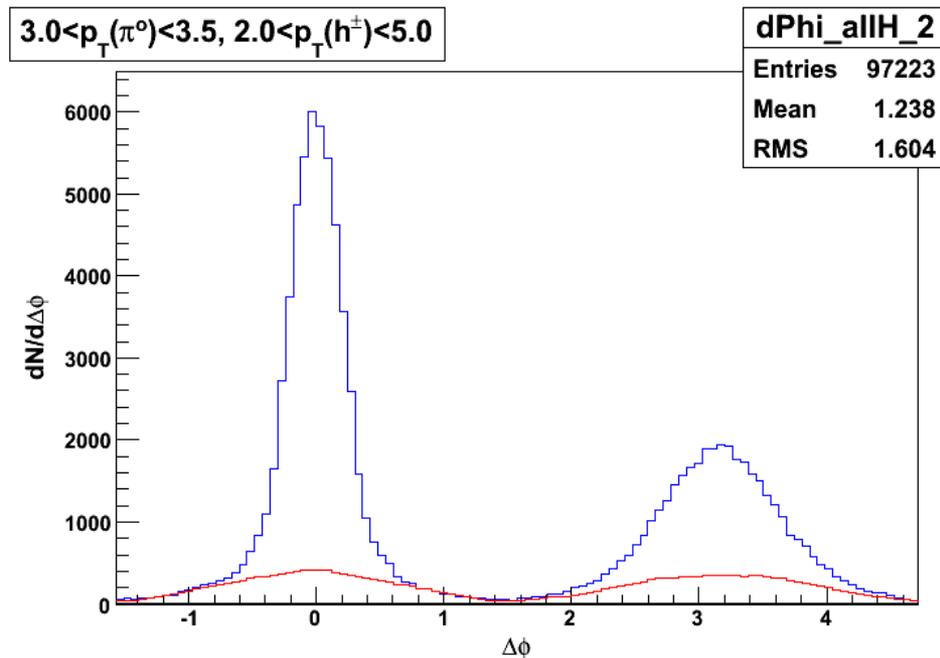
EMCal cluster energy > 1.4 GeV

The mass,  $M_{\pi^0}$ , and the width,  $\sigma_{\pi^0}$ , are determined by fitting the di-photon mass spectrum with *gaus+pol3*.



# Correlation Function

- **Blue** - real (raw  $\Delta\phi$ )
- **Red** - background (mixed events)

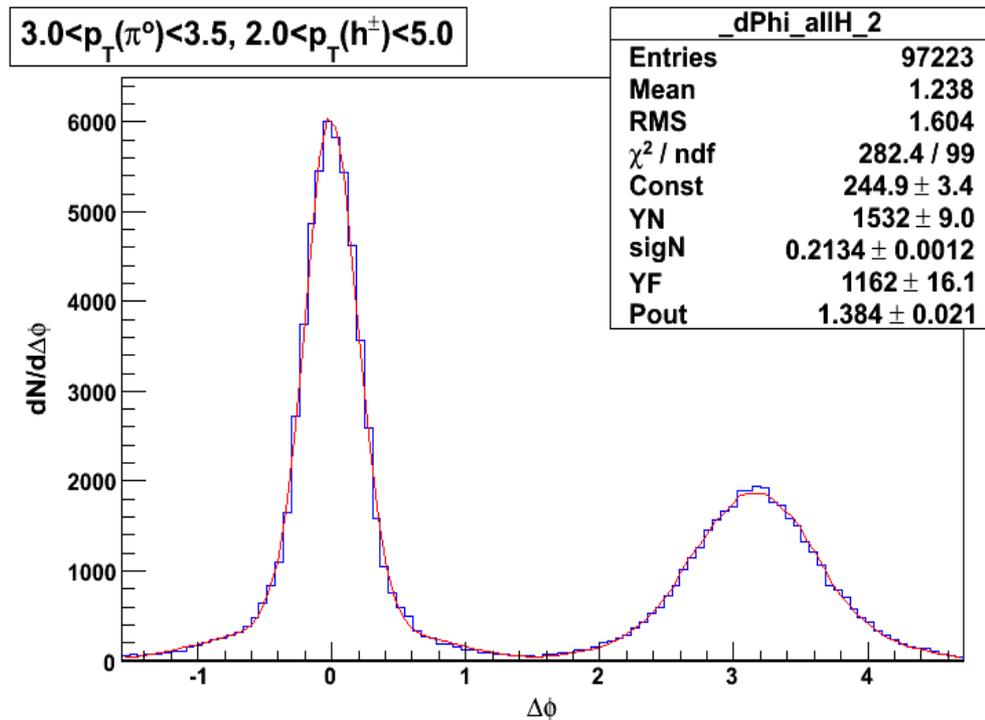


# Fit Function

$$\frac{dN_{real}}{d\Delta\phi} = \frac{1}{N} \frac{dN_{mix}}{d\Delta\phi} \cdot \left( C_o + C_1 \cdot Gaus(0) + C_2 \cdot \frac{dN_{away}}{d\Delta\phi} \Big|_{\pi/2}^{3\pi/2} \right)$$

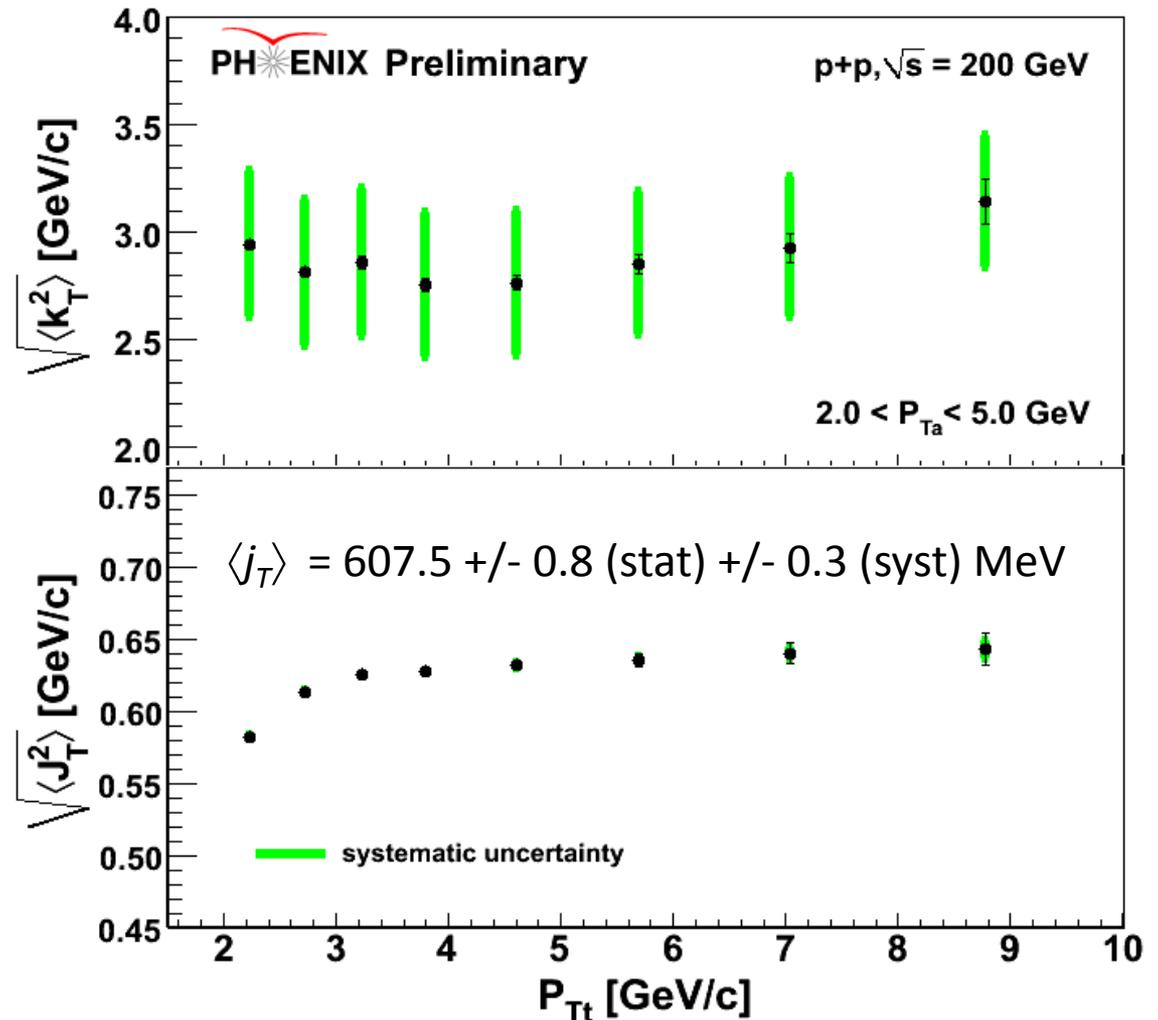
$$\frac{dN_{away}}{d\Delta\phi} \Big|_{\pi/2}^{3\pi/2} = \frac{-p_{Ta} \cos \Delta\phi}{\sqrt{2\pi} \langle p_{out}^2 \rangle \text{Erf} \left( \sqrt{2} p_{Ta} / \sqrt{\langle p_{out}^2 \rangle} \right)} \exp \left( -\frac{p_{Ta}^2 \sin^2 \Delta\phi}{2 \langle p_{out}^2 \rangle} \right)$$

- **Blue** – real (raw  $\Delta\phi$ )
- **Red** – Fit from above equation



# $k_T$ results for unpolarized case

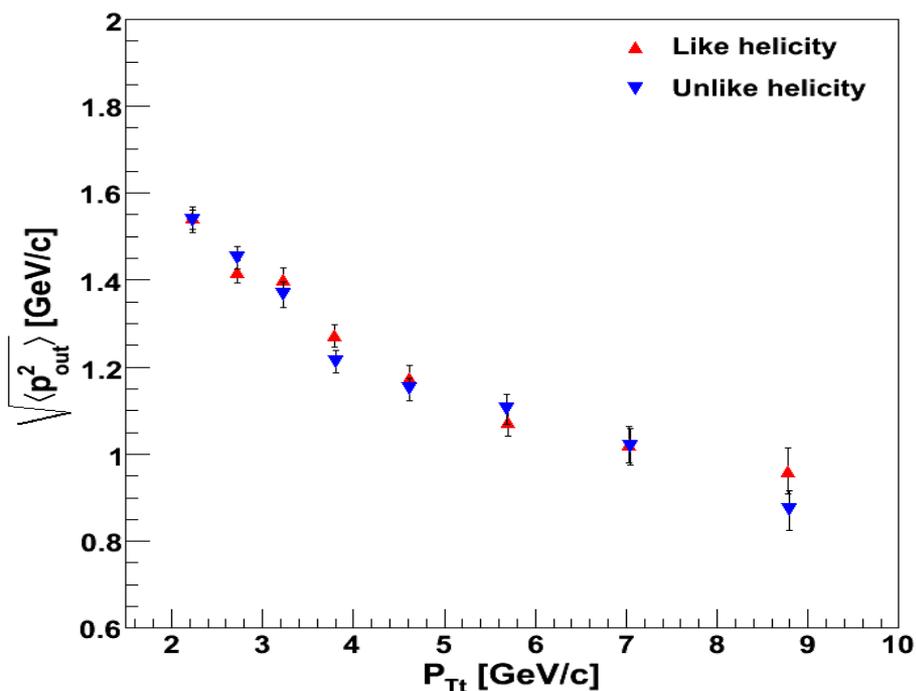
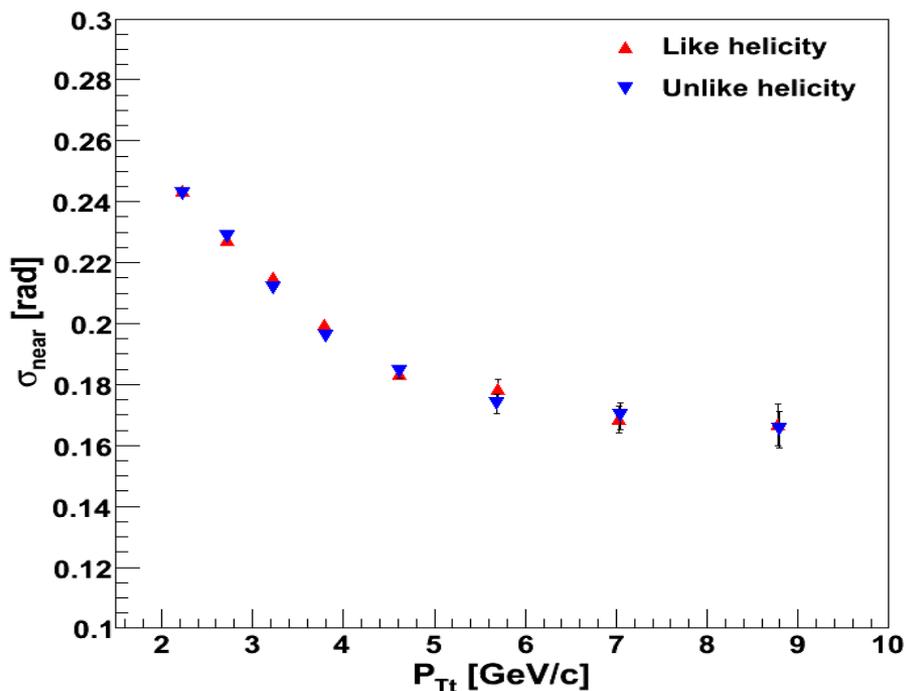
$$\langle k_T \rangle = 2.88 \pm 0.02 \text{ (stat)} \pm 0.12 \text{ (syst)} \text{ GeV}$$



# Spin sorted dihadron correlations

## Fit Results

- Red  $\blacktriangle$  – Like helicity ( $++$  &  $--$ )
- Blue  $\blacktriangledown$  – Unlike helicity ( $+-$  &  $-+$ )

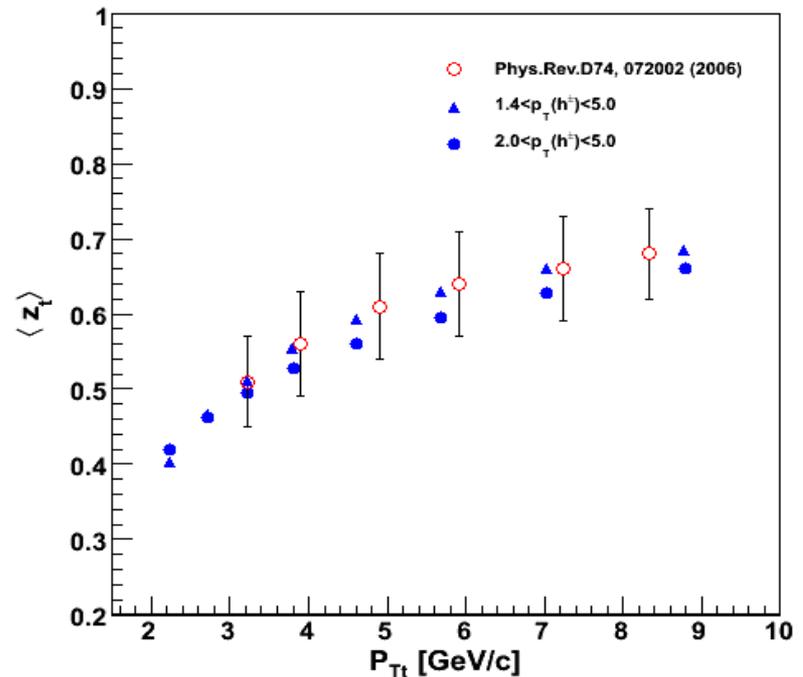
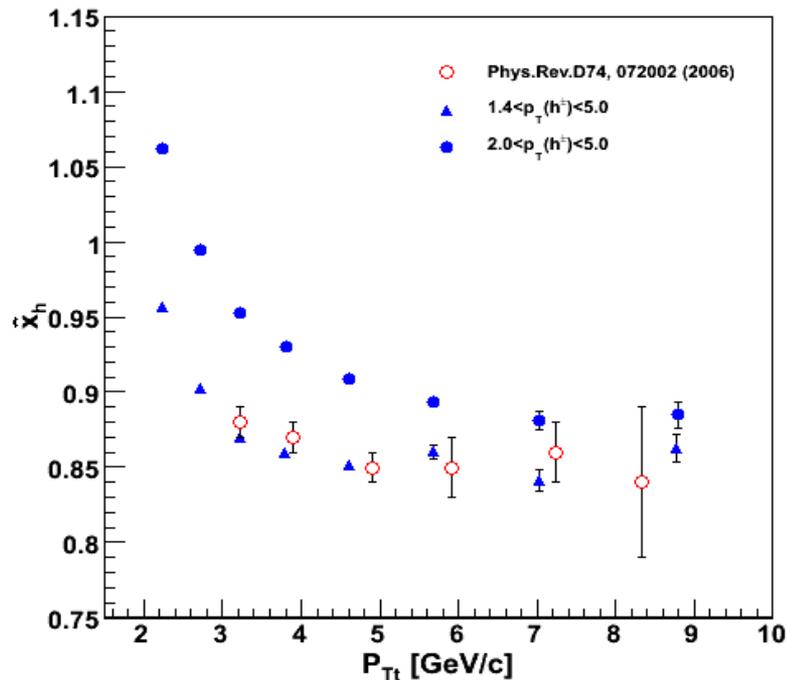


# Need $\langle z_t \rangle$ and $x_h$

$$\frac{\langle z_t(k_T, x_h) \rangle \sqrt{\langle k_T^2 \rangle}}{\hat{x}_h(k_T, x_h)} = \frac{1}{x_h} \sqrt{\langle p_{\text{out}}^2 \rangle - \langle j_{Ty}^2 \rangle (1 + x_h^2)}$$

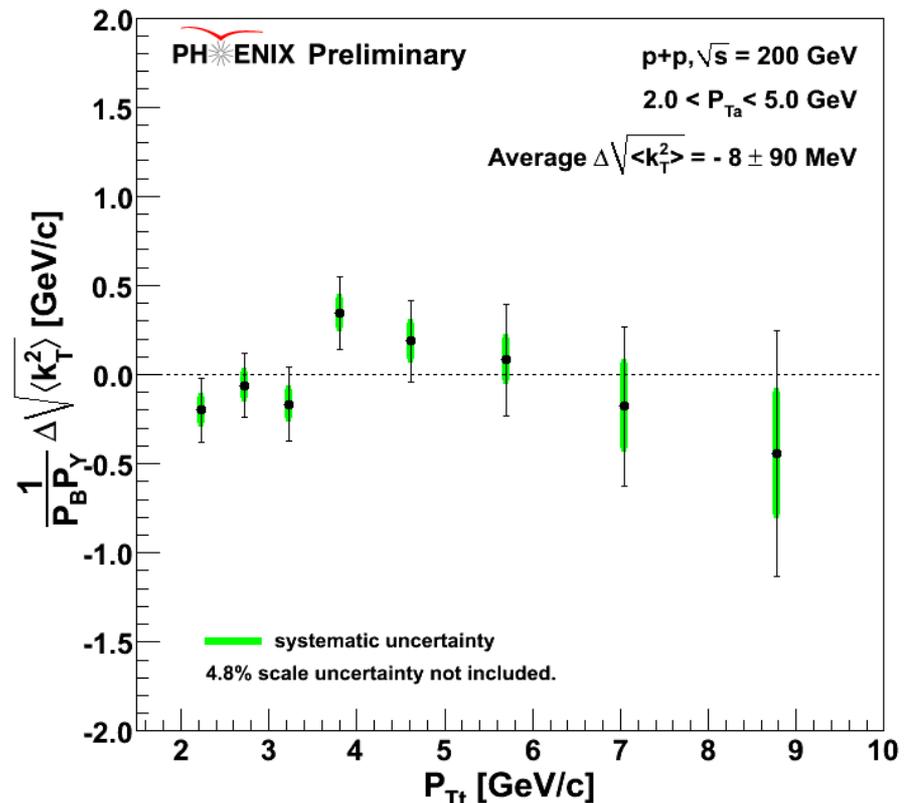
$\hat{x}_h$  and  $\langle z_t \rangle$  are determined through a combined analysis of the  $\pi^0$  inclusive cross section and using jet fragmentation function measured in  $e^+e^-$  collisions. Ref. PRD 74, 072002 (2006)

$\langle k_T^2 \rangle$  asymmetry results are not sensitive to this  $\hat{x}_h$  and  $\langle z_t \rangle$  values – only needed to set the scale.



# Preliminary Results

- Take the difference Like - Unlike helicities
- Normalized by beam polarizations
- $\langle j^2_T \rangle$  asymmetry small ( $-6 \pm 6$  MeV out of  $\sim 630$  MeV for unpolarized)
- $\langle k^2_T \rangle$  asymmetry also small ( $-8 \pm 90$  MeV out of  $\sim 2.8$  GeV for unpolarized).



# Interpretation

$$\Delta \langle k_T^2 \rangle_{initial} \neq 0$$

$$\Delta \langle k_T^2 \rangle_{hard} \neq 0$$

$$\Delta \langle k_T^2 \rangle_{frag} \approx \mathcal{O}\left(\frac{1}{Q}\right)$$

- The small asymmetry in  $j_T$  verifies our assumption that the third term is suppressed.
- The second term can be constrained from current knowledge of the  $\pi^0 A_{LL}$ .

# Interpretation

- The first term can be related in the “Meng conjecture” to partonic transverse momentum:

$$\Delta \langle k_T^2 \rangle_{initial} = \sum_i \sum_j c^{ij} W^{ij} \langle \vec{k}_T^i \cdot \vec{k}_T^j \rangle^{++} - \langle \vec{k}_T^i \cdot \vec{k}_T^j \rangle^{+-}$$

where  $c^{ij}$  give the initial state weights and  $W^{ij}$  give the impact parameter weighting.

- Since PYTHIA studies show that the final state studied here is dominated by gg scattering (~50%),
- We can interpret this as a constraint on the gluon orbital angular momentum.
- Consistent with previous (Sivers) measurements, and complimentary, since one naively needs no initial or final state interactions.

# Summary

- We have an analysis tool - di-hadron azimuthal correlations - that allows us to measure  $k_T$ .
- We are studying this in helicity-sorted collisions to see if there is a spin-dependent, coherent component of  $k_T$ .
- In our kinematic range, we find  $\Delta j_T^{RMS} = -6 \pm 6$  MeV, and  $\Delta k_T^{RMS} = -8 \pm 90$  MeV.
- While this is consistent with other measurements related to gluon OAM, the connection to parton OAM needs theoretical guidance!