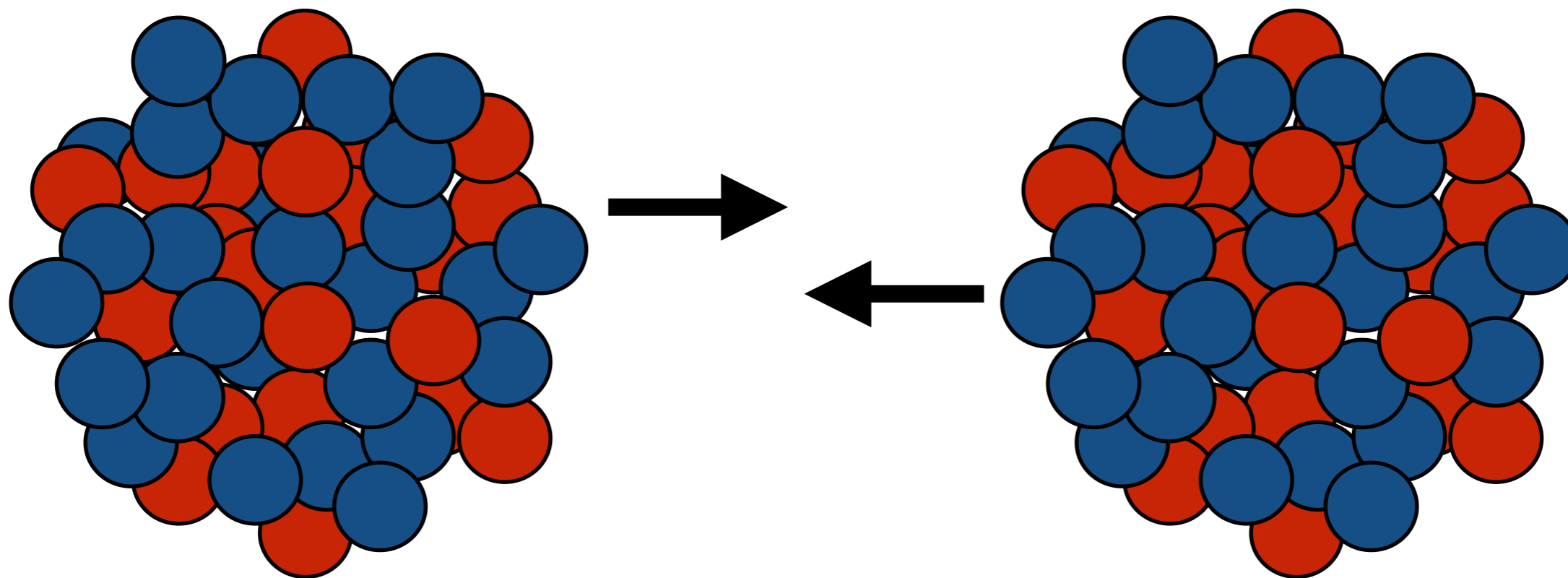


# The Lighter Side of Heavy Ions

News from p/d+A Collisions

Anne M. Sickles, BNL

March 10, 2014

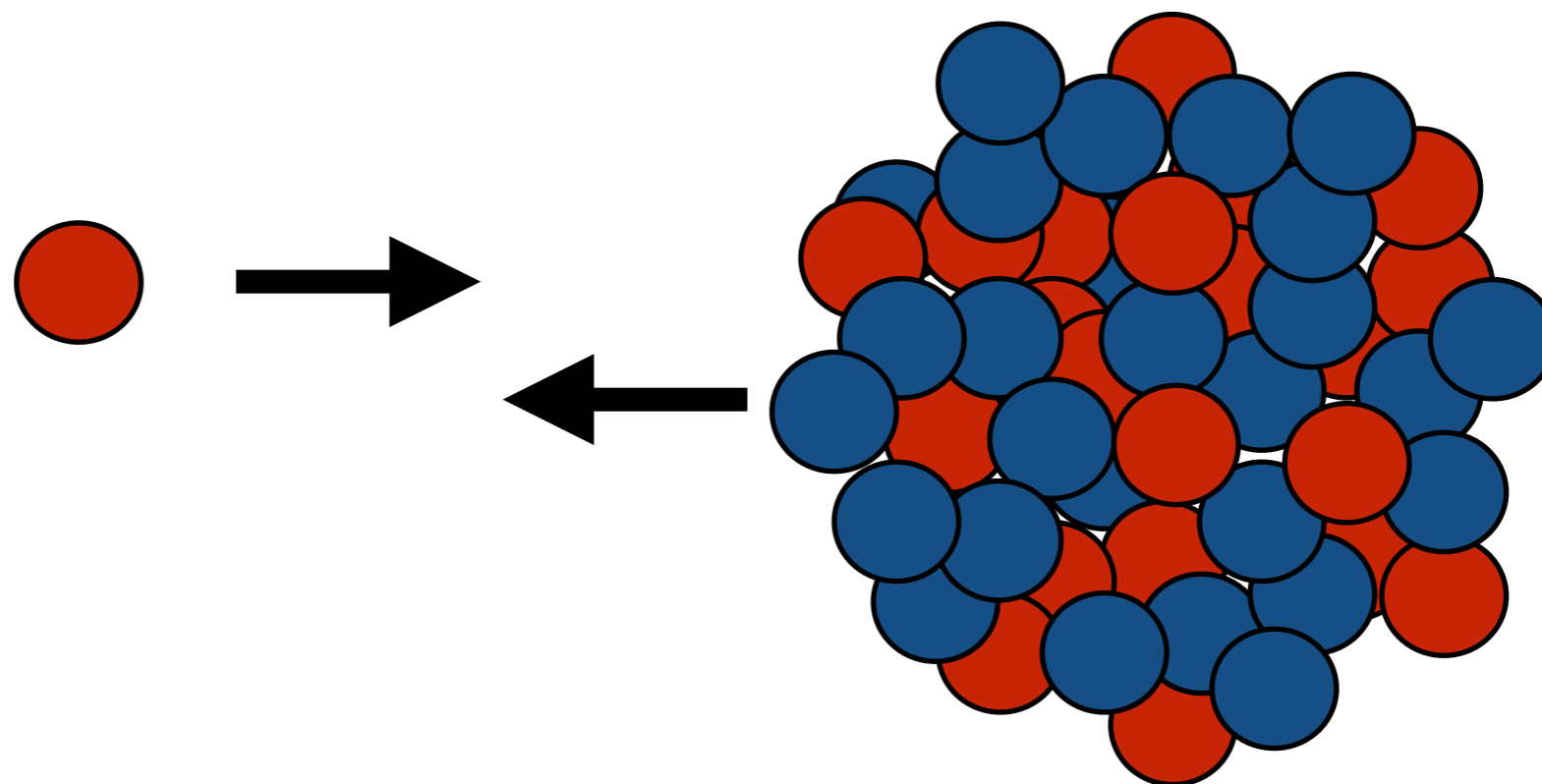


# The Lighter Side of Heavy Ions

News from p/d+A Collisions

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March 10, 2014



# The Lighter Side of Heavy Ions

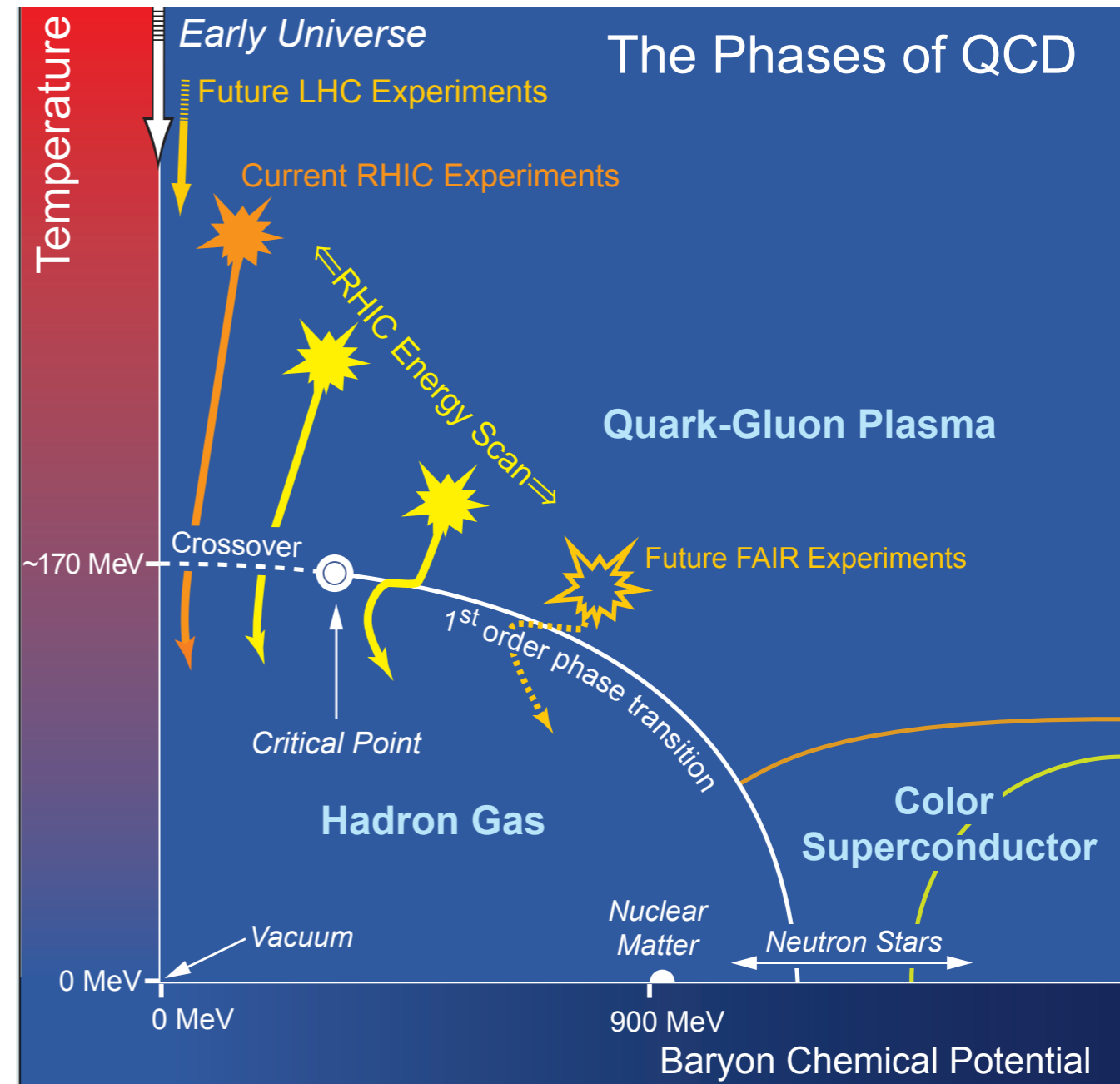
News from p/d+A Collisions

Anne M. Sickles, BNL

March 10, 2014

# why heavy ion collisions?

- study the phase structure of QCD
- here we will focus on the high temperature side, but there are also exciting new investigations at lower temperatures and higher baryon densities
- quantitative understanding of the properties of QCD matter at extreme temperatures



# Heavy Ion Programs at RHIC and LHC

---



2000 - present

7.7-510 GeV collision energy

AuAu, dAu, pp, CuCu, UU, CuAu

strengths: collision system &  
energy versatility and long running  
times



2010 - present

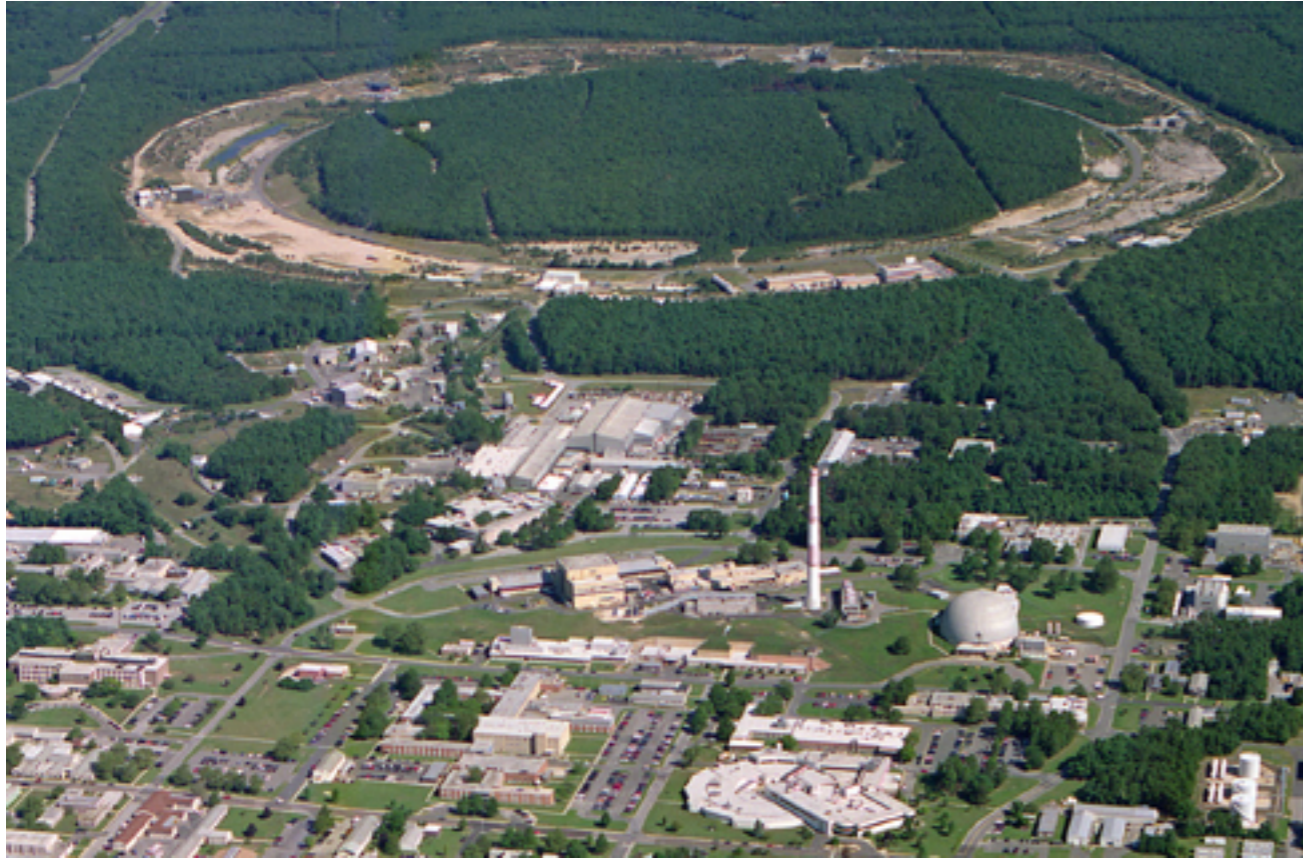
2.76 TeV collision energy PbPb

5.02 TeV pPb

pp @ multiple energies

strengths: excellent detectors and  
very high energy

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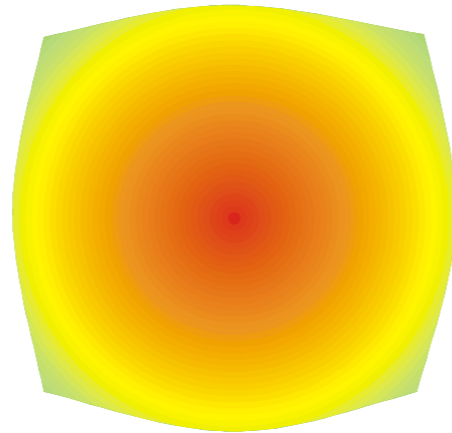
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# relativistic heavy ion collisions

---

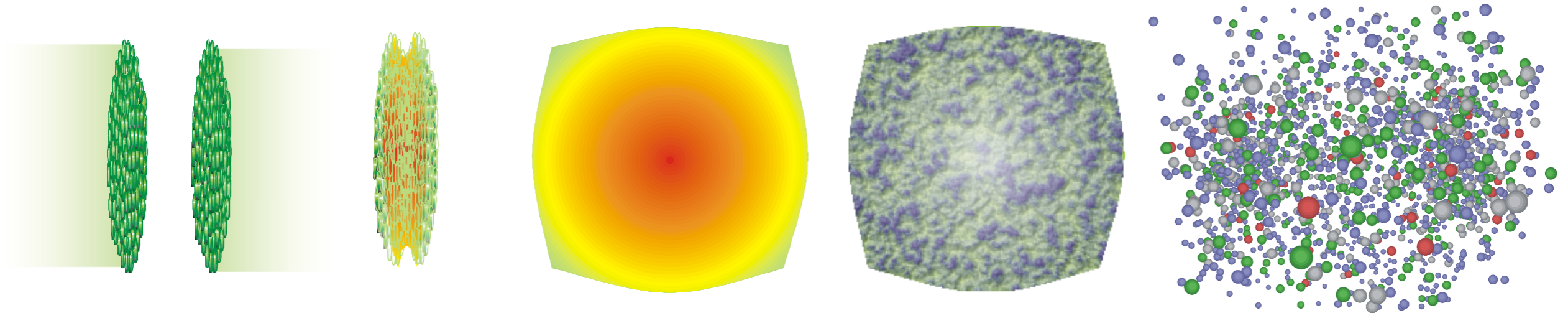
**quark gluon plasma**



# relativistic heavy ion collisions

---

## quark gluon plasma

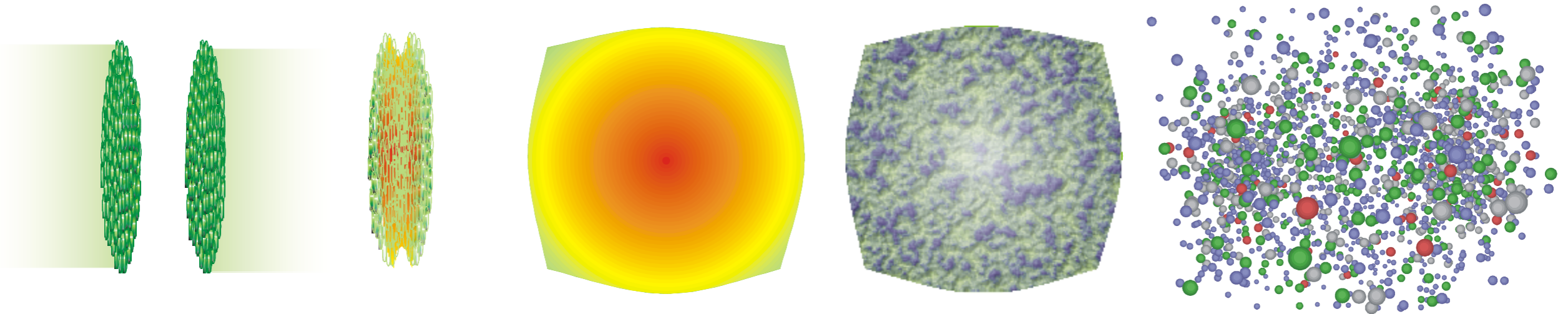




# relativistic heavy ion collisions

---

## quark gluon plasma



want to untangle **QGP** effects from  
effects of initial nucleus and  
hadronic matter

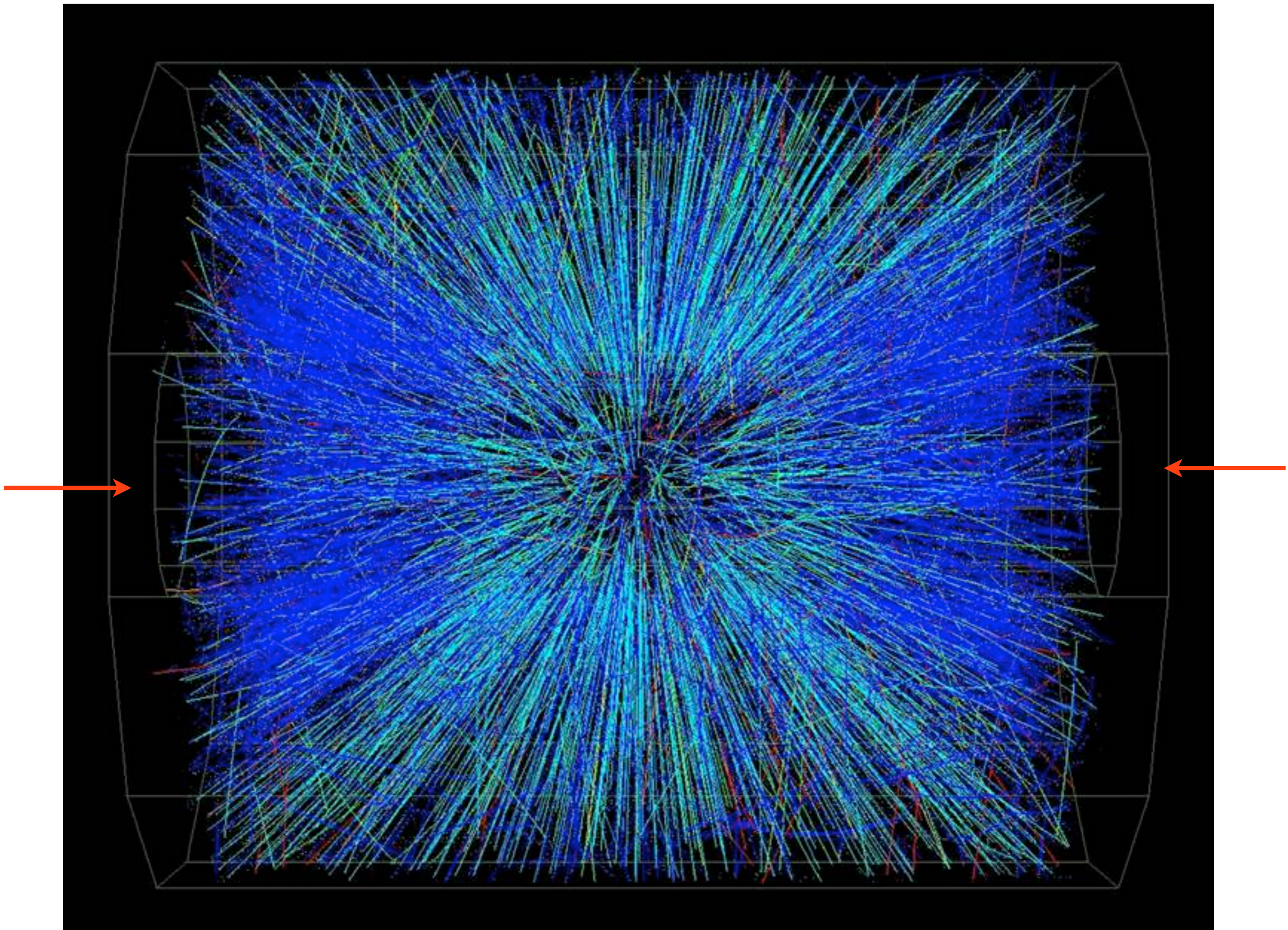
# the aftermath

---

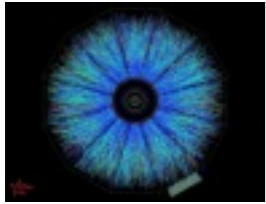


# the aftermath

---

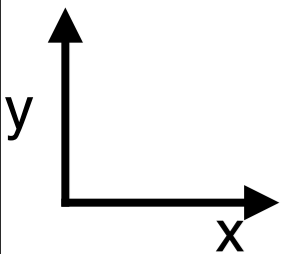
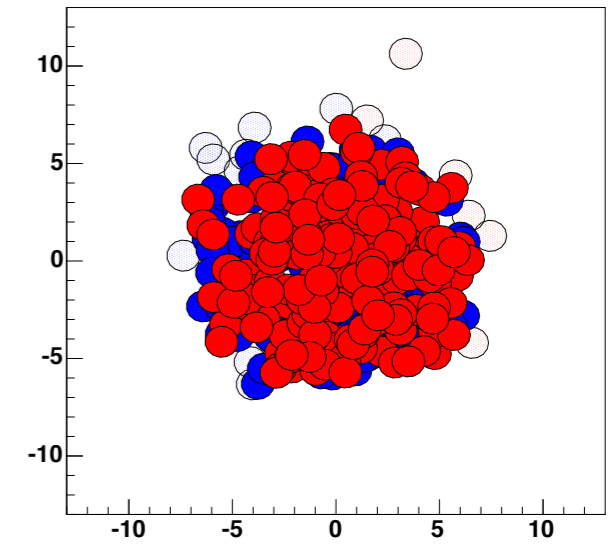
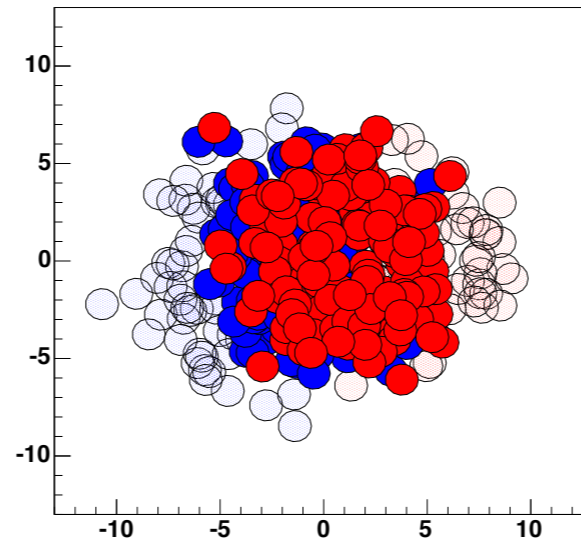
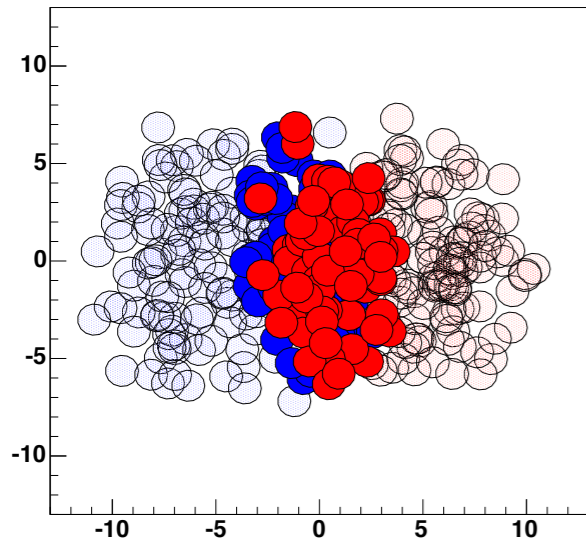


# collision geometry



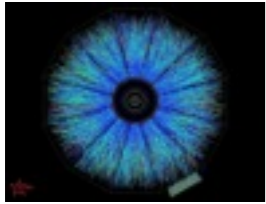
view: one nuclei going into the screen and one coming out

nucleon positions for the colliding nuclei for three different collisions



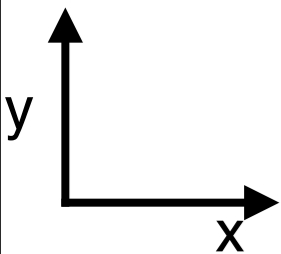
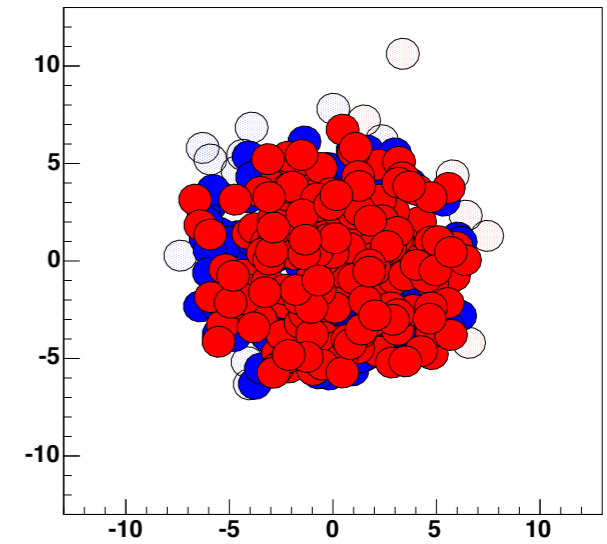
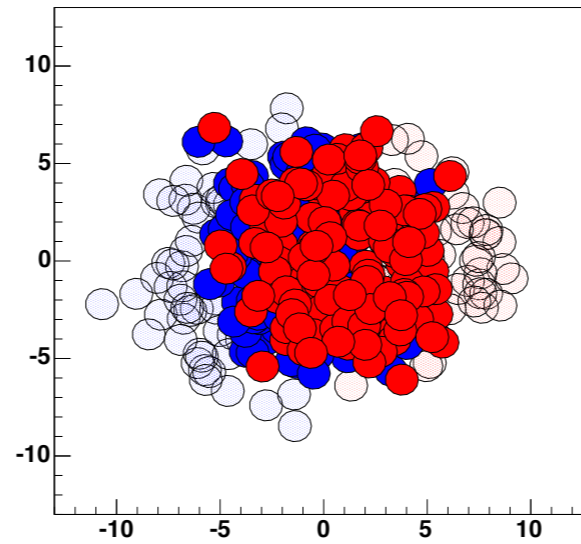
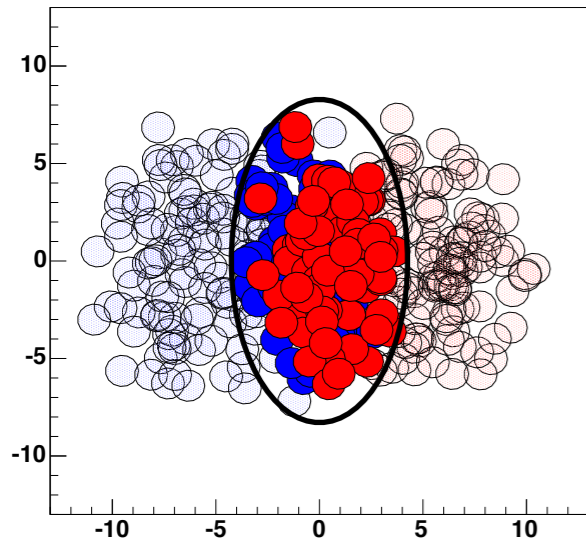
varying the distance between the nuclei, changes the shape and size of the region where the nuclei overlap

# collision geometry



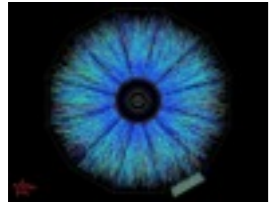
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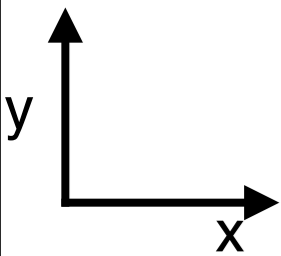
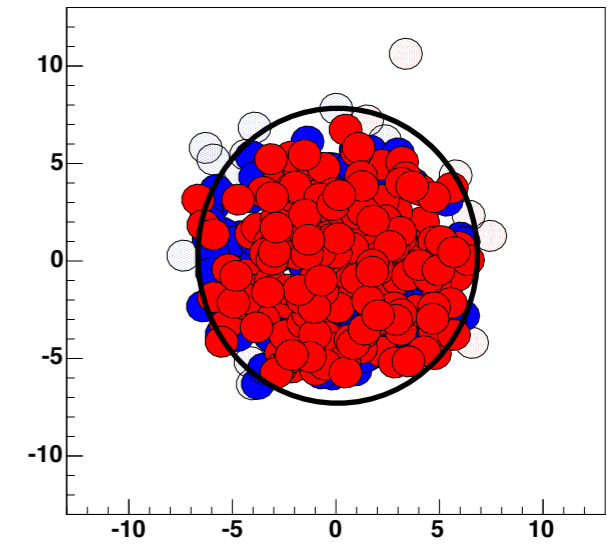
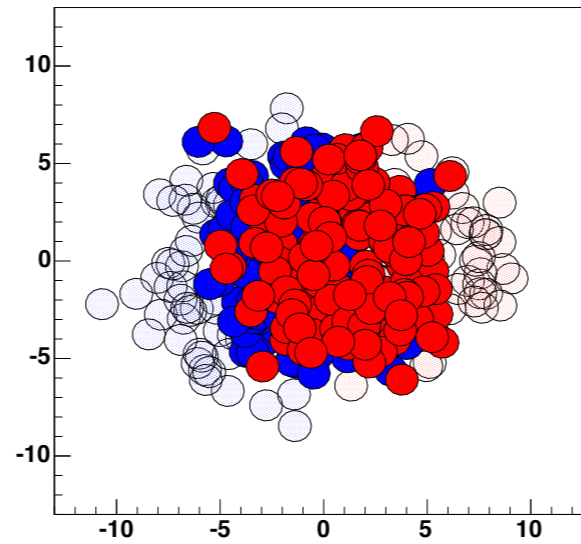
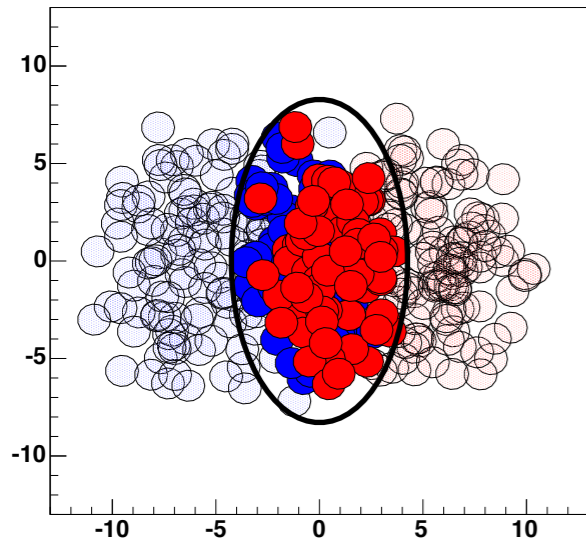
varying the distance between the nuclei, changes the shape and size of the region where the nuclei overlap

# collision geometry



view: one nuclei going into the screen and one coming out

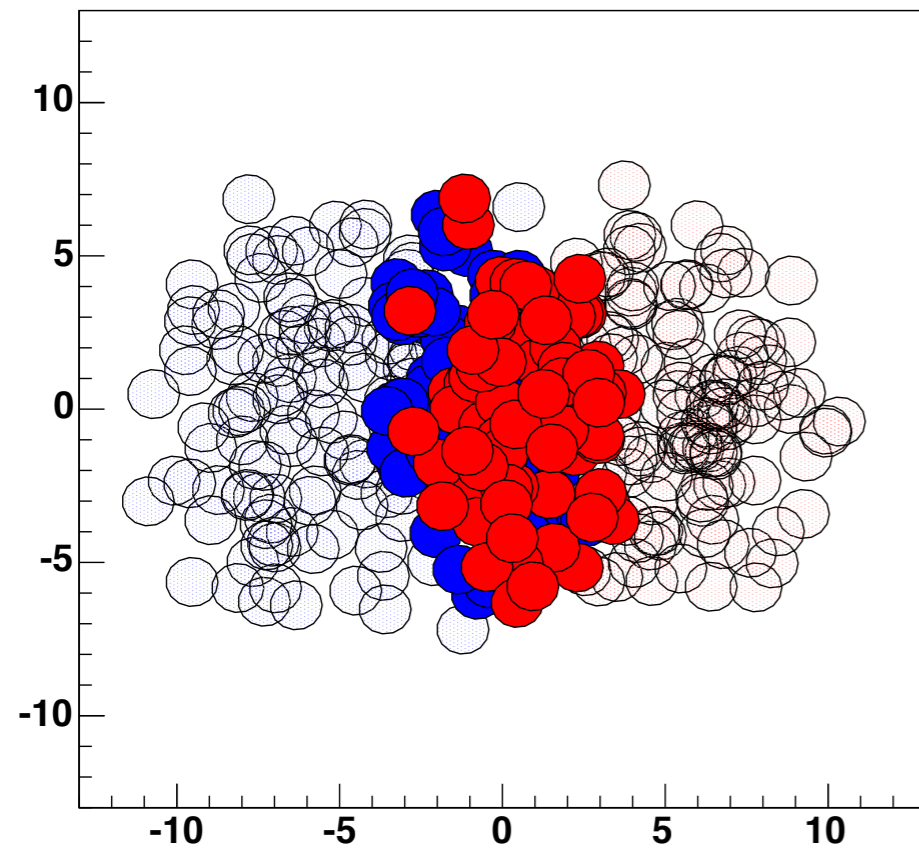
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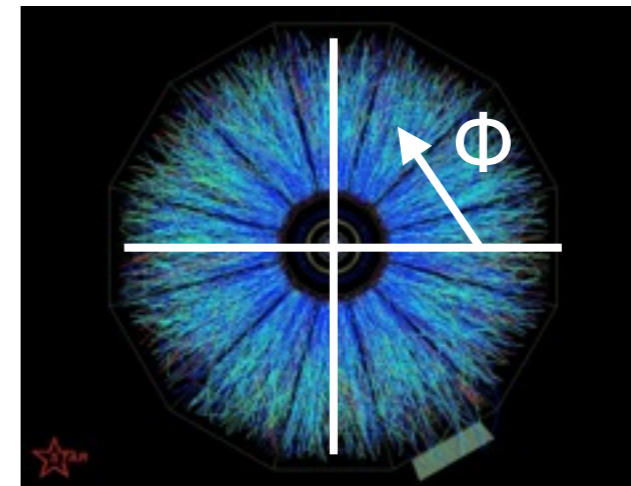
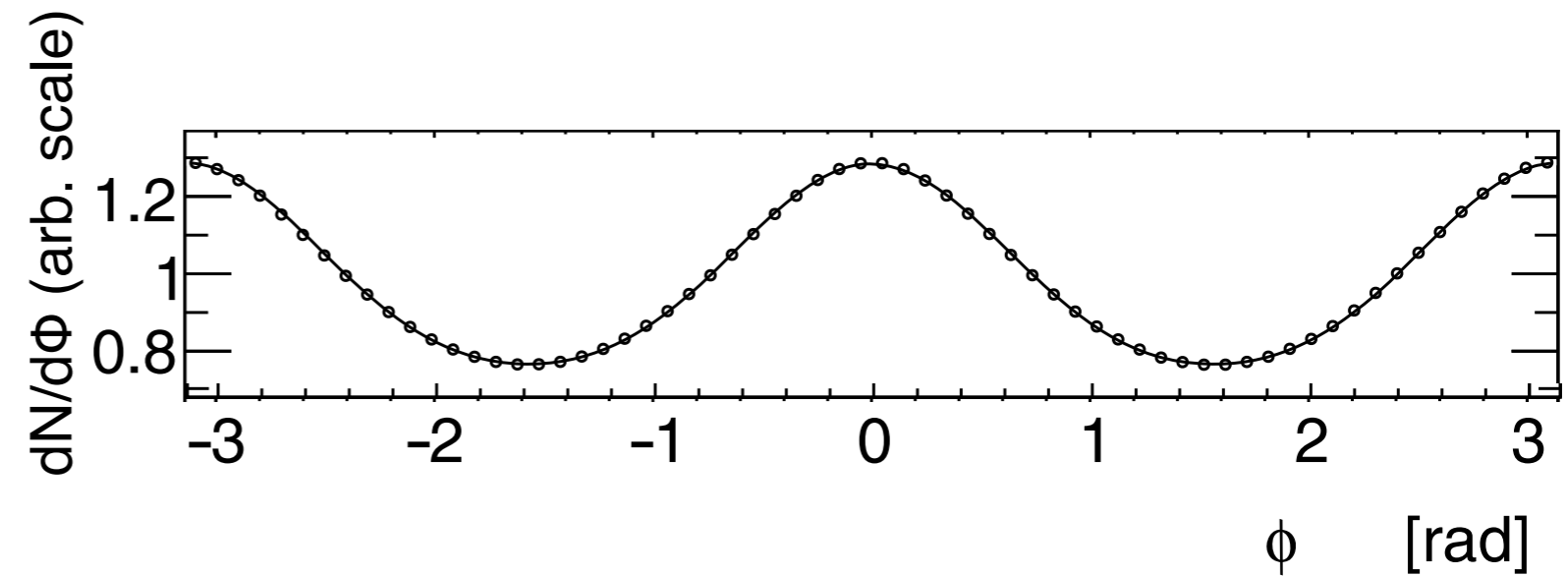
varying the distance between the nuclei, changes the shape and size of the region where the nuclei overlap

# collision geometry $\rightarrow$ measured particles

## initial collision geometry

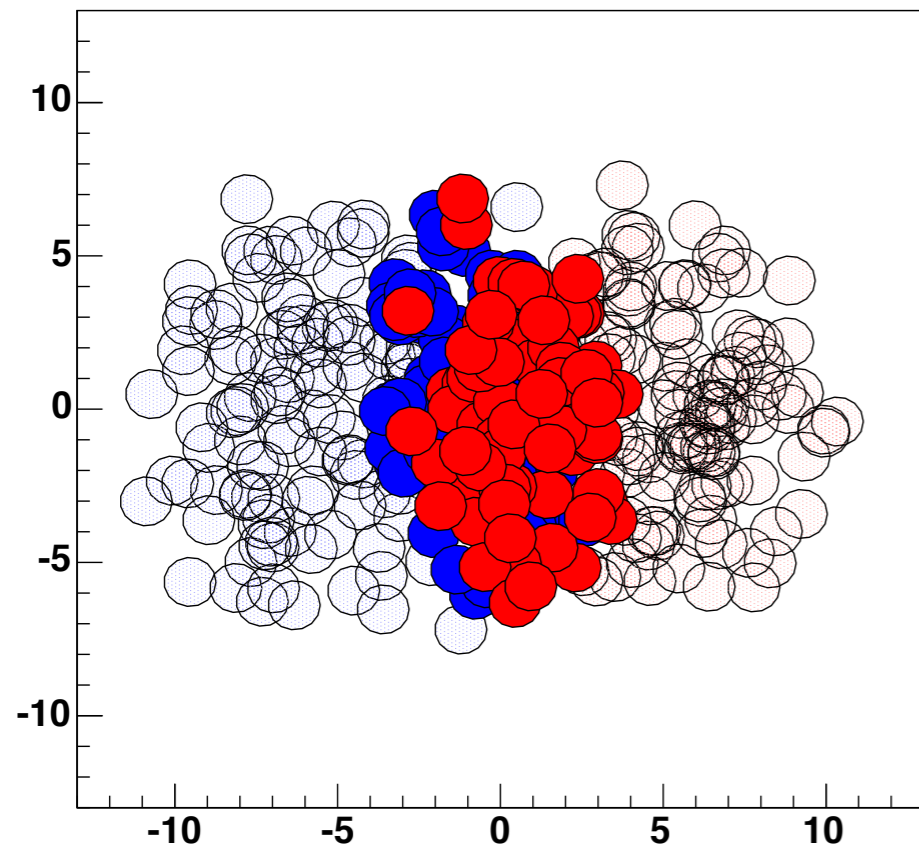


## measured hadron distributions

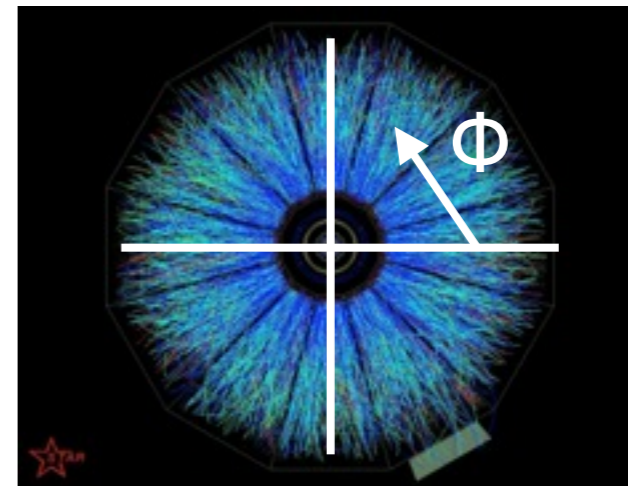
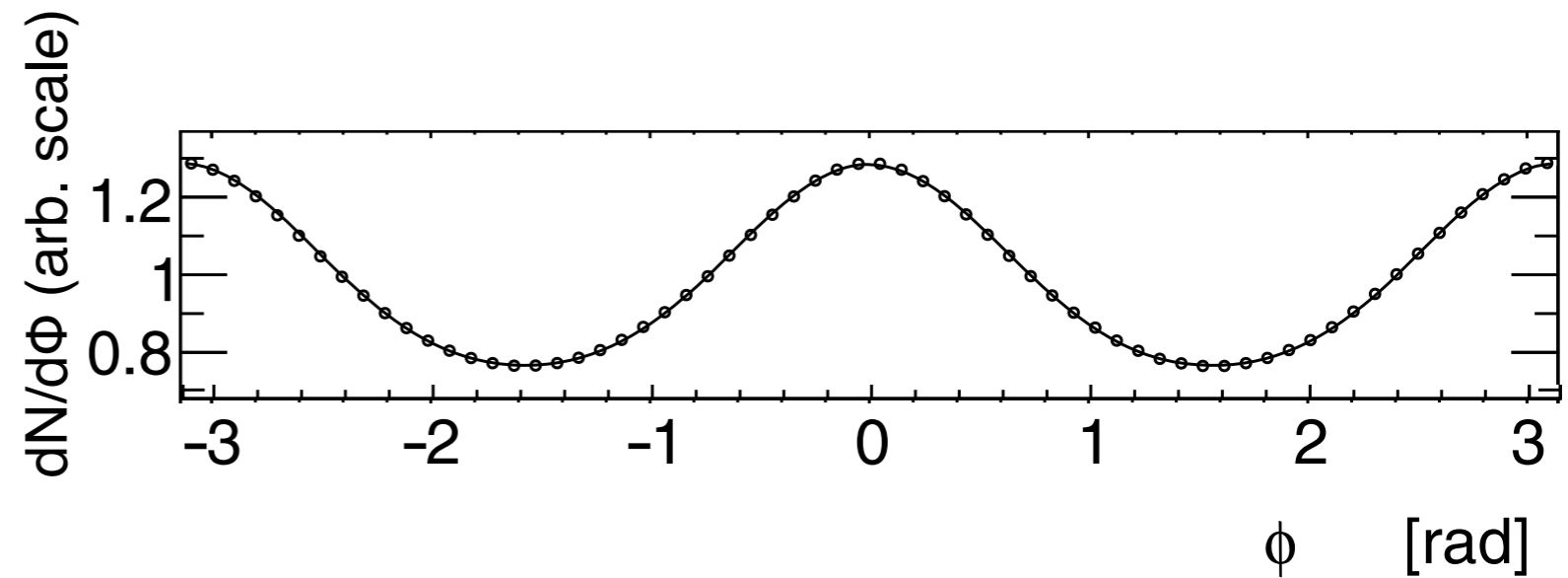


# collision geometry $\rightarrow$ measured particles

## initial collision geometry



## measured hadron distributions

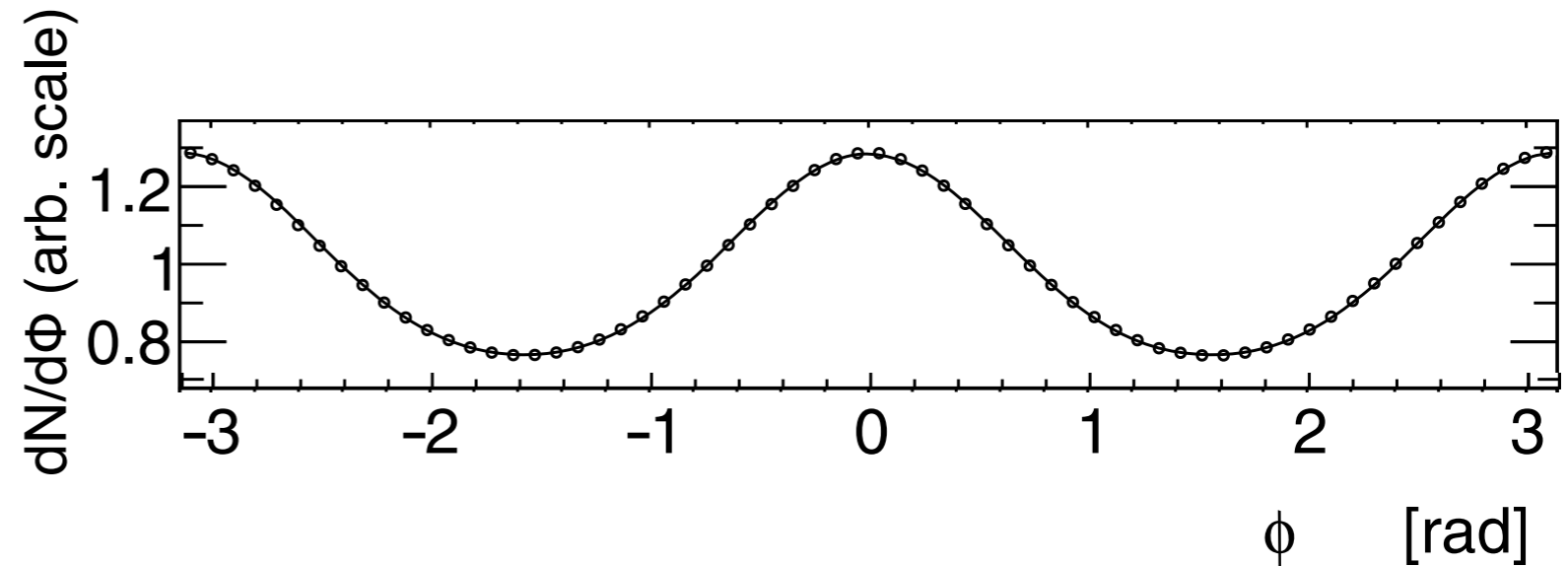
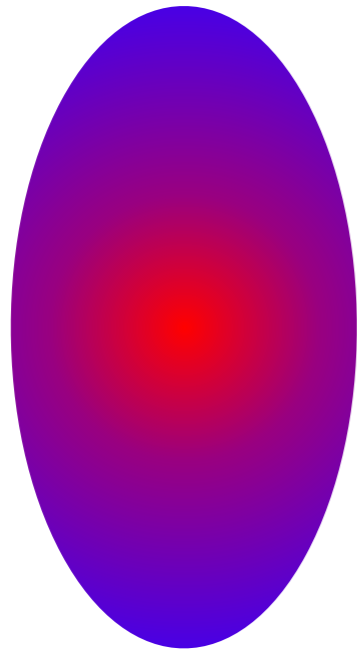


**the shape of the initial collision geometry is imprinted on the final particle distributions**



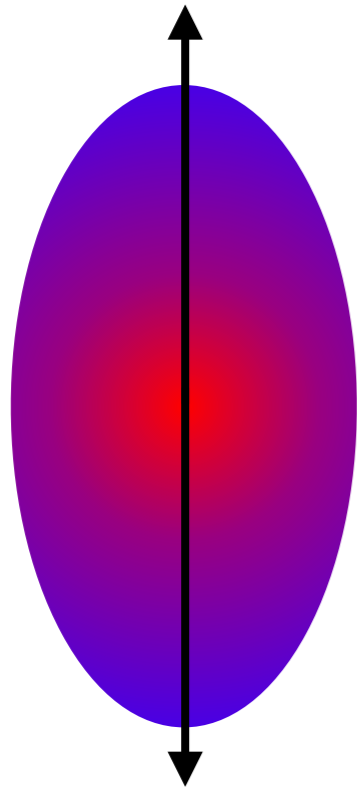
# strong interactions

---

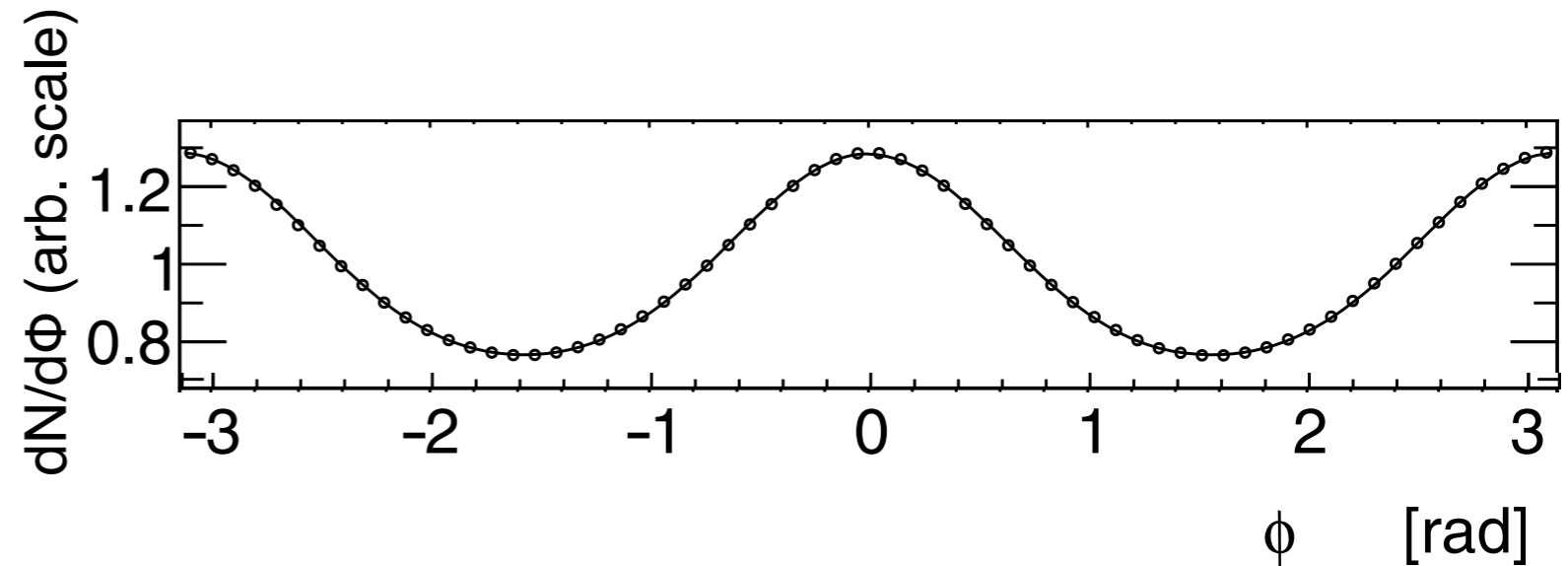


- large observed anisotropies  $\rightarrow$  strong interactions:
  - **fluid behavior, hydrodynamics**
- larger pressure gradients push more particles out in the x direction than in y

# strong interactions

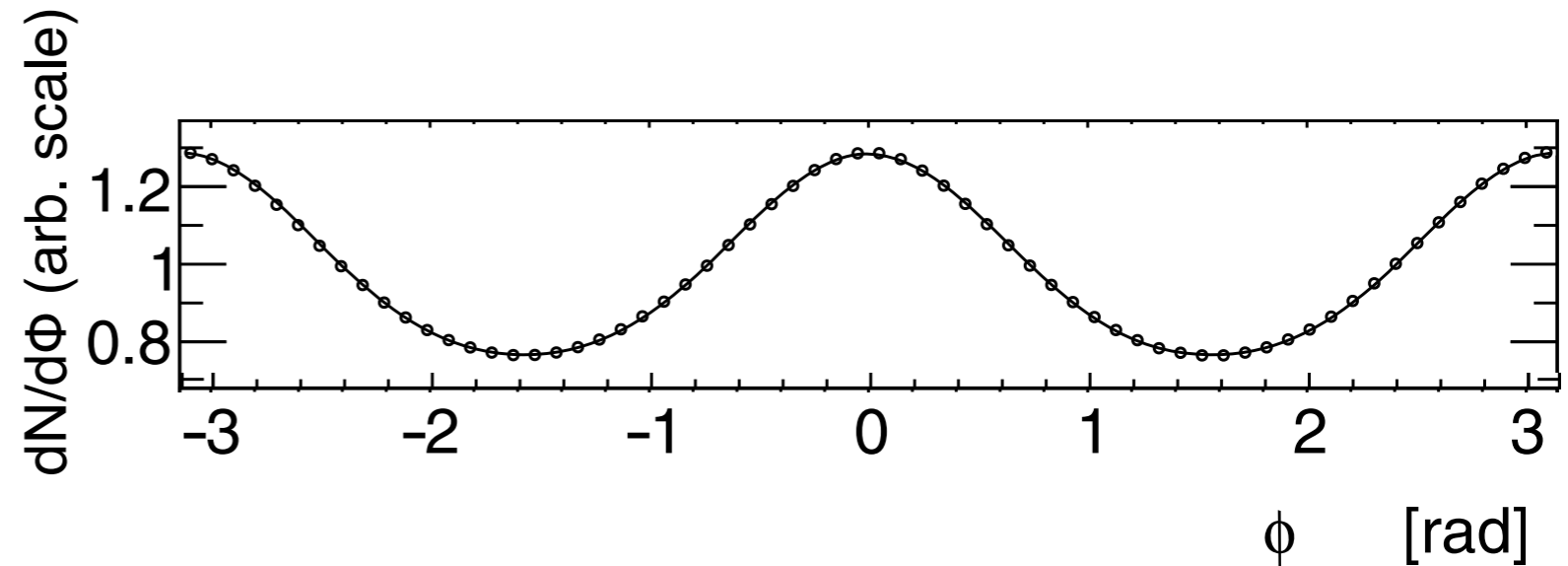
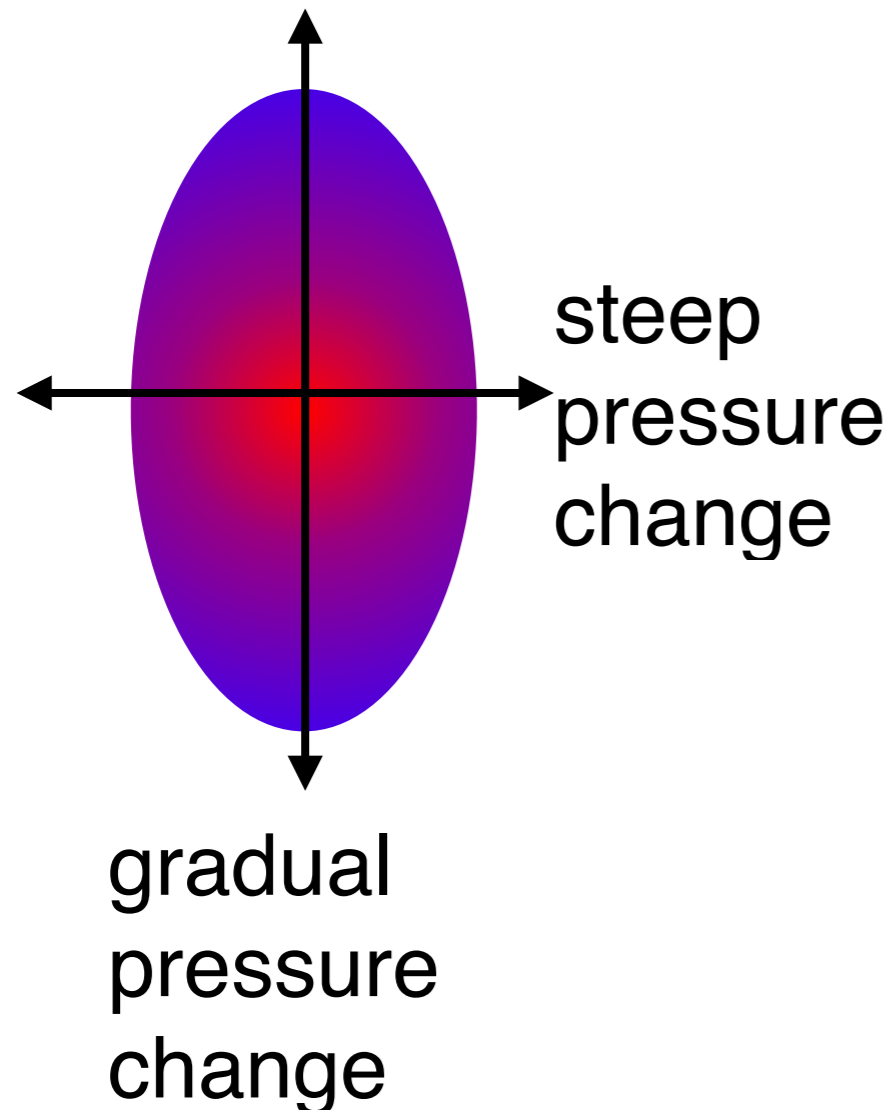


gradual  
pressure  
change



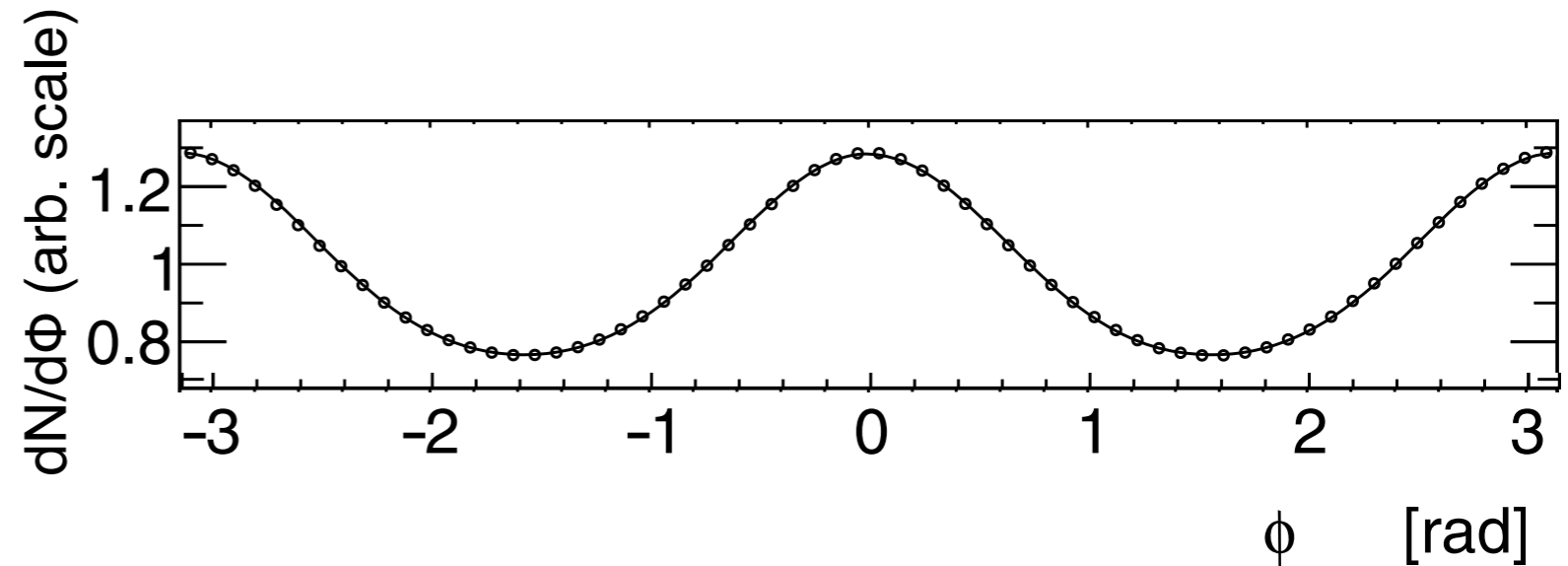
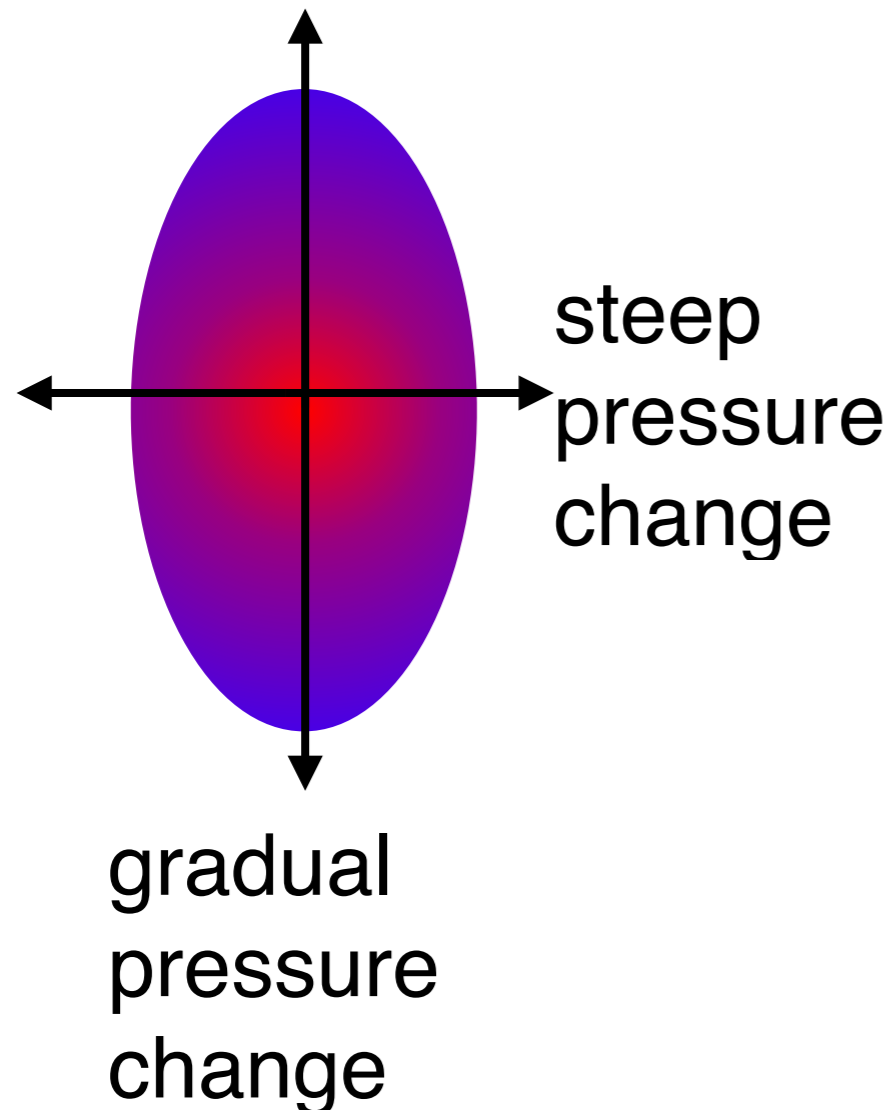
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# strong interactions



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# strong interactions

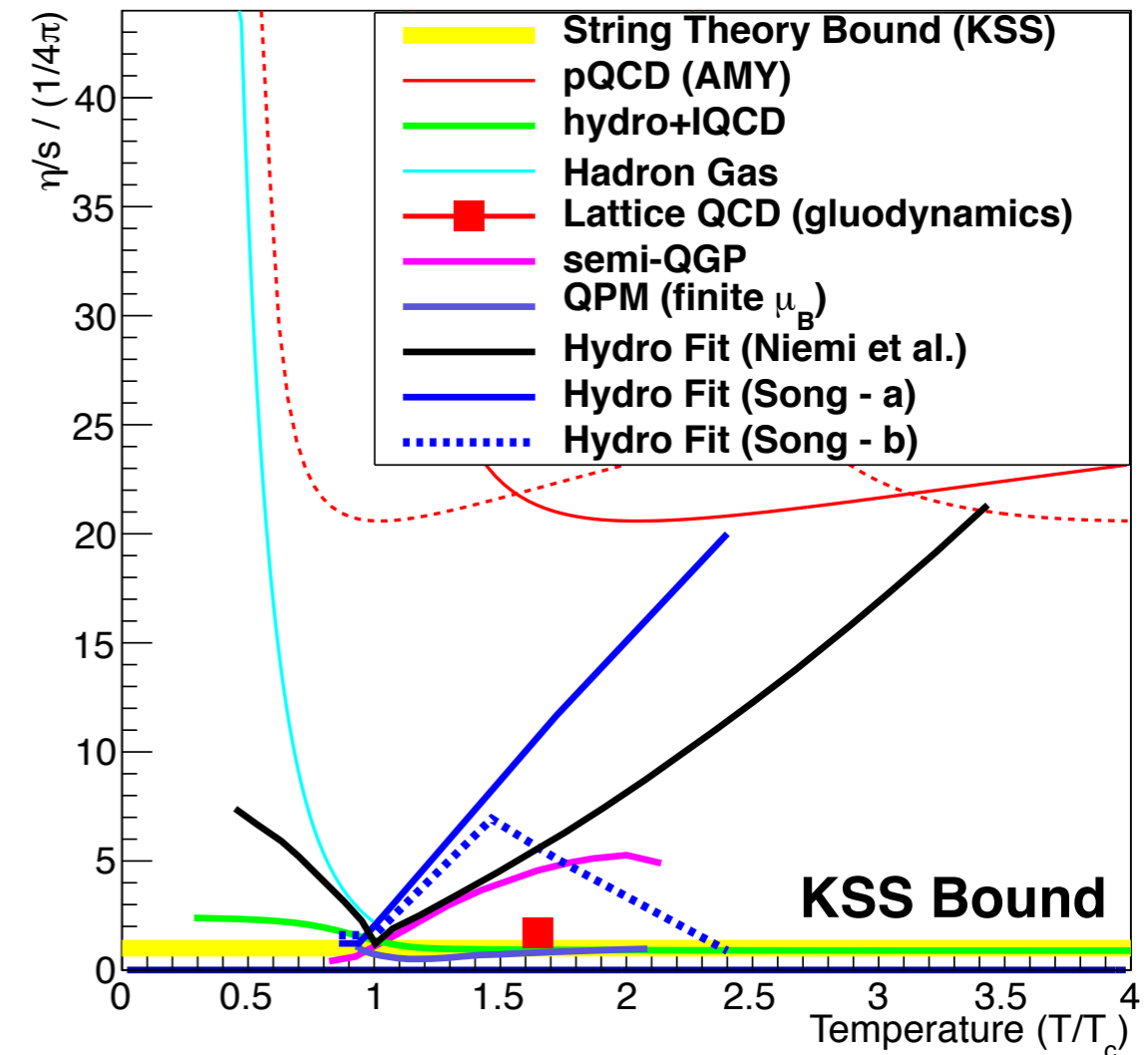


$$\frac{dN}{d\phi} = 1 + 2v_2 \cos 2\phi$$

- large observed anisotropies → strong interactions:
  - **fluid behavior, hydrodynamics**
- larger pressure gradients push more particles out in the x direction than in y

# the viscosity of the QGP

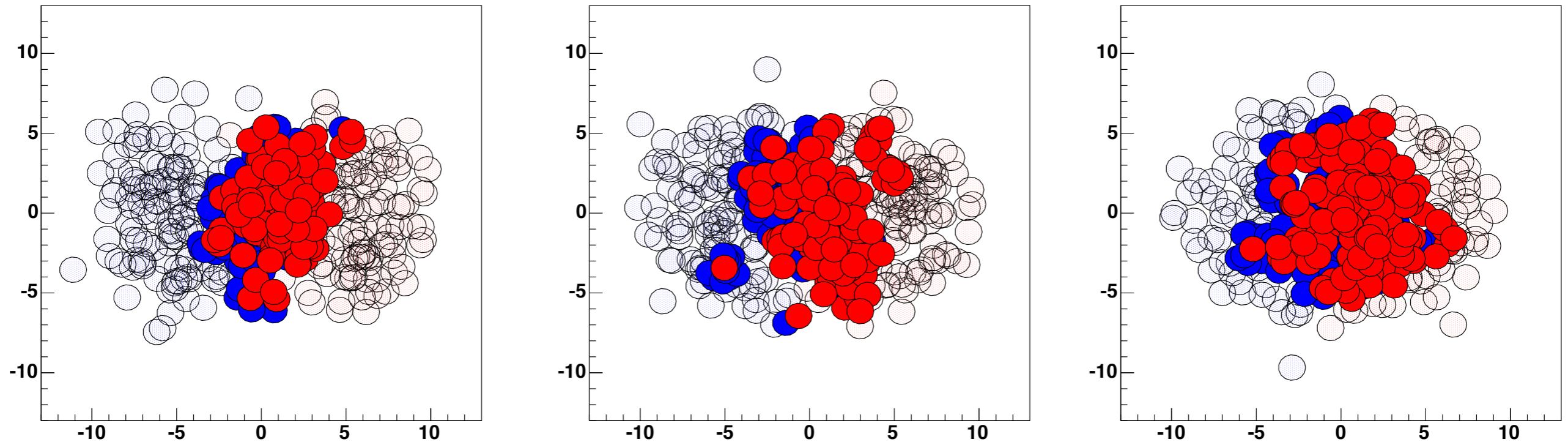
- what kind of fluid is the QGP?
- more like water or honey?
- characterize by ratio of shear viscosity to entropy density:  $\eta/s$
- we know that  $\eta/s(\text{QGP})$  is very small near the critical T
- but how does that change with temperature?



# each collision is unique

---

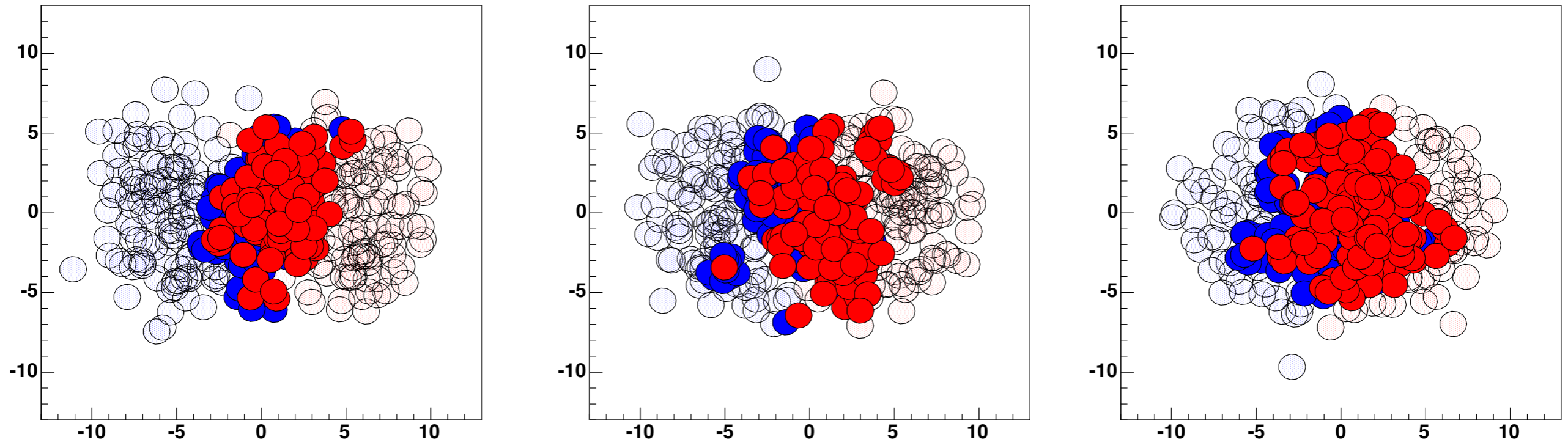
nucleon distributions for 3 single collisions (xy-plane)



each collision evolves in isolation without knowing what the “typical” collision is

# each collision is unique

nucleon distributions for 3 single collisions (xy-plane)



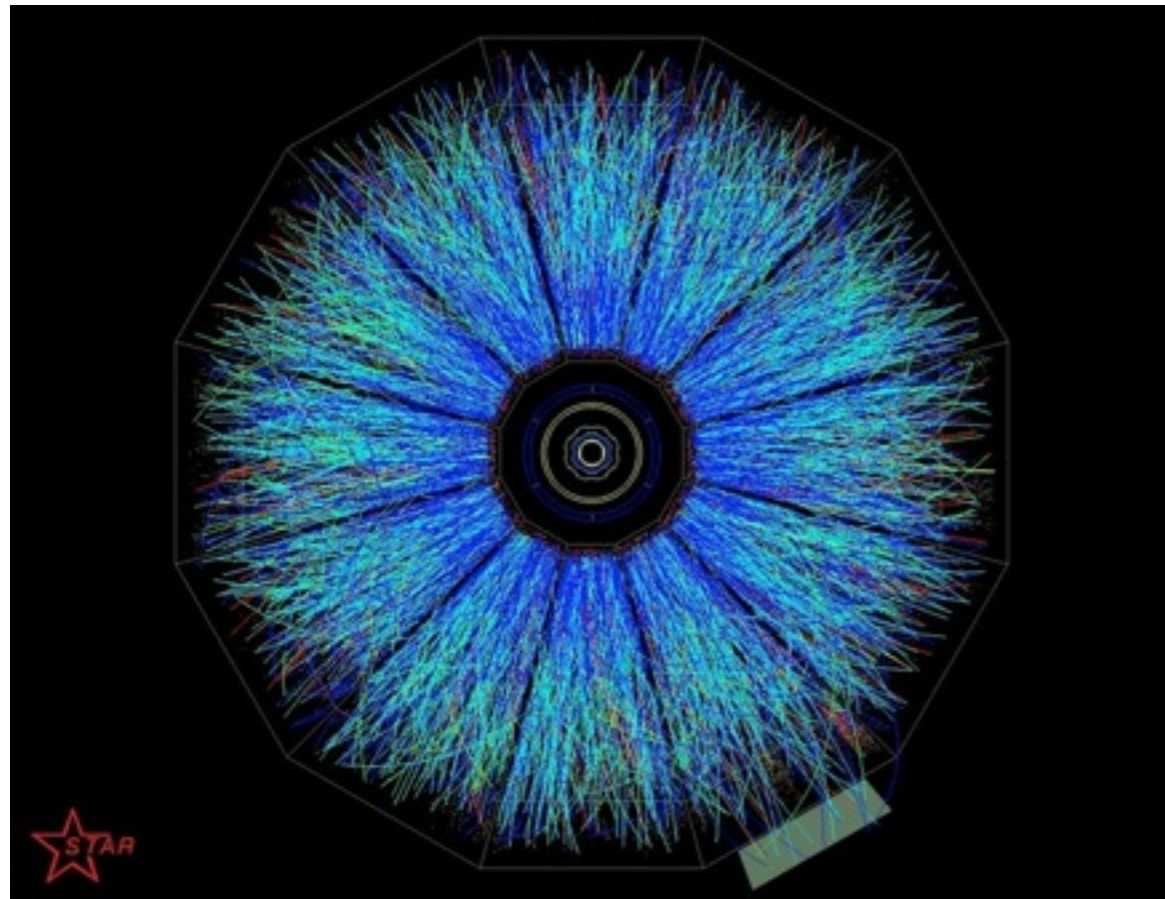
each collision evolves in isolation without knowing what the “typical” collision is

not just  $v_2$  describing  $\cos 2\Phi$ , but  $v_n$ :

$$\frac{dN}{d\phi} \propto 1 + \sum_n 2v_n \cos n(\phi - \Psi_n)$$

# two particle correlations

---



$$\frac{dN}{d\phi} \propto 1 + \sum^n 2v_n \cos n (\phi - \Psi_n)$$

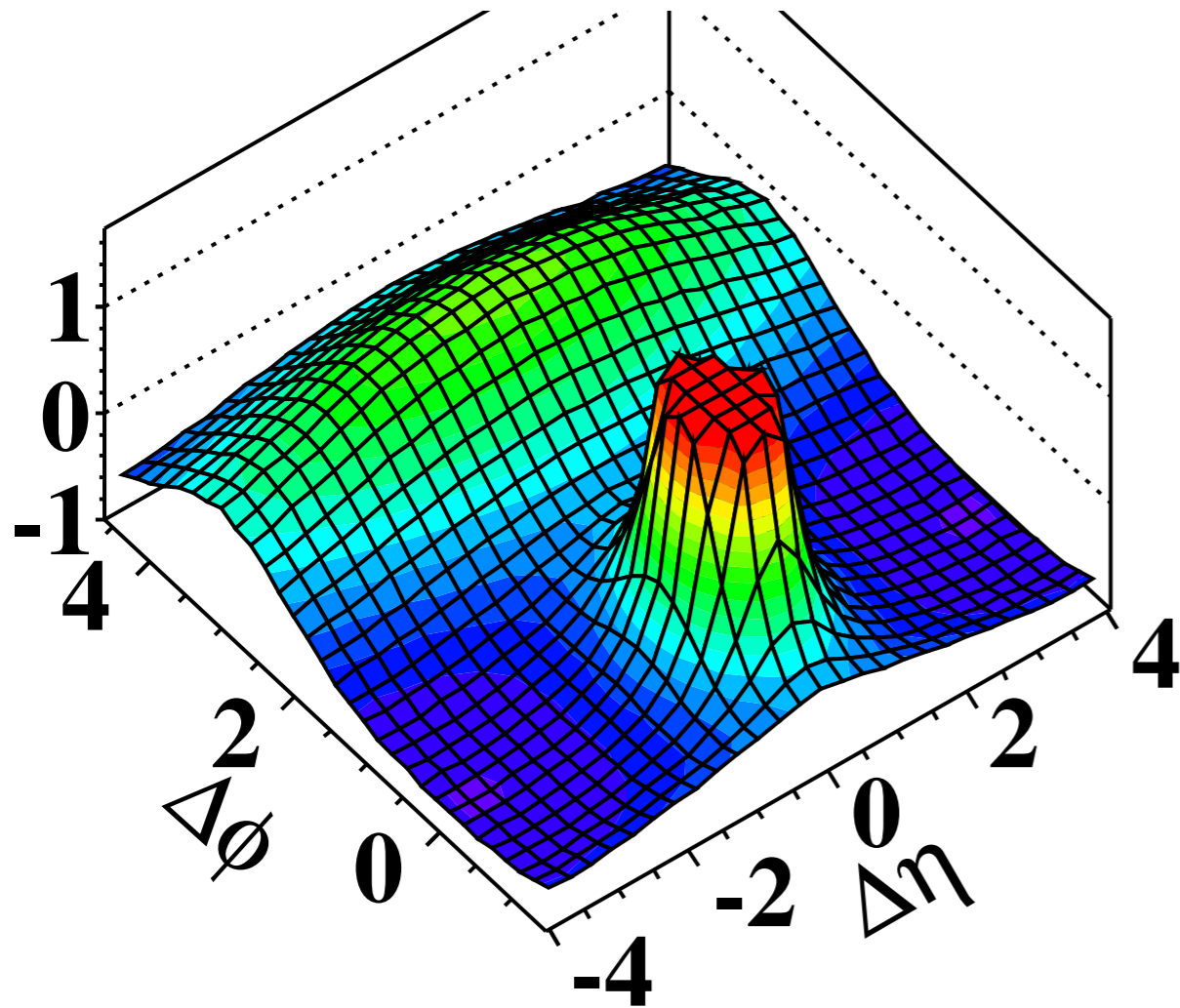
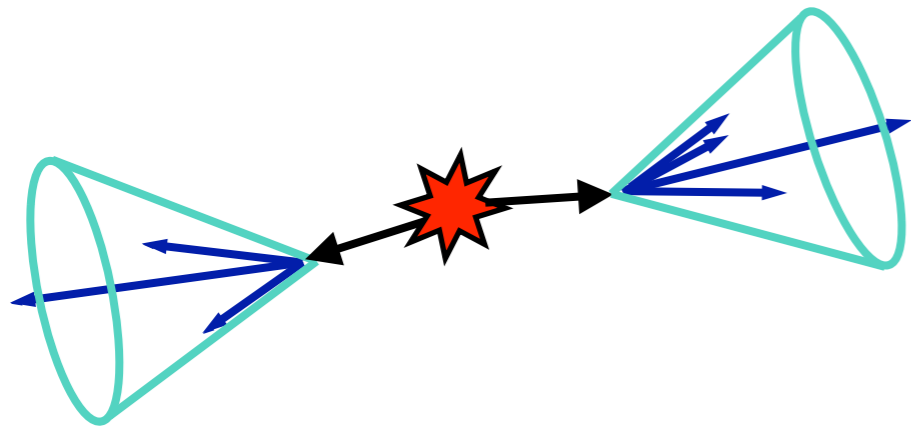


$$\frac{dN_{AB}}{d\Delta\phi} \propto 1 + \sum^n 2v_{n,A}v_{n,B} \cos (n\Delta\phi)$$



# two particle correlations

jets in pp collisions



flow

single particles

$$\frac{dN}{d\phi} \propto 1 + \sum_n 2v_n \cos n(\phi - \Psi_n)$$

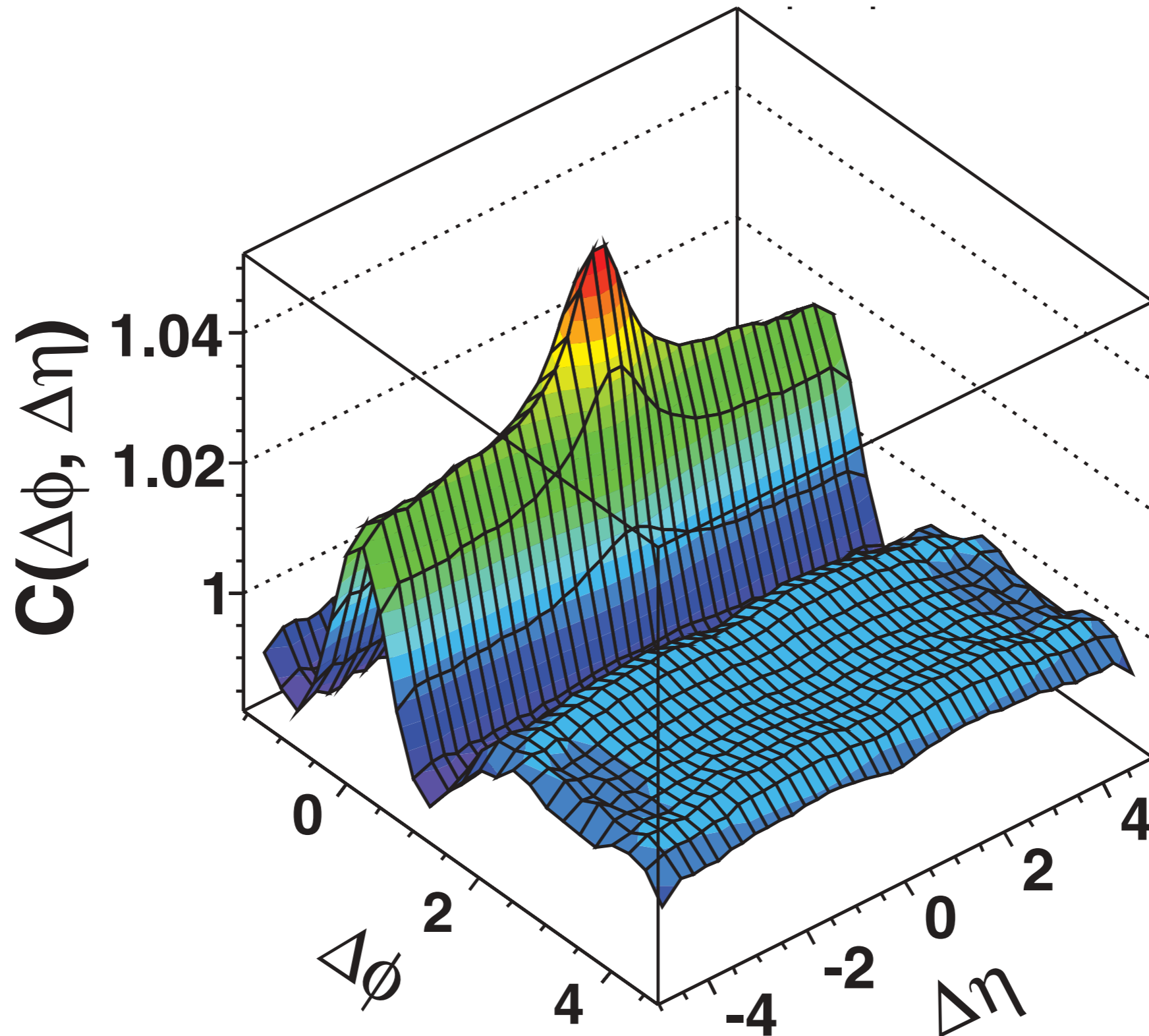
pairs of particles

$$\frac{dN_{AB}}{d\Delta\phi} \propto 1 + \sum_n 2v_{n,A}v_{n,B} \cos(n\Delta\phi)$$

flow correlations should be long range  $\eta$

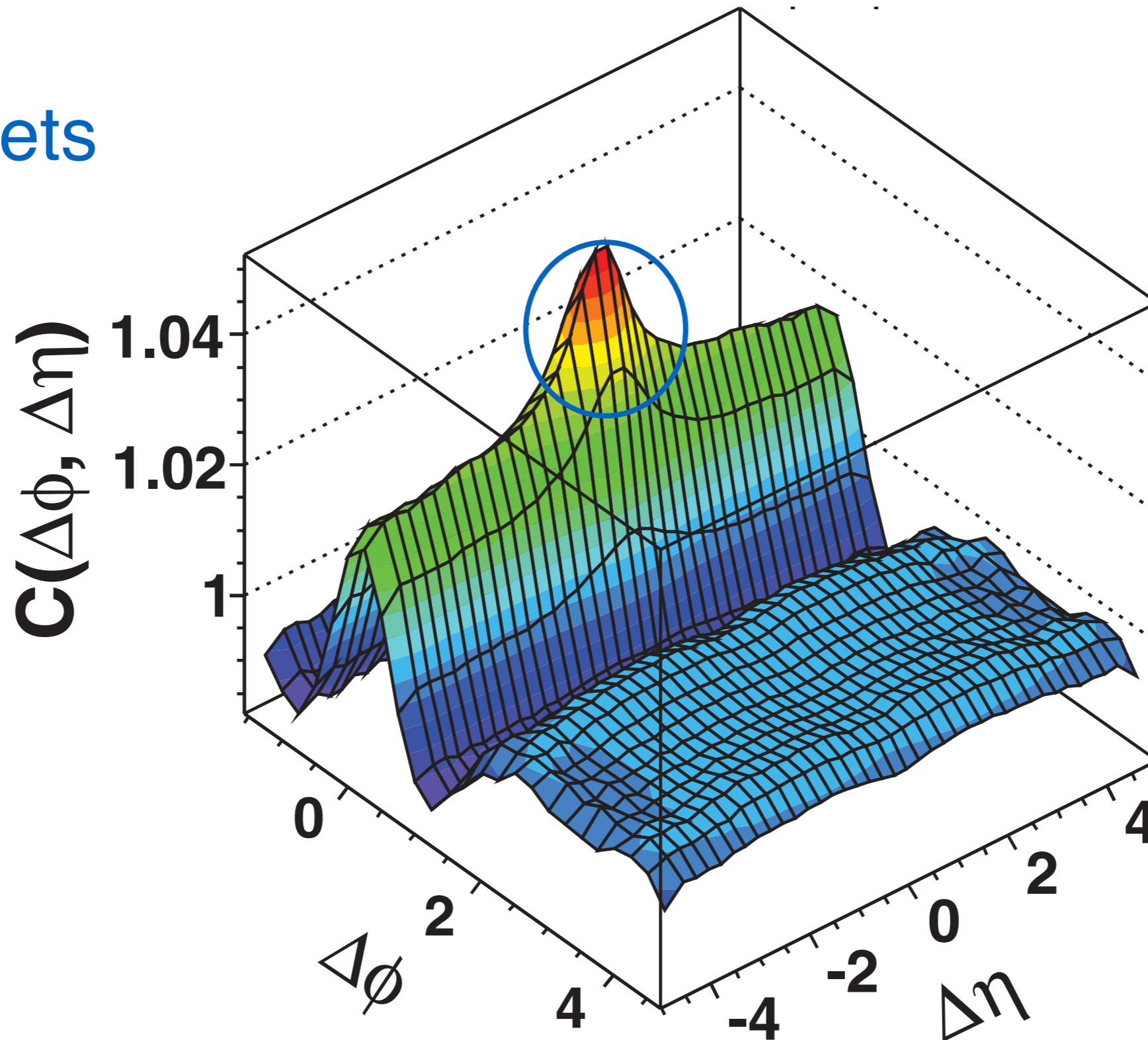
# correlations in PbPb

---



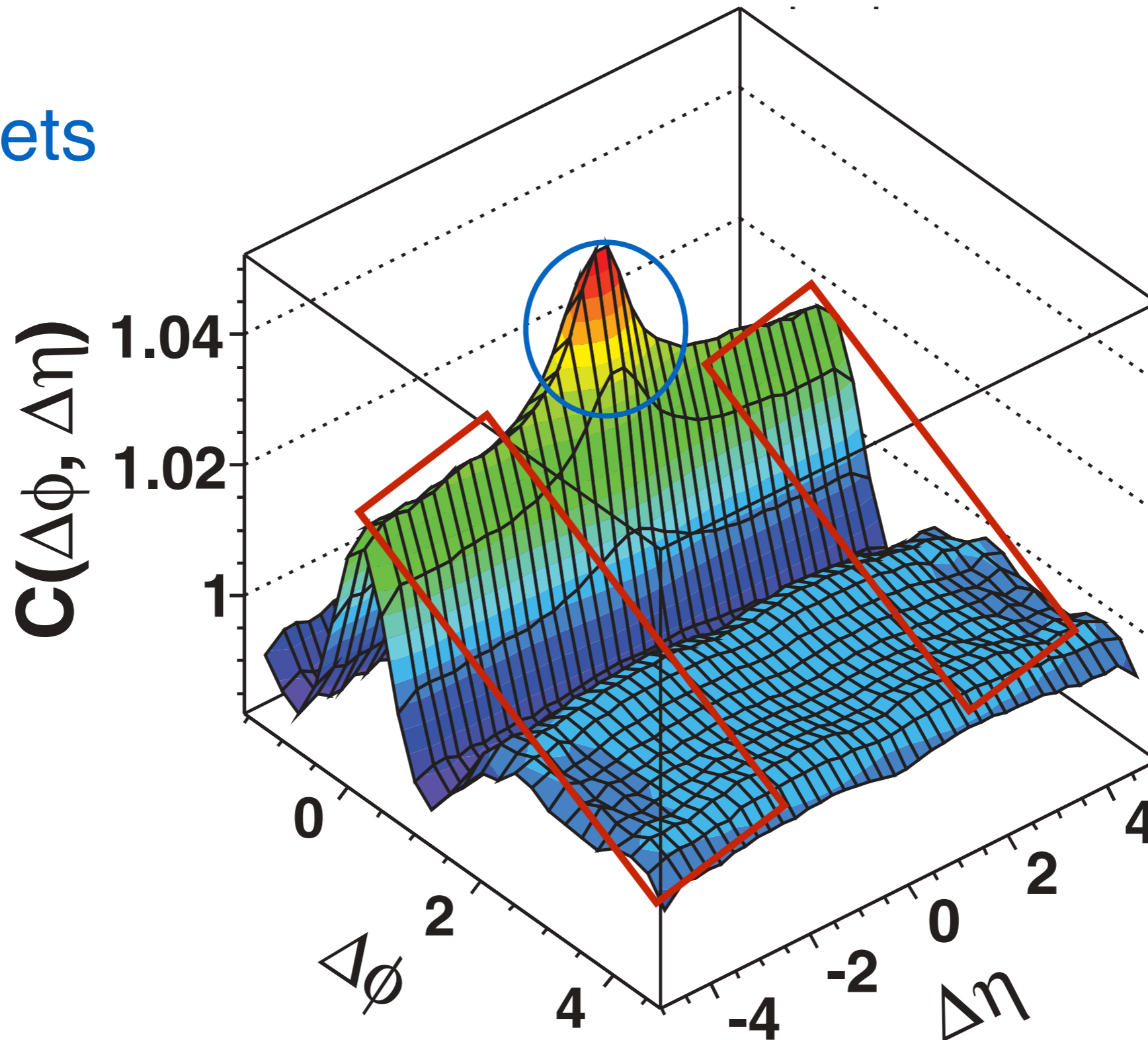
# correlations in PbPb

jets



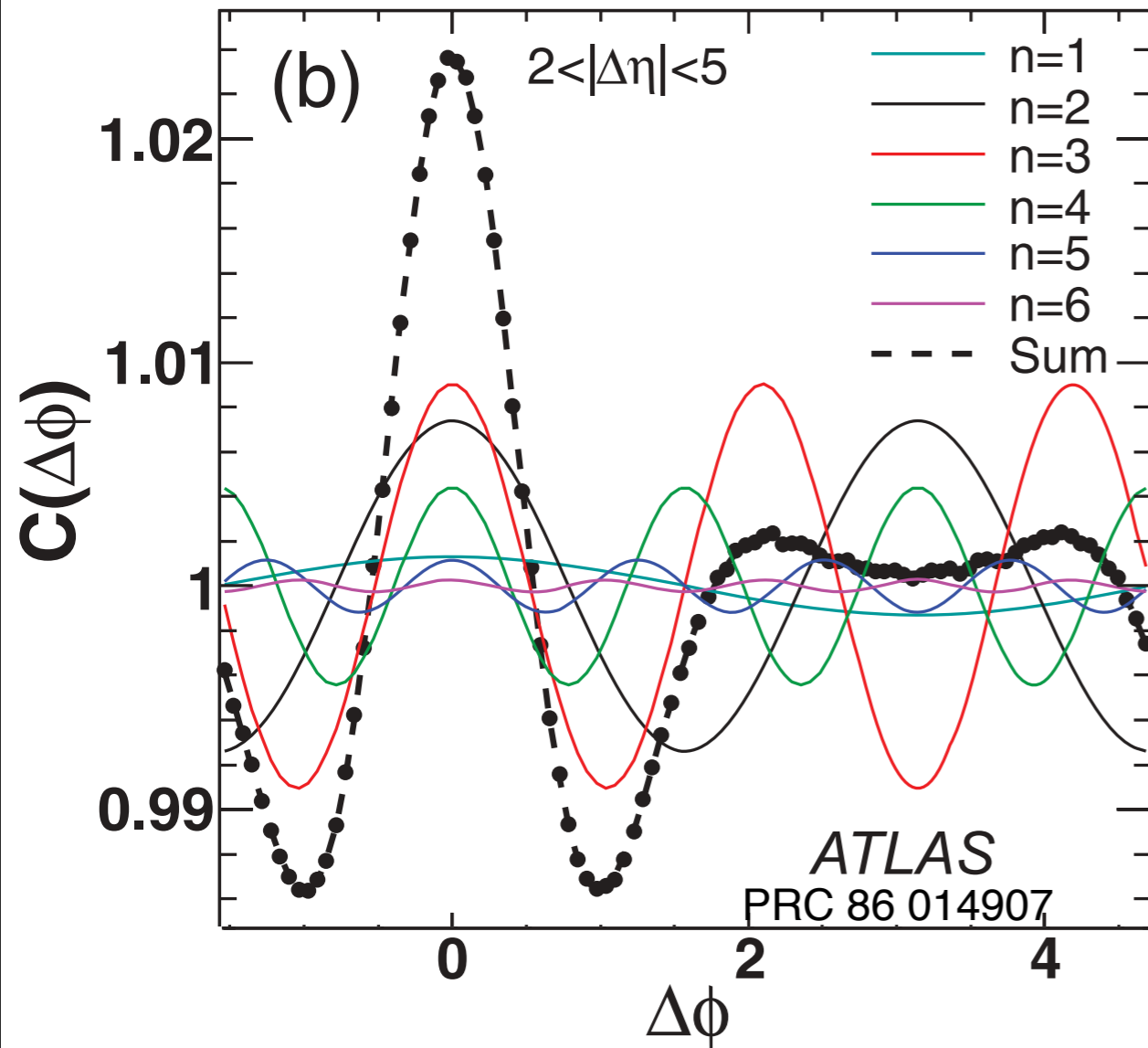
# correlations in PbPb

jets



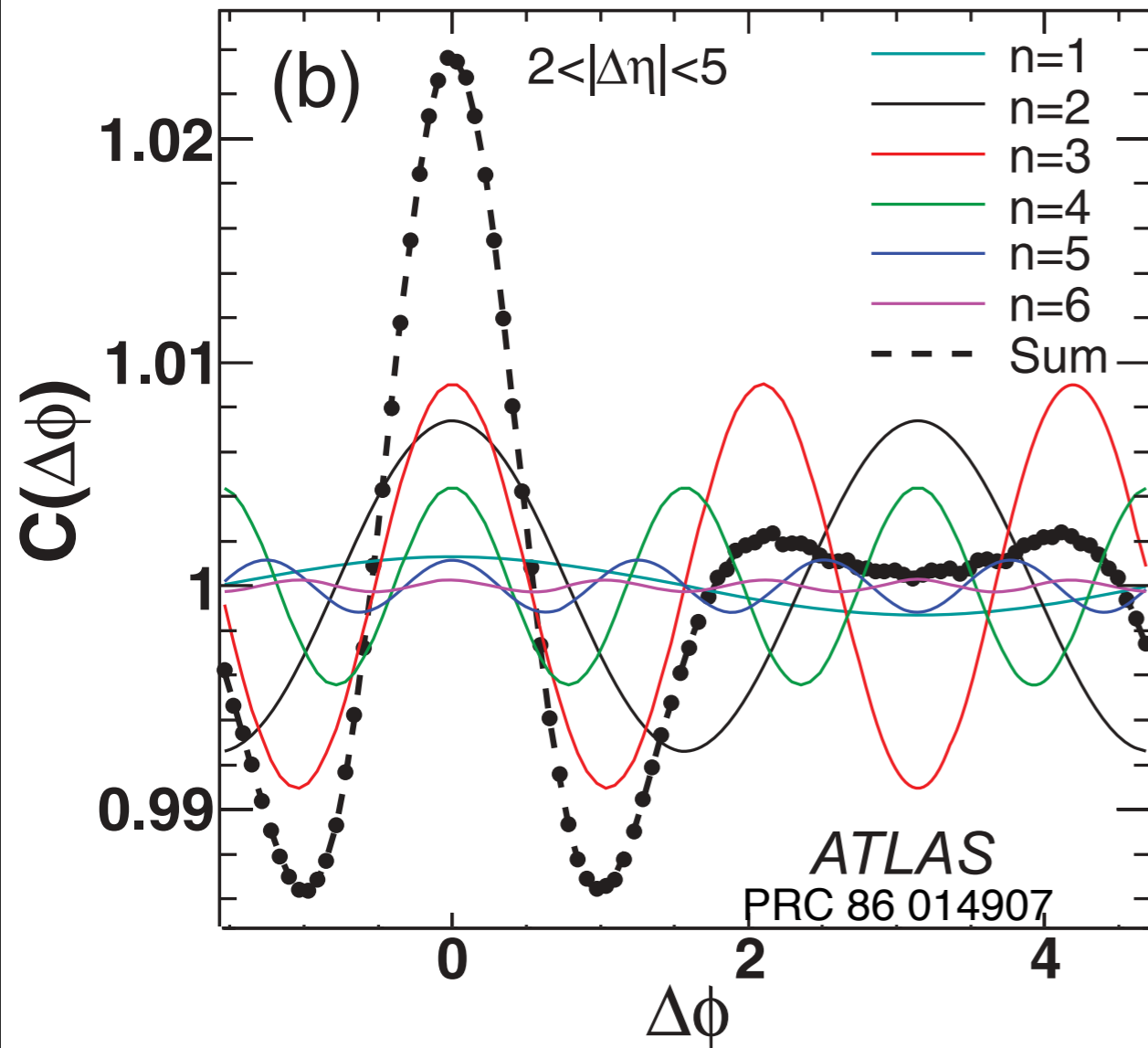
flow

# ridge: $v_N$ & two particle correlations



$$\frac{dN_{AB}}{d\Delta\phi} \propto 1 + \sum^n 2v_{n,A}v_{n,B} \cos(n\Delta\phi)$$

# ridge: $v_N$ & two particle correlations

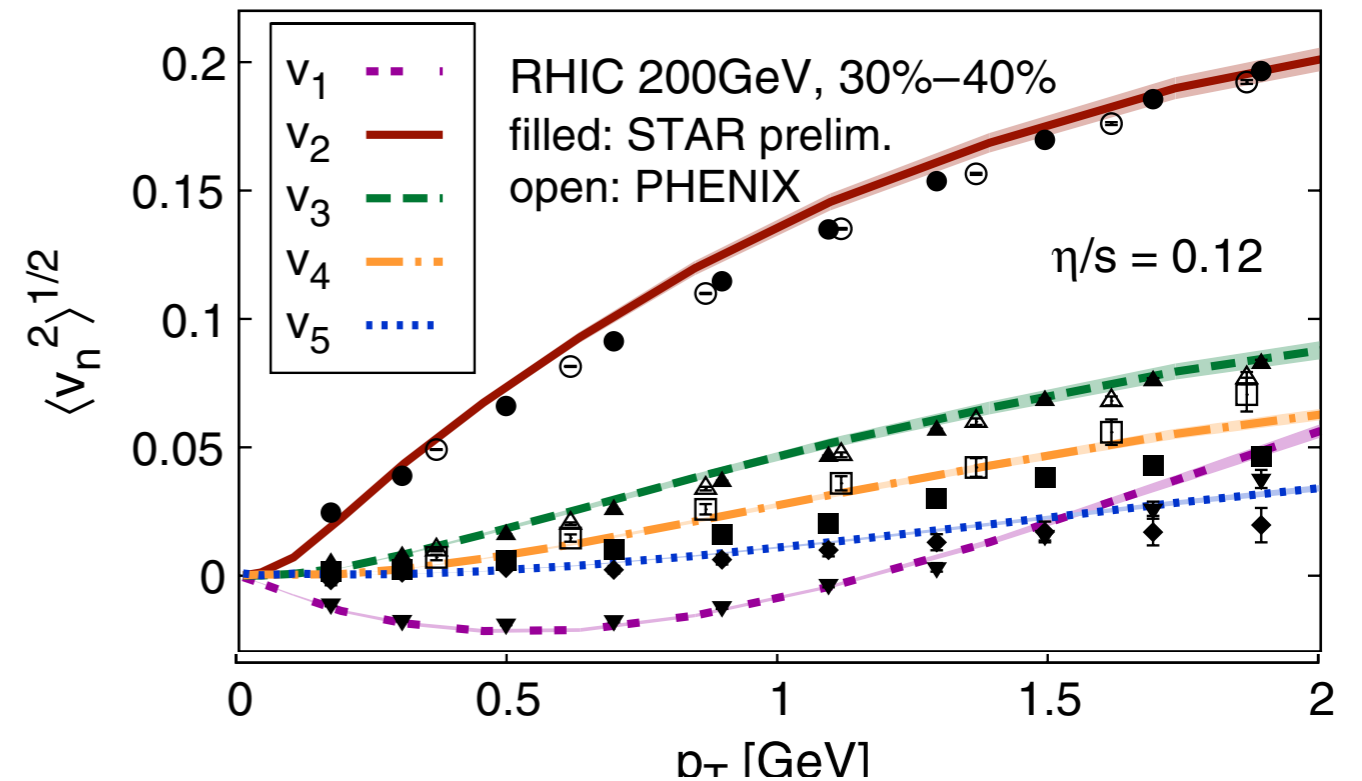
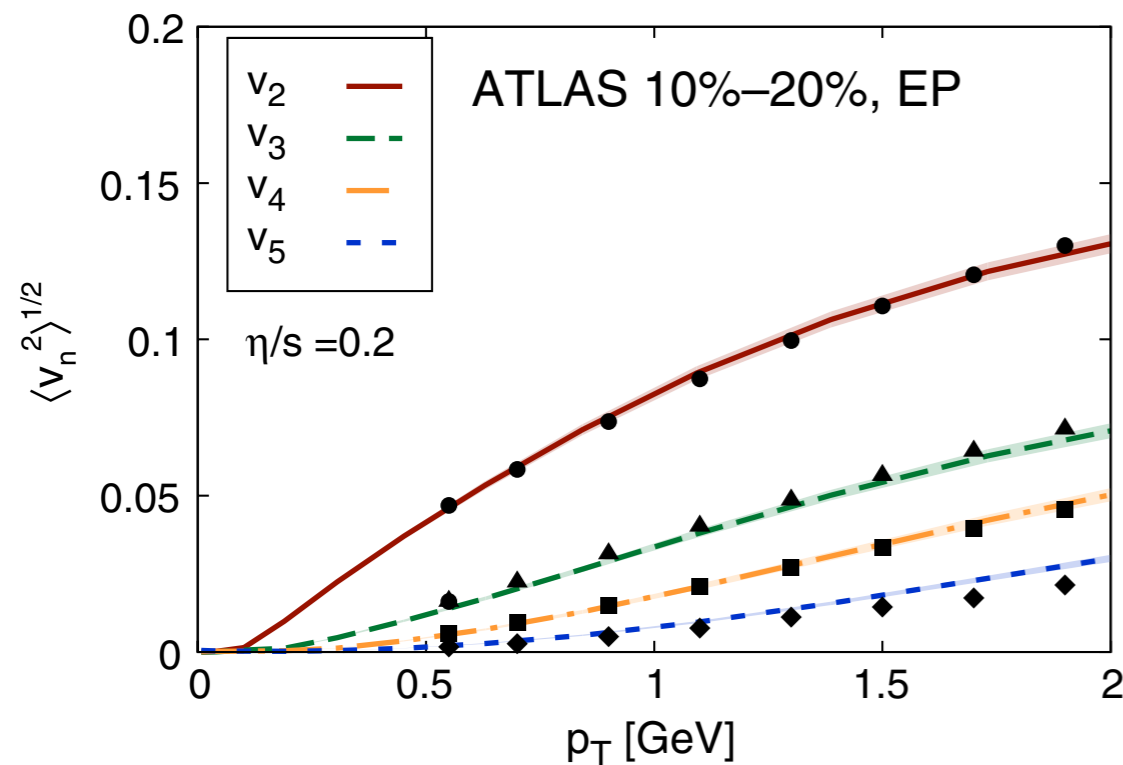


$$\frac{dN_{AB}}{d\Delta\phi} \propto 1 + \sum_n 2v_{n,A}v_{n,B} \cos(n\Delta\phi)$$

**evidence for many higher order terms in particle correlations**

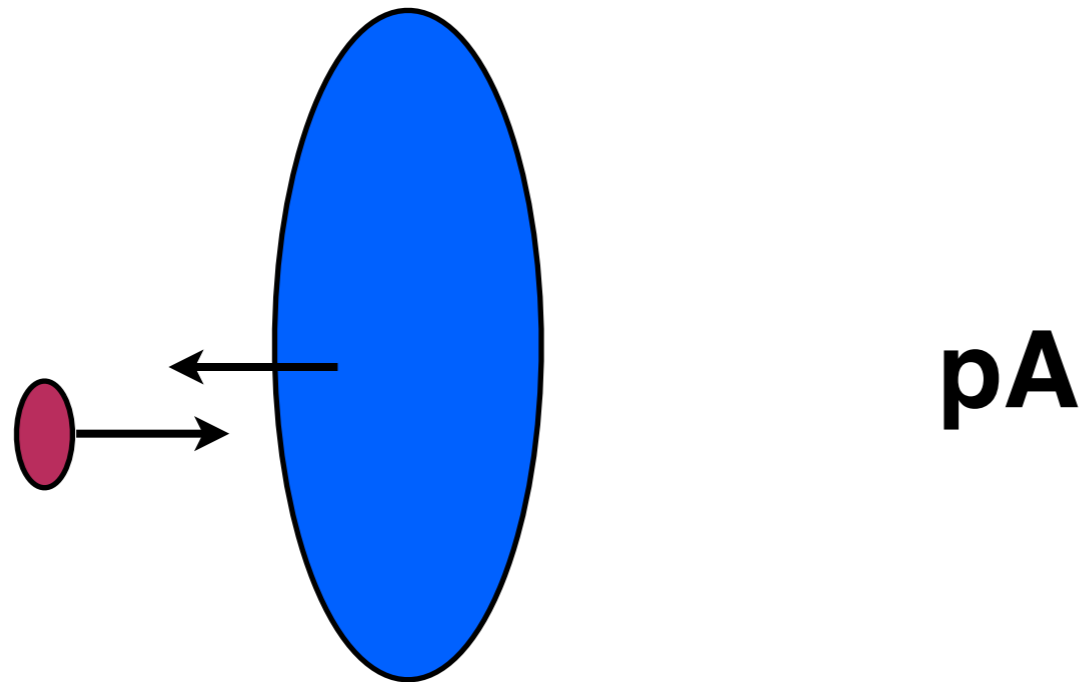
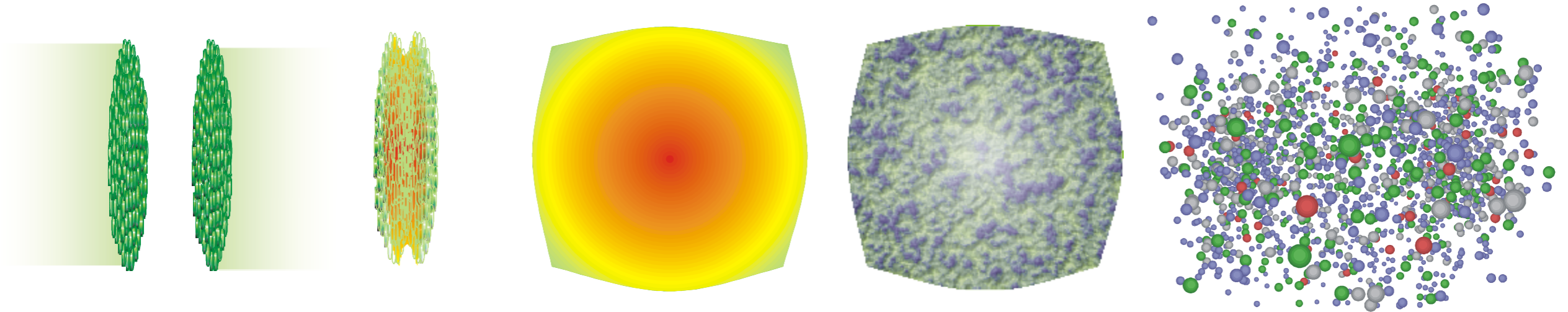
# state of the art hydrodynamic calculations

## 3 +1d viscous hydrodynamics



quantitative description of  $v_1 - v_5$  at both RHIC and LHC  
sensitivity to  $\eta/s$

# pA physics

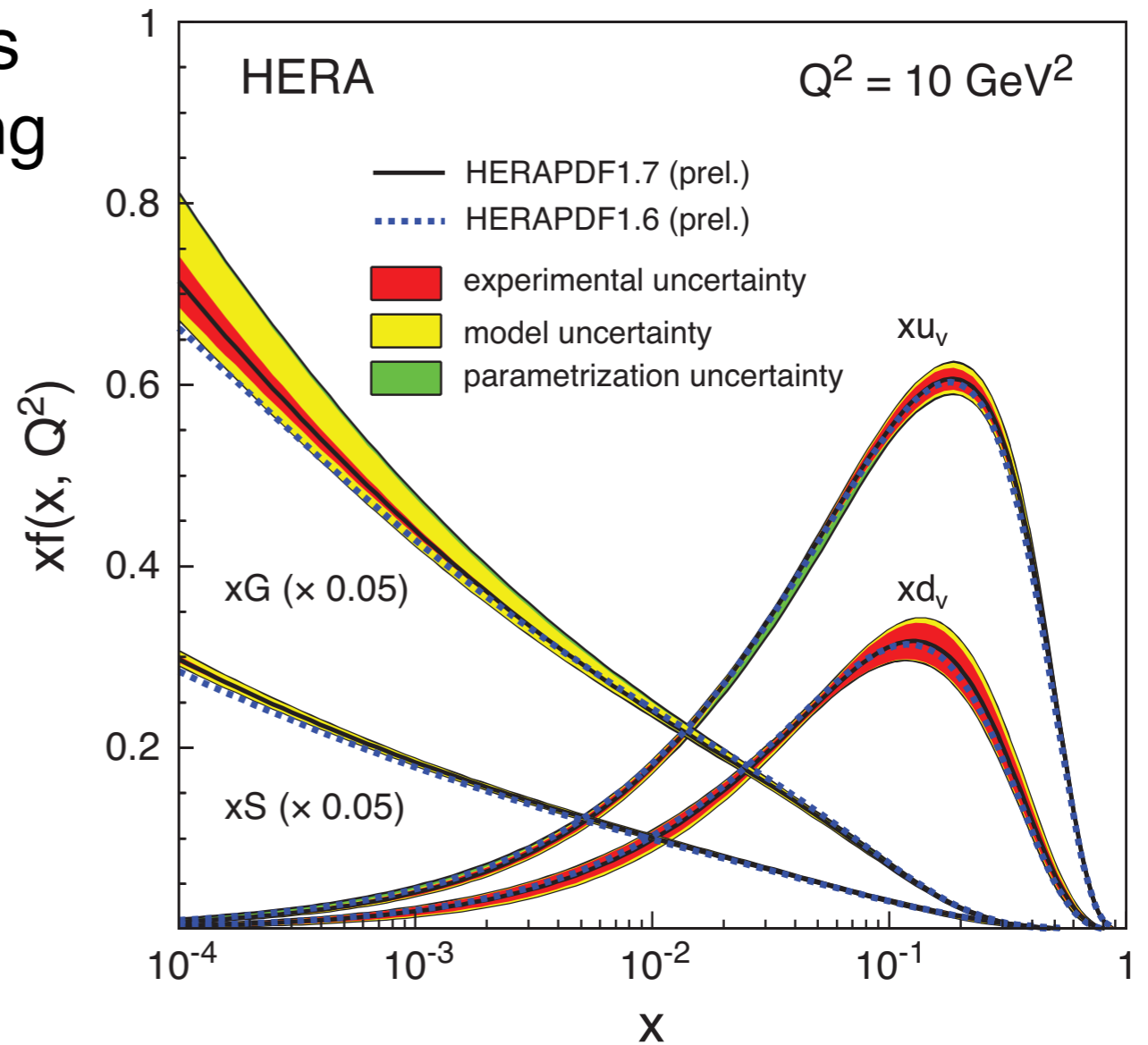
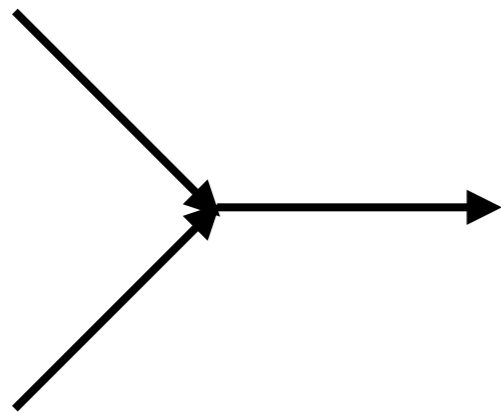


- isolate QGP effects from something present in the incoming nuclei



# saturation of low x gluons

- basic idea: the number of gluons increases quickly with decreasing  $x$ . At some point there are so many gluons that the recombination rate becomes significant, saturating the distribution

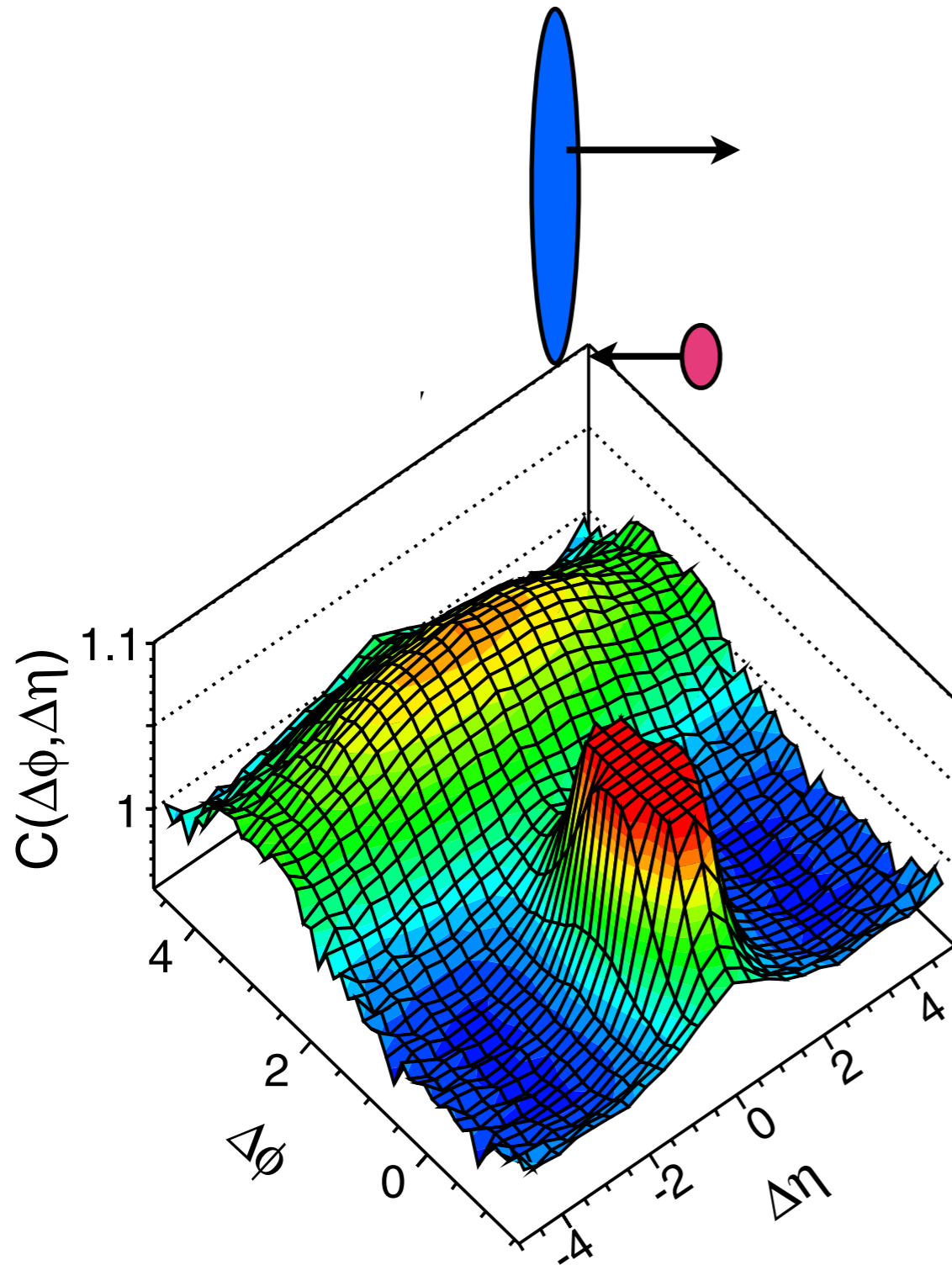


in a large nucleus in high energy collisions, this happens more readily because the nucleons overlap, increasing the density

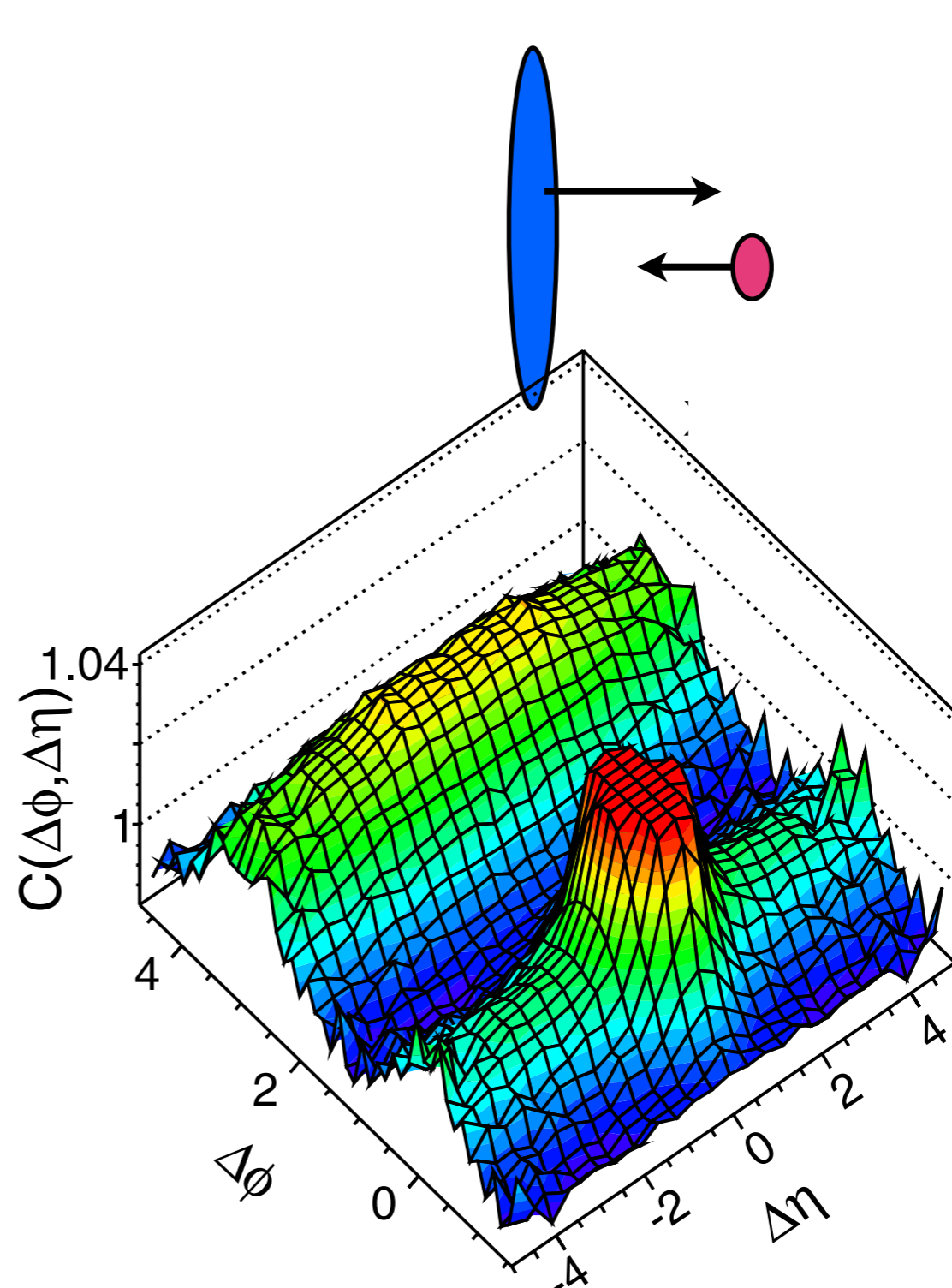
# a closer look at pPb

peripheral collisions

central collisions

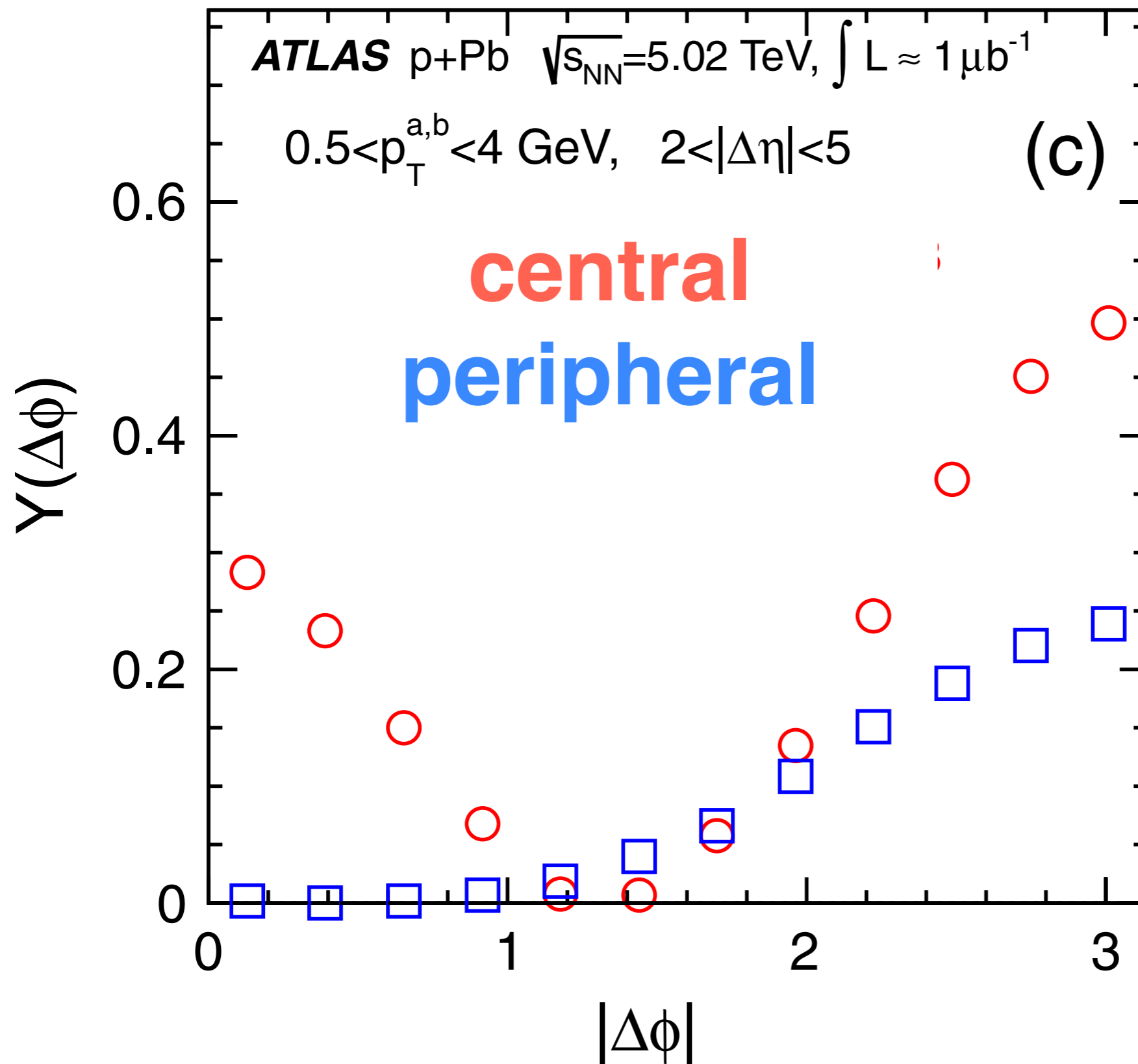


jets

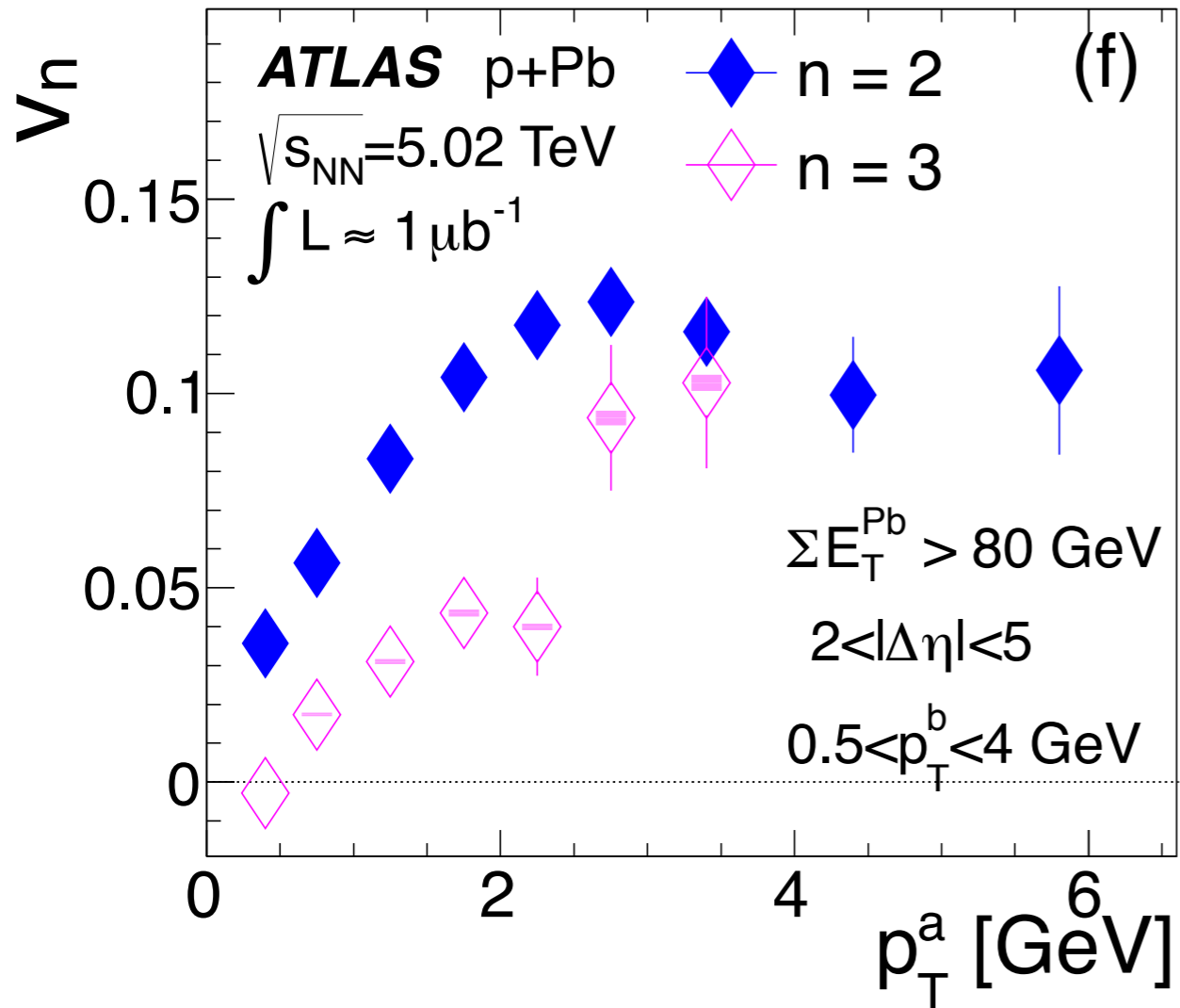


jets + flow

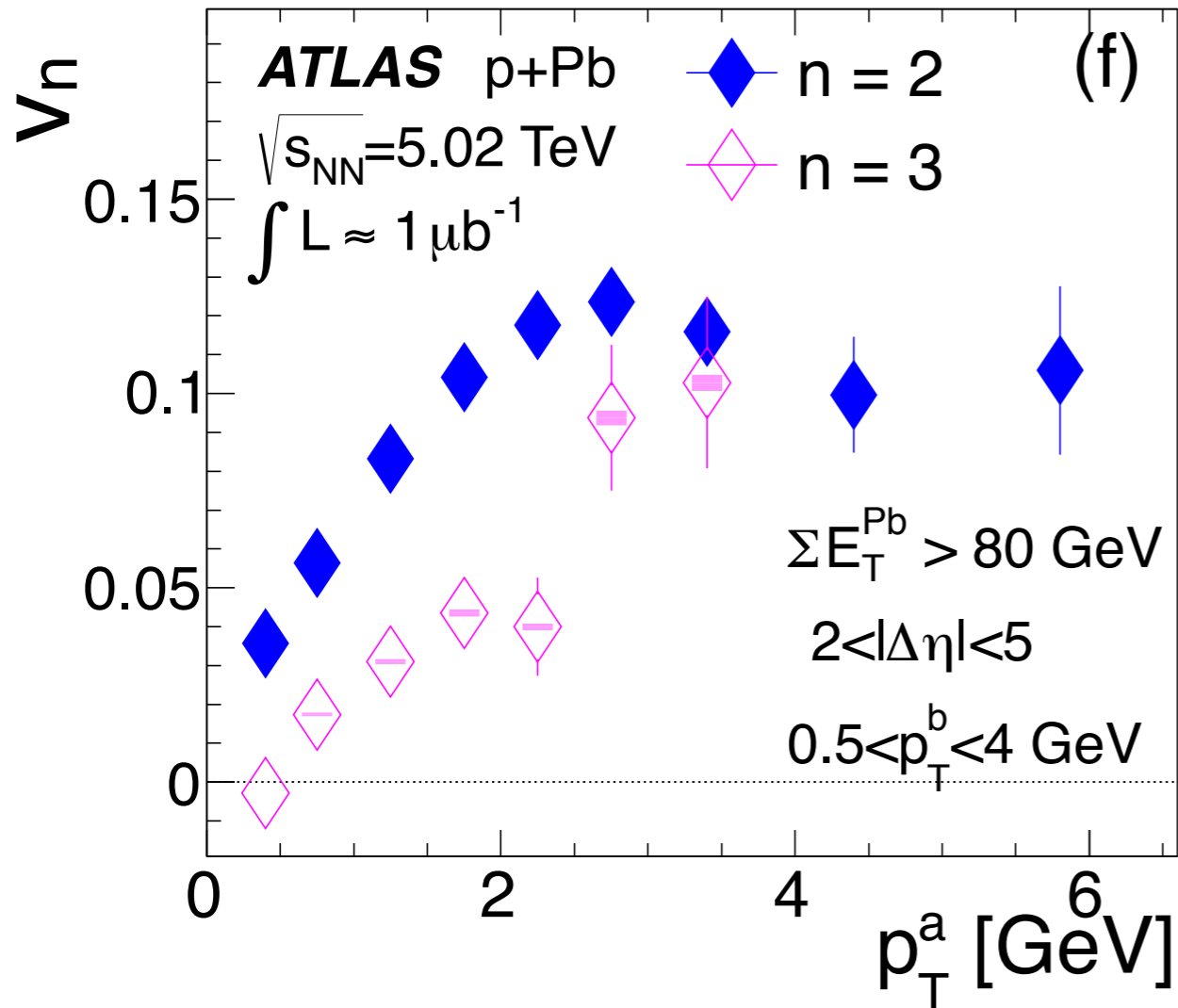
# a closer look at pPb



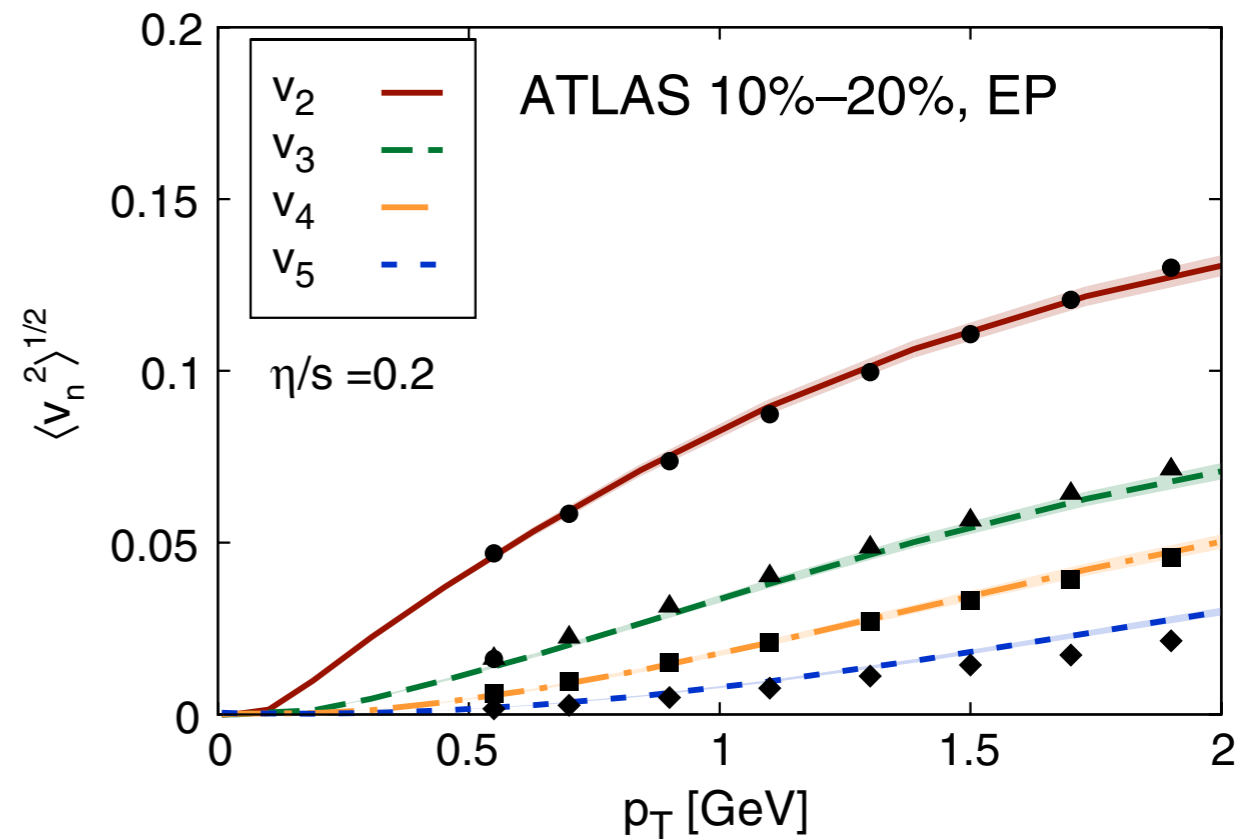
# $v_2$ & $v_3$ in pPb collisions



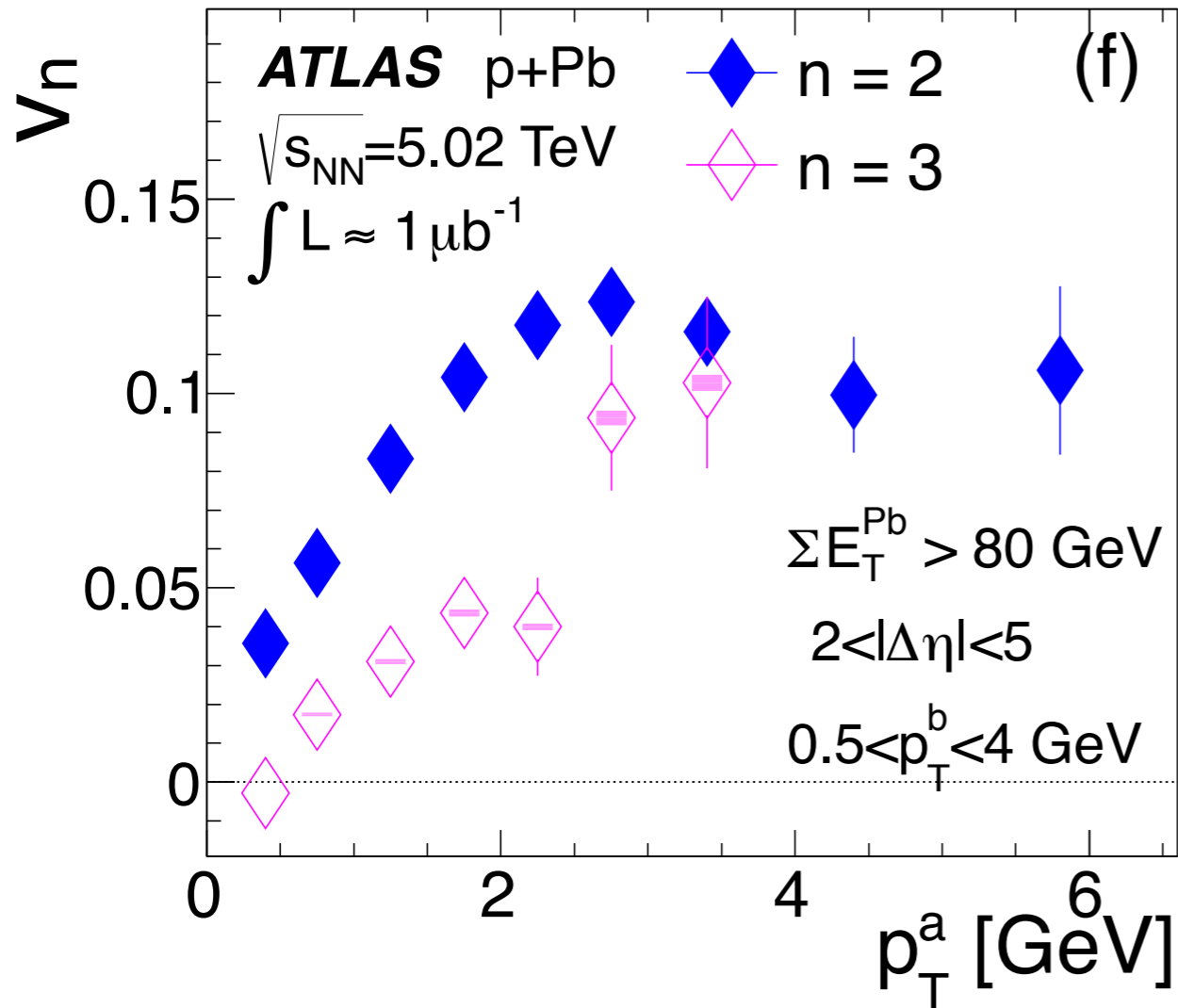
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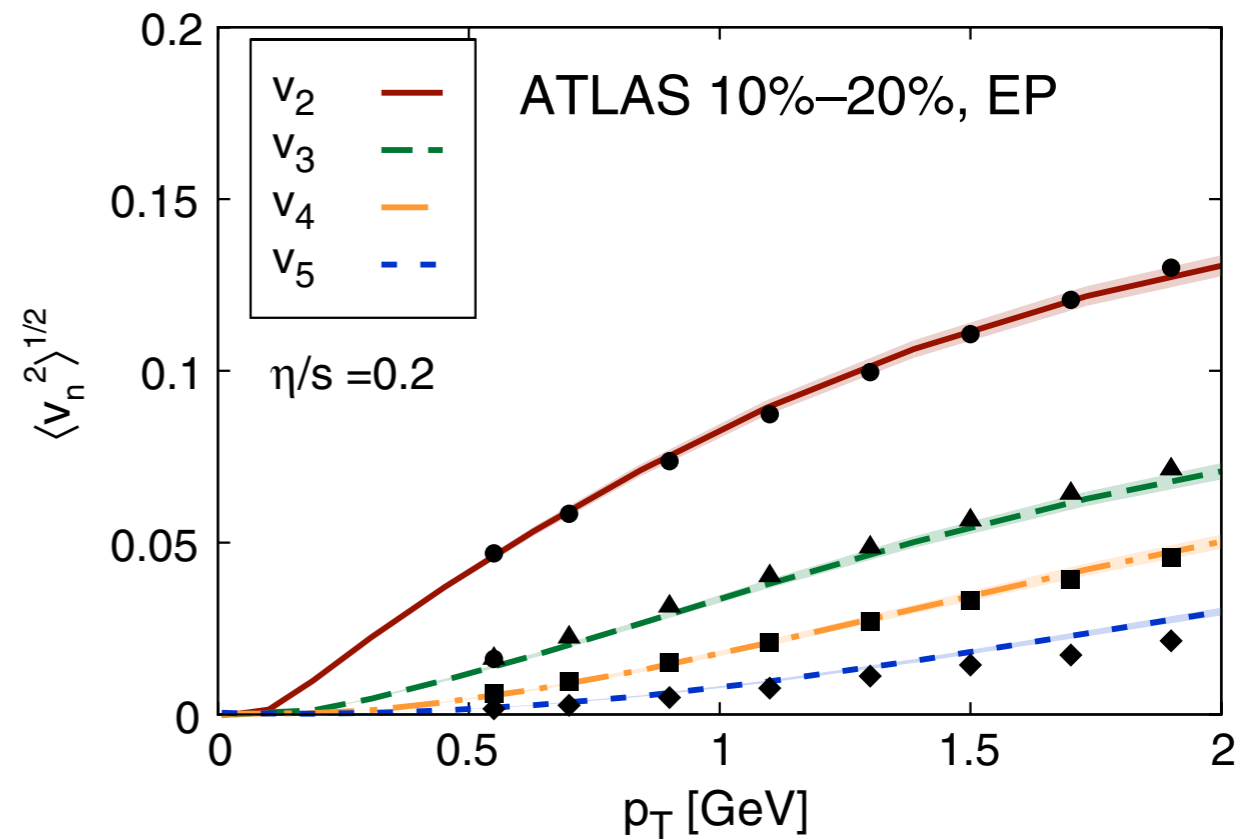
very similar to AA results



# $v_2$ & $v_3$ in pPb collisions



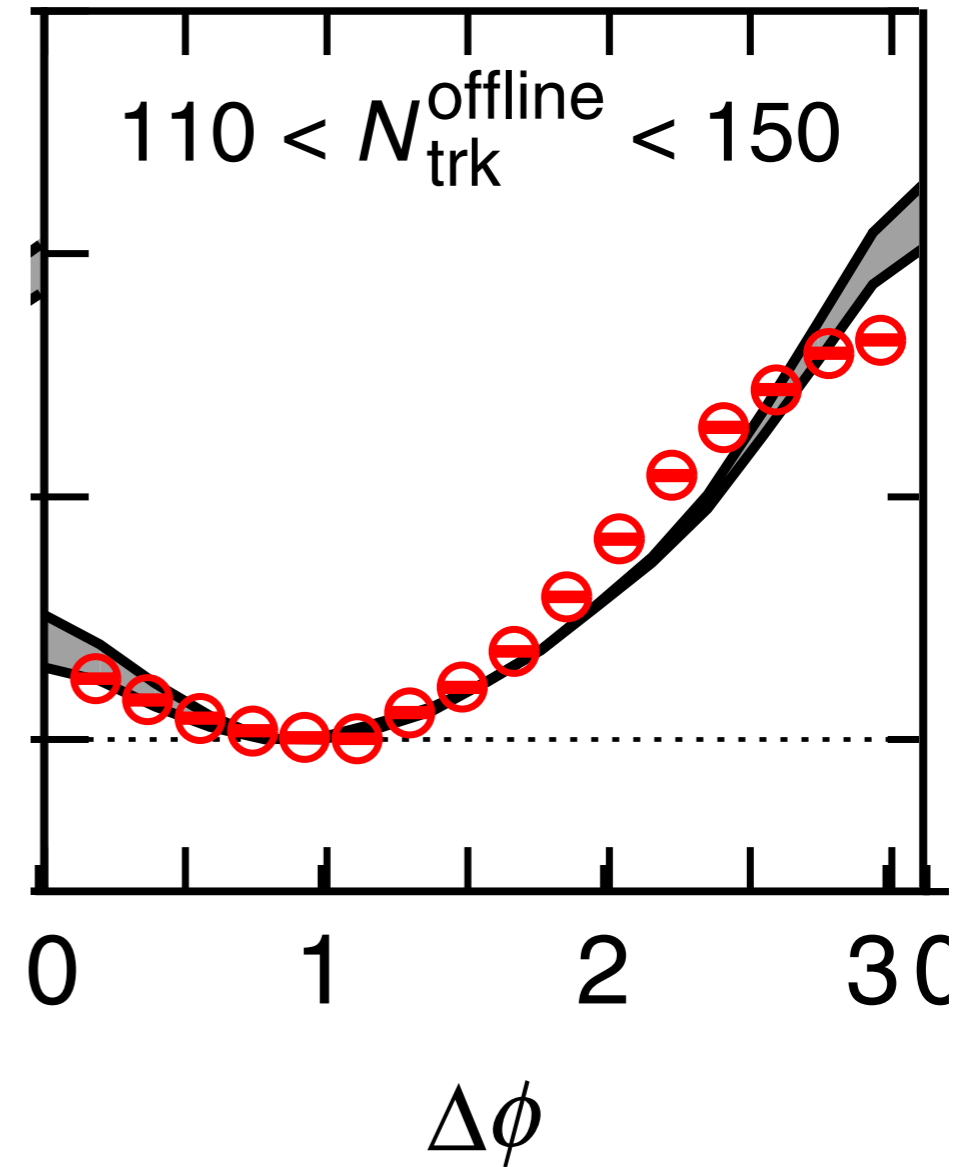
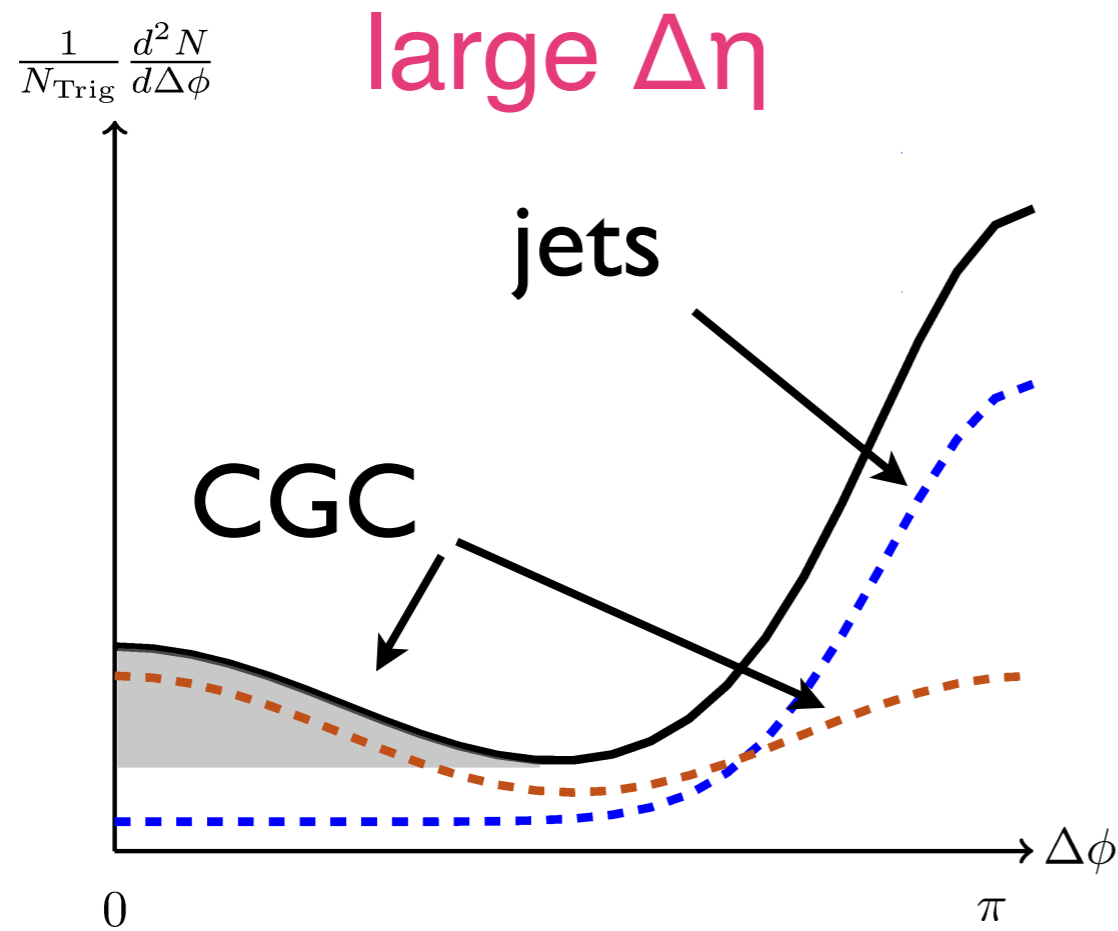
very similar to AA results



are the pA and AA  $v_2$  related to the same physics?

# ridge in pp/pPb from color glass condensate?

Color Glass Condensate: calculational framework for saturation



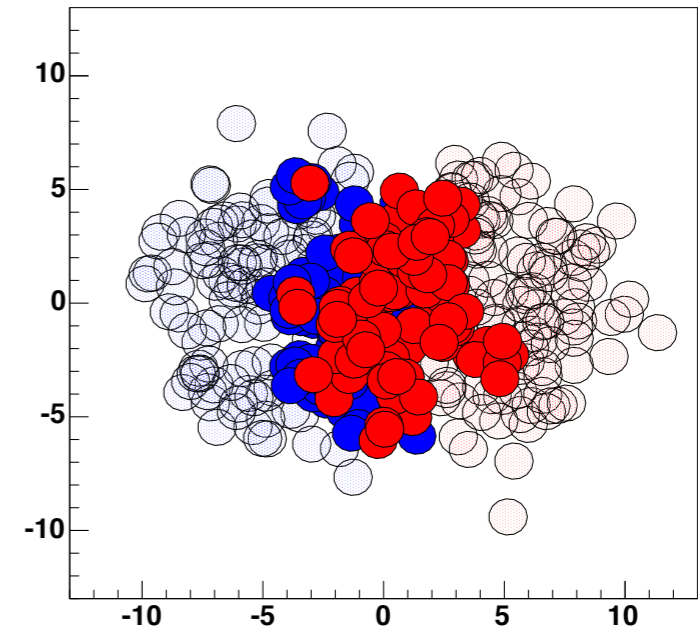
good description of the data in pPb

# geometry in AA & pA

---

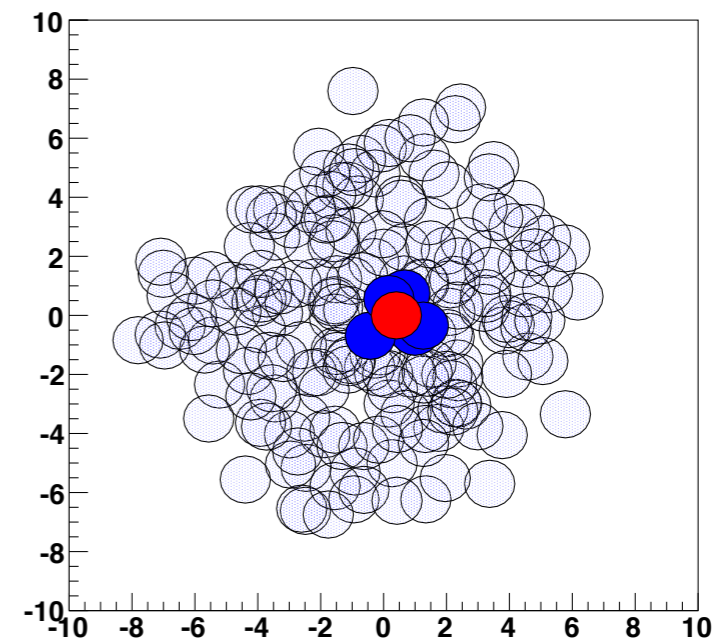
**AA**

**impact parameter  
+ fluctuations**



**pA**

**fluctuations**

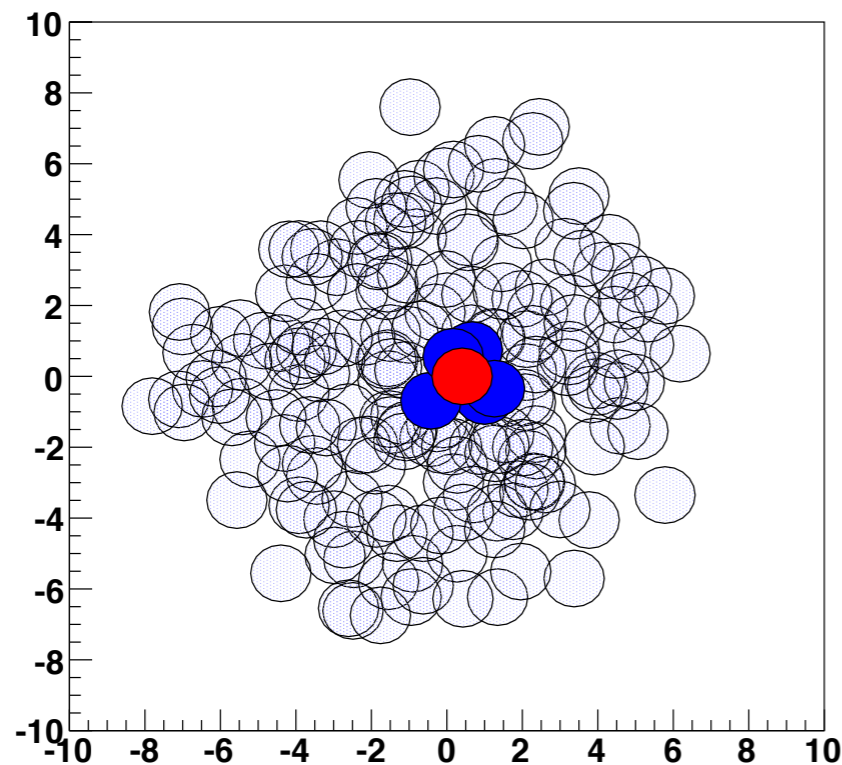




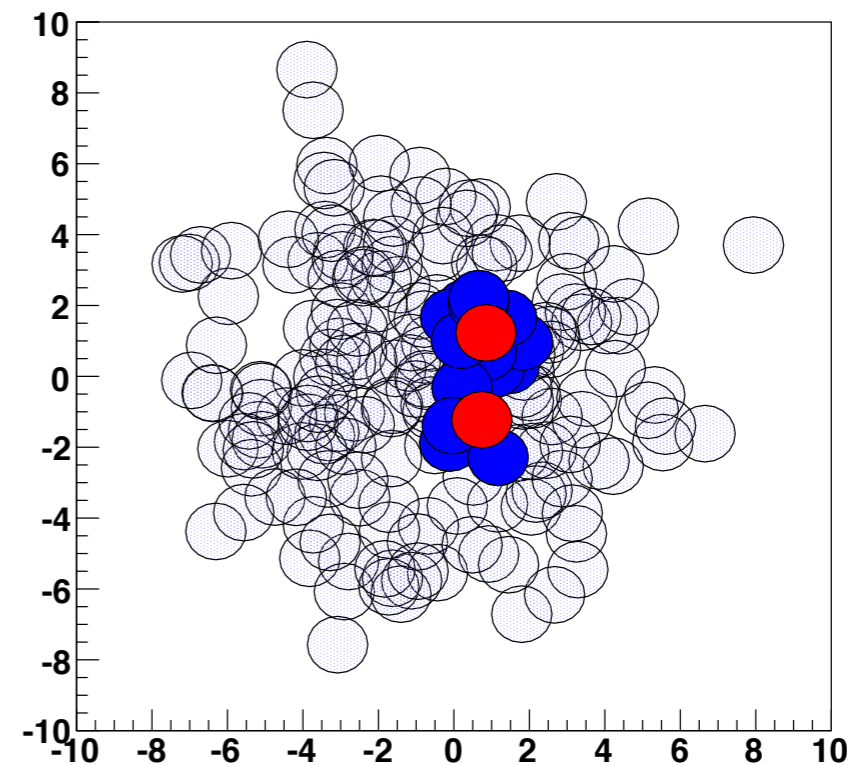
# variation of the small nucleus

---

**pA**



**dA**

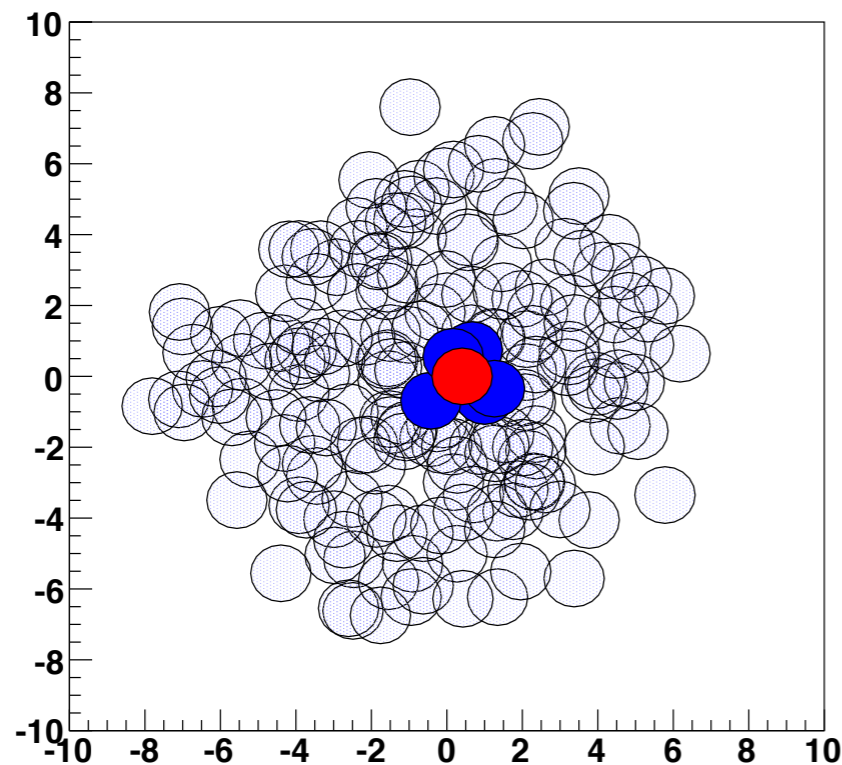


**control the collision geometry by varying the small nucleus**

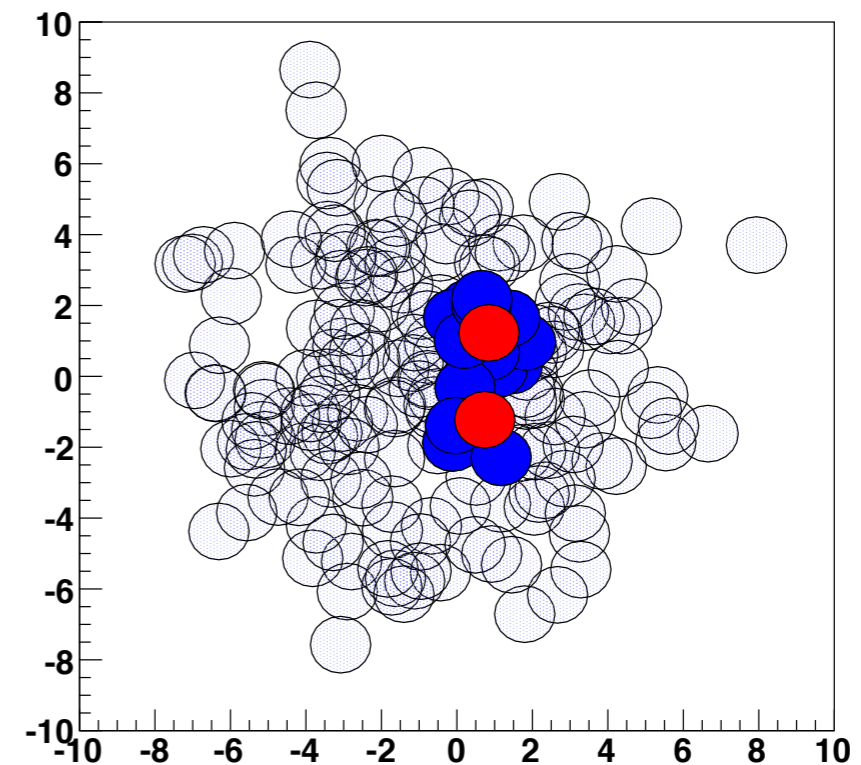
# variation of the small nucleus

---

**pA**



**dA**



**control the collision geometry by varying the small nucleus**

does  $v_2$  reflect the geometry of the initial state in p/d+A as in A+A?

# what can RHIC add?

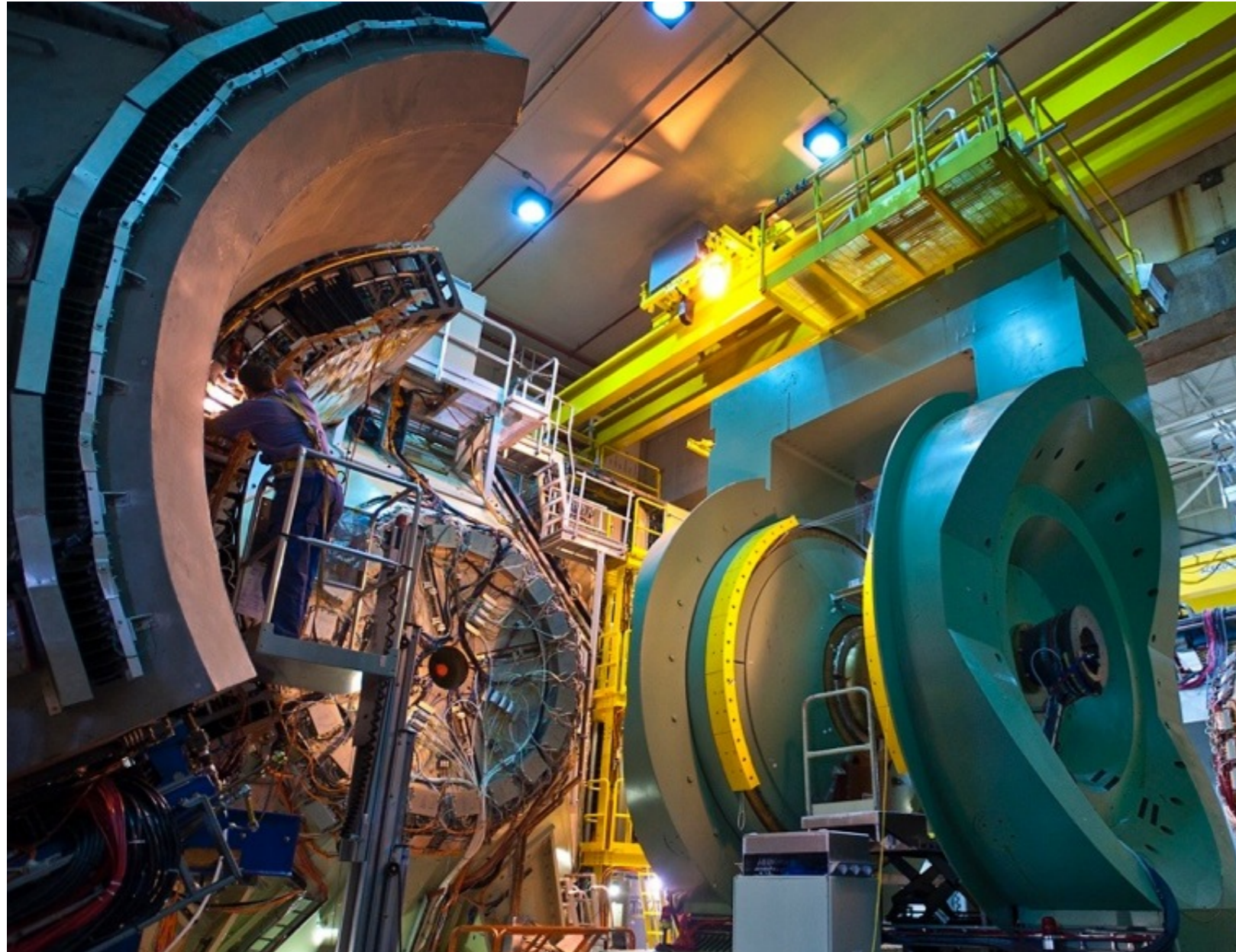
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RHIC had huge d+Au sample  
25x smaller collision energy than the LHC

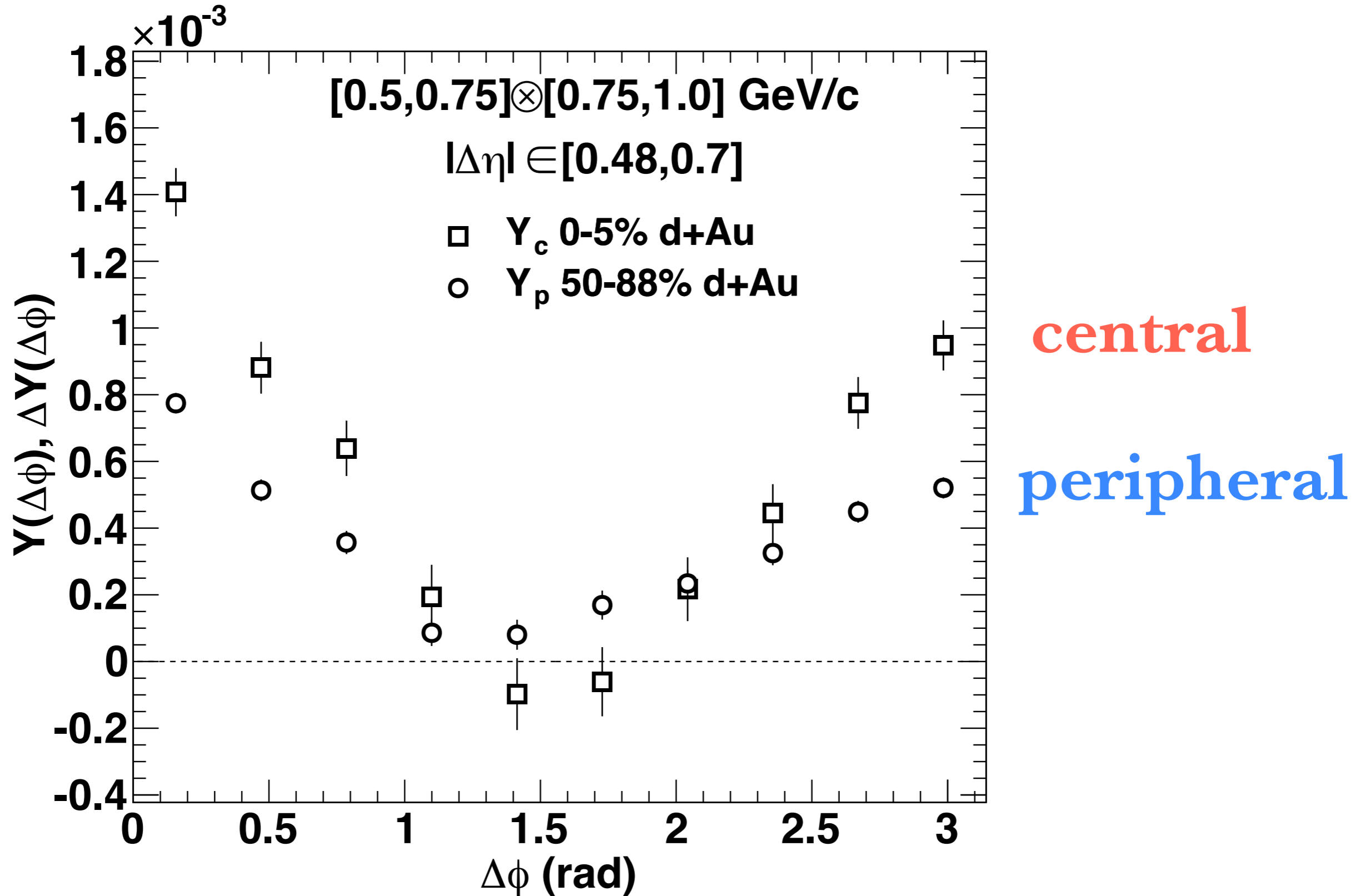
# PHENIX

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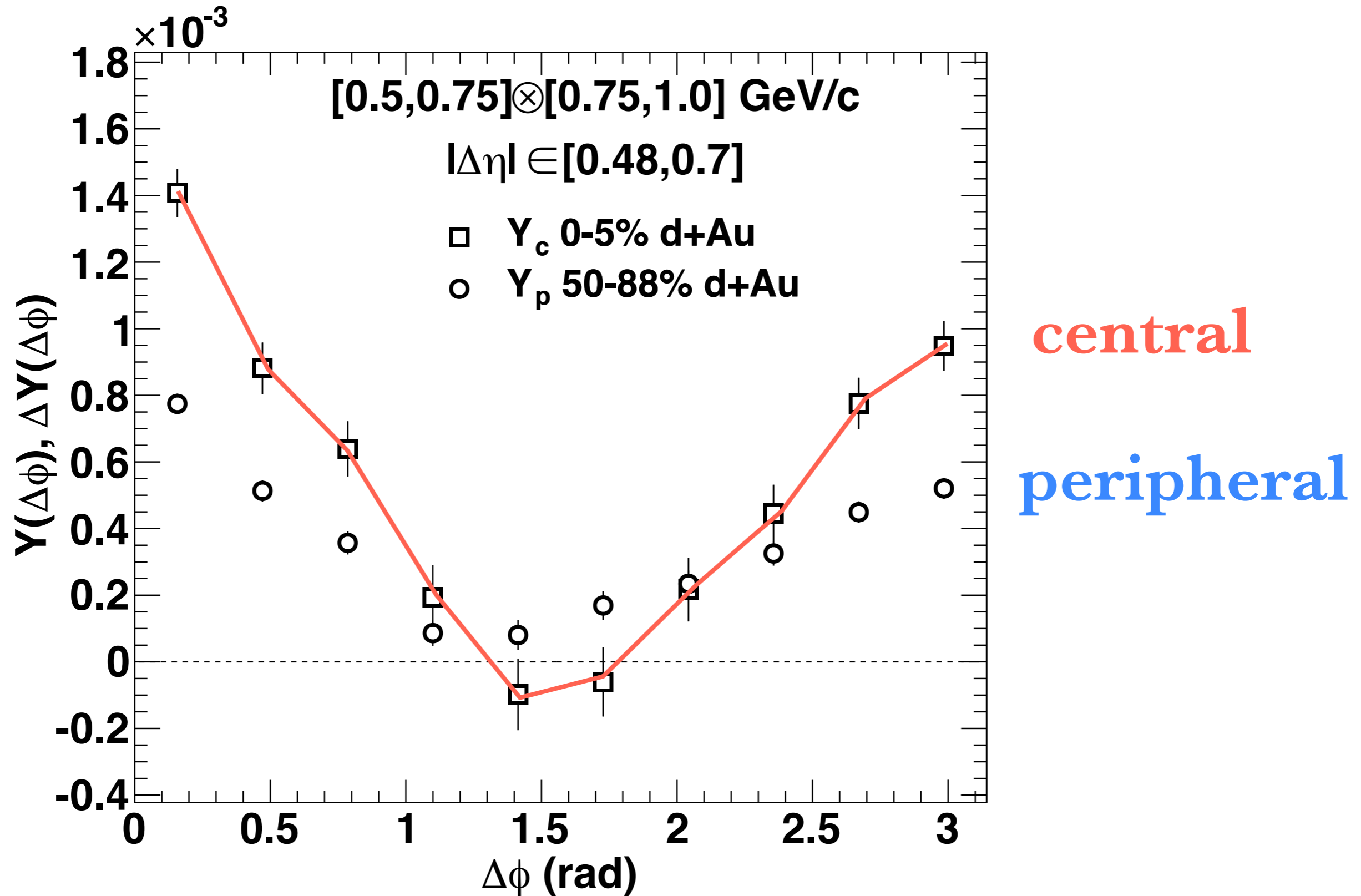


- charged hadrons
- $|\eta| < 0.35$
- $|\Delta\eta| < 0.7$
- centrality determined by charged particles in the Au going direction:  $3 < |\eta| < 4$
- 1.6B minimum bias events

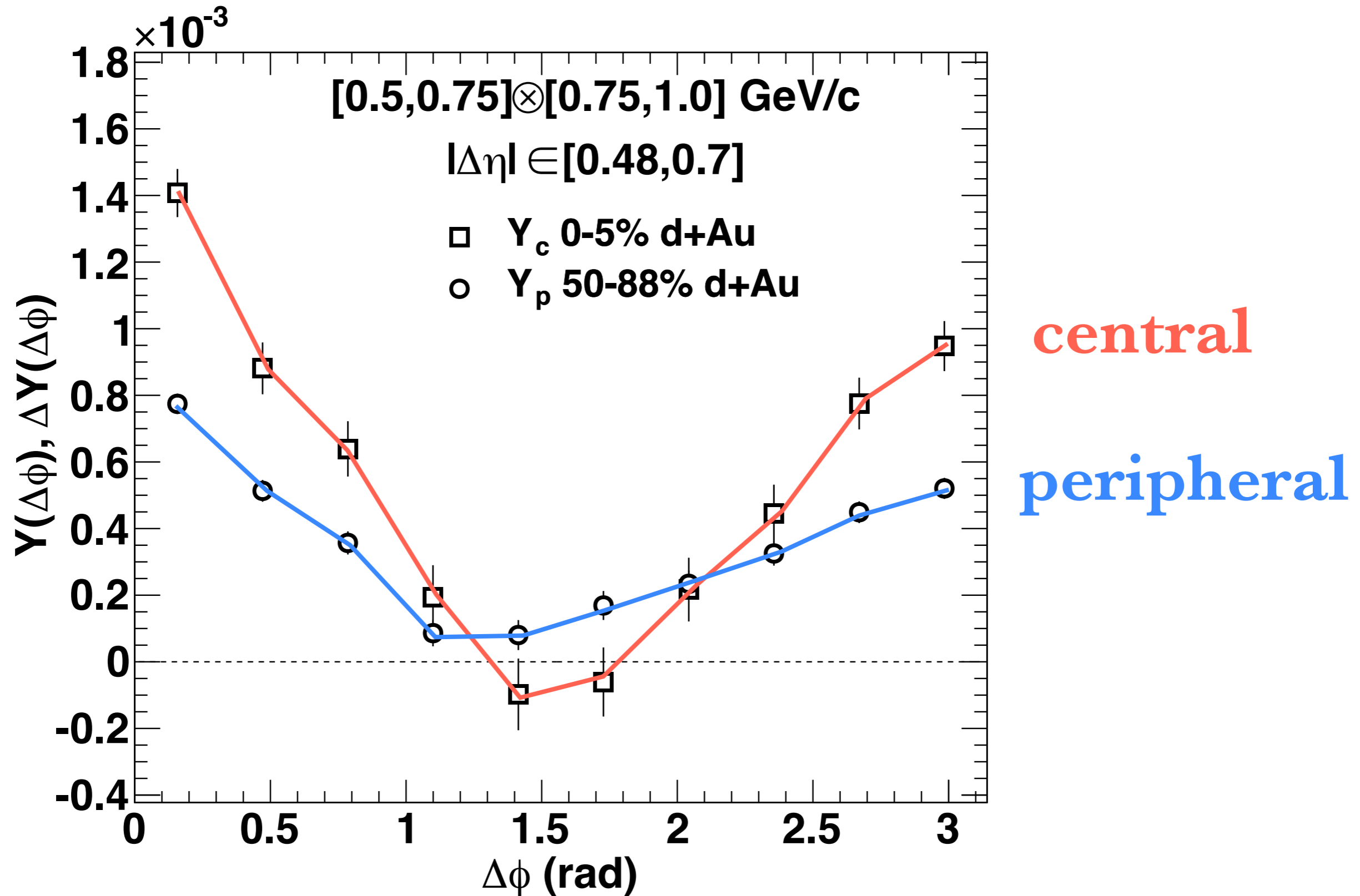
# two particle correlations in dAu



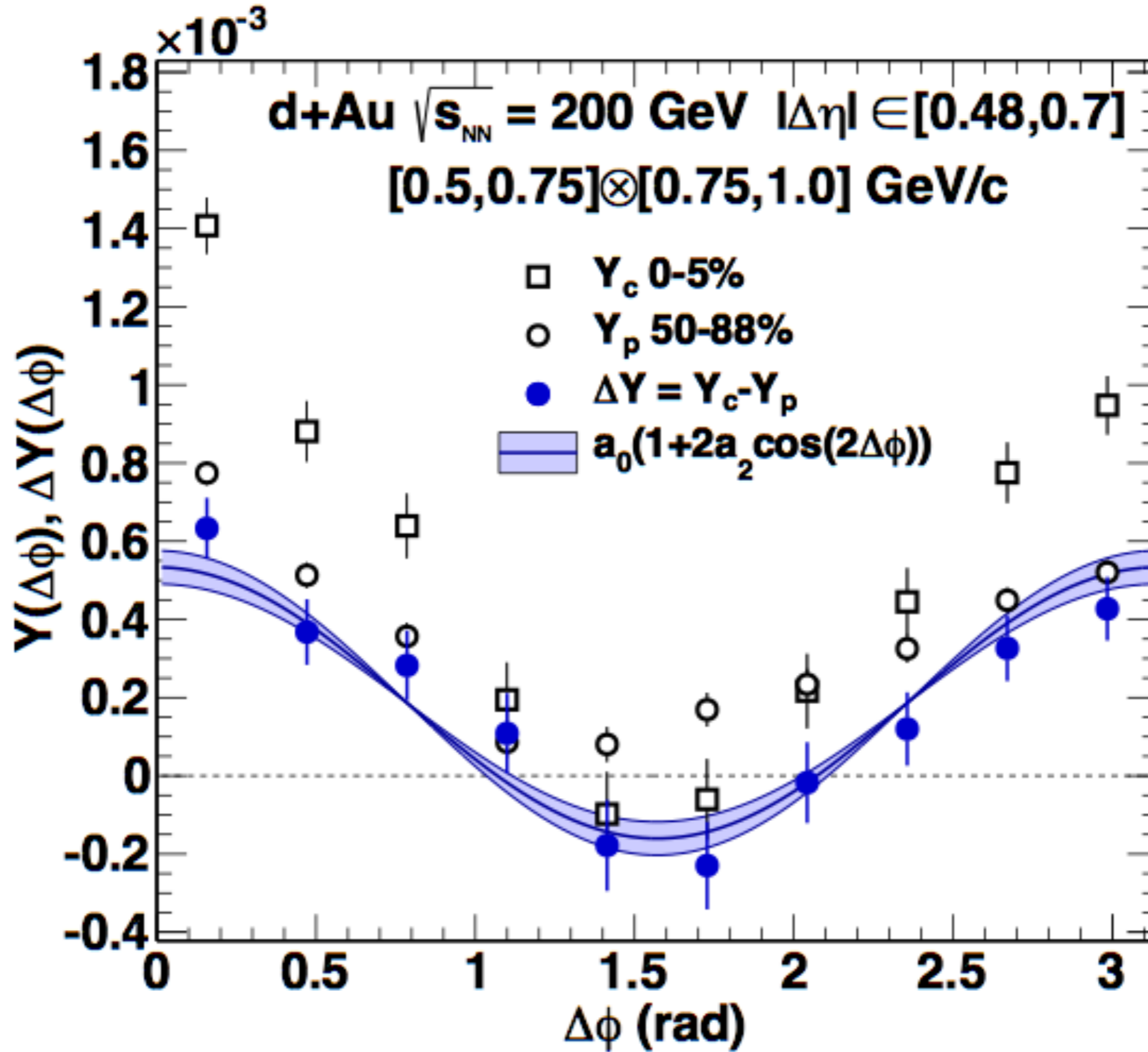
# two particle correlations in dAu



# two particle correlations in dAu

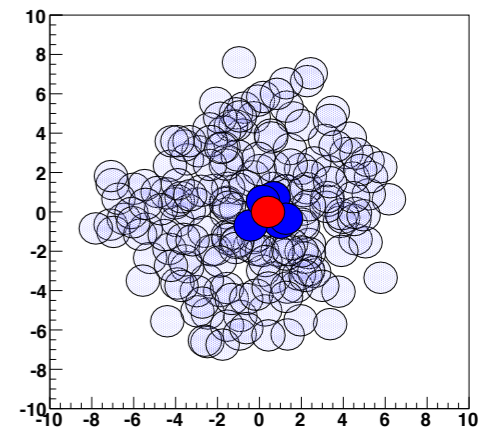
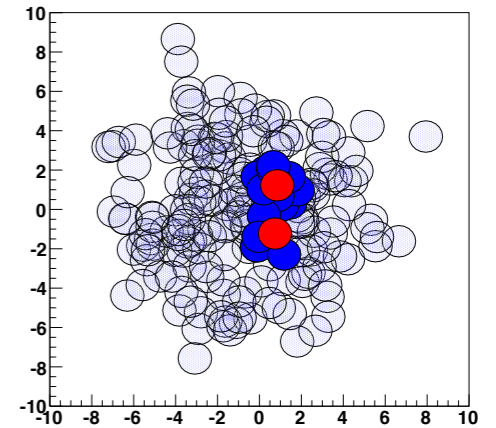
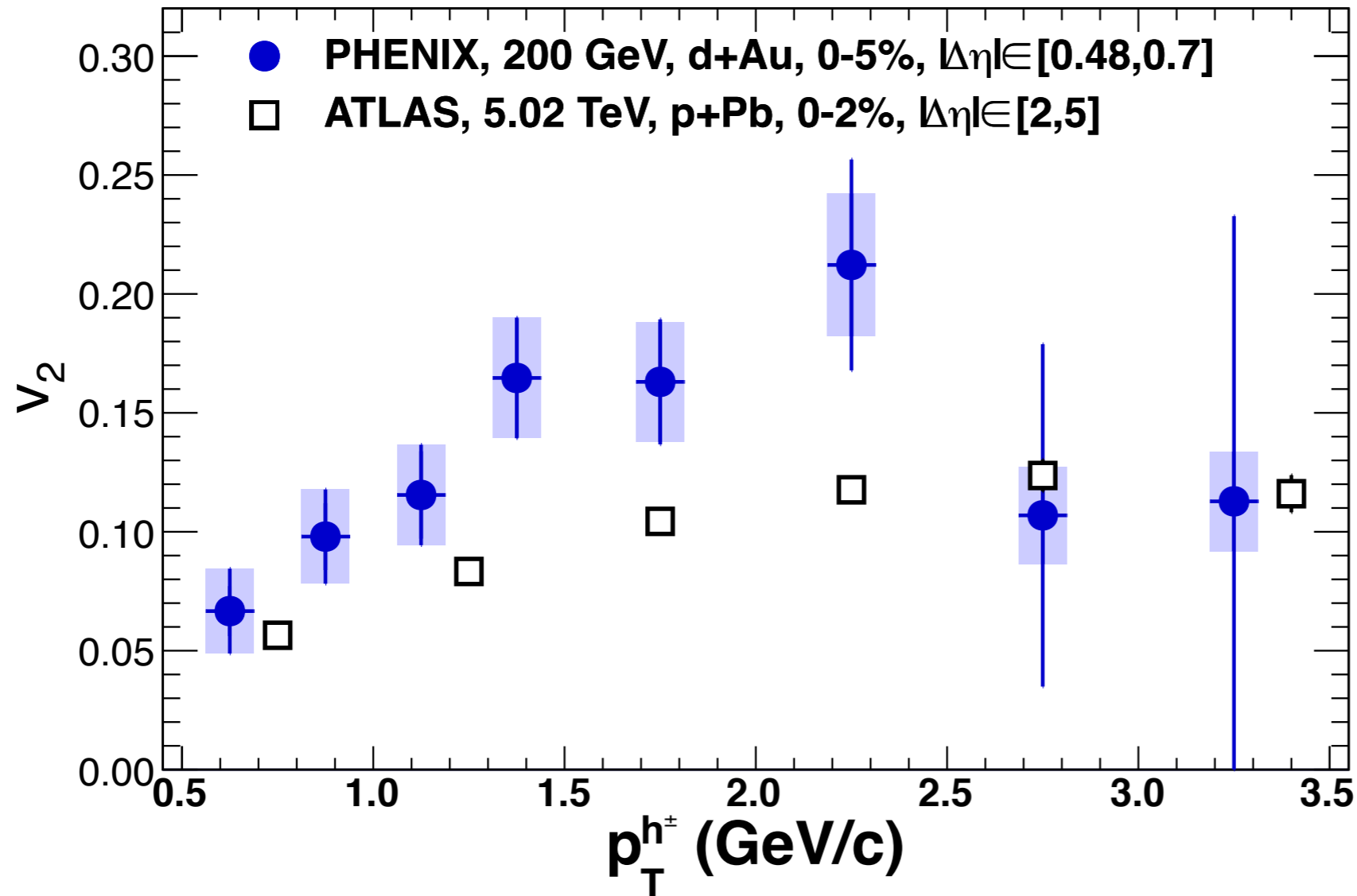


# centrality dependence





# v2: pPb & dAu

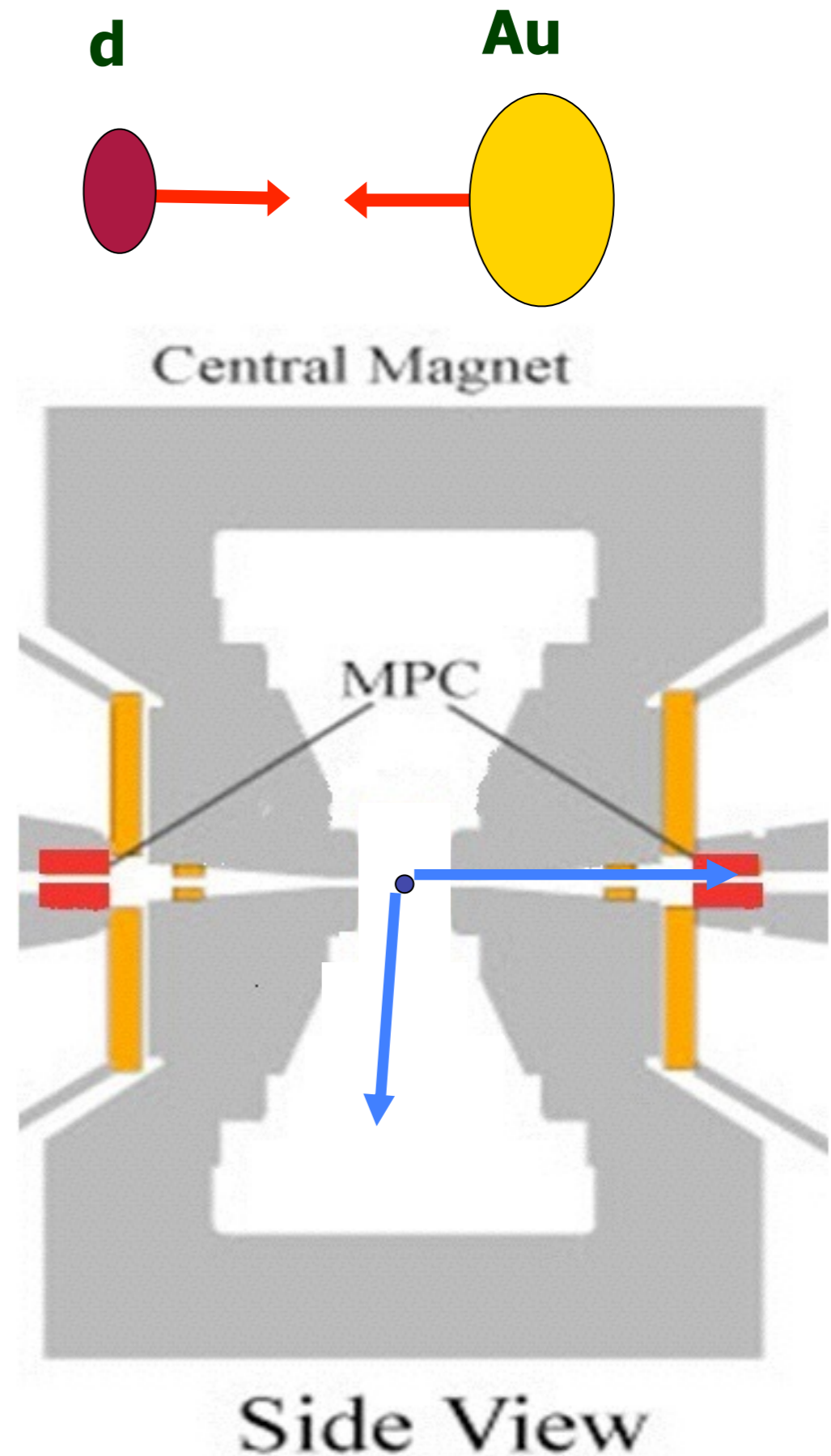


# rapidity separated correlations

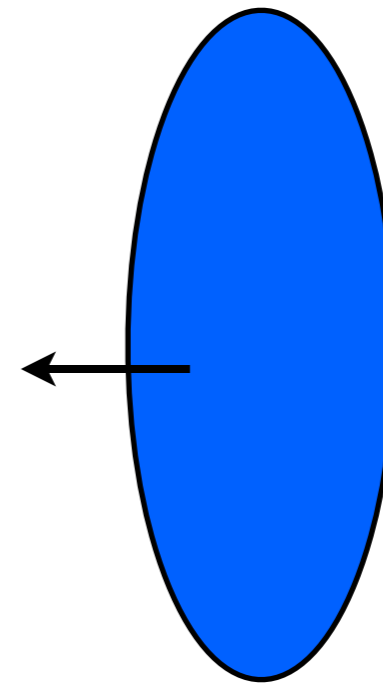
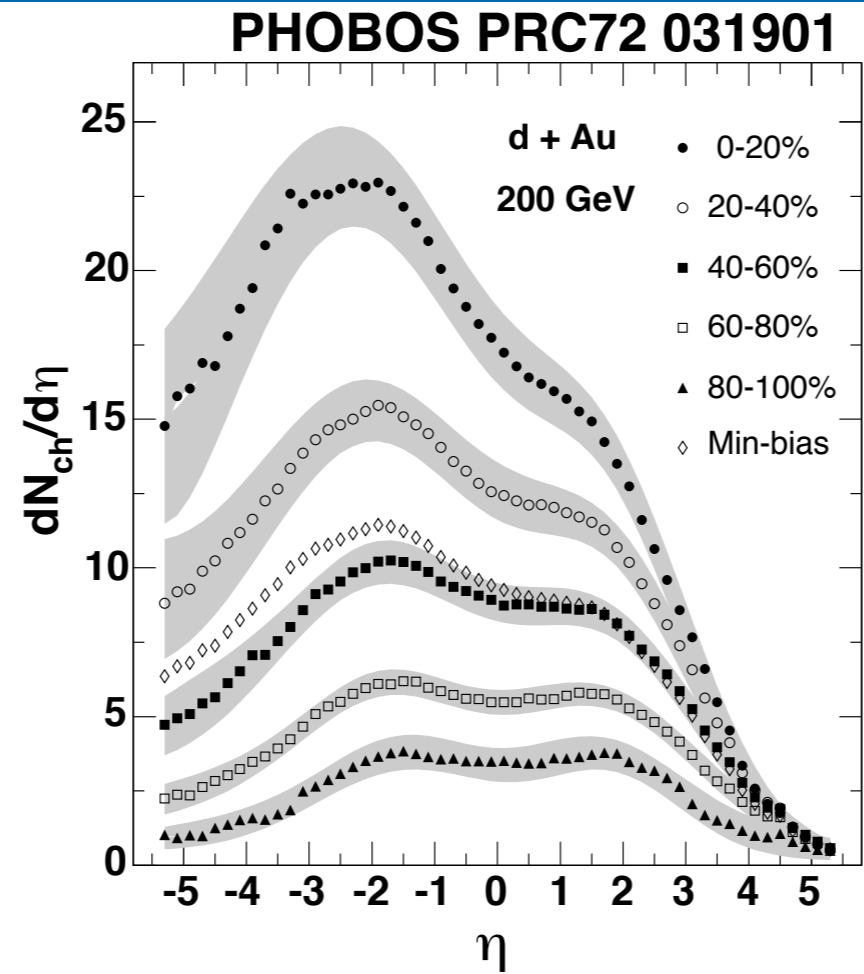
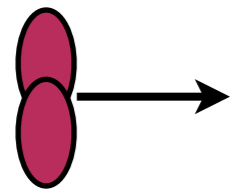
## Muon Piston Calorimeters

both d-going & Au-  
going directions

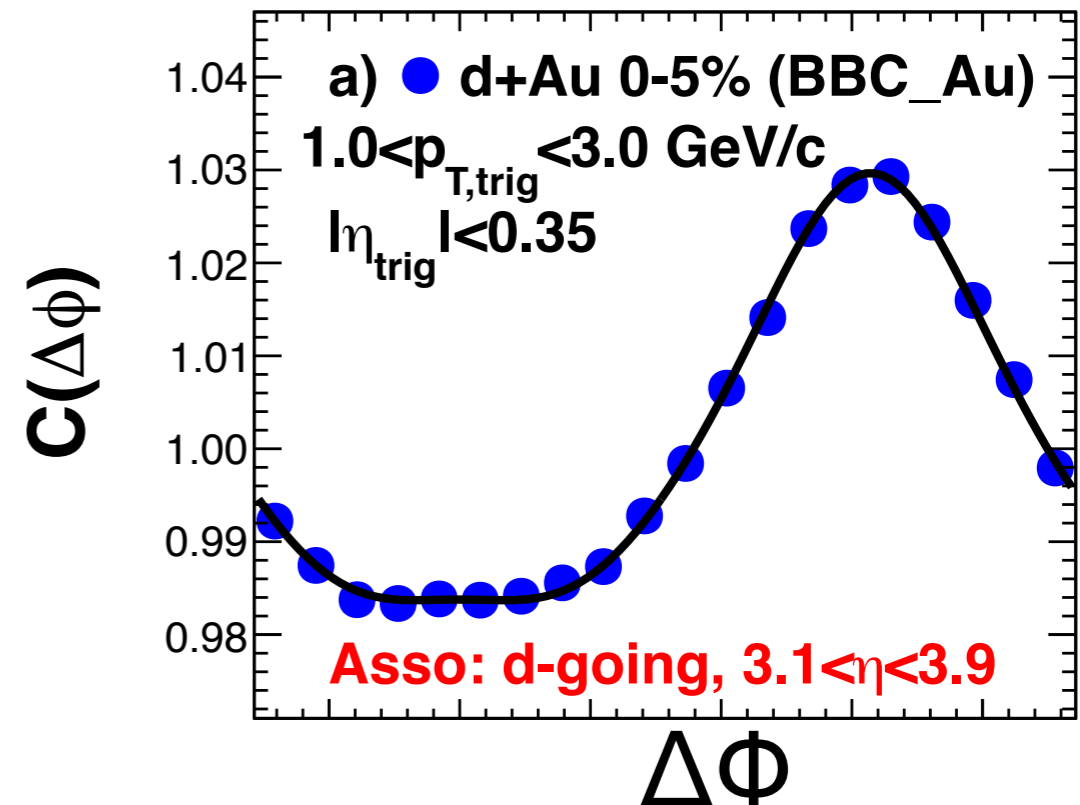
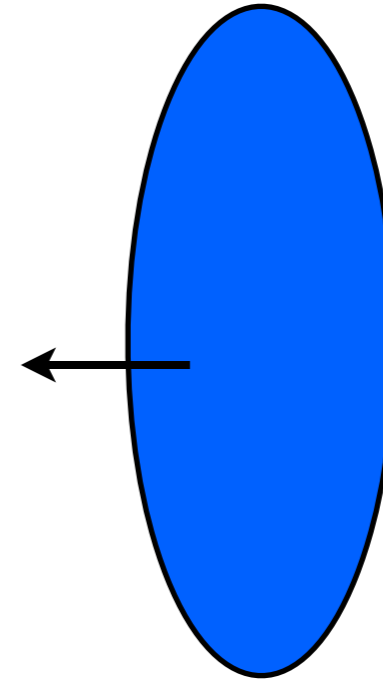
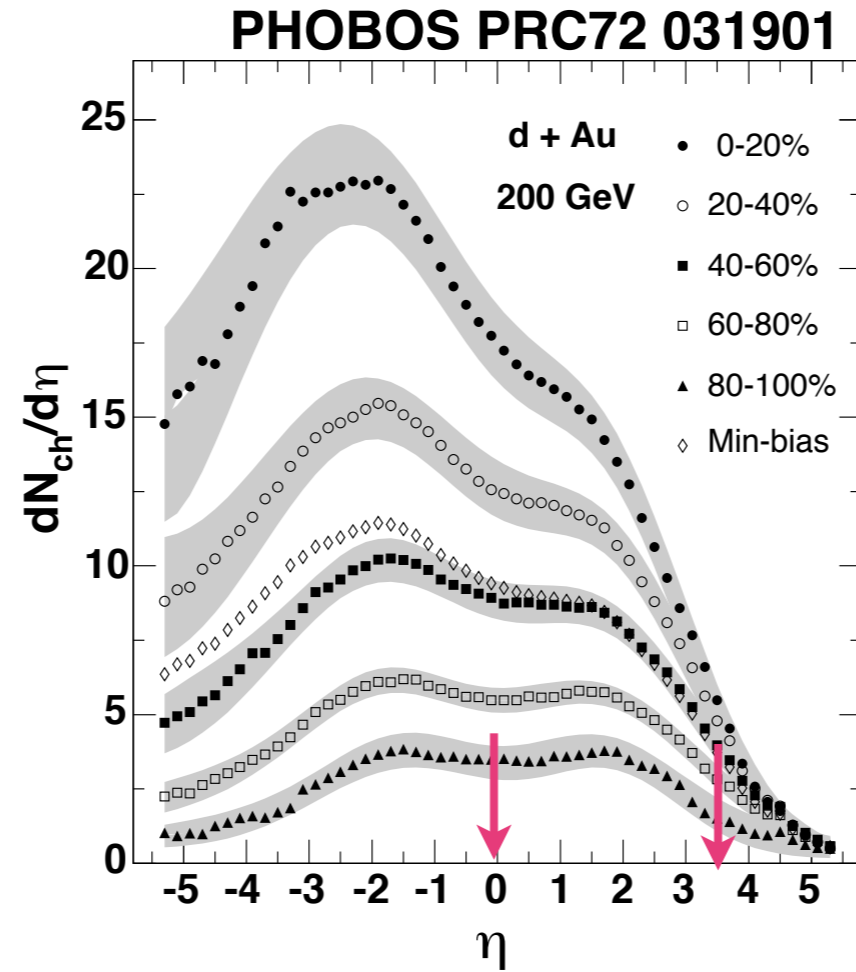
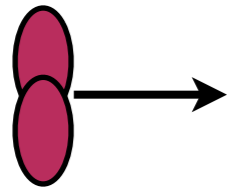
$$3 < |\eta| < 4$$



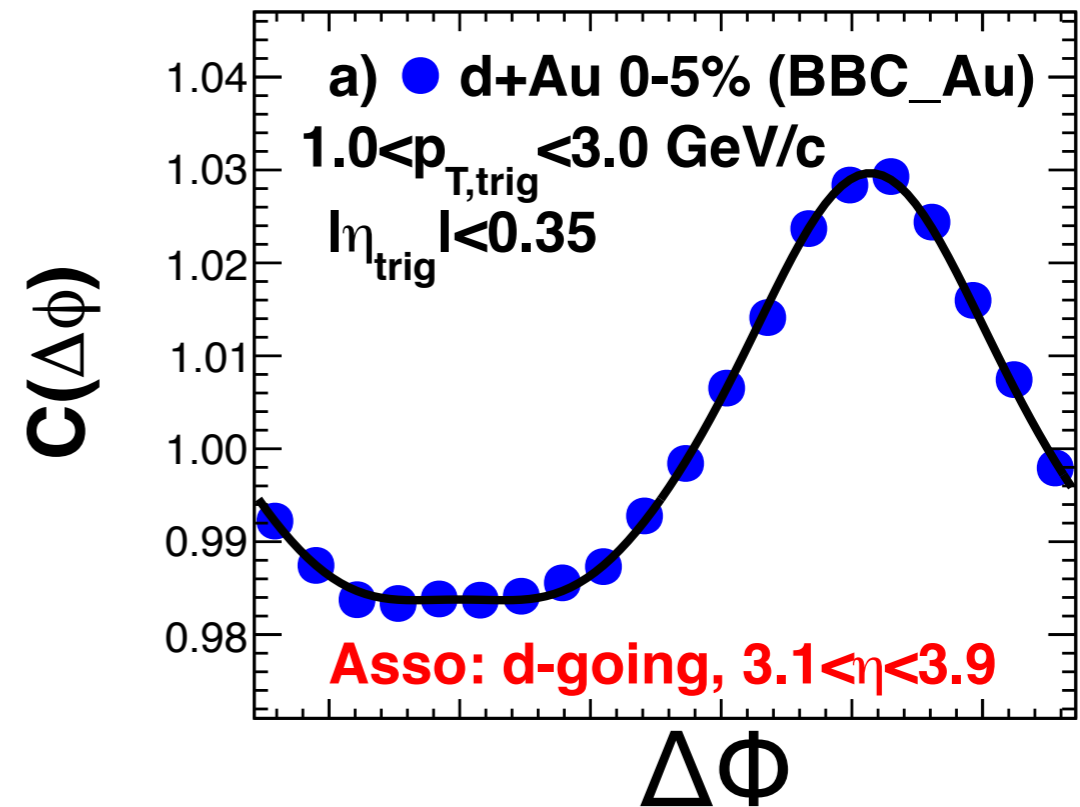
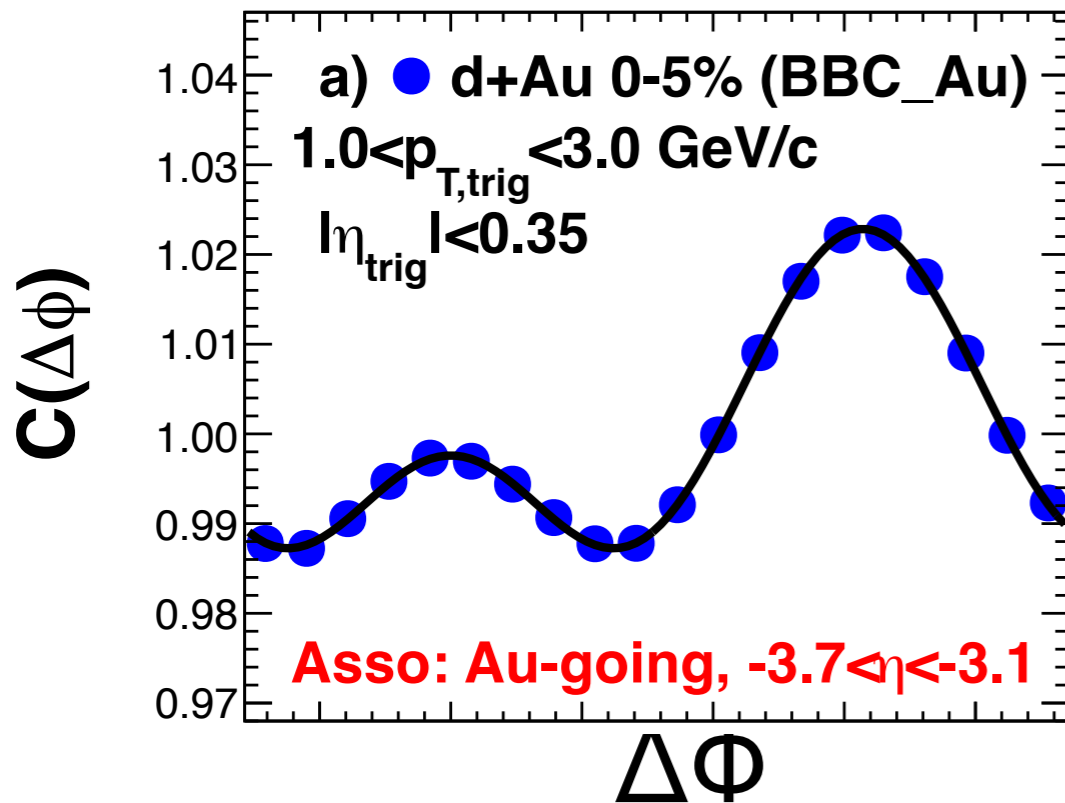
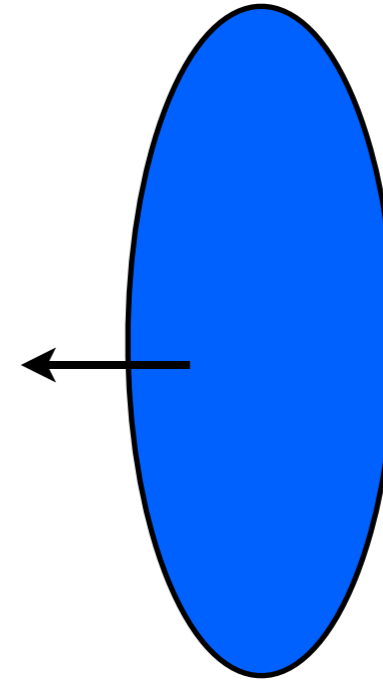
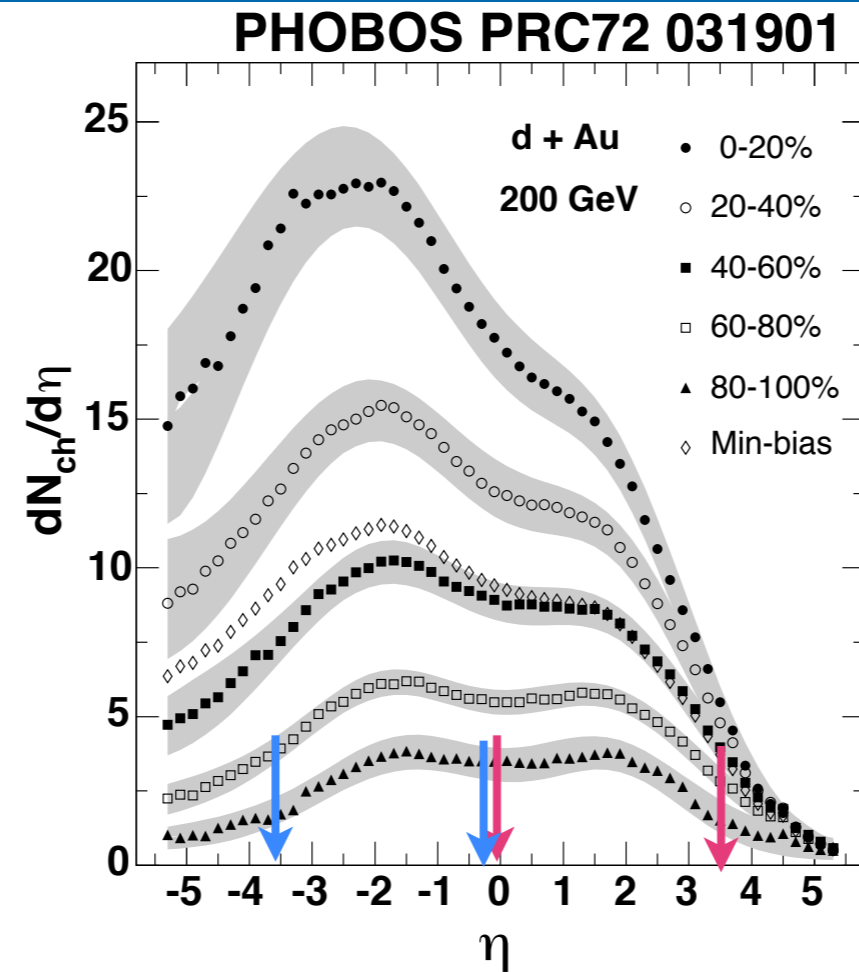
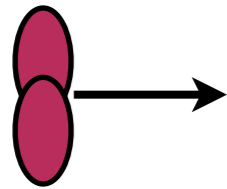
# long range correlations in dAu



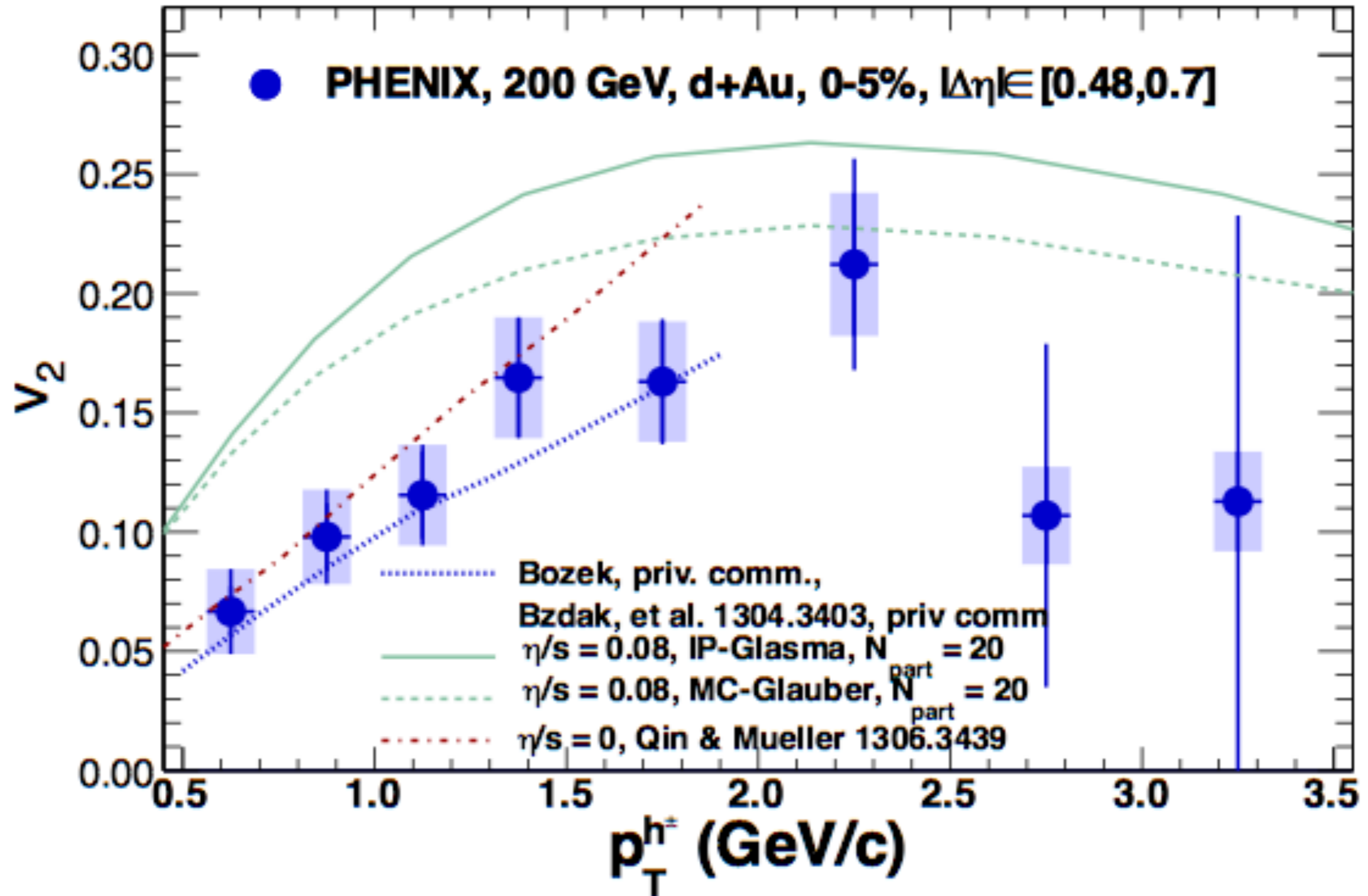
# long range correlations in dAu



# long range correlations in dAu



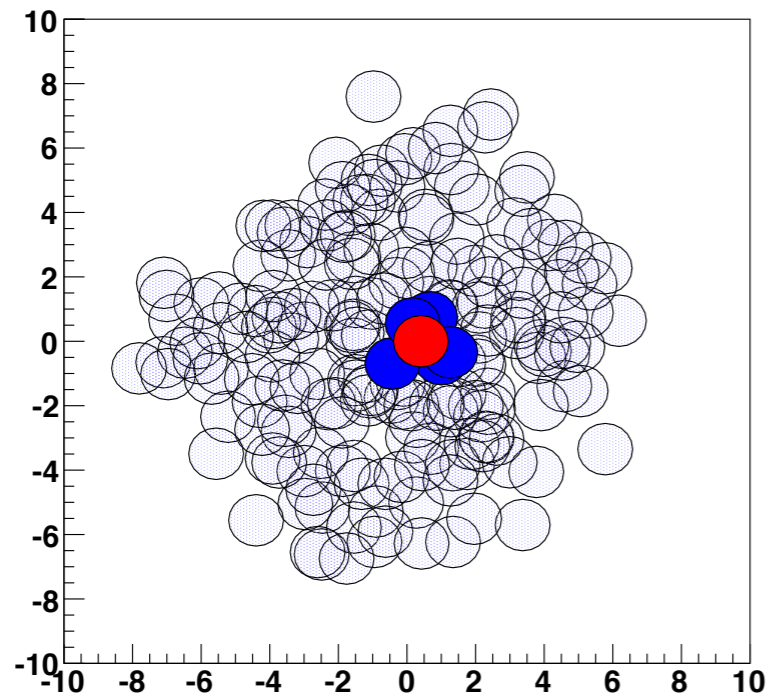
# $v_2$ in dAu compared to hydro. calculations



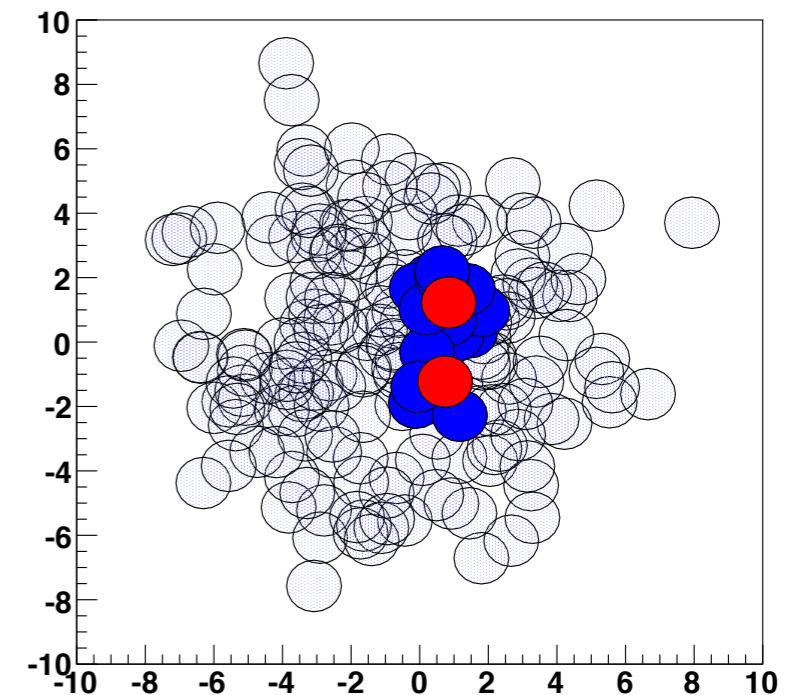
# shapes of pA & dA

---

pA, small  $\varepsilon_2$

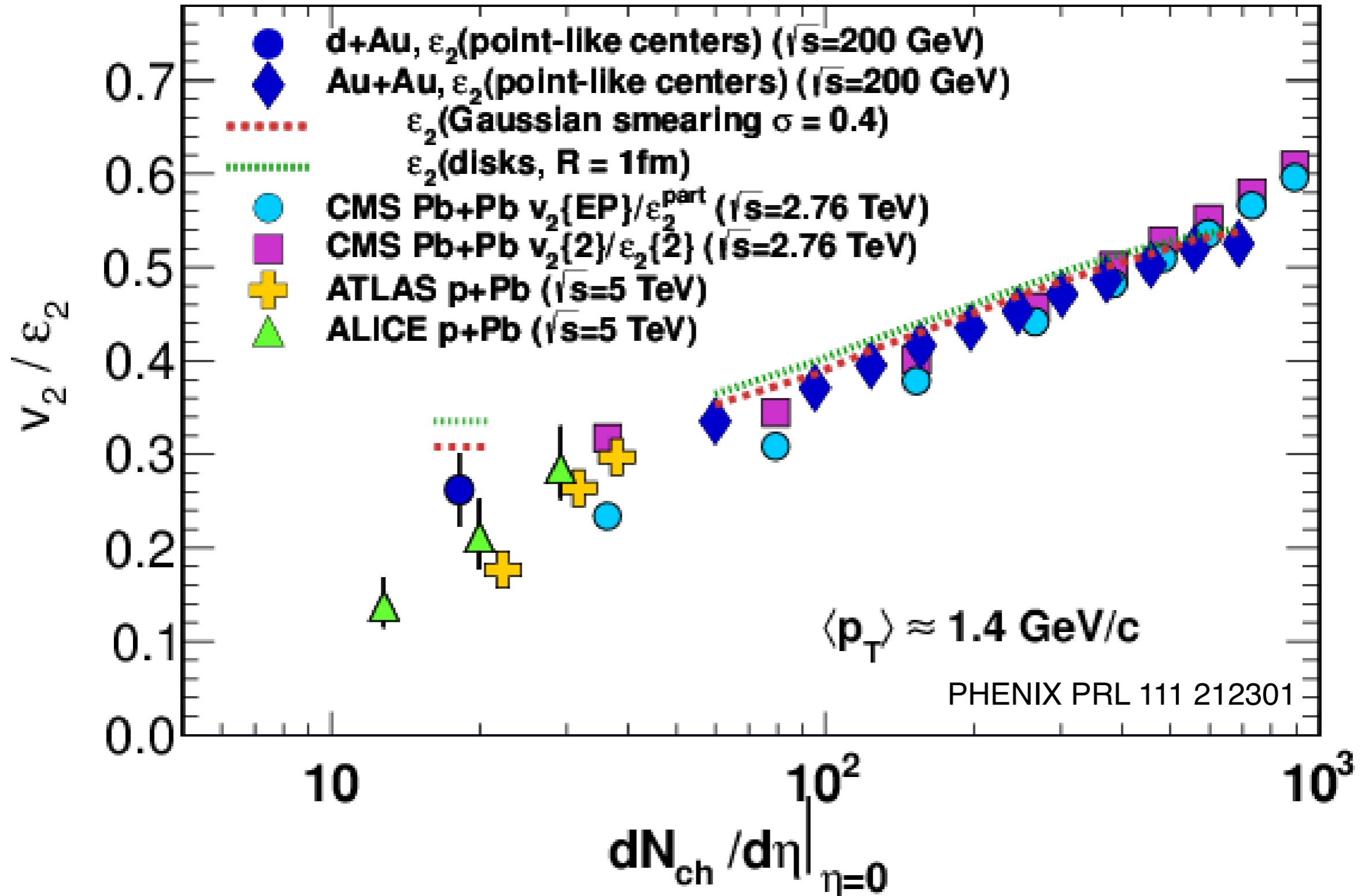


dA, large  $\varepsilon_2$



Glauber Monte Carlo used to generate single event initial energy density distributions  
used to determine  $\langle \varepsilon_n \rangle$  values for event selections

# dAu, pPb, AuAu & PbPb

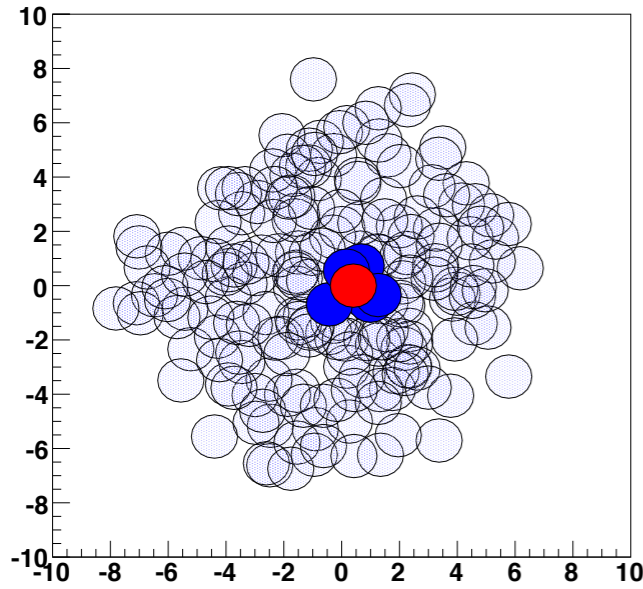


single trend, AA data understood as initial geometry + hydrodynamics



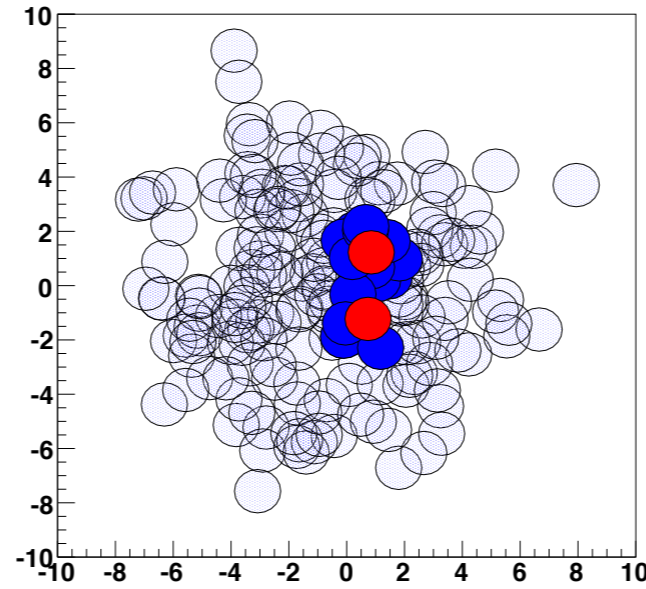
# variation of the small nucleus

**pA**



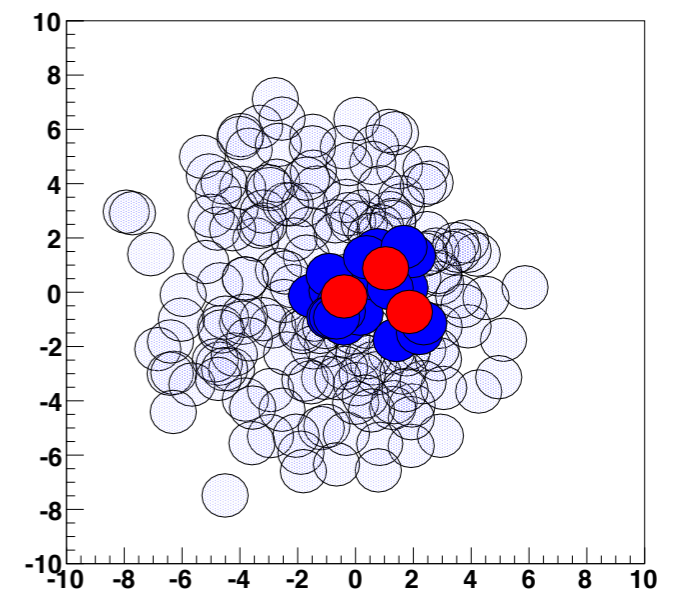
small  $\varepsilon_2$

**dA**



large  $\varepsilon_2$   
small  $\varepsilon_3$

**<sup>3</sup>HeA**

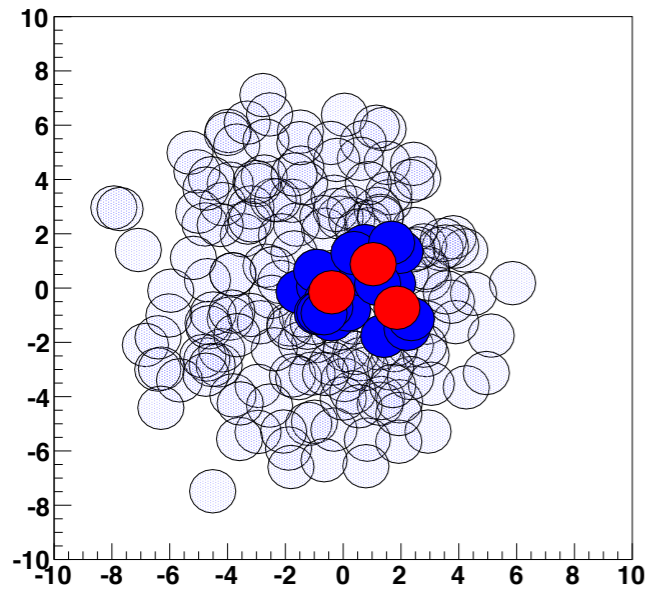


large  $\varepsilon_3$

$$\varepsilon_n = \frac{\sqrt{\langle r^2 \cos n\phi \rangle^2 + \langle r^2 \sin n\phi \rangle^2}}{\langle r^2 \rangle}$$

**control the collision geometry by varying the small nucleus**

# importance of $v_3$

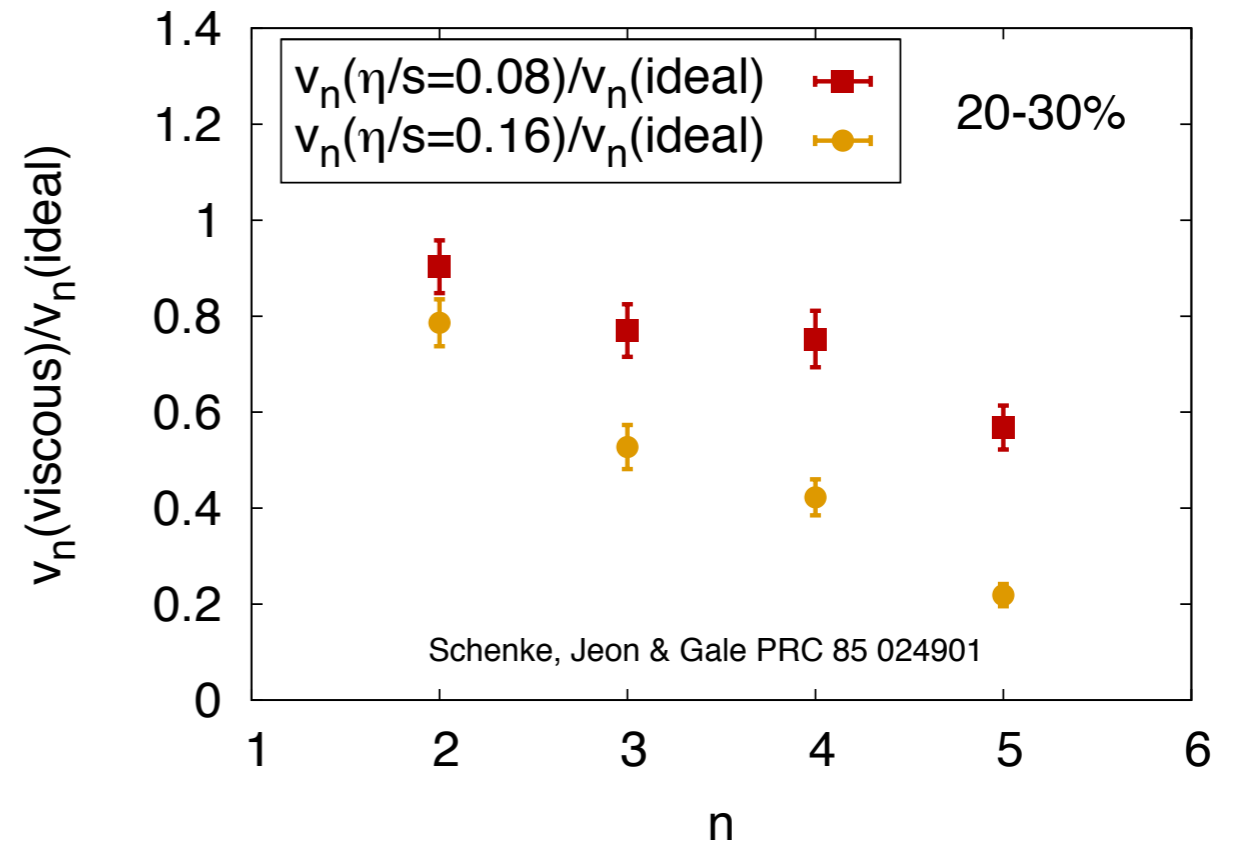


if:  $\varepsilon_3 \rightarrow \cos 3\Delta\Phi$  modulation

direct confirmation of hydrodynamic behavior in small systems

new handle on viscosity

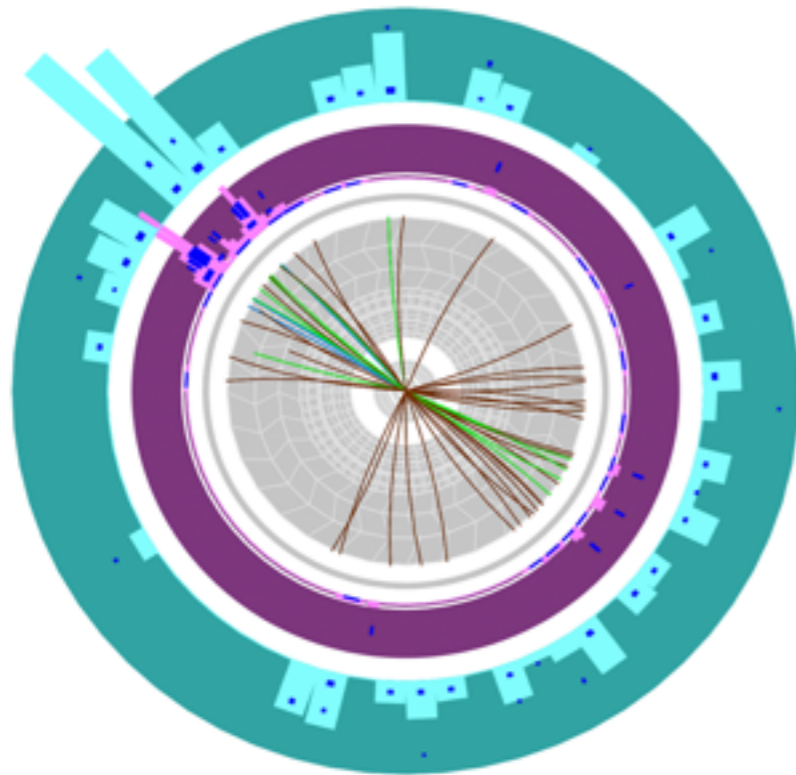
higher moments,  
more sensitive to viscous effects



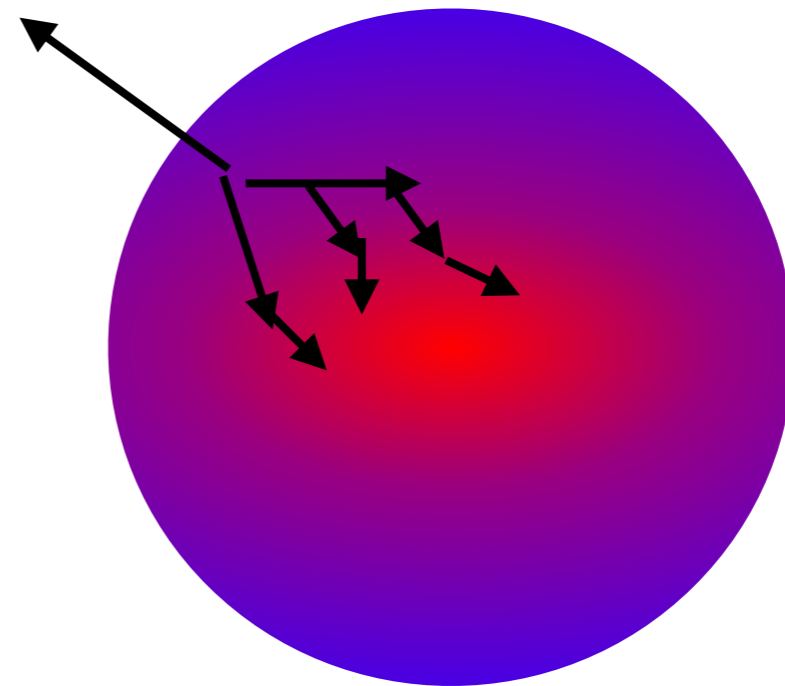
# jet quenching

---

jets act as an external probe of the QGP and lose energy  
as they go through the matter  
quenching is sensitive to the **matter** itself and how **long**  
the jet is in the matter



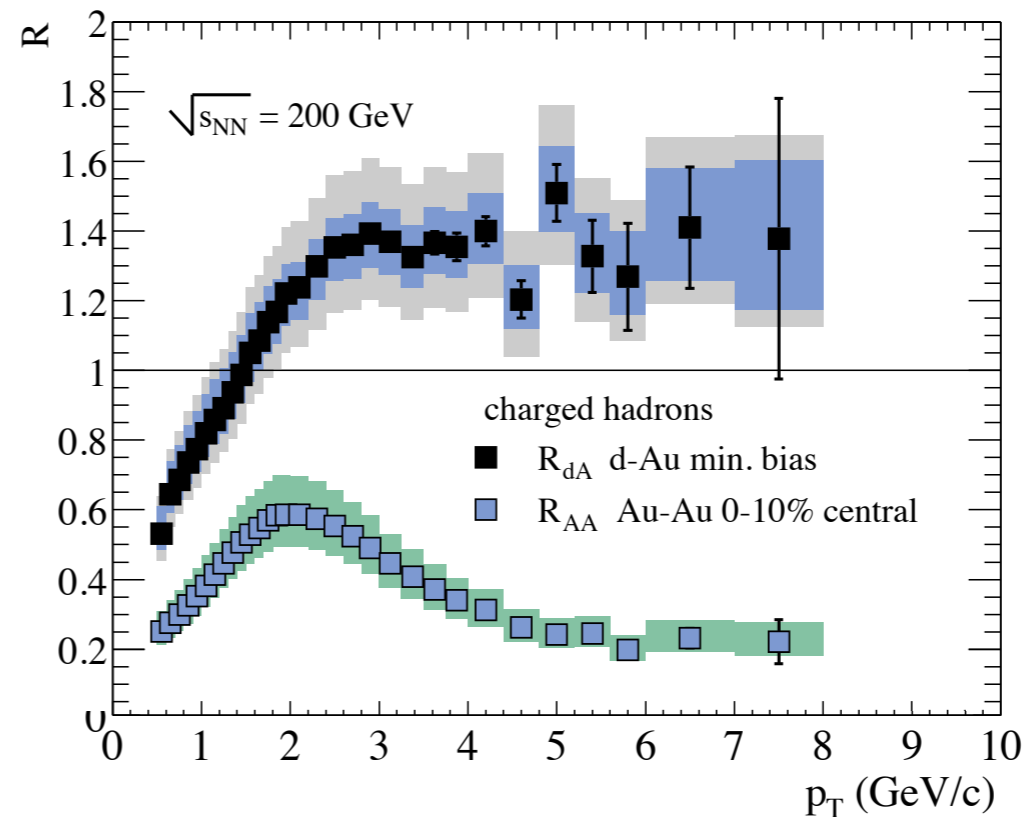
ATLAS PRL 105 252303 (2011)



# jet quenching

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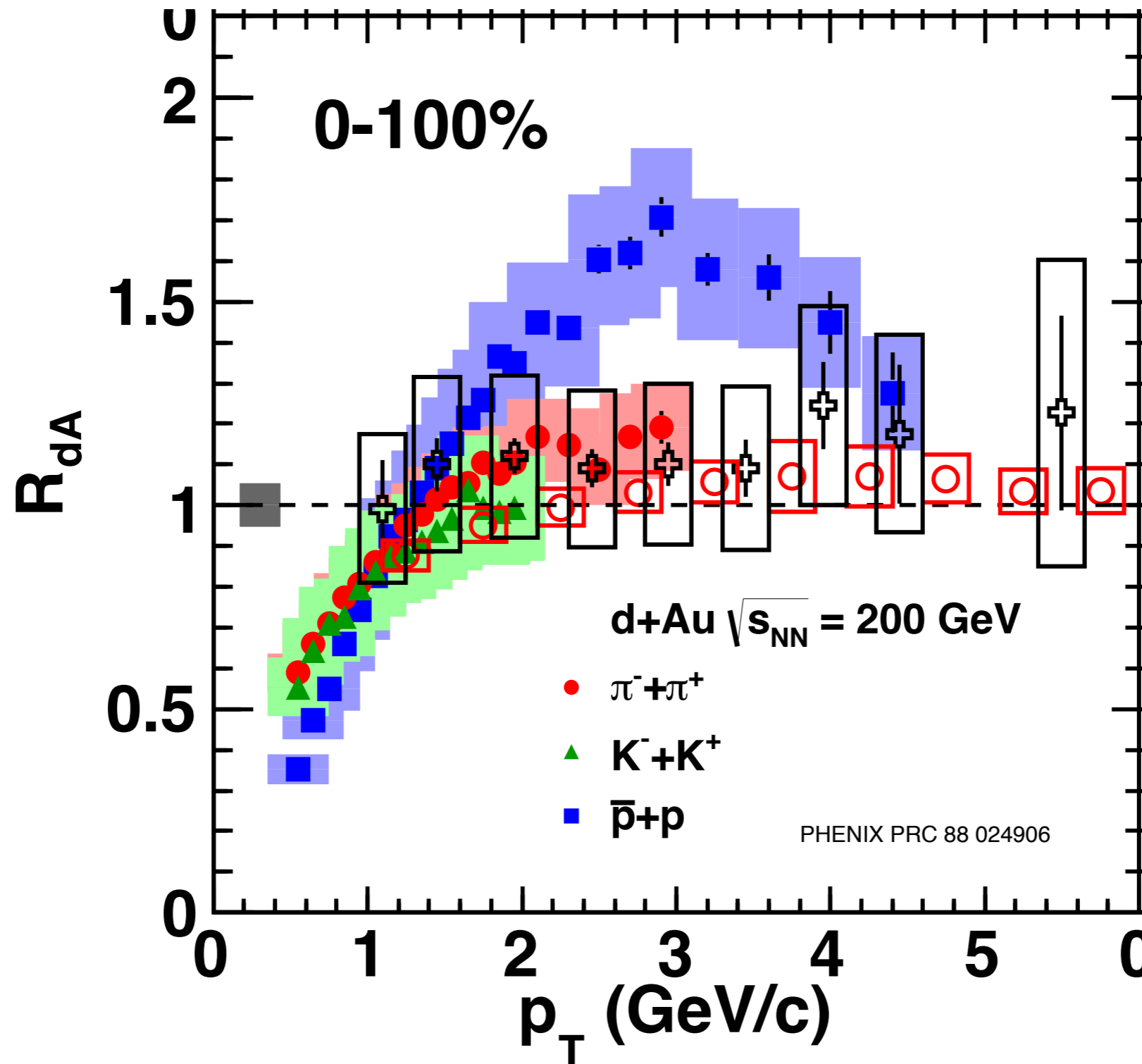
$$R_{AA} = \frac{\text{observed in AA}}{\text{expectation from pp}}$$



dAu:  
no quenching

AuAu: quenching

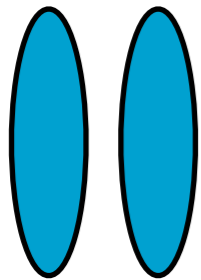
# particle species dependence



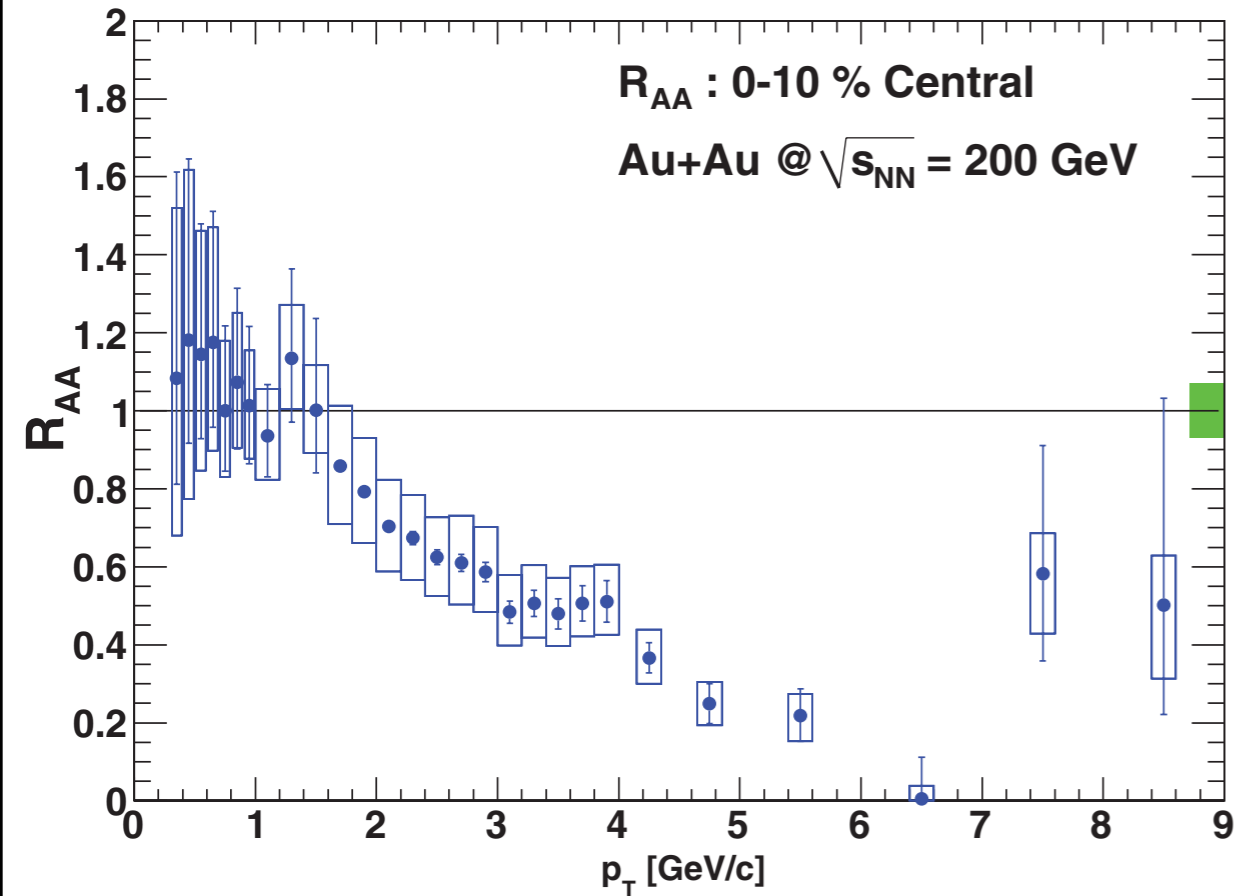
charged hadron enhancement in protons

# ...and heavy flavor

## electrons from heavy flavor decays

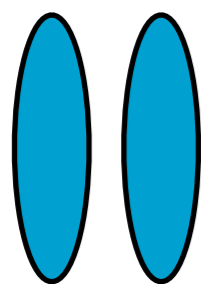


Au+Au



# ...and heavy flavor

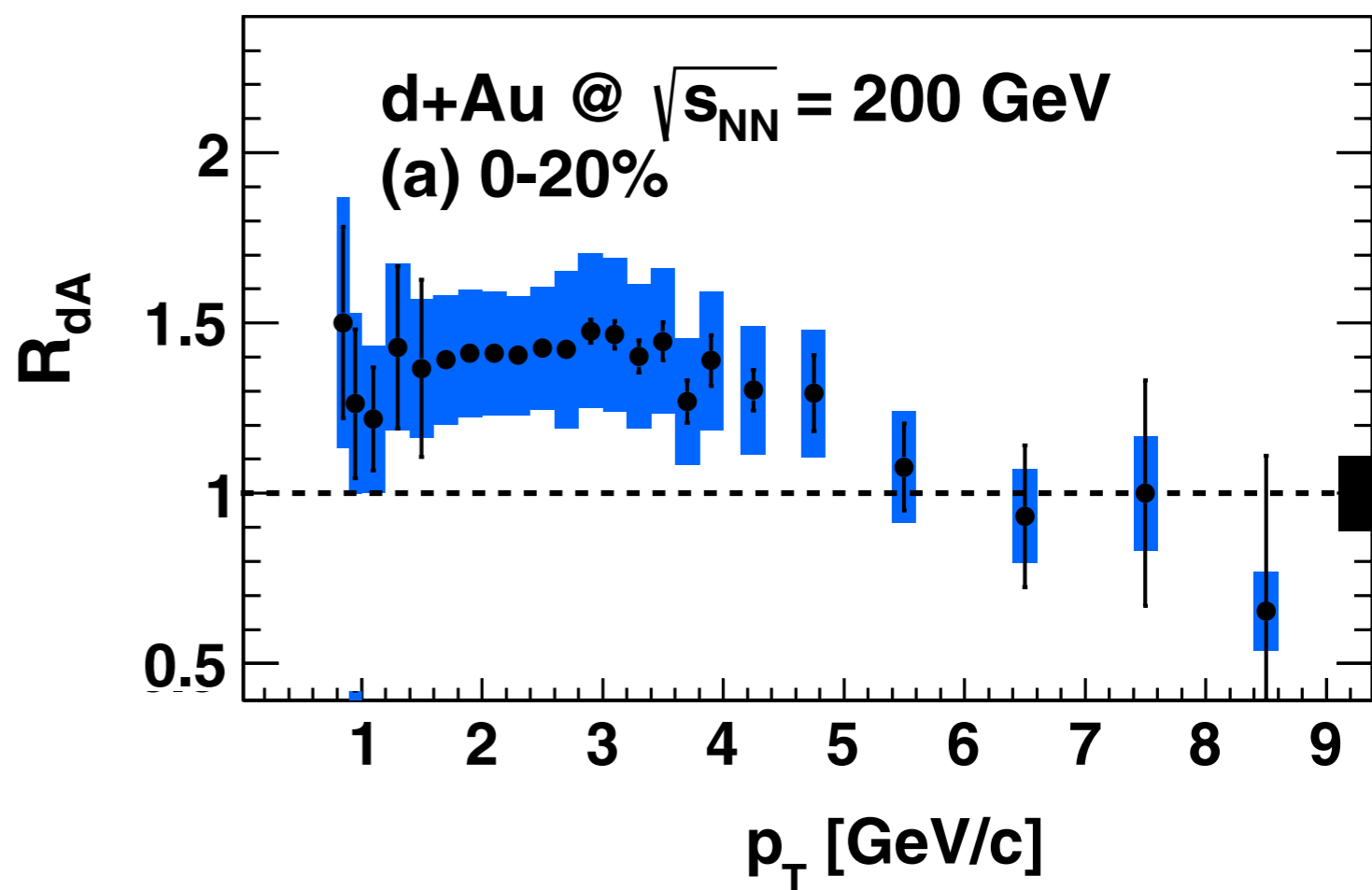
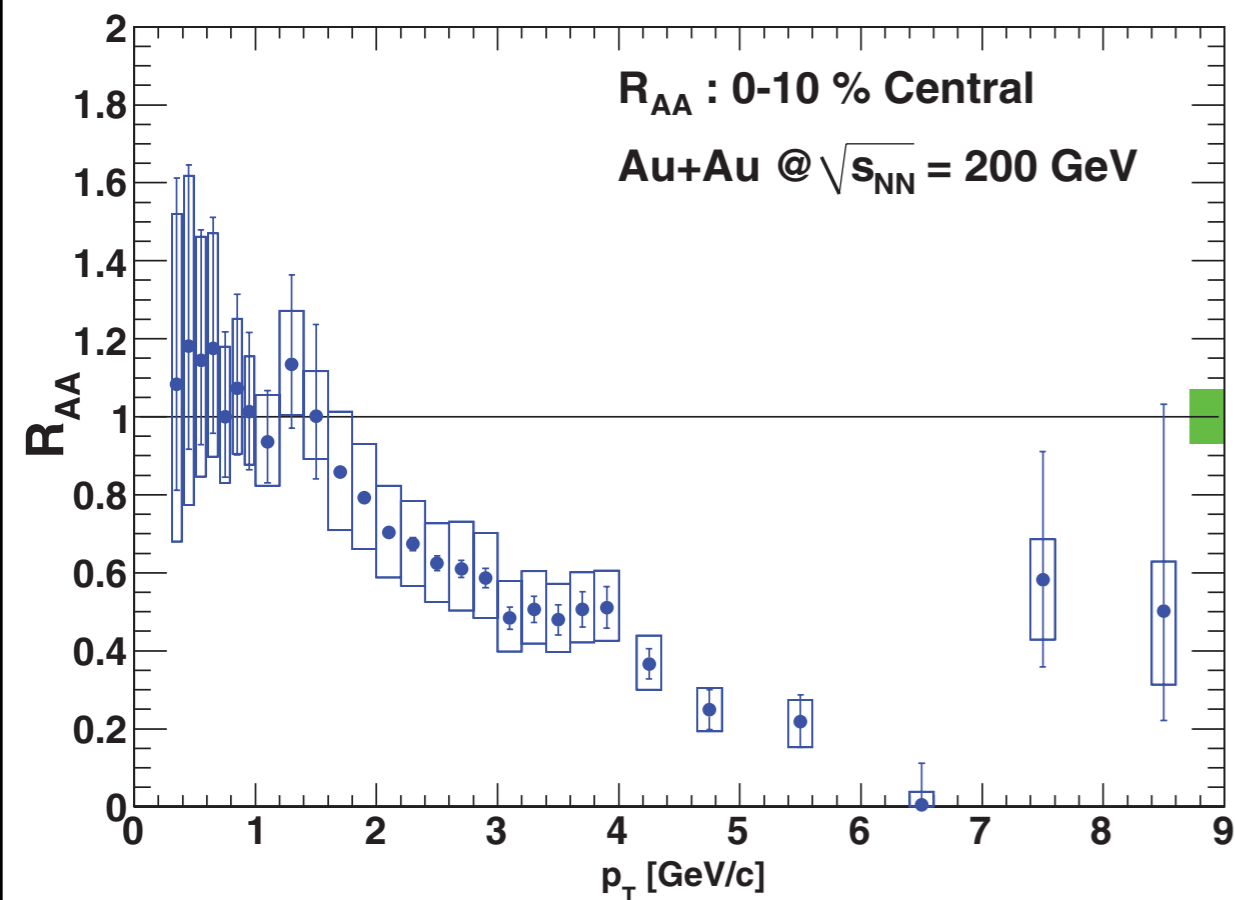
electrons from heavy flavor decays



Au+Au

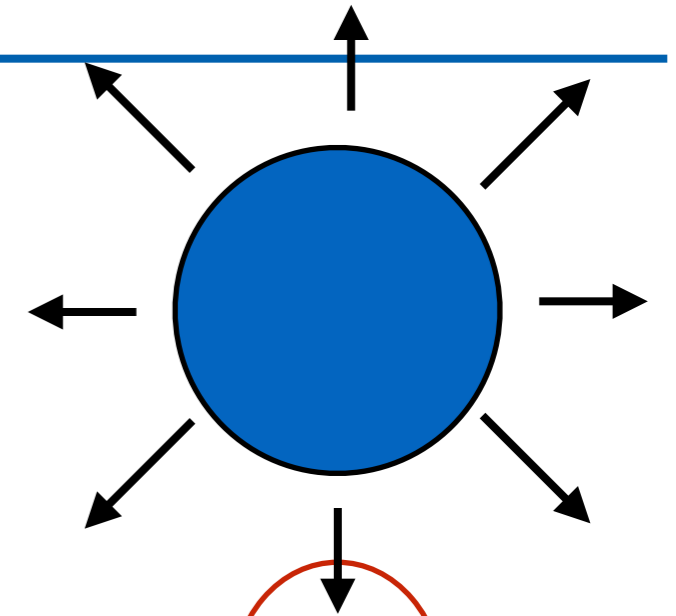


d+Au



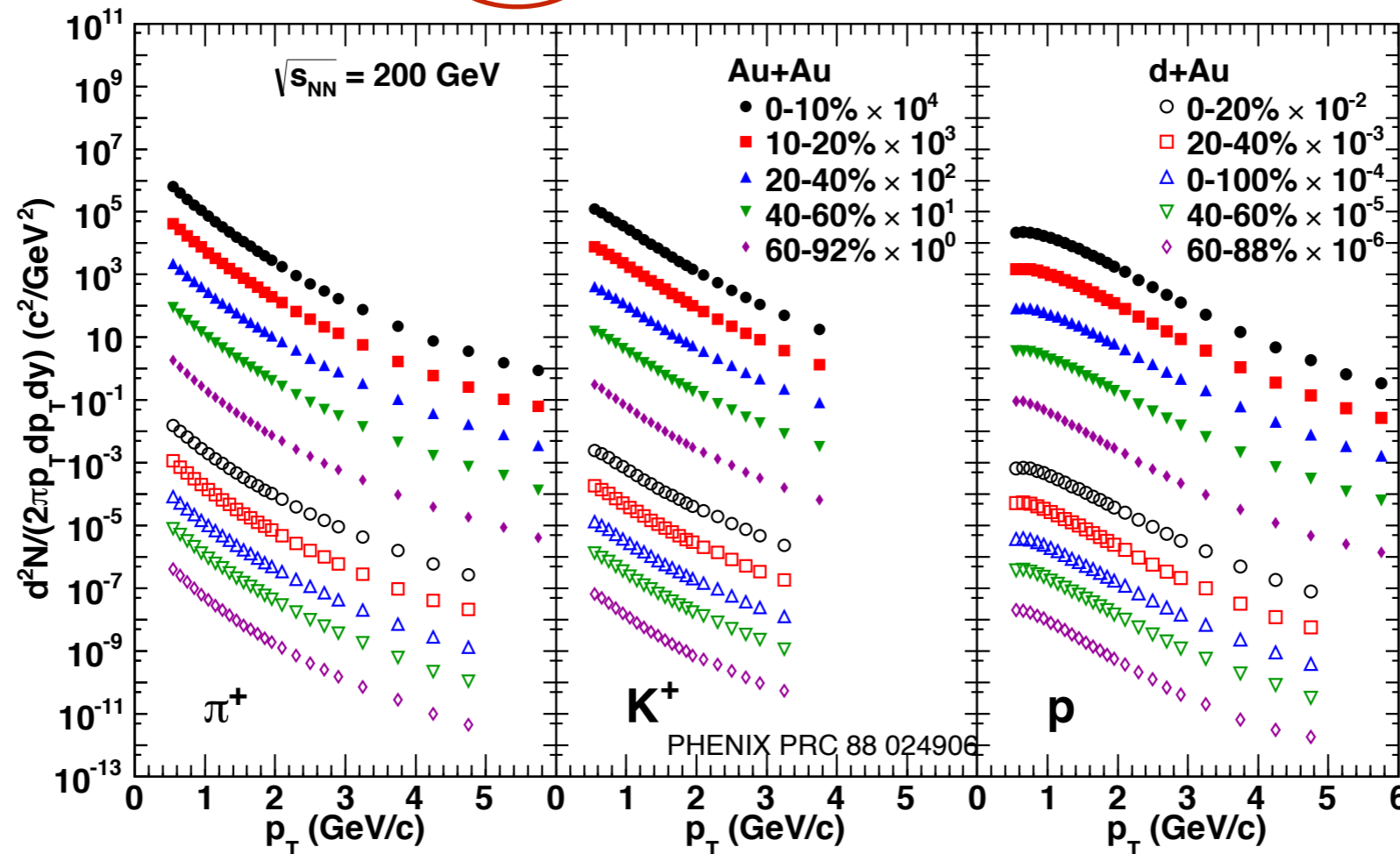
# what about radial flow?

**the Blast-Wave:** outward velocity boost, from a hydrodynamic source



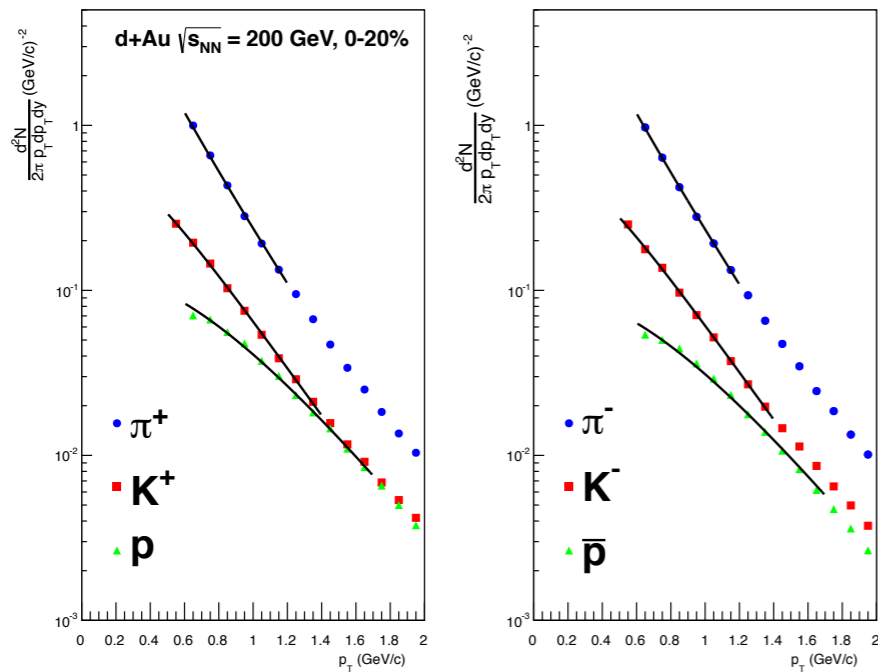
$$\frac{1}{p_T} \frac{dN}{dp_T} \propto \int_0^R r dr m_T I_0 \left( \frac{p_T \sinh \rho}{T_{fo}} \right) K_1 \left( \frac{m_T \cosh \rho}{T_{fo}} \right) \quad \rho = \tanh^{-1} (\beta_{max} (r/R)^n)$$

Schnedermann, Sollfrank, & Heinz PRC48 2462 (1993)

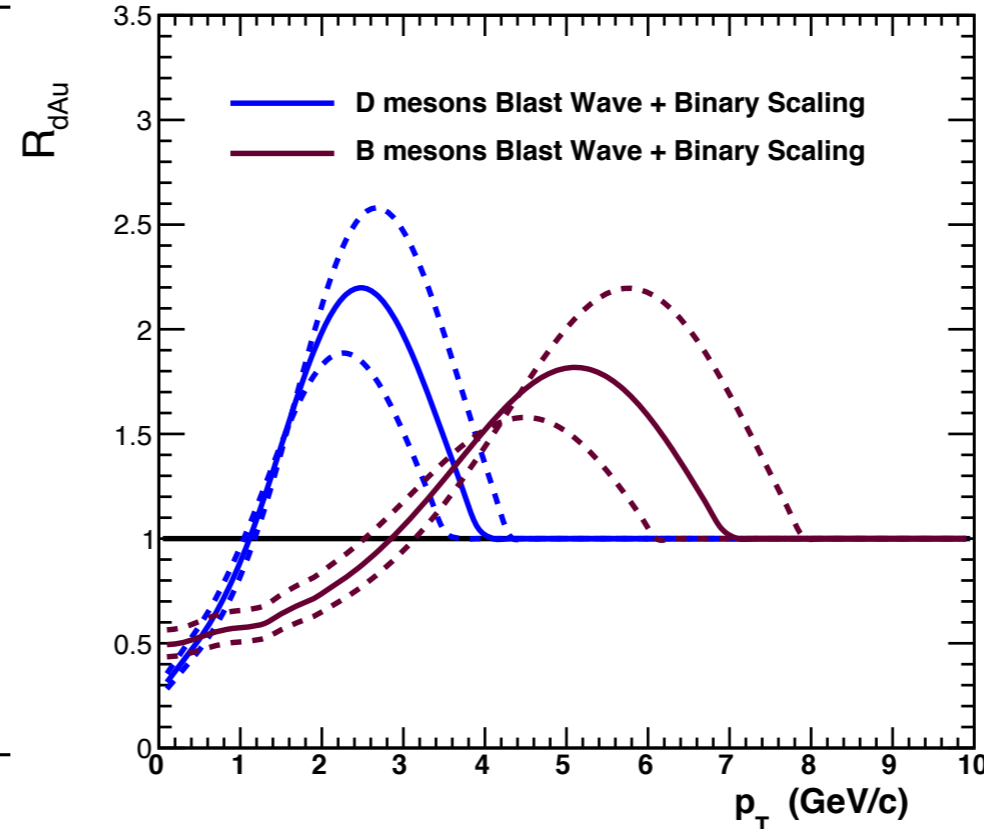
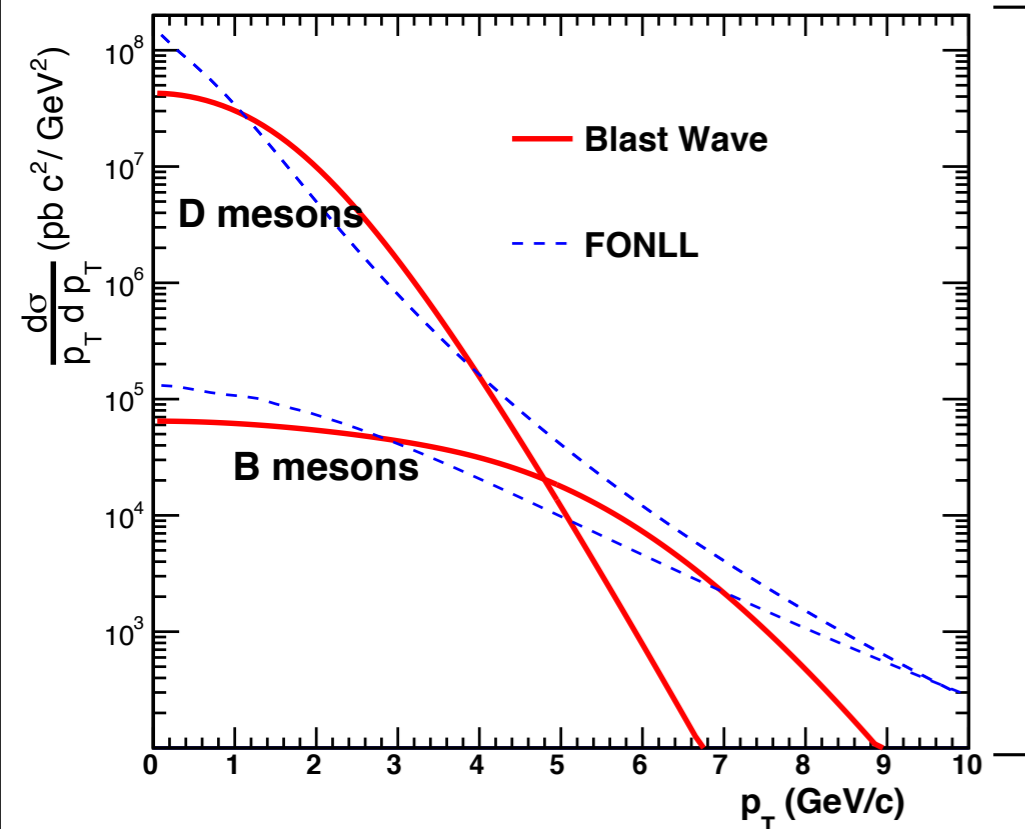




# blast-wave fit to dAu data

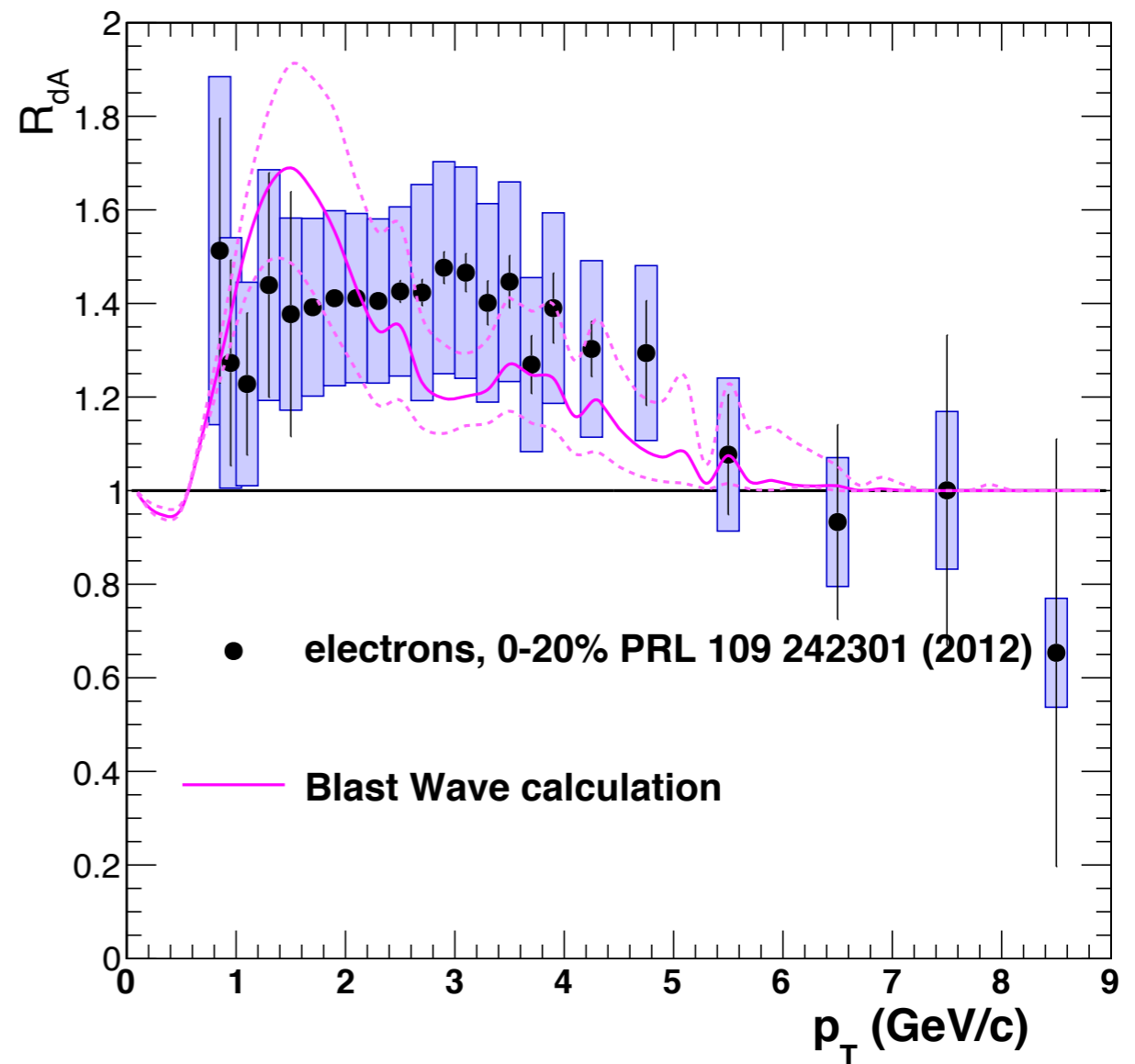
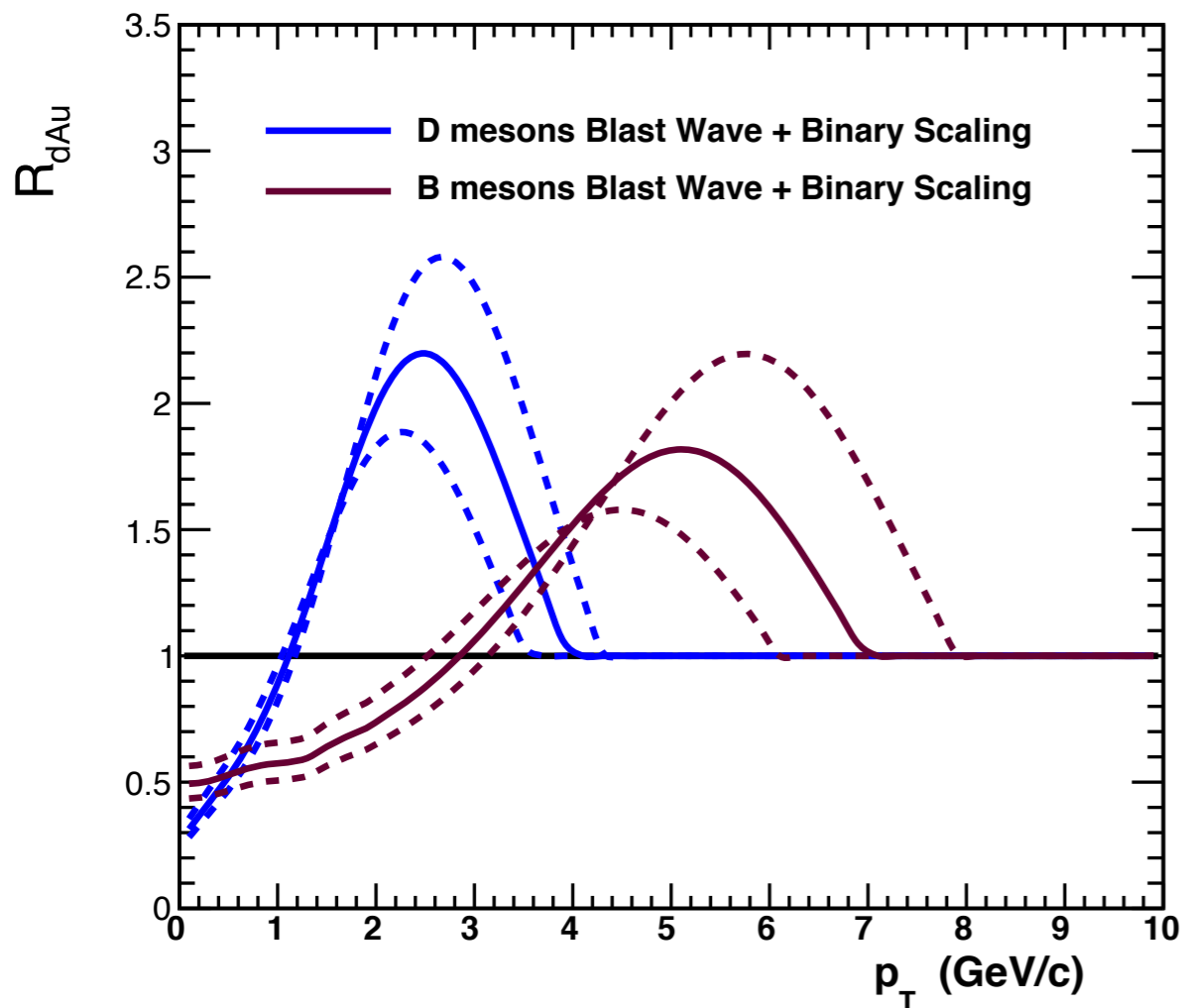


0-20% d+Au  
simultaneous fit to  $\pi$ ,  $K$ ,  $p$   
 $\beta_{\max} = 0.70$   
 $T_{fo} = 139\text{MeV}$



large enhancement of heavy mesons!

# and for the electrons?



another flow effect?

charm and bottom separated measurements key to clarifying

# jets in dAu

---

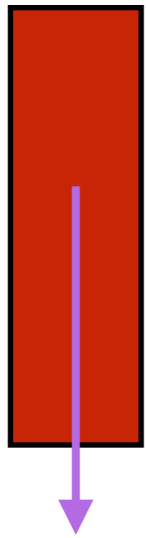
**AA collisions: quenching  
depends on L**



# jets in dAu

---

**AA collisions: quenching  
depends on L**

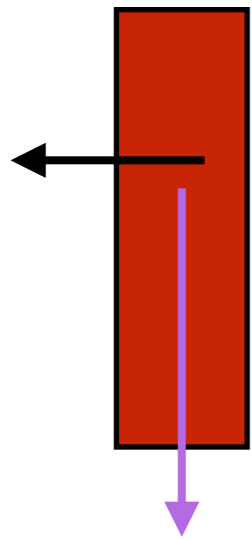


a lot of quenching

# jets in dAu

---

**AA collisions: quenching depends on L**

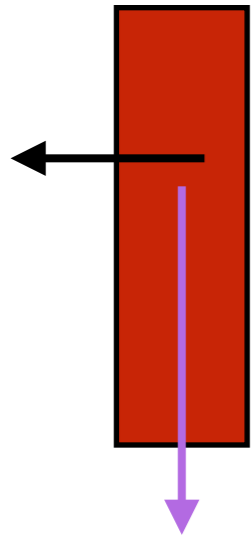


a little quenching

a lot of quenching

# jets in dAu

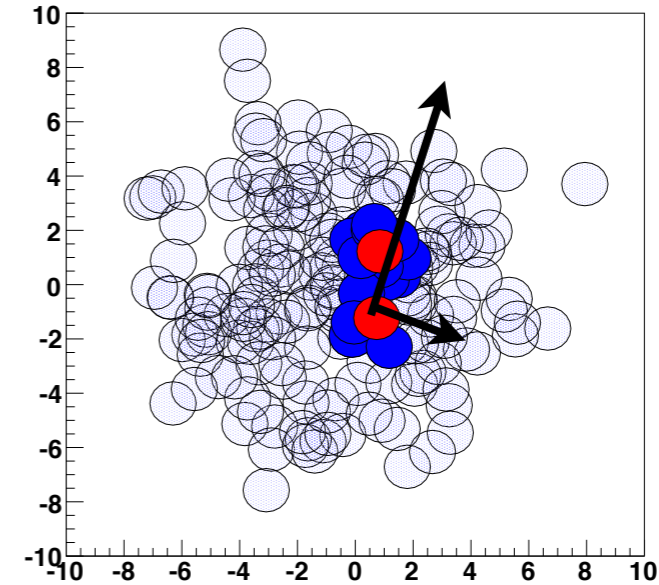
**AA collisions: quenching depends on L**



a little quenching

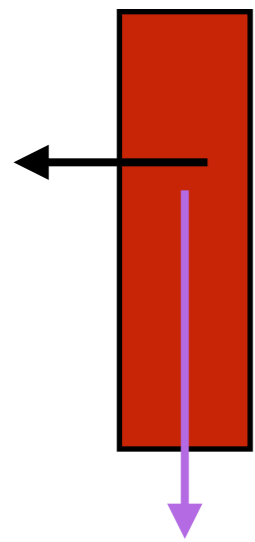
a lot of quenching

could something similar happen in dA?



# jets in dAu

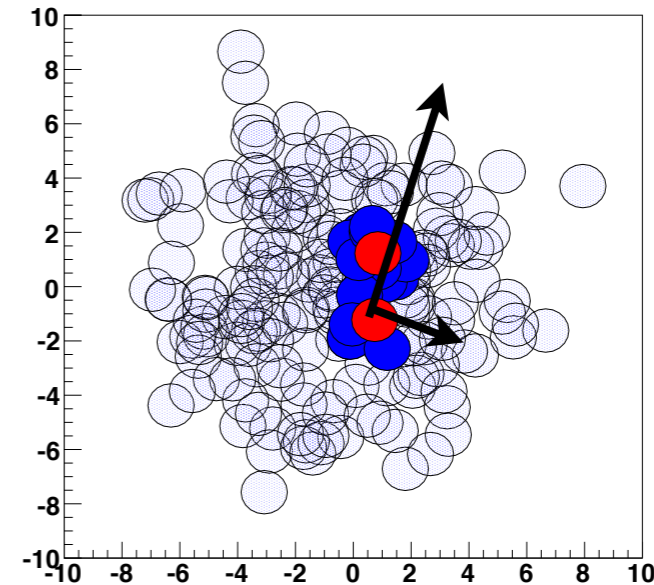
**AA collisions: quenching depends on L**



a little quenching

a lot of quenching

could something similar happen in dA?

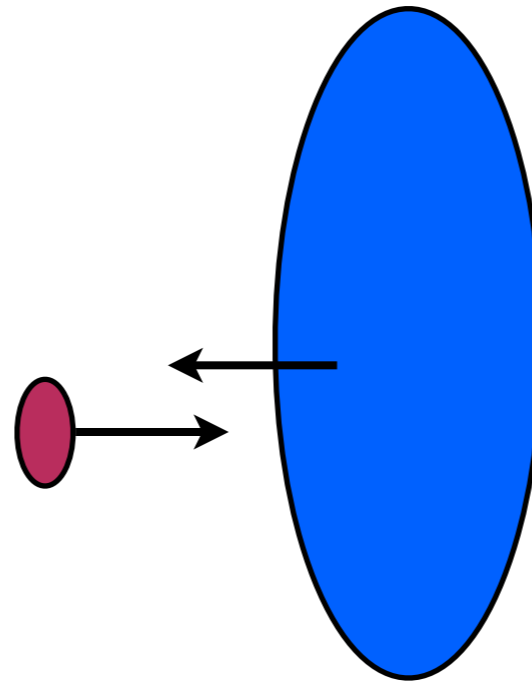


geometrical dependence might be observable even though we know the overall level of quenching is small in dAu

recent calculations (Zhang & Liao, 1311.5463), show **~10x** bigger effect in dAu than pPb

# investigating initial state of the nucleus?

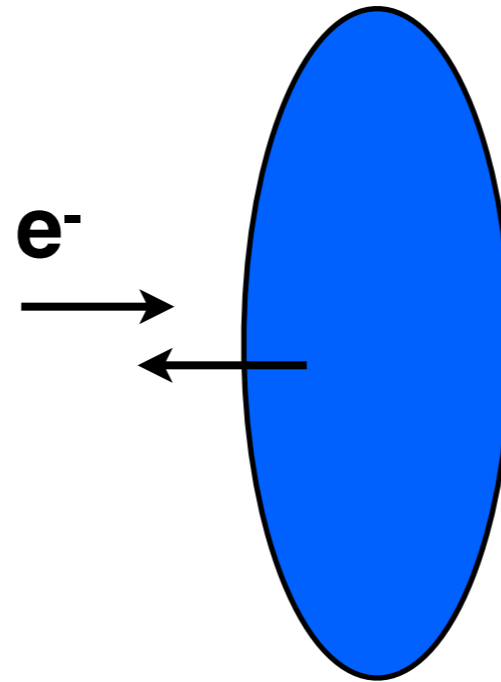
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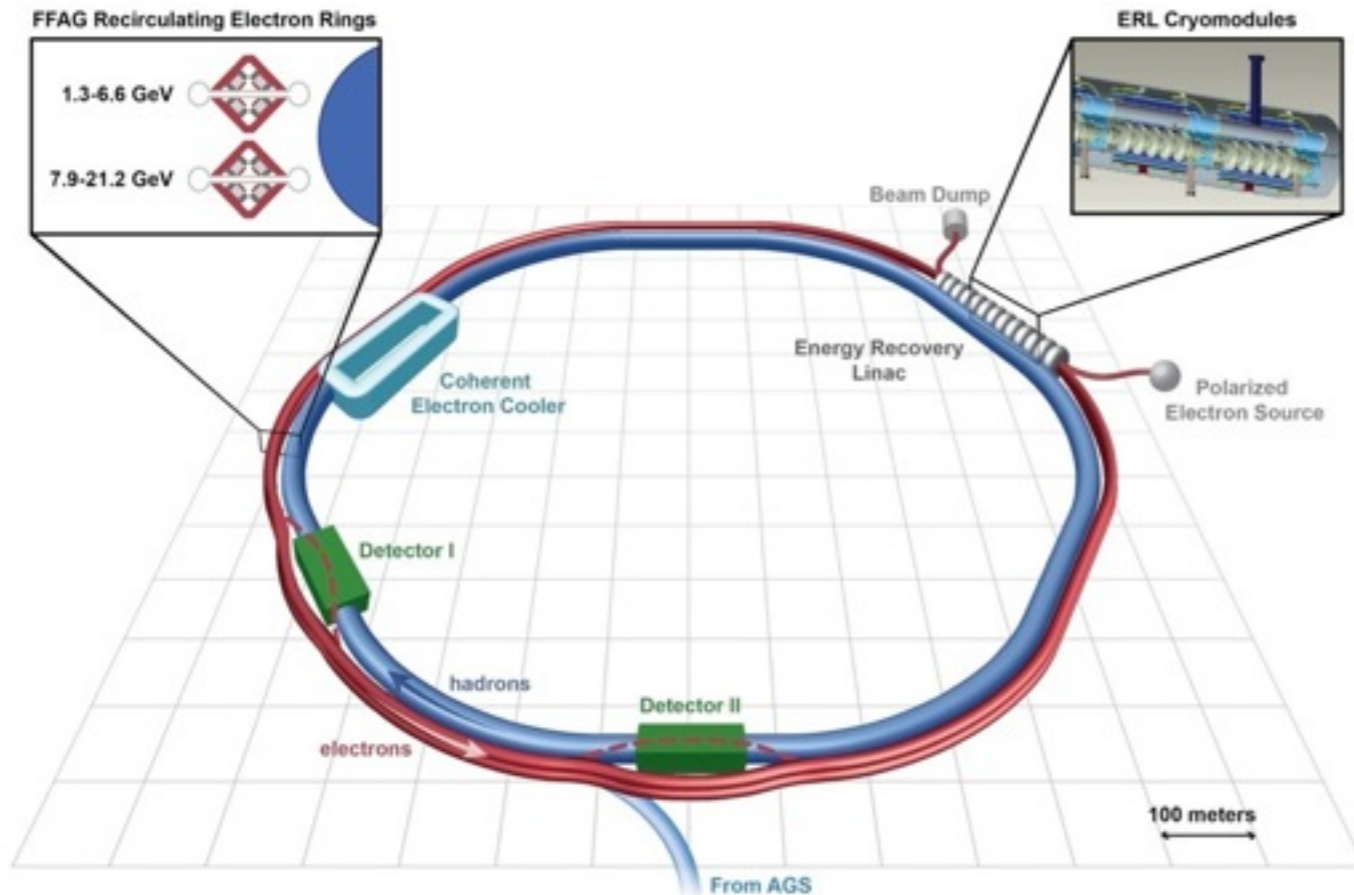
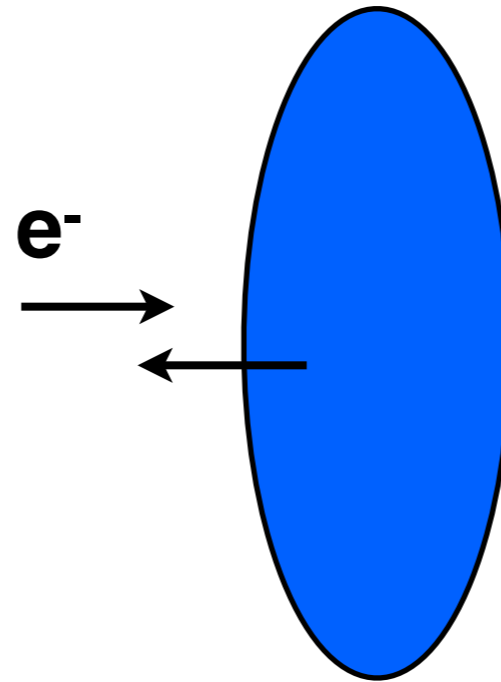


# investigating initial state of the nucleus?

---



# investigating initial state of the nucleus?

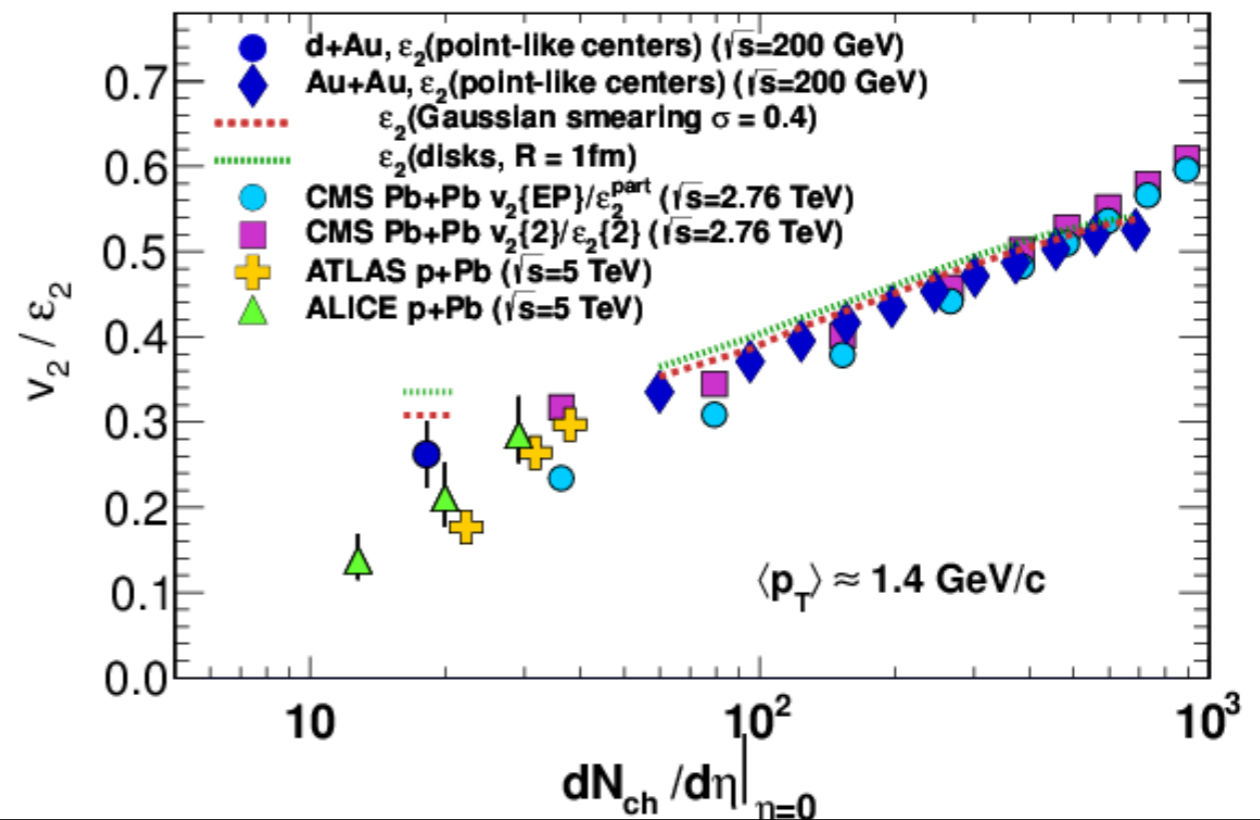
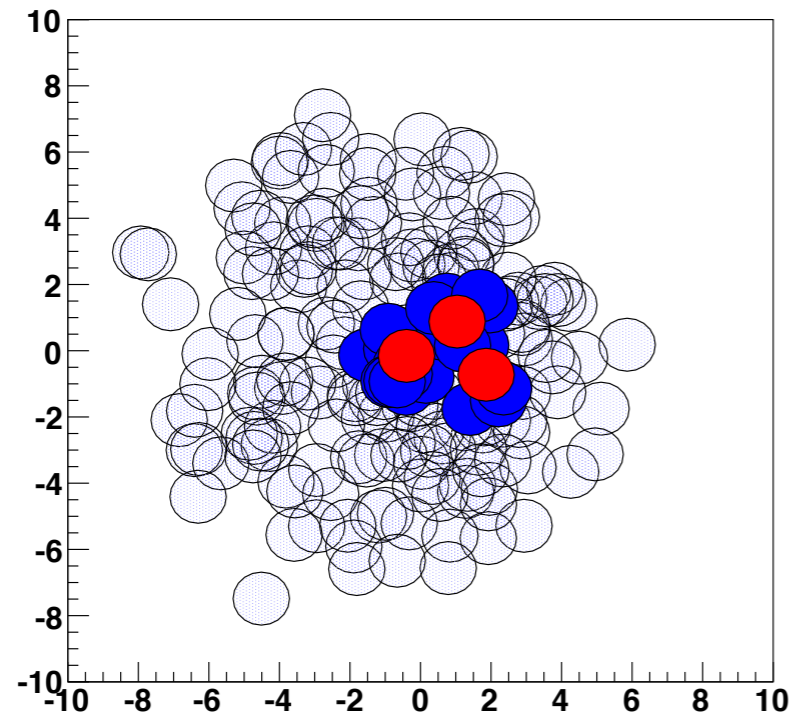


## eRHIC

upgrade to allow  
electrons at RHIC  
timescale ~ 2025

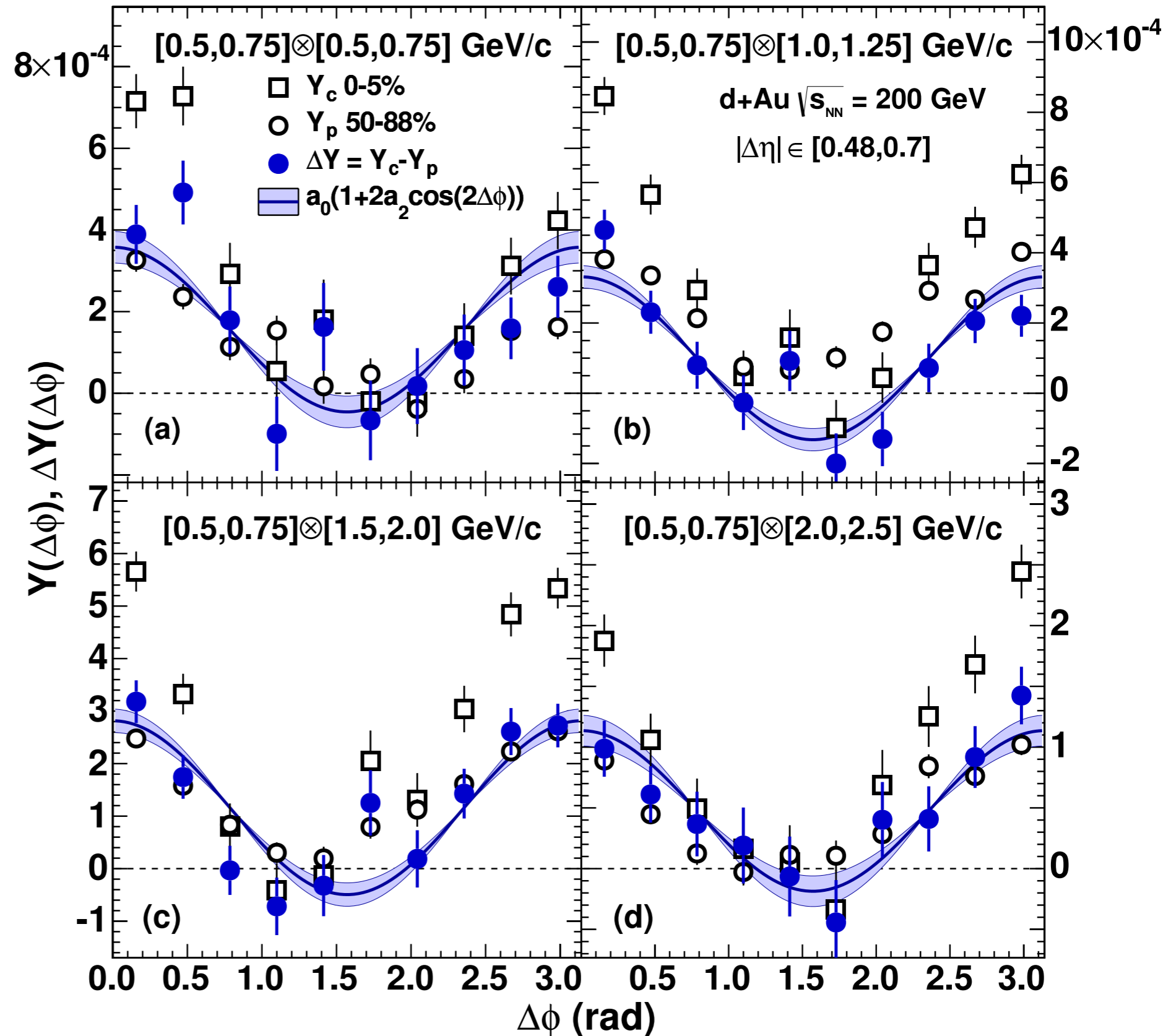
# pushing the limits of the QGP

- RHIC and the LHC are pushing the size limits of the quark gluon plasma
- suggestive of evolution, rather than a transition, from big to small systems
- looking forward to new measurements very soon to support/challenge this interpretation and quantitative understanding



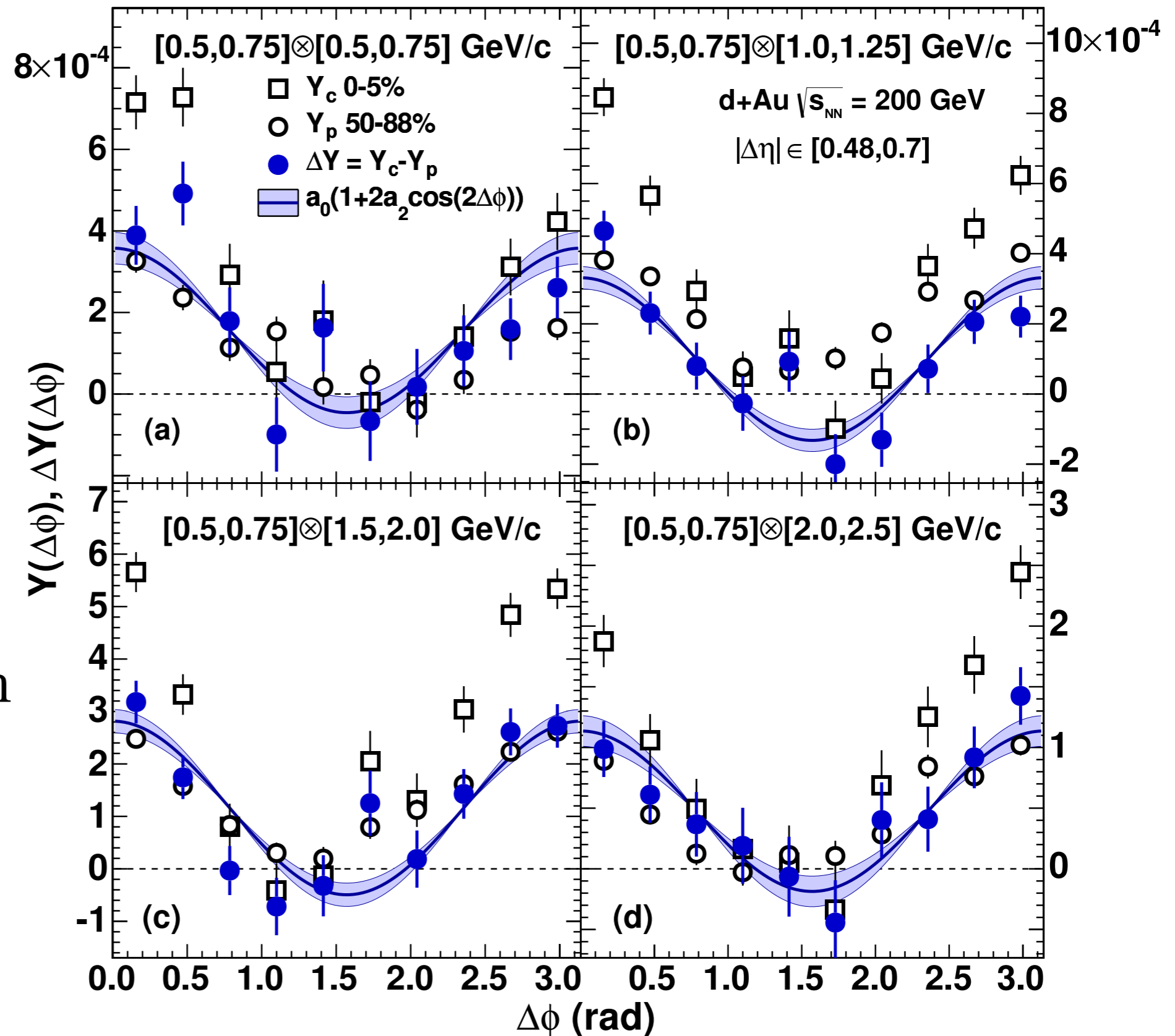
- 
- backups

centrality  
dependence  
consistently  
described by  
 $\cos 2\Delta\phi$  shape  
evidence for  
double ridge



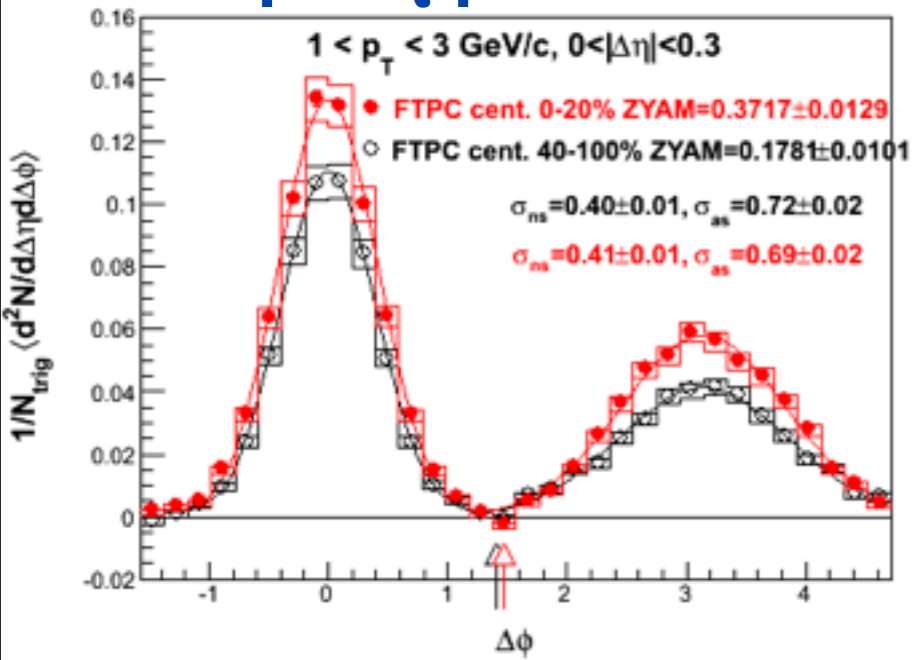
centrality  
dependence  
consistently  
described by  
 $\cos 2\Delta\phi$  shape  
evidence for  
double ridge

but is this just an  
artifact of the  
small  $|\Delta\eta|$   
acceptance?



# results from STAR

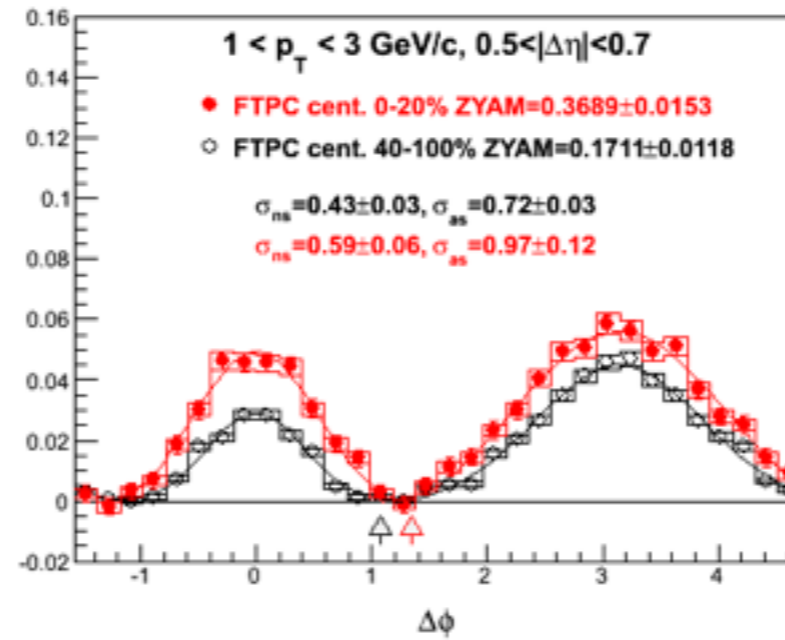
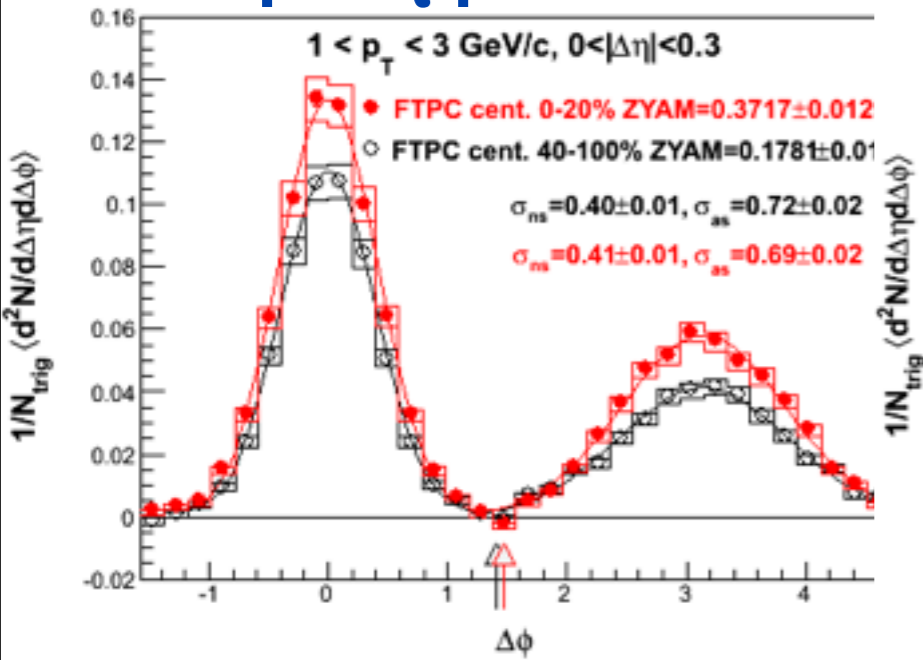
$|\Delta\eta| < 0.3$



# results from STAR

$|\Delta\eta| < 0.3$

$0.5 < |\Delta\eta| < 0.7$



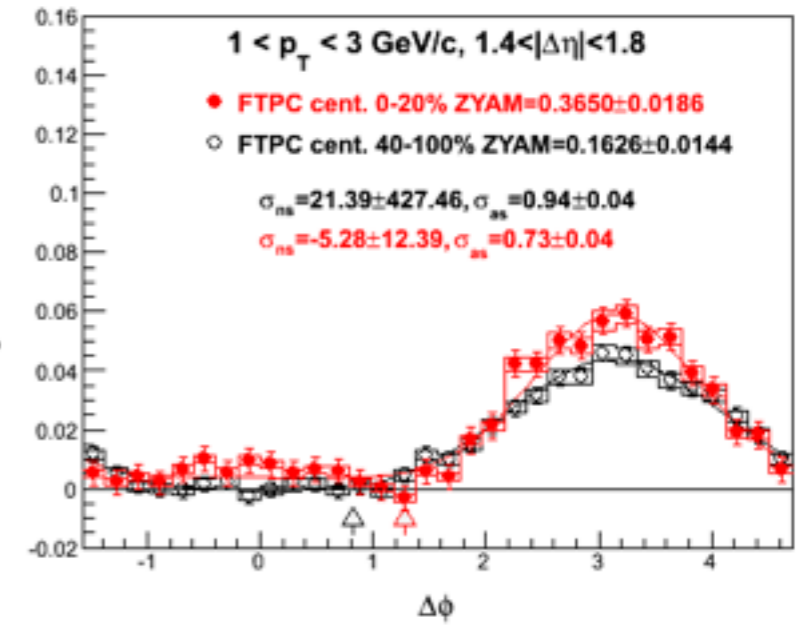
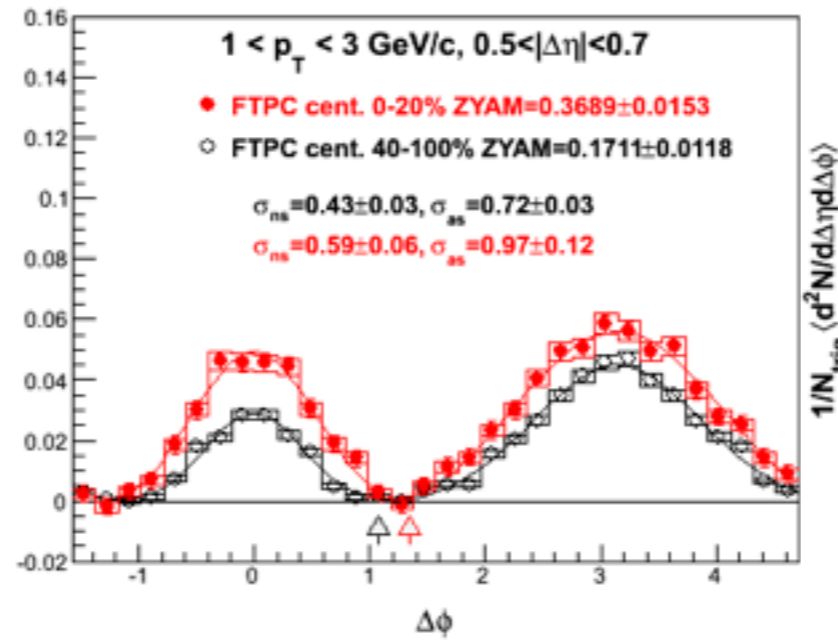
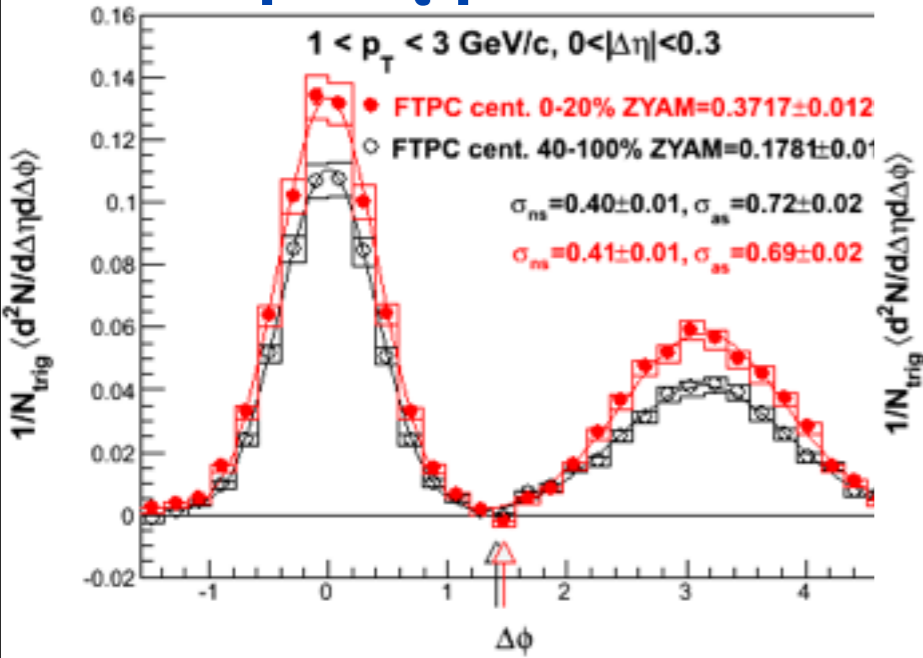


# results from STAR

$|\Delta\eta| < 0.3$

$0.5 < |\Delta\eta| < 0.7$

$1.4 < |\Delta\eta| < 1.8$

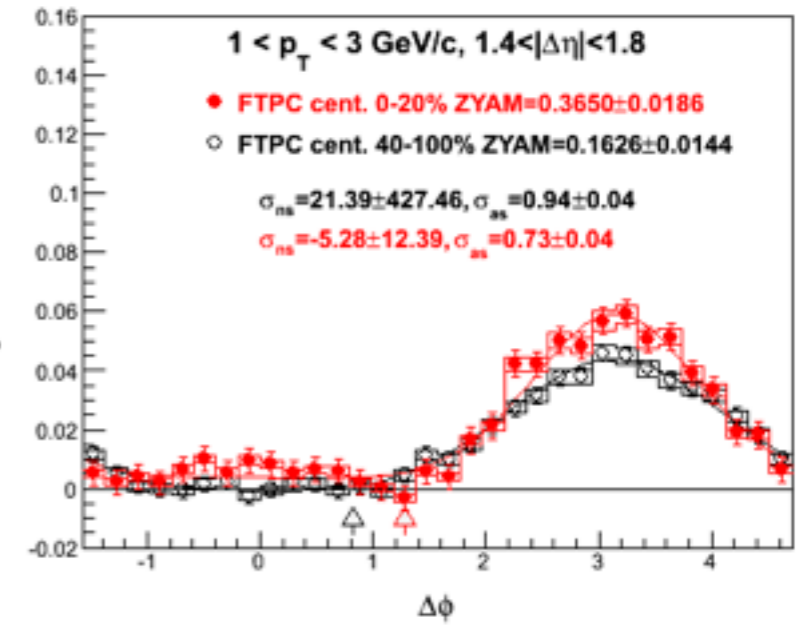
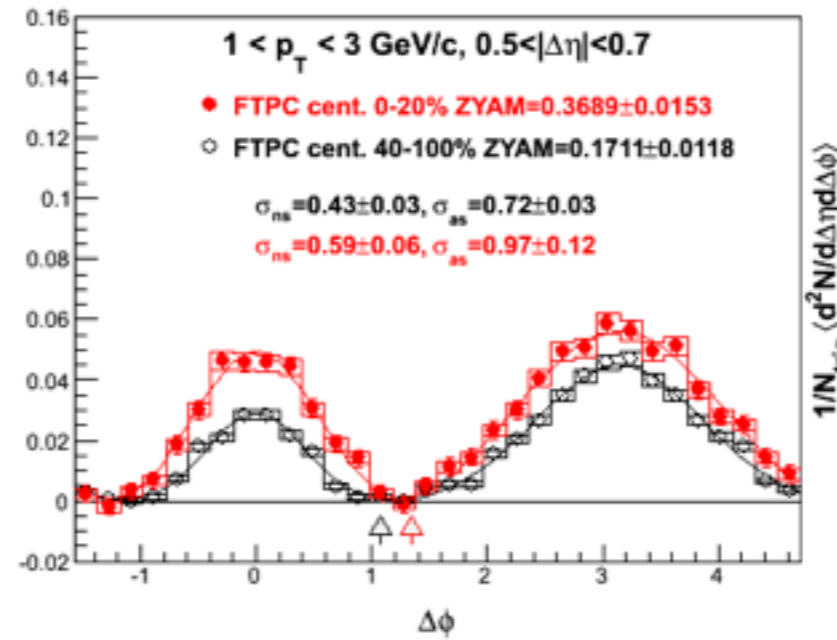
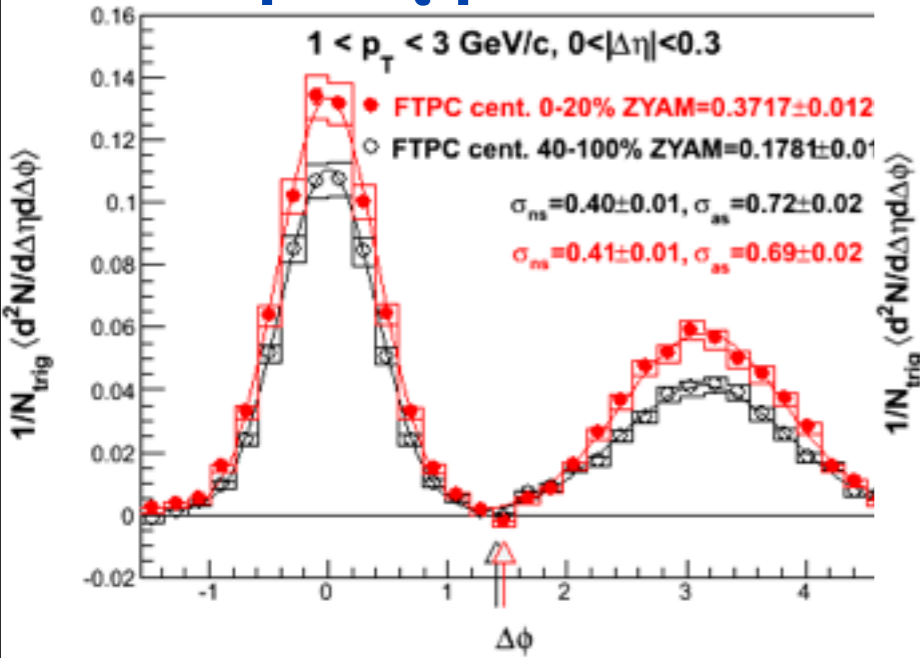


# results from STAR

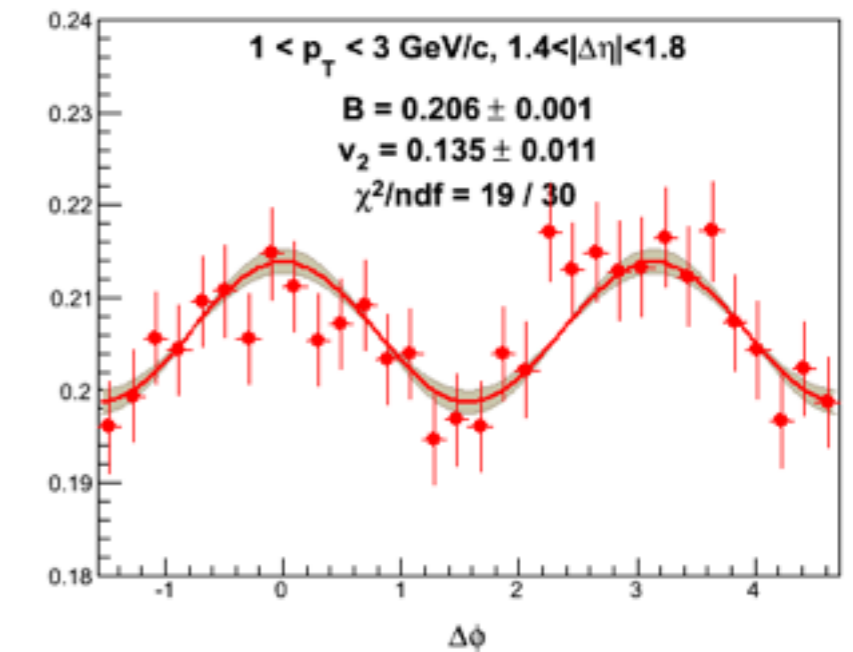
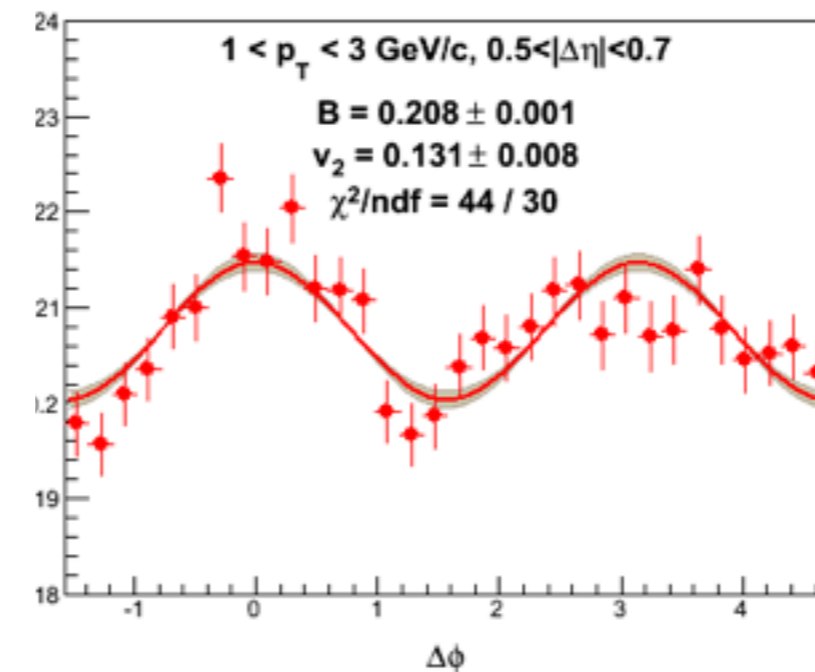
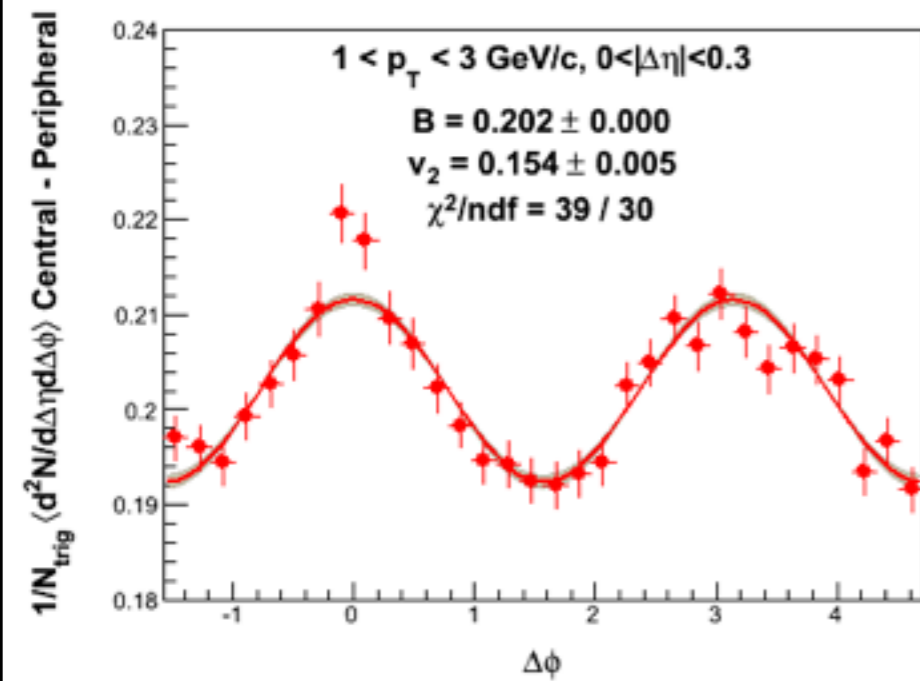
$|\Delta\eta| < 0.3$

$0.5 < |\Delta\eta| < 0.7$

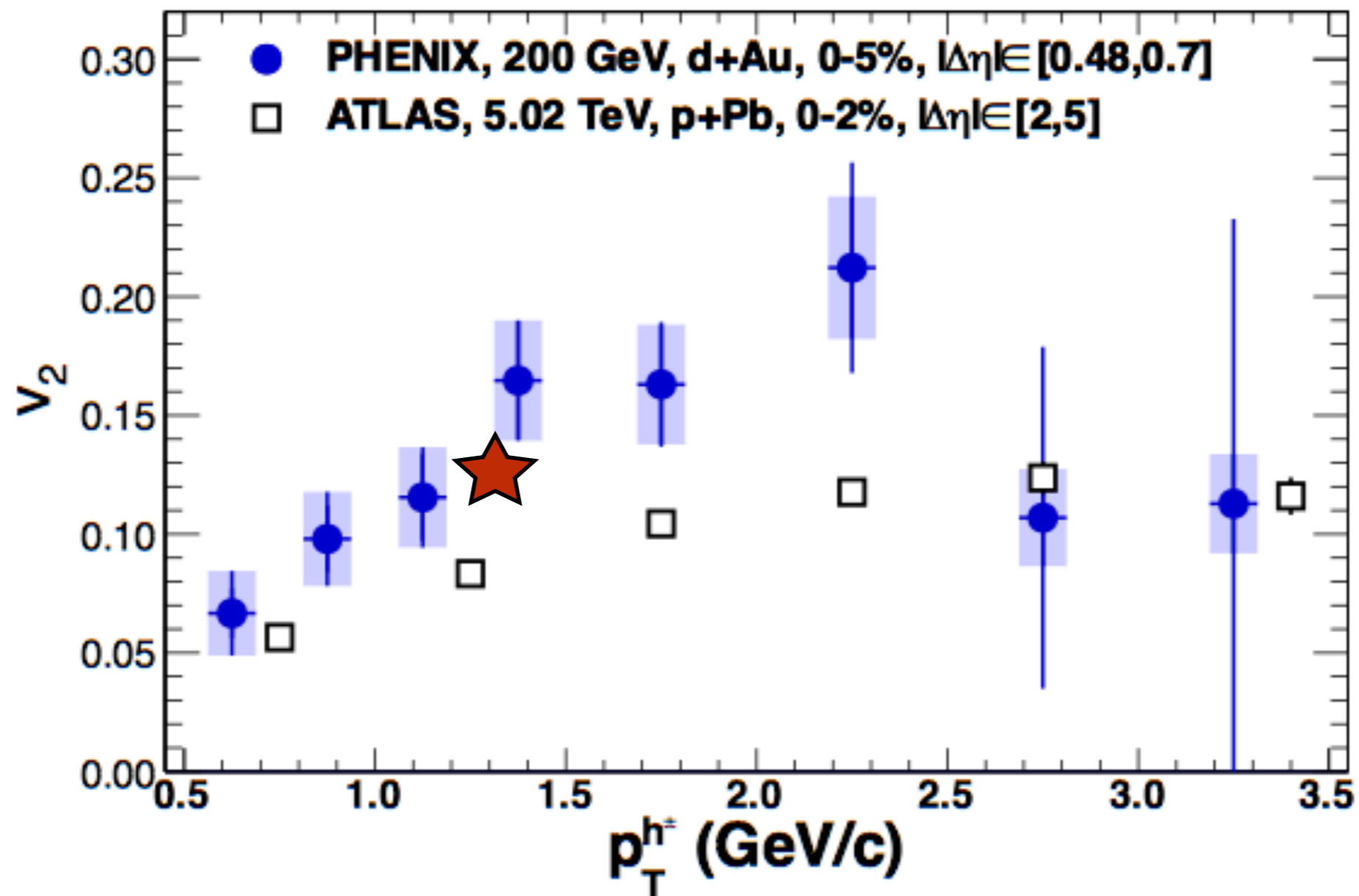
$1.4 < |\Delta\eta| < 1.8$



central - peripheral

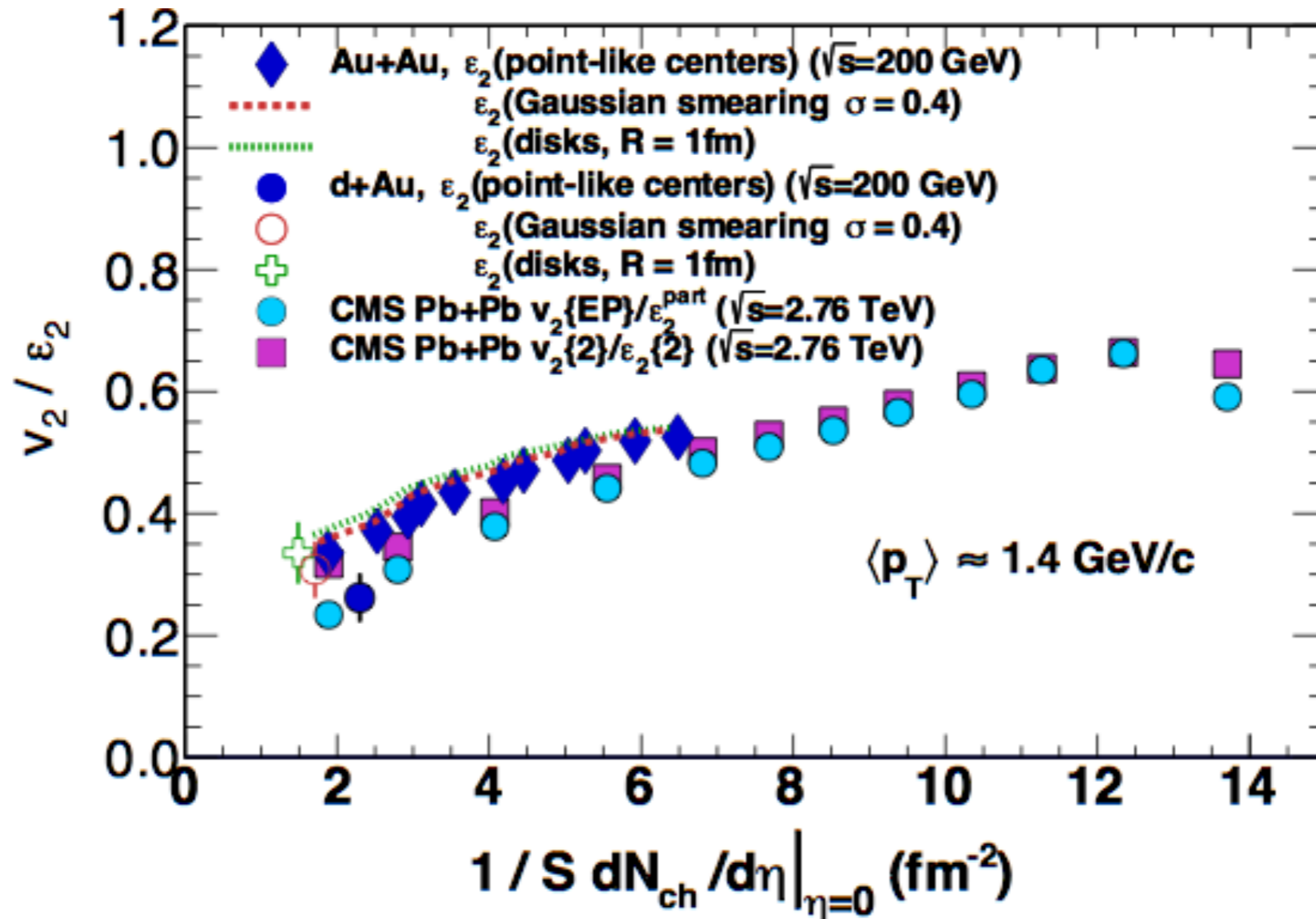


# RHIC comparisons



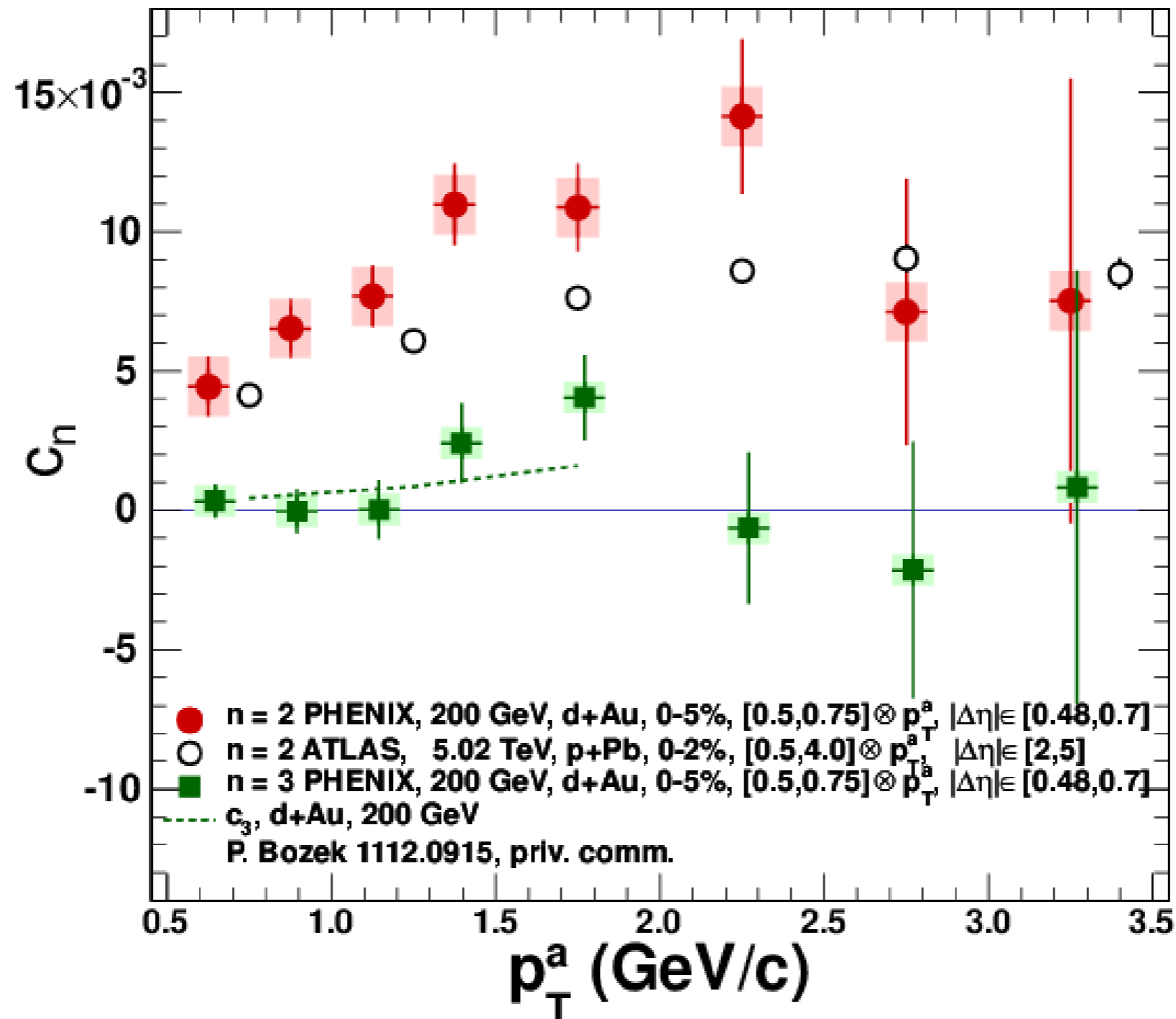
STAR  $v_2$ :  $\sim 13 \pm 1\%$   $1 < p_T < 3 \text{ GeV}/c$   
good consistency at RHIC!

# scaling with overlap area?



- approximate scaling with  $1/S \, dN_{\text{ch}}/d\eta$
- significant uncertainties due to nucleon representations in d+Au
- n.b. not directly comparable to other  $1/S$  plots, here  $v_2$  at fixed  $p_T$ !

# $v_3$ at RHIC?



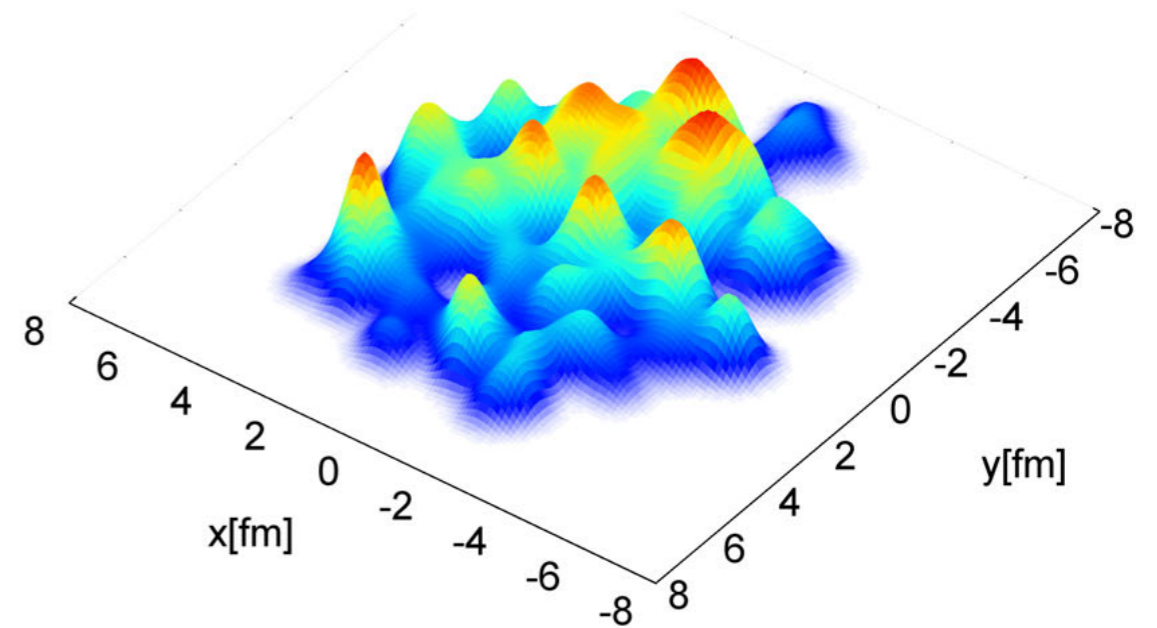
**no evidence for significant  $v_3$ , consistent with hydro expectations**

# nucleon positions to energy density

---

single event initial energy density

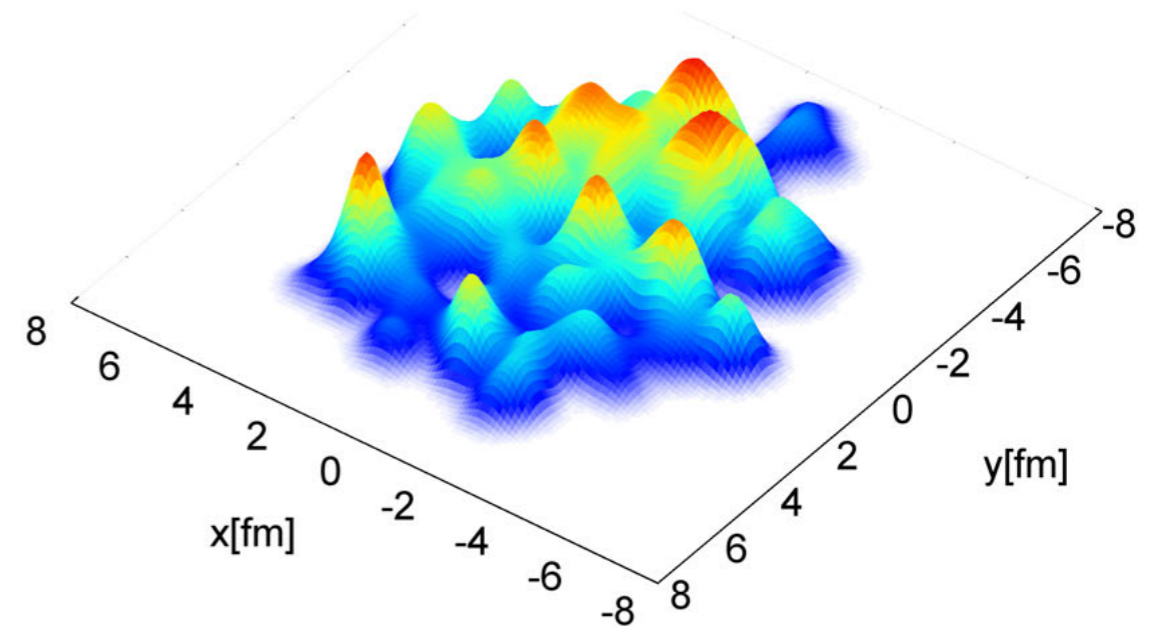
nucleons: Gaussians,  
 $\sigma = 0.4\text{fm}$



# nucleon positions to energy density

single event initial energy density

nucleons: Gaussians,  
 $\sigma = 0.4\text{fm}$



subnucleonic fluctuations:  
IP-Glasma model

