

## Where Feynman, Field and Fox failed and how we fixed it at RHIC

M. J. Tannenbaum<sup>a</sup>

<sup>a</sup>Physics Department, 510c, Brookhaven National Laboratory  
Upton, N.Y., 11973-5000, U.S.A., *mjt@bnl.gov*

Hard-scattering of point-like constituents (or partons) in p-p collisions was discovered at the CERN-ISR in 1972 by measurements utilizing inclusive single or pairs of hadrons with large transverse momentum ( $p_T$ ). Due to the steeply falling power-law  $p_T$  spectrum of the hard-scattered partons, the inclusive single particle (e.g.  $\pi^0$ )  $p_{T_t}$  spectrum from jet fragmentation of a parton with  $\hat{p}_{T_t}$  is dominated by trigger fragments with large  $\langle z_t \rangle \sim 0.7 - 0.8$ , where  $z_t = p_{T_t}/\hat{p}_{T_t}$  is the fragmentation variable. It was generally assumed, following Feynman, Field and Fox [1], as shown by data from the CERN-ISR experiments, that the  $p_{T_a}$  distribution of away side hadrons from a single particle trigger [with  $p_{T_t}$ ], corrected for  $\langle z_t \rangle$ , would be the same as that from a jet-trigger and follow the same fragmentation function as observed in  $e^+e^-$  or DIS. PHENIX [2] attempted to measure the fragmentation function from the away side  $x_E \sim p_{T_a}/p_{T_t}$  distribution of charged particles triggered by a  $\pi^0$  in p-p collisions and showed by explicit calculation that the  $x_E$  distribution is actually quite insensitive to the fragmentation function. A new formula for the distribution of an associated away-side particle with transverse momentum  $p_{T_a}$ , which is presumed to be a fragment of an away-jet with  $\hat{p}_{T_a}$ , with exponential fragmentation function  $D(z) = Be^{-bz}$ , triggered by a particle with transverse momentum  $p_{T_t}$ , presumably from a trigger-side jet with  $\hat{p}_{T_t}$ , invariant cross section,  $Ed^3\sigma/dp^3 = A\hat{p}_{T_t}^{-n}$ , was given [2]:  $dP_{p_{T_a}}/dx_E|_{p_{T_t}} \approx \frac{\langle m \rangle}{\hat{x}_h} \frac{(n-1)}{(1+x_E/\hat{x}_h)^n}$ , which relates  $x_E$ , the ratio of the transverse momenta of the measured particles, to  $\hat{x}_h = \hat{p}_{T_a}/\hat{p}_{T_t}$ , the ratio of the transverse momenta of the away-side to trigger-side jets, where  $\langle m \rangle$  is the mean multiplicity of particles in the away jet. Many analyses of the away-jet  $p_{T_a}$  distributions in Au+Au collisions are available; but these tend to describe the effect of the medium with the variable  $I_{AA}(x_E)$ , the ratio of the  $x_E$  distribution in A+A collisions to that in p-p collisions, which typically shows an enhancement at low values of  $x_E$  and a suppression at higher values of  $x_E$ . Such behavior could be explained as a decrease in  $\hat{x}_h$  in A+A collisions due to energy loss of the away jet in the medium. Fits of the above formula to the available data will be presented to establish whether: a) the away-jets simply lose energy; b) some of the away-jets lose energy, others punch-through without losing energy; etc.

### References

- [1] R. P. Feynman, R. D. Field, and G. C. Fox, Nucl. Phys. B **128** (1977) 1–65.
- [2] S. S. Adler, et al., PHENIX Collaboration, Phys. Rev. D **74** (2006) 072002.