

Nuclear physics at small x and (s)PHENIX

D. Kharzeev

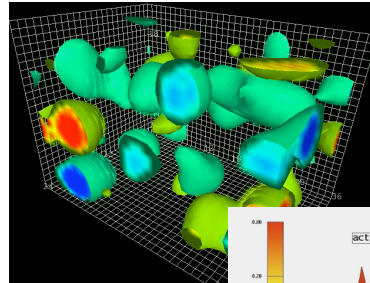


From particles to fields in QCD

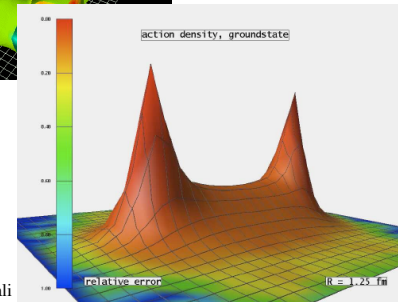
Particles

| | | | |
|---------------------------|------------------------|-------------------------------|---------------------------------|
| Quarks | u up | c charm | t top |
| | d down | s strange | b bottom |
| Leptons | ν_e e- Neutrino | ν_μ μ - Neutrino | ν_τ τ - Neutrino |
| | e electron | μ muon | τ tau |
| I II III | | | |
| The Generations of Matter | | | |

Fields (Geometry)



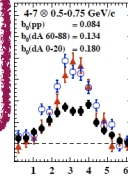
D. Leinweber



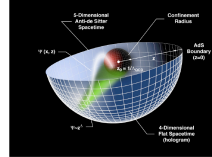
G. Bali

From particles to fields: collective phenomena as the essence of QCD

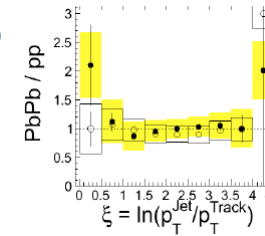
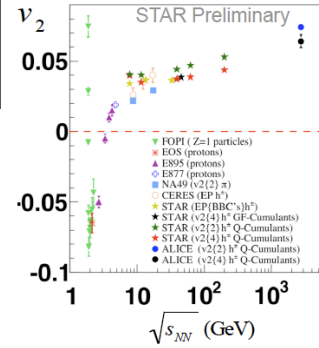
- From partons to strong color fields:
nuclear wave functions at small x



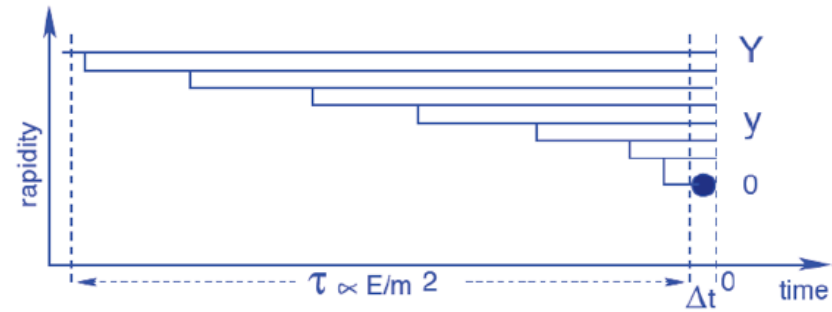
- Hydrodynamics and collective behavior:
transport properties; phase diagram &
fluctuations; anomalies and
chiral magnetic effect



- Jets: the flow of energy and momentum in QCD
- The probes: heavy quarks, dileptons, ..



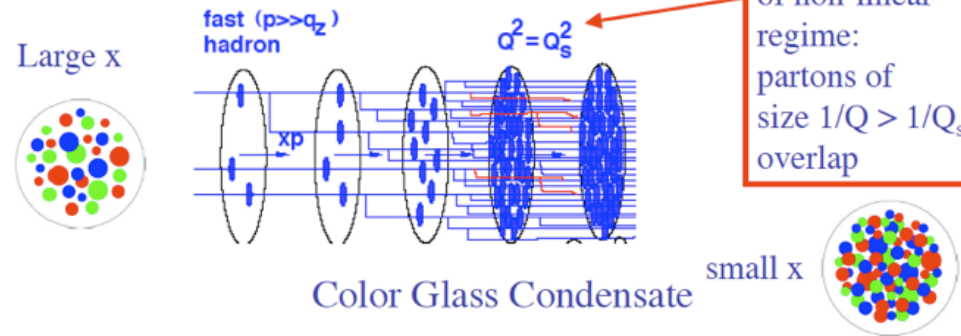
The space-time picture of high-energy interactions in QCD



1. Fast (large y) partons live for a long time;
2. Parton splitting probability is $\sim \alpha_s y$ - not small!

From partons to fields at small x

Bjorken x : the fraction of hadron's momentum carried by a parton; high energies s open access to small $x = Q^2/s$



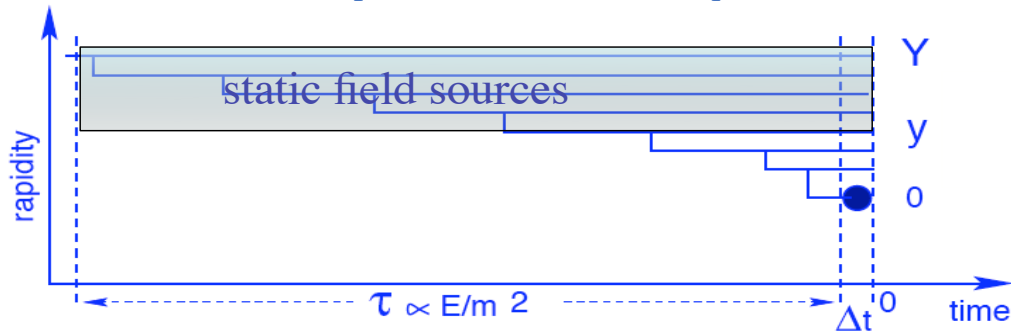
Because the probability to emit an extra gluon is $\sim \alpha_s \ln(1/x) \sim 1$, the number of gluons at small x grows; the transverse area is limited

→ transverse density becomes large

L.Gribov, E.Levin, M.Ryskin;
A.Mueller, J.Qiu

The origin of classical background field

Coherent field with occupation number $\sim \frac{1}{\alpha_s(Q_s)}$
suppression of hard processes at small x ;
depletion of back-to-back (quantum) correlations



Gluons with large rapidity and large occupation number
act as a background field for the production of slower gluons

QCD at small x: the parameters

$$\mathcal{L} = -\frac{1}{4}G_{\mu\nu}^a G_{\mu\nu}^a + \sum_f \bar{q}_f^a (i\gamma_\mu D_\mu - m_f) q_f^a;$$

Parameters of QCD lagrangian:
quark masses m_f and Λ_Q

New parameter at small x:
saturation momentum Q_s

is related to \hat{q} by $Q_s^2 = q L$ ⁷

Renormalization group and the effective action

RG constraints the form of the effective action:

$$\mathcal{L}_{\text{eff}} = -\frac{1}{4\bar{g}^2(t)} G^2, \quad t \equiv \ln \left(\frac{G^2}{\Lambda^4} \right)$$

the coupling is defined through $t = \int_g^{\bar{g}(t)} \frac{dg}{\beta(g)}$:

At large t (strong color field),

$$\frac{1}{\bar{g}^2(t)} \sim t + \dots \quad \text{and} \quad \mathcal{L}_{\text{eff}} \sim G^2 \ln \left(\frac{G^2}{\Lambda^4} \right) \quad 10$$

$$\rightarrow \frac{1}{N_{\text{part}}} \frac{dN}{d\eta} \sim \frac{1}{\alpha_s(Q_s^2)} \quad \text{KLN} \quad 8$$

QCD evolution at small x: nuclear modification of hard processes

New parameter at small x:
saturation momentum

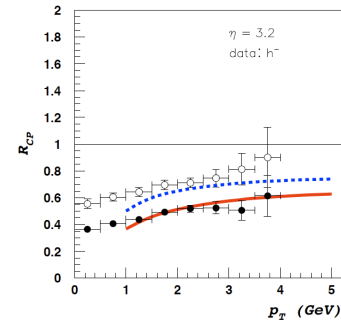
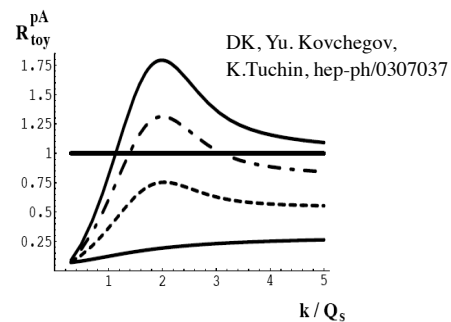
Therefore, instead of $\ln(Q/\Lambda_Q)$,

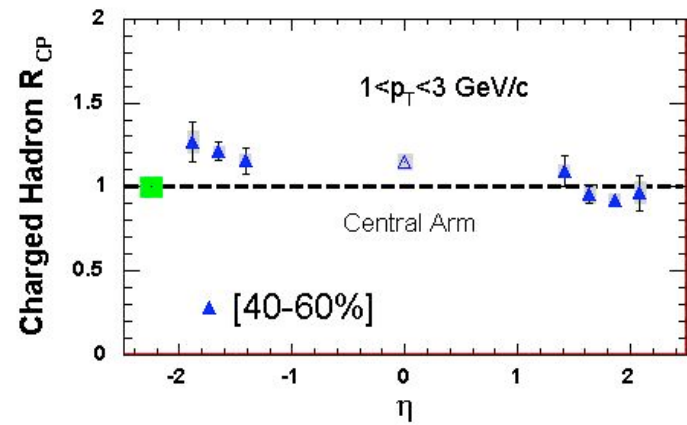
we get $\ln(Q/Q_s)$

Resummation of these logs yields
“anomalous dimensions” -
and results in nuclear suppression of hard
processes at small x

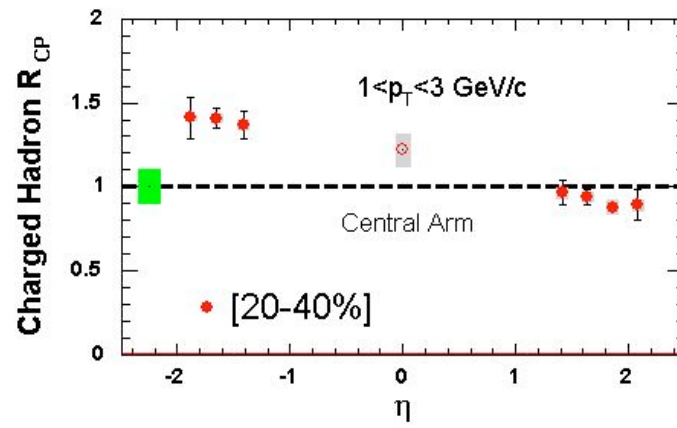
Nuclear suppression of hard processes at small x: an intuitive explanation

Saturation slows down the growth of gluon densities at small x; this effect is stronger in the nuclear wave function than in the proton wave function - thus R_{pA} decreases

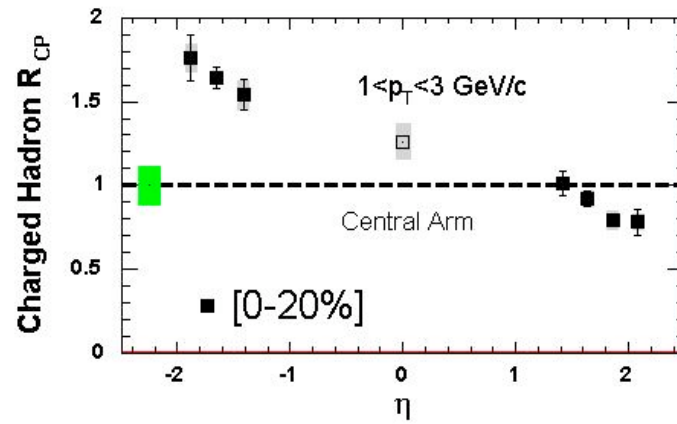




PHENIX data



PHENIX data



PHENIX data

Recent theoretical studies:
A.Dumitru, J. Jalilian-Marian,
C.Marquet, K.Tuchin, ¹³.

Baryons vs mesons at small x

At forward rapidities at RHIC, the nuclear modification of baryon production is linked to the mechanism of baryon number transport (diquarks? gluon junctions? Skyrmions?)

Measurement: production cross sections of PID hadrons: baryons (protons, hyperons), and mesons; compare mid- and forward-(+backward?) rapidities

Quarks vs gluons at small x

**Quarks and gluons interact differently
because of different color charge**

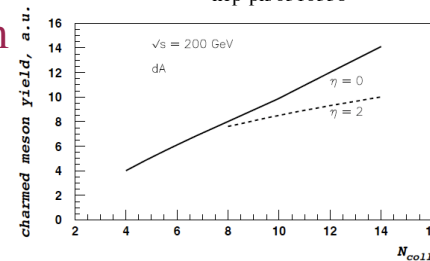
Measurement: production cross sections of photons
and dileptons at forward rapidities

Flavor dependence of small x suppression

At forward rapidities at RHIC, the saturation momentum exceeds the charm quark mass
(always the case at LHC!)
so charm becomes “light” at small x -
a sensitive probe of saturation

Measurement: identified charm production cross section, compare mid- and forward- (+backward?) rapidities

DK, K.Tuchin,
hep-ph/0310358



Charmonium suppression at small x

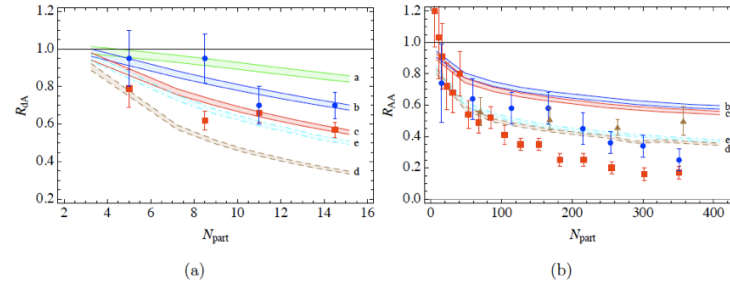
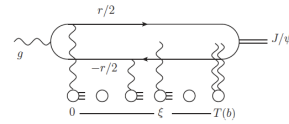
In the entire kinematical range of RHIC, the saturation momentum exceeds the inverse radius of charmonium, $(\alpha_s M_c)^{-1}$, so charmonium production is a sensitive probe of saturation

Measurement: different charmonia production (χ , ψ' , ...bottomonium?) compare mid- and forward- (+backward?) rapidities

Gluon saturation effects on the color singlet J/ψ production in high energy dA and AA collisions

F. Dominguez,¹ D.E. Kharzeev,^{2,3} E.M. Levin,^{4,5} A.H. Mueller,¹ and K. Tuchin⁶

arXiv:11091250 (Sep 6, 2011)

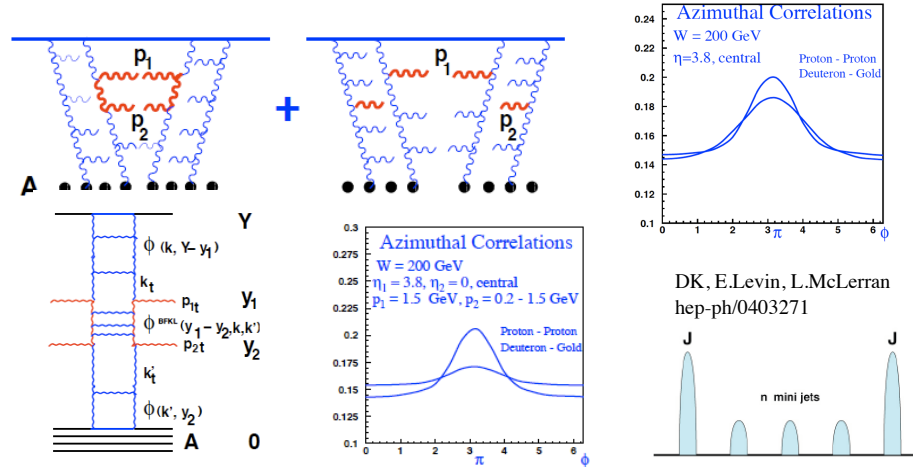


no
factorization!

AA is **NOT**
reproduced,
neither at
RHIC nor
at LHC

FIG. 3: Nuclear modification factor vs N_{part} in (a) dAu and (b) AA collisions using the DHJ model [13]. Band 'a' (green) represents rapidity $y = -1.7$ at $\sqrt{s} = 200$ GeV, 'b' (blue): $y = 0$, $\sqrt{s} = 200$ GeV, 'c' (red): $y = 1.7$, $\sqrt{s} = 200$ GeV, 'd' (brown): $y = 3.25$, $\sqrt{s} = 2.76$ TeV, 'e' (cyan): $y = 0$, $\sqrt{s} = 5.5$ TeV. $m = 1.5$ GeV, $C = 1$. Experimental data [16–19] is represented by (blue) circles in 'b', by (red) squares in 'c' and by (brown) triangles in 'd'. (Color online).

Disappearance of di-jet correlations at small x and at large rapidity separations

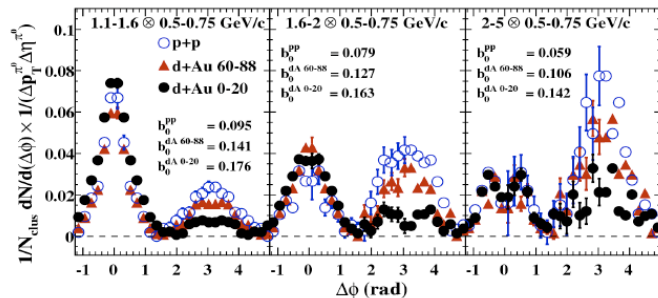


Measurement: di-jet(leading hadron) correlations with large rapidity separation (forward-backward)

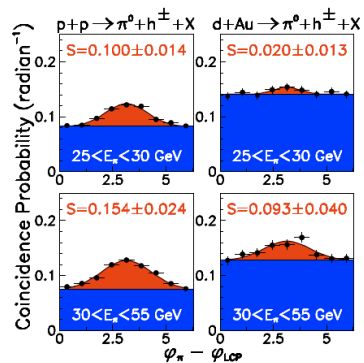
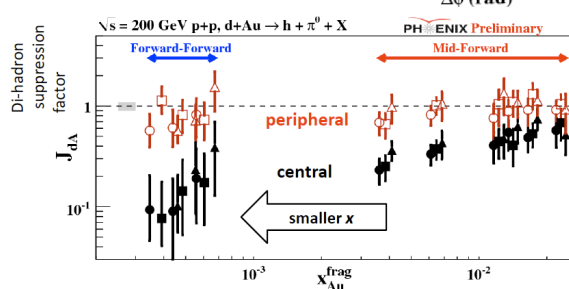
Di-hadron correlations at small x

$\sqrt{s_{NN}} = 200 \text{ GeV}$, d+Au, p+p \rightarrow Cluster + π^0 ; $3.0 < \eta_{\text{clus}}, \eta_{\pi^0} < 3.8$

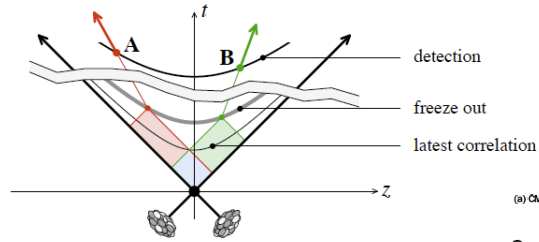
PHENIX Coll,
QM 2011
arXiv:1105.5112



STAR Coll,
arXiv:0602011 (PRL)



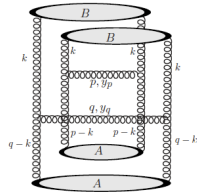
Probes of coherent gluon fields at small x: long-range correlations in rapidity



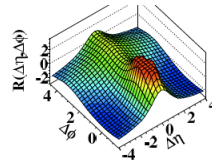
causality:

$$\tau_{\text{init.}} = \tau_{\text{f.o.}} \exp\left(-\frac{1}{2}\Delta y\right)$$

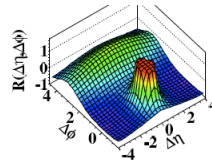
A.Dumitru, K. Dusling, F. Gelis,
J. Jalilian-Marian, T. Lappi, R. Venugopalan,
arXiv: 1009.5295



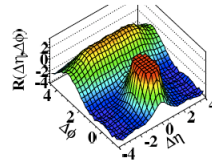
(a) CMS MinBias, $p_T > 0.1 \text{ GeV}/c$



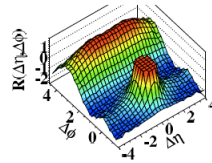
(b) CMS MinBias, $1.0 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$



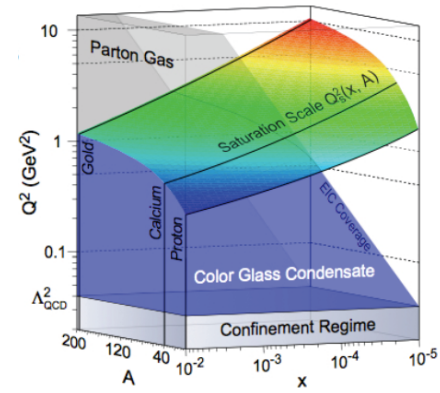
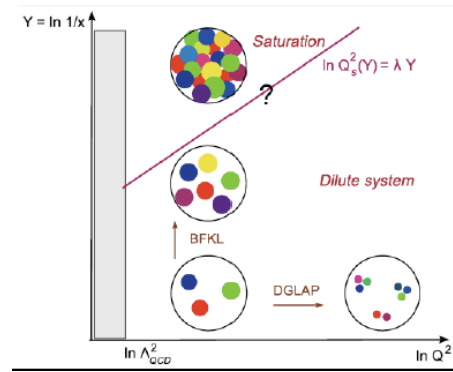
(c) CMS N = 110, $p_T > 0.1 \text{ GeV}/c$



(d) CMS N = 110, $1.0 \text{ GeV}/c < p_T < 3.0 \text{ GeV}/c$



The path to eRHIC



Some speculative ideas:

- Polarized protons on nuclei?
(spin asymmetries are due to quark interactions with gluon fields; study the dependence on the gluon field strength)
- Double-(single-)diffractive production of hadrons (light and charmonia) at mid-(forward-)rapidity in polarized pp collisions (was never done before?) - spin contents at small x; glueballs; sphalerons; mechanism of charmonium production; ...²³