



Particle Correlation Measurements from

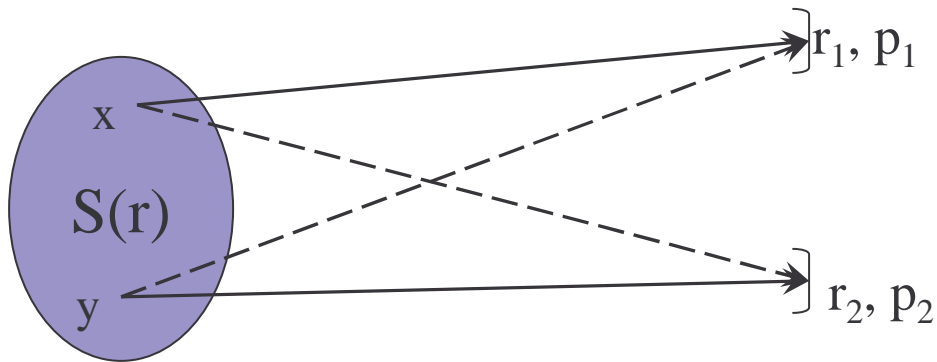


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for the PHENIX Collaboration

HBT results presented here:
submitted to PRL: 16 Jan `02
nucl-ex/0201008



The Physics of Hanbury-Brown Twiss



n Consider a source $S(r)$ of identical bosons (γ or π) whose wave functions can be described as plane waves.

n Assume:

- Production amplitudes independent of momentum
- Mutually incoherent

Amplitude for this diagram:



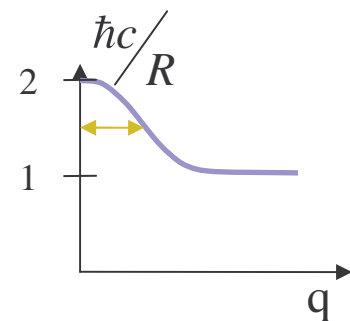
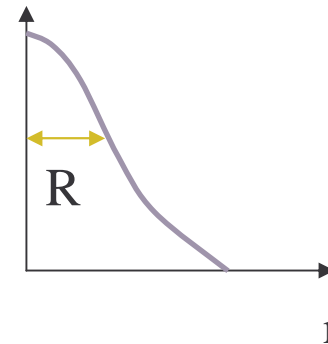
$$A = \frac{1}{\sqrt{2}} \left[e^{ip_1(r_1-x)} e^{ip_2(r_2-y)} + e^{ip_1(r_1-y)} e^{ip_2(r_2-x)} \right]$$

$$A^2 = 1 + e^{i(p_1-p_2)(y-x)}$$

$$\approx 1 + \cos(\Delta p \Delta x)$$

Corresponding normalized probability:

$$P = \int dx dy |A|^2 S(x) S(y) = 1 + \left| \tilde{S}(q) \right|^2 \equiv C_2(q)$$





Some details:

In principle:

$$C_2 = \frac{\langle n \rangle^2}{\langle n(n-1) \rangle} \frac{dn/dp_1 dp_2}{\left(\frac{dn}{dp_1} \right) \left(\frac{dn}{dp_2} \right)}$$

In practice:

$$C_2 = \frac{A(q)}{B(q)}$$

Pairs from same event

n One dimension works fine when measuring stars:

- Static
- Isotropic emission (no position momentum correlations)

n Heavy Ion Collisions are anything but

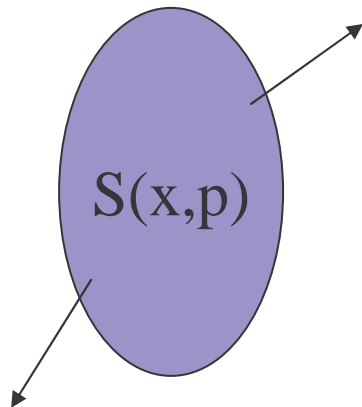
- Consider a more complicated source

Pairs from 'mixed' events

n S() – effective single particle Wigner phase space density

- Often replaced by a classical phase-space density in practical calculations

n Note: due to mass shell constraints, the function is non-invertible (model assumptions).



$$E \frac{dN}{d^3 p} = \int d^4 x S(x, p)$$

$$\begin{aligned} q &= p_1 - p_2 \\ k &= p_1 + p_2 / 2 \end{aligned}$$

$$C_2(q, k) = 1 + \frac{\left| \int d^4 x S(x, k) e^{iq \cdot x} \right|^2}{\int d^4 x S(x, k + \frac{1}{2} q) \int d^4 y S(y, k - \frac{1}{2} q)}$$



Kinematic Variables

- Currently the field tends to consider projections of the momentum difference into:

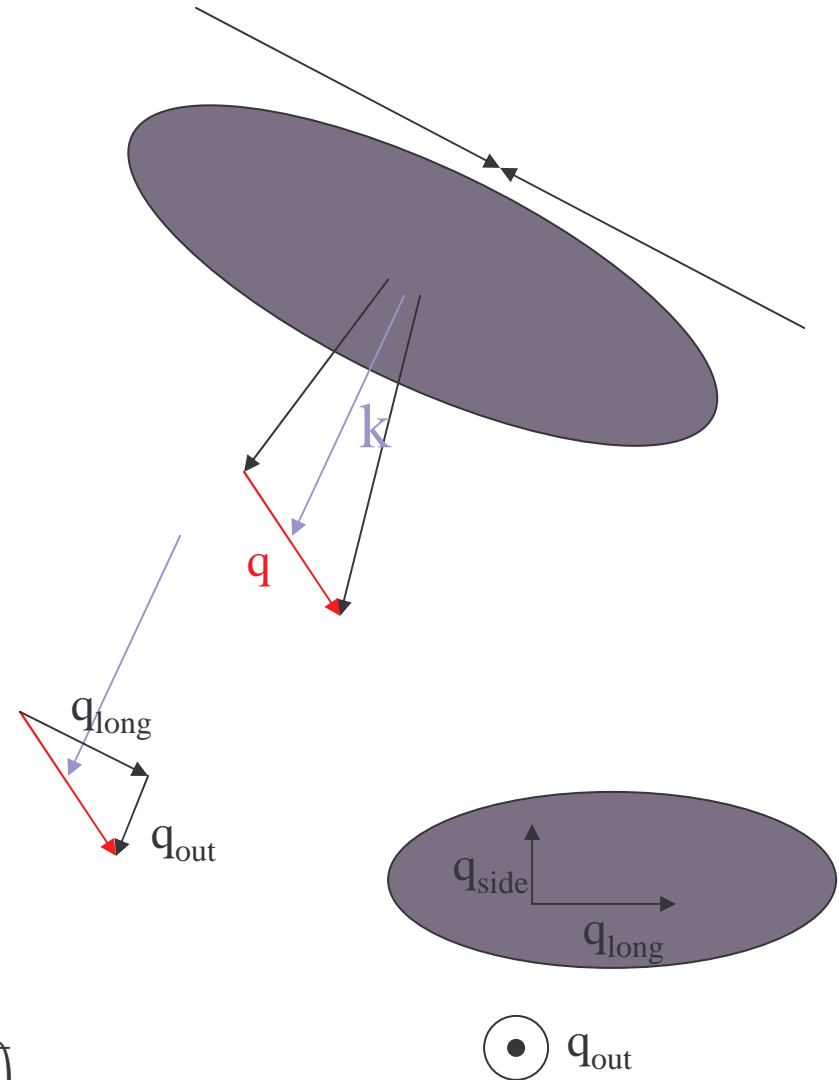
$$\square \quad q_{\text{long}}, q_{\text{side}}, q_{\text{out}}$$

- $q_{\text{long}} \propto R_{\text{long}} \sim \text{longitudinal extent}$

- $q_{\text{side}}/q_{\text{out}} \propto R_{\text{side}}/R_{\text{out}} \sim \text{transverse extent}$

$$\begin{aligned} q &= p_1 - p_2 \\ k &= p_1 + p_2 / 2 \end{aligned}$$

$$C_2(q, k) = 1 + \frac{\left| \int d^4 x S(x, k) e^{iq \cdot x} \right|^2}{\int d^4 x S(x, k + \frac{1}{2} q) \int d^4 y S(y, k - \frac{1}{2} q)}$$

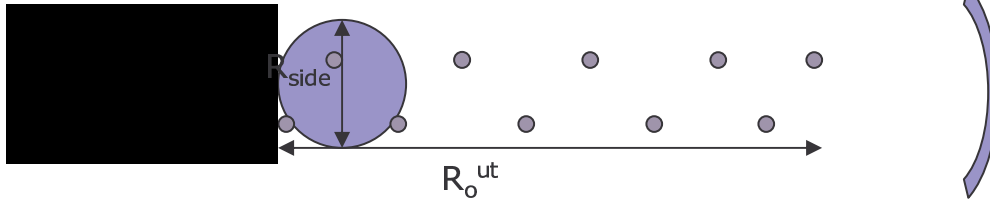




HBT and the QGP

n “Naïve” picture:

$$\square R_{\text{out}}^2 = R_{\text{side}}^2 + (\beta_{\text{pair}} \tau)^2$$



n Concrete predictions are few:

- Pratt PRD 1314 ('86): fireball and EOS $\tau \sim 90$ fm/c
- Bertsch NPA 173 (89) QGP + cascade $\tau \sim 12$ fm/c
- Hydro calculation of Rischke & Gyulassy expects $R_{\text{out}}/R_{\text{side}} \sim 2$ - >4 @ $k_t = 350$ MeV.
- Result robust to T_{freeze} , d_Q/d_H , 1st order vs. rapid cross-over.

n Response: can hadronic rescattering mask this prediction?

D.H. Rischke, M. Gyulassy/Nuclear Physics A 608 (1996) 479–512

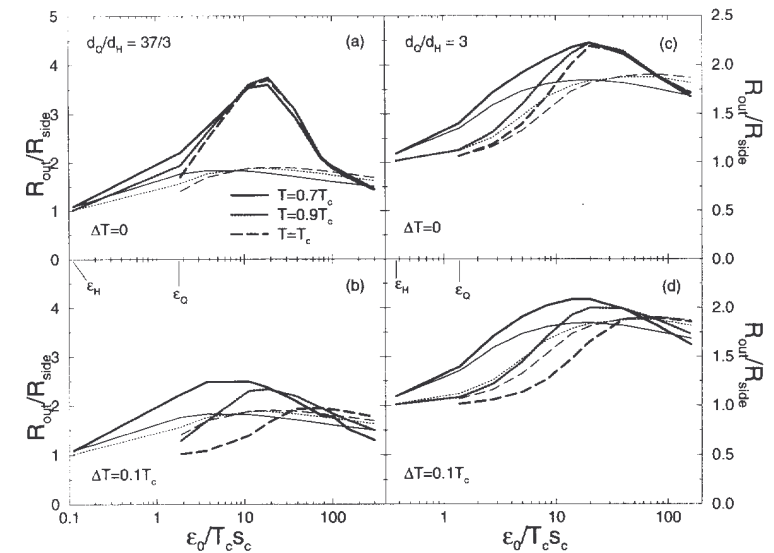
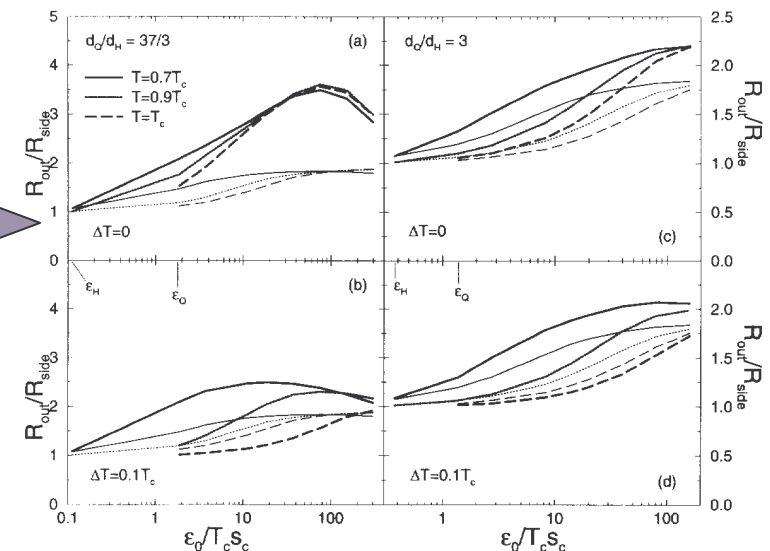


Fig. 17. The same as in Fig. 12, but for the ratio $R_{\text{out}}/R_{\text{side}}$.



One step closer

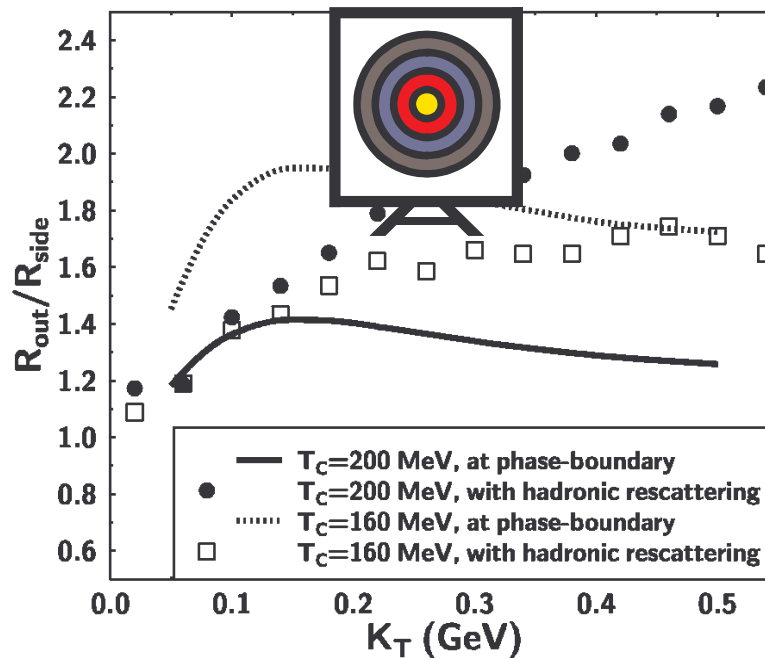


FIG. 3. R_{out}/R_{side} for RHIC initial conditions, as a function of K_T at freeze-out (symbols) and at hadronization (lines).

- n Soff, Bass, Dumitru (PRL 86)
 - Couple microscopic transport to hydro with phase transition
 - Still expect $R_{out}/R_{side} > 1$ \perp measurements at high k_t are very interesting.
- n Note:
 - Hydro: $R_o/R_s(200) < R_o/R_s(160)$
 - n Longer time at phase transition
 - Transport: $R_o/R_s(200) > R_o/R_s(160)$
 - n Longer time rescattering

PHENIX – Year 1 Configuration

- n Both arms provide hadron PID (contrary to popular belief)
- n East:
 - ▣ DC + TOF (~100 ps)
 - ▣ π/K separation to 2 GeV/c
- n West:
 - ▣ DC + PbSc (~600 ps)
 - ▣ π/K separation to 1 GeV/c
- n B-field + geometry limits lower k_t bound to 200 MeV.

$$\frac{\partial p}{p} = 0.6\% \oplus 3.6\% p$$

I use 'electron' acceptance for this hadron

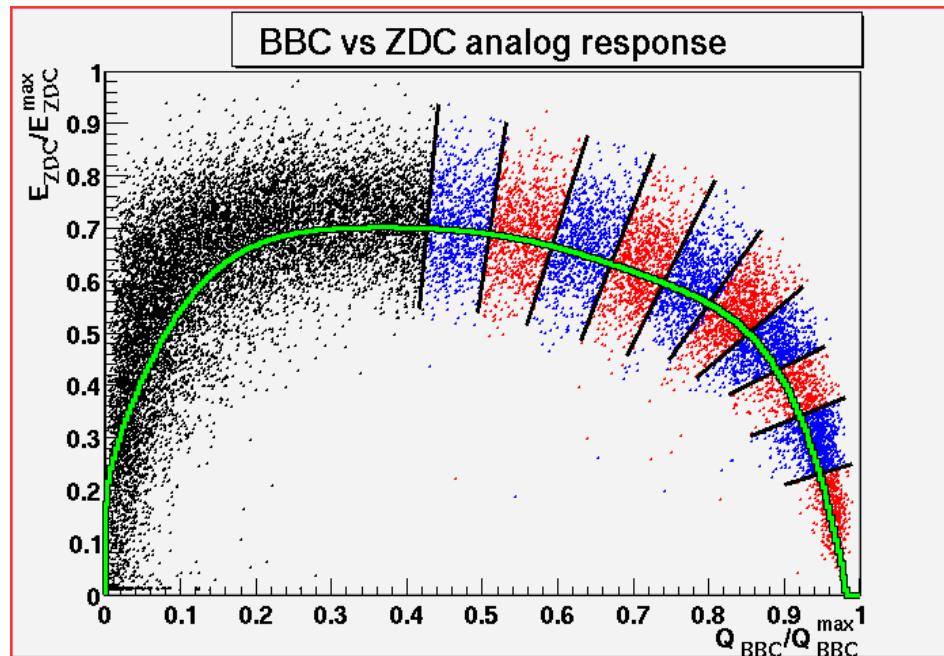


-2.0 0.0 2.0

Rapidity



Centrality definition and sample

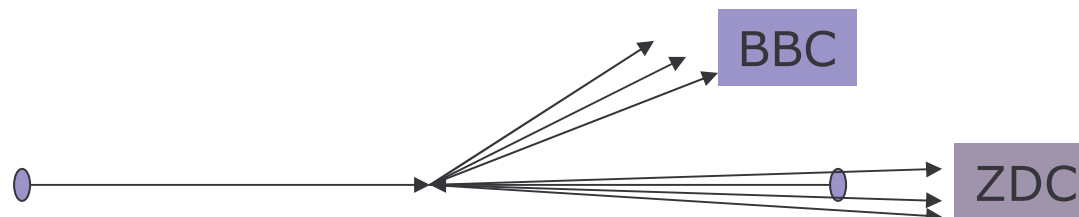


This data sample uses the 30% most central collisions

- $\langle \text{cent}_{\text{pairs}} \rangle = 10\%$
- 493k events

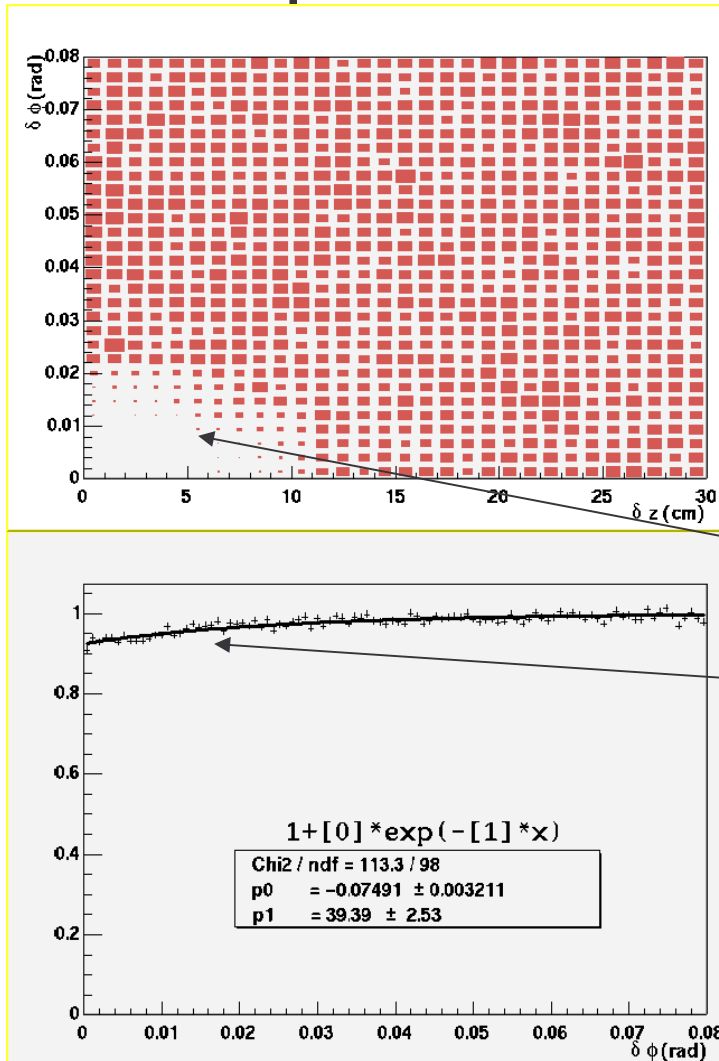
After all analysis cuts:

- 3.1M π^+ pairs
- 3.3M π^- pairs





Pair acceptance and corrections



- n π definition $\pm < 1.5\sigma$ from π peak & $> 2.5\sigma$ from K peak
- n Require pairs from mixed events to have reconstructed vertices within 1cm
 - Acceptance varies as a function of vertex position
- n Remove pairs within 2cm in drift chamber
 - Ghosting
- n Remove tracks with EMC clusters within 12cm of each other in both real and mixed sample
 - Shower + tower size in EMC
- n Correct for two track inefficiencies at low relative ϕ in the drift chamber.
- n Full Coulomb Correction modified for momentum smearing
 - Partial correction changes radii results marginally
- n Residual correlations in event mixed background $\pm < 2\%$ error
- n Momentum smearing correction to correlation function

Systematic Errors:

R_{long} & R_{side} : 8%
 R_{out} : 4.5%



Results:

Assume $T = 125 \text{ MeV}$, $\beta_f = .69/\eta_f = .85$
From fits to singles spectra in centrality region 5-15%
(J. Burward-Hoy)

n Theoretical hydro-inspired fits.

$$\square R_{side}^2(m_T) = \frac{R_{geom}^2}{1 + \beta_f^2 \left(\frac{m_T}{T} \right)} \Rightarrow R_{geom} = 6.7 \pm 0.2 \text{ fm}$$

$$m_t^2 = k_t^2 + m_\pi^2$$

□ Chapman, Nix, Heinz PRC 52, 2694

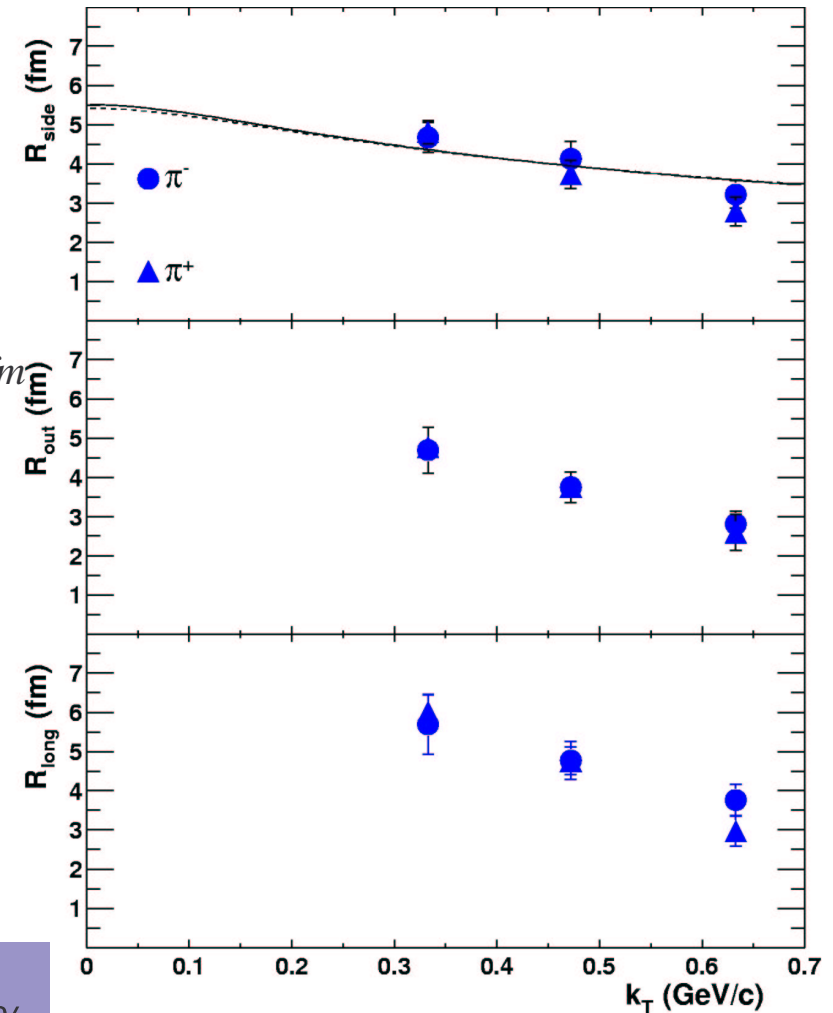
$$\square R_{side}^2(m_T) = \frac{R_{geom}^2}{1 + \eta_f^2 \left(\frac{1}{2} + \frac{m_T}{T} \right)} \Rightarrow R_{geom} = 8.1 \pm 0.3 \text{ fm}$$

□ Wiedemann, Scotto, Heinz PRC 53, 918

n However, hydro calculations predict R_o than data, R_s , smaller.

n Much larger than comparable 1D RMS Au radius of 3.07 fm

n k_t dependence suggest larger β_f/T (η_f/T) than fits to singles

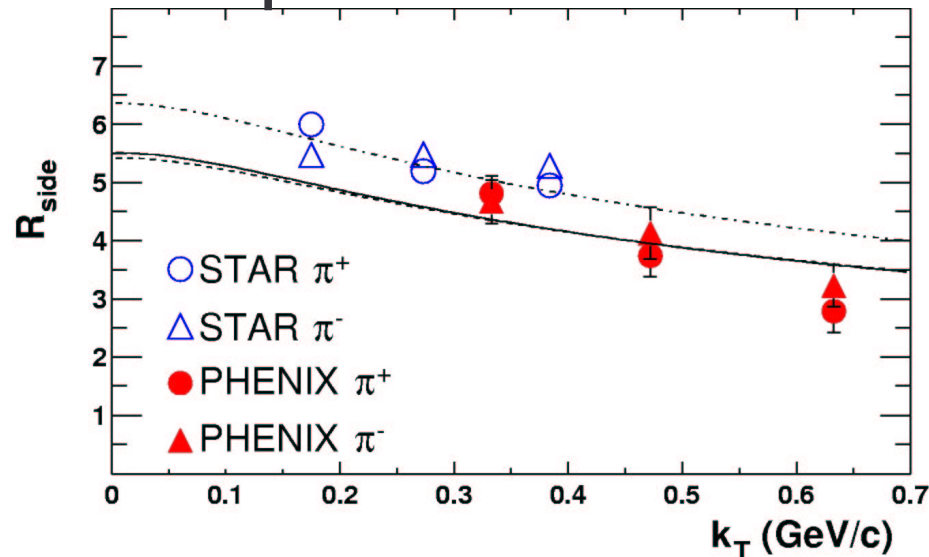


Systematic Errors:

R_{long} & R_{side} : 8%
 R_{out} : 4.5%



Compared to STAR



- n Well described by hydro model if consider datasets separately.
 - But they indicate a much higher flow/temp ratio when taken together
 - Need to be careful about systematics between different measures.
 - Both experiments should be able to sort this out in the next dataset.
- n Inconsistent with models of QGP that include an hadronic rescattering phase.



Energy dependence ...

- n R_{out} and R_{side} are energy independent within error bars.
 - n Smooth energy dependence in R_{long}
 - n No immediate indication of very different physics
 - n Fit R_{long} to:
- $$\frac{A}{\sqrt{m_T}}$$
- n AGS: $A = 2.19 \pm 0.05$
 - n SPS: $A = 2.90 \pm 0.10$
 - n RHIC: $A = 3.32 \pm 0.03$

$A = \tau_0 T$ in 1st order T/m_T calculation

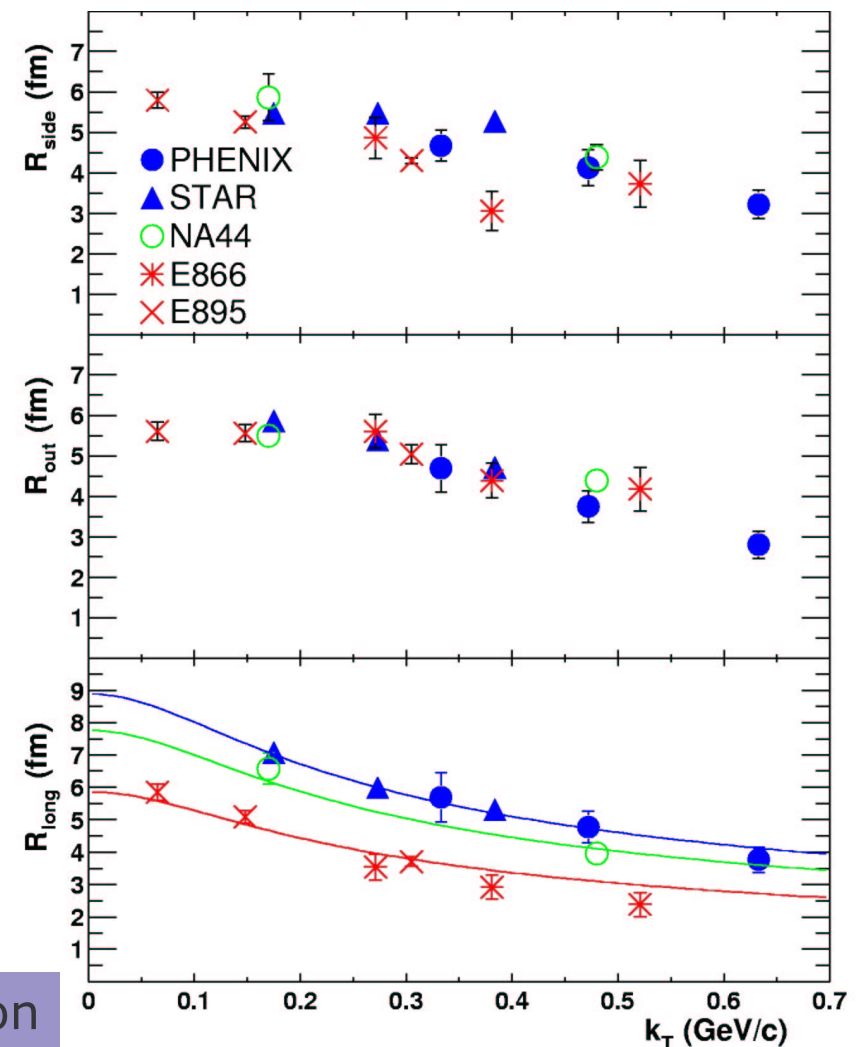
τ_0 = average freeze-out proper time

M. Lisa *et al.*, PRL **84**, 2798 (2000)

R. Soltz *et al.*, to be sub PRC

C. Ader *et al.*, PRL **87**, 082301

I.G. Bearden *et al.*, EJP **C18**, 317 (2000)

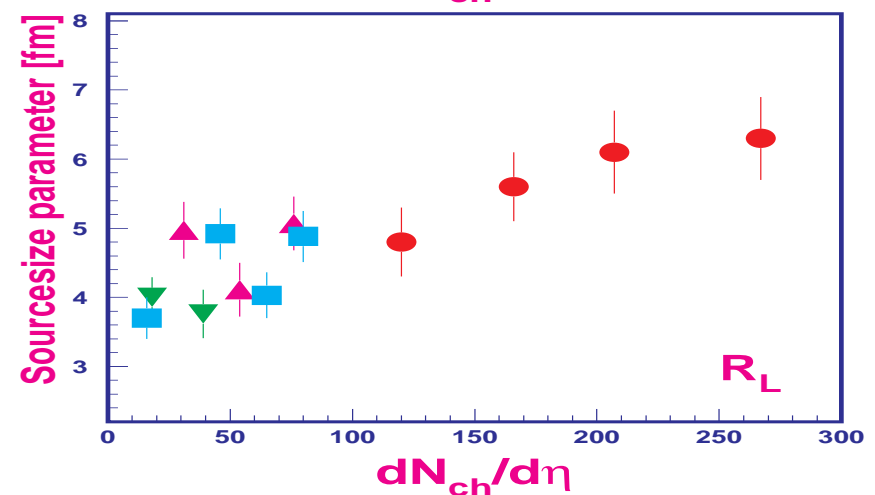
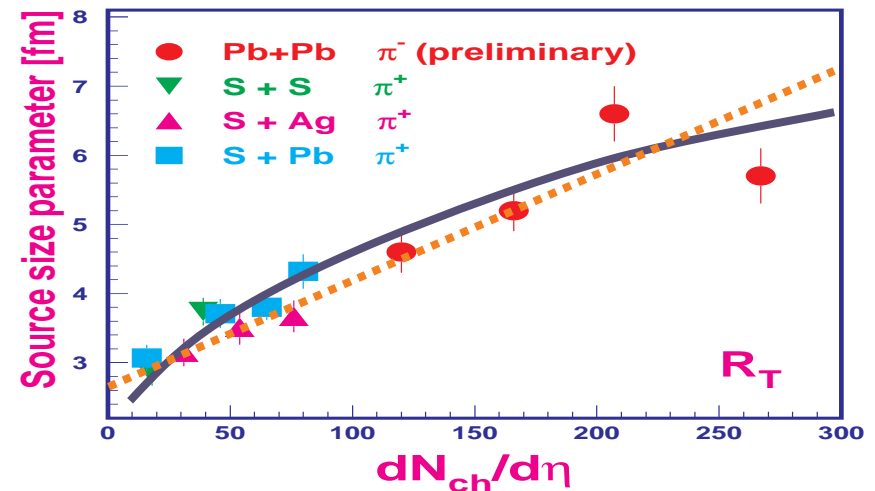




Even an experimentalist can't predict this

- n “A prognostication”
 - Local PHENIX HBT expert keeps his predictions on the web.
- n Zajc at Nucleus-Nucleus 97 showed this slide
- n Extrapolation to RHIC multiplicities?
 - $(dN/dy)^{1/3} \Rightarrow R \sim 9 \text{ fm}$
 - $(dN/dy) \Rightarrow R \sim 20 \text{ fm}$
- n Neither.
- n What is happening here?
- n **

Multiplicity dependence of pion source size



** To be fair: Bill's radii predictions are way off. His guess at the charged particle multiplicity in the extrapolations...bang on

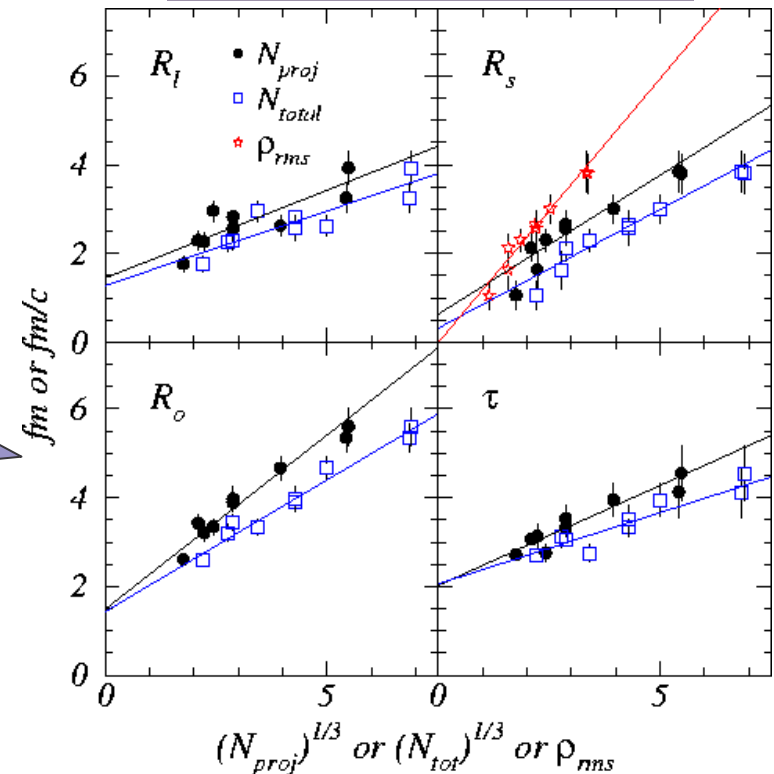
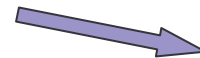


Q & A

Q: Why are R_{out} and R_{side} ~identical over an order of magnitude of beam energy?

There is ample evidence at AGS and SPS of dependence of radii on # target+projectile nucleons

- Even though the larger nuclei have larger flow the radii follow a simple scaling
- Why no multiplicity/energy dependence?



Q: Flow higher at RHIC leads to smaller radii?

- HBT results suggest high flow but spectra imply flow comparable to SPS **
- k_t dependence of radii is similar for different energies (competing mechanisms that create similar β_T/T ?)

Q: Higher opacity at RHIC energies?

Why would opacity be higher at RHIC than at AGS?

- $\pi n \rightarrow \Delta \rightarrow \pi n$ & $\pi \pi \rightarrow \rho \rightarrow \pi \pi$
- And why would size and opacity effects identically cancel?
- Red herring? – You can't create opaque source with smaller lifetime.

Q: Why is it only changing in the longitudinal direction?

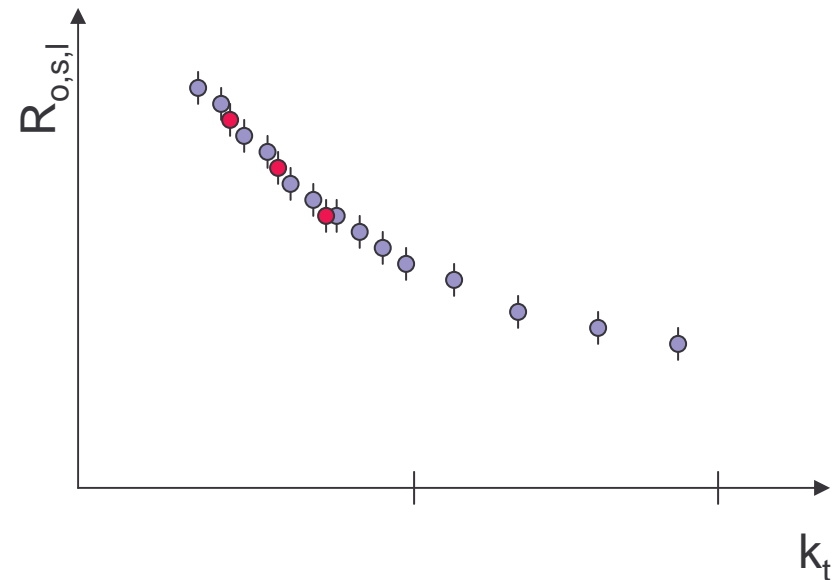
- Some surprise among theorists that the difference is so small.
- Is there a quantitative expectation of the R_{long} dependence?

** To my knowledge no one has fit both SPS and RHIC results to same hydro model



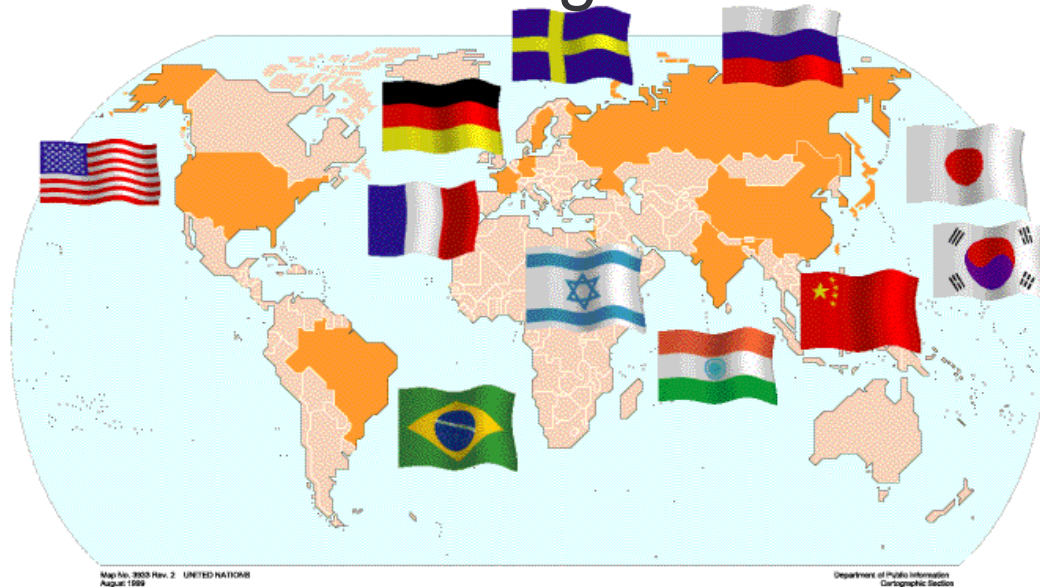
A taste of the future

- n For year-1 we wrote ~5 million min. bias events.
 - 1/2 million events in this analysis after all cuts
- n In the past run we wrote ~90 million min. bias events (+14 million rare event triggers)
- n Therefore, for the year-2 pion correlations:
 - easily 10 high statistics bins in k_t from .2 to 1.0 GeV + a few bins from 1.0 to 2.0 GeV
 - n Capability to exclude detailed theories
 - n Systematic errors start to become the real problem
- n Beyond pions:
 - 1D proton and kaon correlations
 - Non-identical correlations (πK , πp , etc.)
 - π^0
 - n maybe possible in year 2 ... year 3?
 - n Probe of very high k_t .
 - anti-neutron correlations
 - n Better than proton correlations \pm lack of coulomb.





The scientists that make me look good*



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*or at least better than I normally do

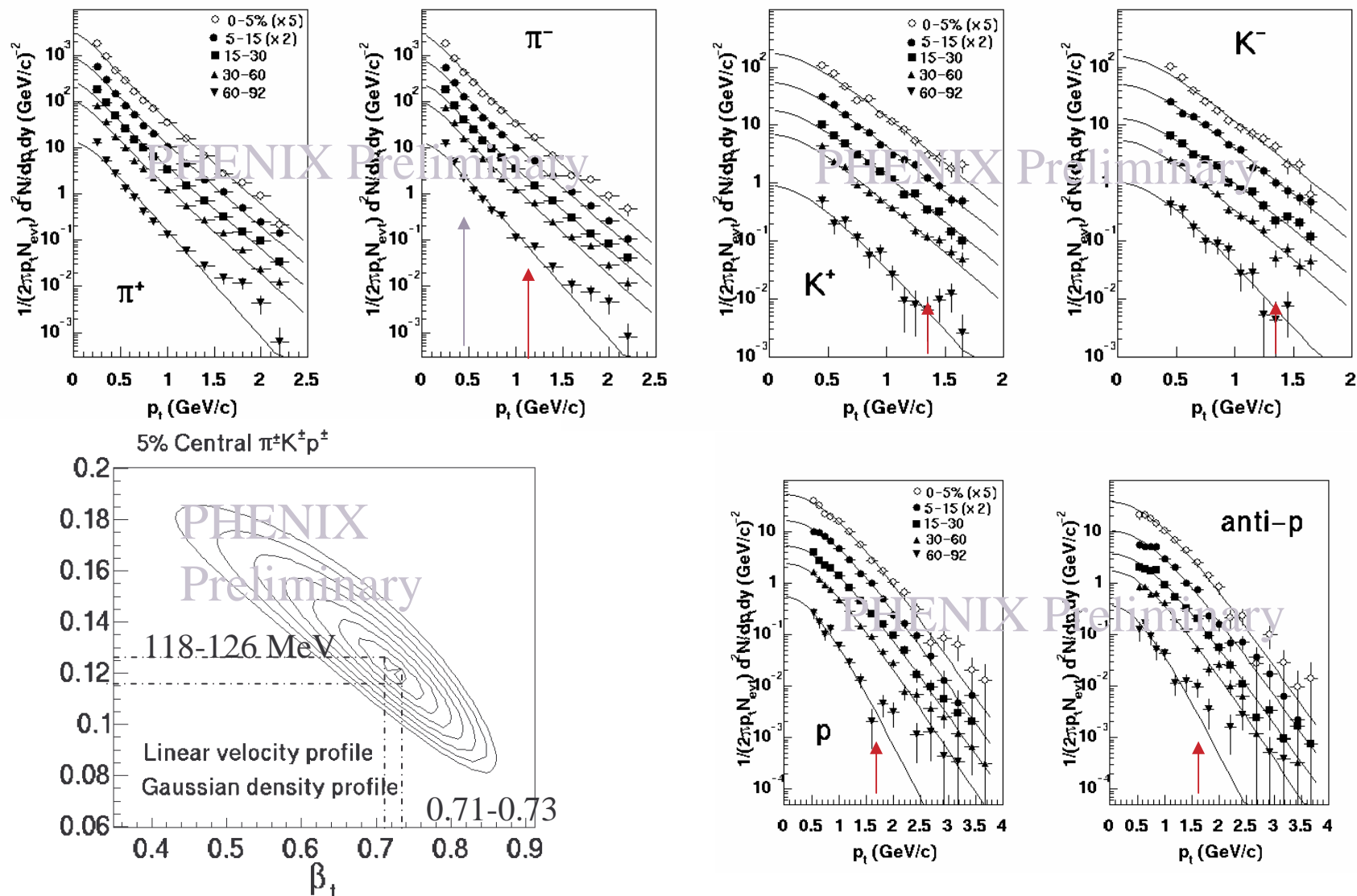


Supporting slides:

Fitting the Single Particle Spectra:

Jane Burward-Hoy (LLNL)

$$1/m_t \, dN/dm_t = A \int f(\xi) \, \xi \, d\xi \, m_T \, K_1(m_T/T_{fo} \cosh \rho) \, I_0(p_T/T_{fo} \sinh \rho)$$





Switching frames

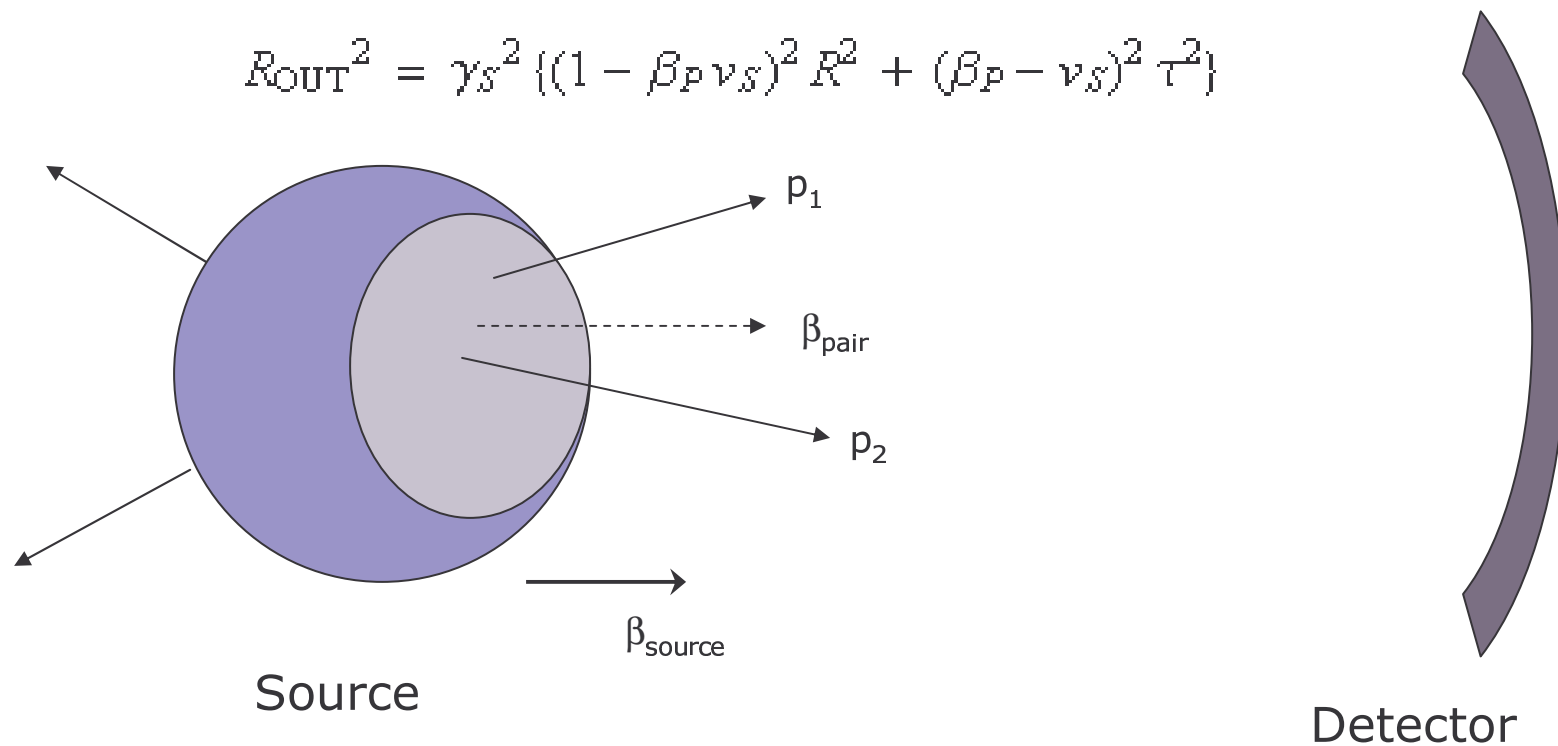
- n Problem: radii extracted depend on the frame in which you measure
- n Historically we have used:
 - Q_{inv} (1D pair rest frame)
 - Collision Center of Mass
 - n AGS
 - LCMS (longitudinally co-moving frame)
 - n Longitudinally boost invariant sources
 - n SPS(?), RHIC
 - Jet frame
 - n p-p collisions.
 - PCMS (pair center-of-mass – 3D version of Q_{inv})
 - ...
- n What is the 'correct' frame?? (= the frame in which the source is not moving)
- n Can we bound it experimentally?



Our source??

2π HBT measurement is simultaneous in the PCMS frame

$$R_{\text{OUT}}^2 = \gamma_s^2 \{ (1 - \beta_P v_s)^2 R^2 + (\beta_P - v_s)^2 \tau^2 \}$$



Pair rest frame \neq source rest frame

What is the β_{source} ?

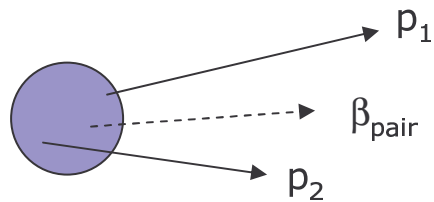
We don't know, but we can bound it:

> 0

$< \beta_{\text{pair}}$

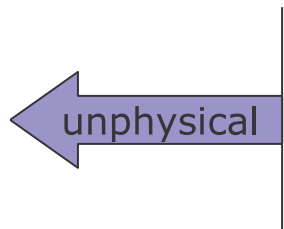


LCMS vs. PCMS

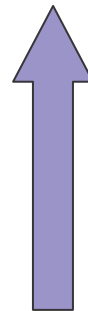


n LCMS

- A Lorentz contracted measurement of the frame

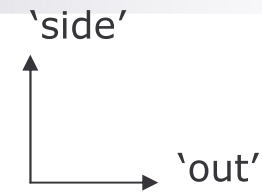
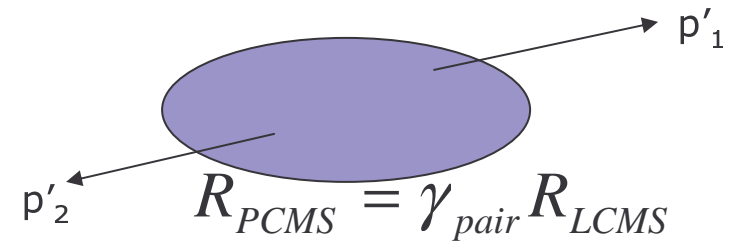
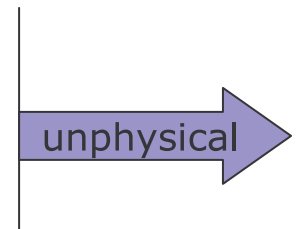


True frame



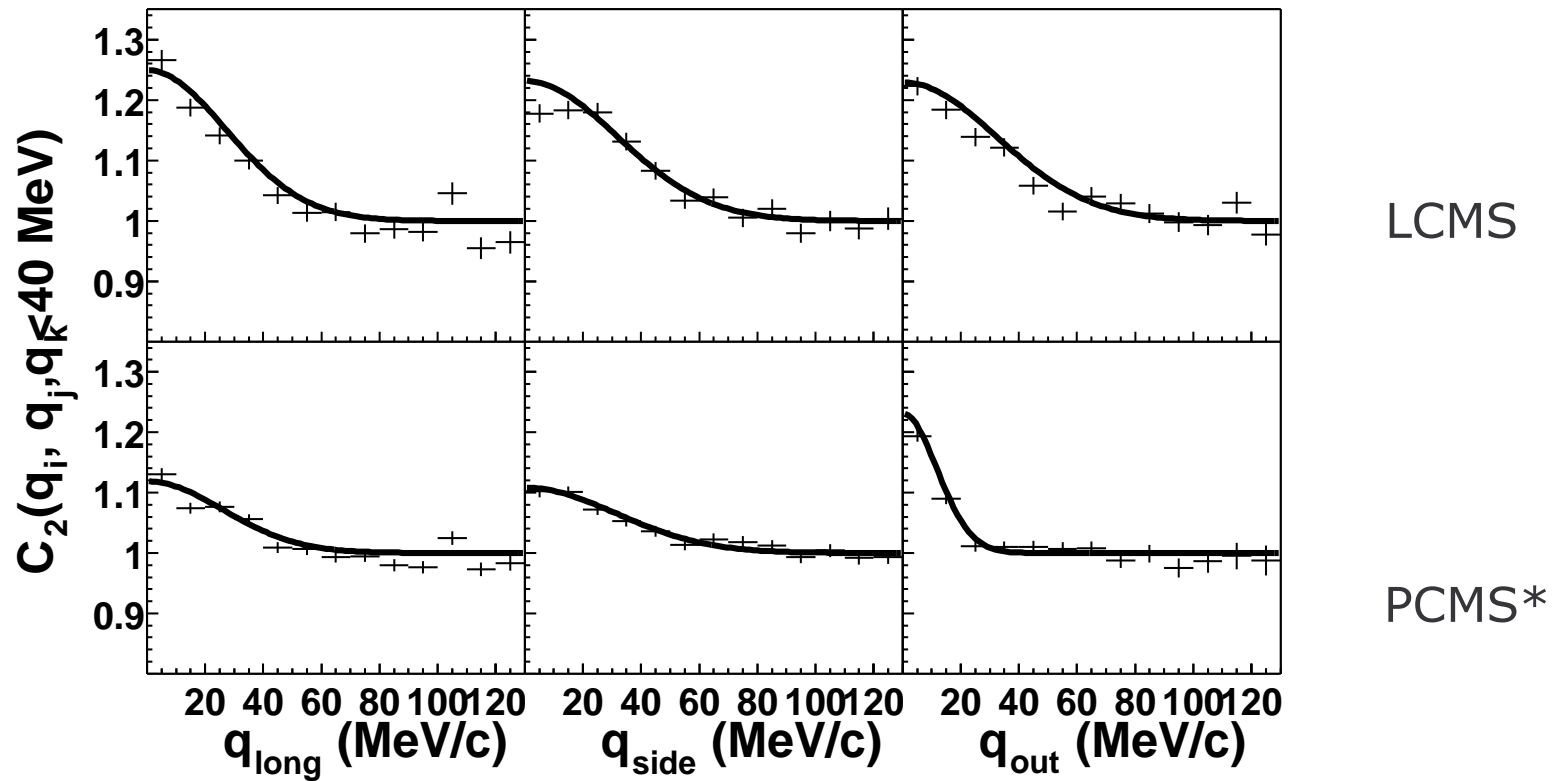
n PCMS

- A measurement of the length of the train in the train frame





Fits to entire dataset (π^-)



*PCMS = Pair Center of Mass System



PCMS results

- n R_{side} and R_{long} unchanged
 - And not plotted
- n $R_{\text{out}}^{\text{PCMS}}$ differs by $\langle \gamma_{\text{pair}} \rangle$ from $R_{\text{out}}^{\text{LCMS}}$
- n Be careful about deducing a lifetime from $R_{\text{out}}^2 = R_{\text{side}}^2 + (\beta_{\text{pair}} \tau)^2$
 - $\beta_{\text{pair}} = 0$

