

Factorized Drell-Yan cross section

□ TMD factorization ($q_{\perp} \ll Q$):

$$\frac{d\sigma_{AB}}{d^4q} = \sigma_0 \int d^2k_{a\perp} d^2k_{b\perp} d^2k_{s\perp} \delta^2(q_{\perp} - k_{a\perp} - k_{b\perp} - k_{s\perp}) \mathcal{F}_{a/A}(x_A, k_{a\perp}) \mathcal{F}_{b/B}(x_B, k_{b\perp}) \mathcal{S}(k_{s\perp}) + \mathcal{O}(q_{\perp}/Q)$$

$$x_A = \frac{Q}{\sqrt{s}} e^y \quad x_B = \frac{Q}{\sqrt{s}} e^{-y}$$

The soft factor, \mathcal{S} , is universal, could be absorbed into the definition of TMD parton distribution

□ Collinear factorization ($q_{\perp} \sim Q, q_{\perp} \gg Q$):

$$\frac{d\sigma_{AB}}{d^4q} = \int dx_a f_{a/A}(x_a, \mu) \int dx_b f_{b/B}(x_b, \mu) \frac{d\hat{\sigma}_{ab}}{d^4q}(x_a, x_b, \alpha_s(\mu), \mu)$$

□ Spin dependence:

The factorization arguments are independent of the spin states of the colliding hadrons

➡ same formula with different distributions for $\gamma^*, W/Z, H^0 \dots$

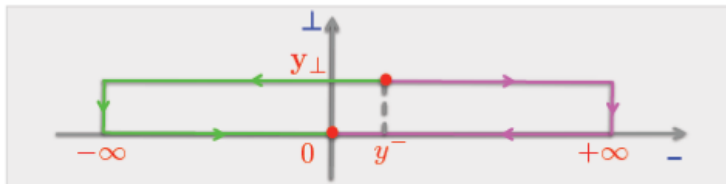
The sign change of Sivers function

□ TMD quark distribution:

See talks by Collins, Kang, ...

$$f_{q/h\uparrow}(x, \mathbf{k}_{\perp}, \vec{S}) = \int \frac{dy^- d^2y_{\perp}}{(2\pi)^3} e^{iyp^+ y^- - i\mathbf{k}_{\perp} \cdot \mathbf{y}_{\perp}} \langle p, \vec{S} | \bar{\psi}(0^-, \mathbf{0}_{\perp}) \text{Gauge link} \frac{\gamma^+}{2} \psi(y^-, \mathbf{y}_{\perp}) | p, \vec{S} \rangle$$

- SIDIS: $\Phi_n^{\dagger}(\{+\infty, 0\}, \mathbf{0}_{\perp}) \Phi_n^{\dagger}(+\infty, \{\mathbf{y}_{\perp}, \mathbf{0}_{\perp}\}) \Phi_n(\{+\infty, y^-\}, \mathbf{y}_{\perp})$
- DY: $\Phi_n^{\dagger}(\{-\infty, 0\}, \mathbf{0}_{\perp}) \Phi_n^{\dagger}(-\infty, \{\mathbf{y}_{\perp}, \mathbf{0}_{\perp}\}) \Phi_n(\{-\infty, y^-\}, \mathbf{y}_{\perp})$



- For a fixed spin state:

$$f_{q/h\uparrow}^{\text{SIDIS}}(x, \mathbf{k}_{\perp}, \vec{S}) \neq f_{q/h\uparrow}^{\text{DY}}(x, \mathbf{k}_{\perp}, \vec{S})$$

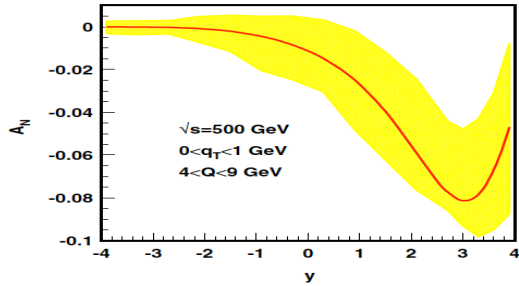
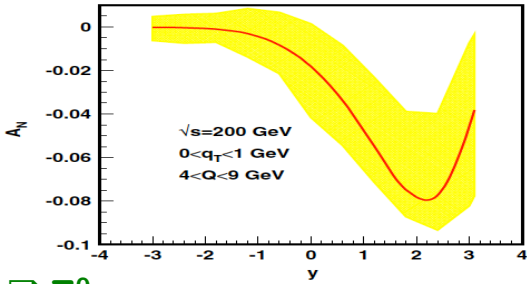
□ Parity + Time-reversal invariance:

$$\longrightarrow f_{q/h\uparrow}^{\text{Sivers}}(x, k_{\perp})^{\text{SIDIS}} = -f_{q/h\uparrow}^{\text{Sivers}}(x, k_{\perp})^{\text{DY}}$$

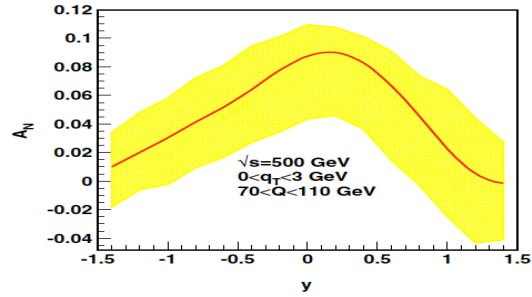
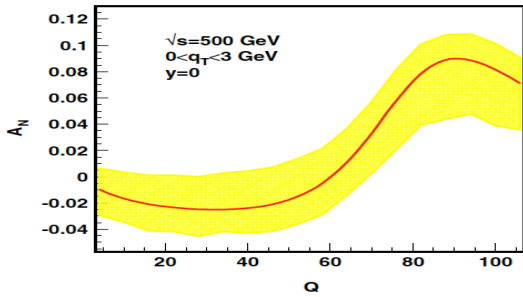
It is a critical test of TMD factorization approach

Test of the modified universality

□ **Drell-Yan:** $A_N^{\sin(\phi-\phi_s)} = -A_N$ See talks by Anselmino, Rogers, ...



□ **Z⁰:**



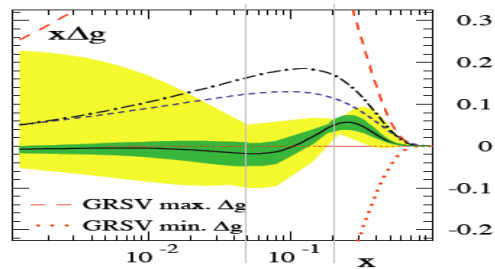
Theoretical challenge: Q-dependence of Sivers function?

The sign “mismatch”

□ **Asymmetry could have a node:** See talks by Kang

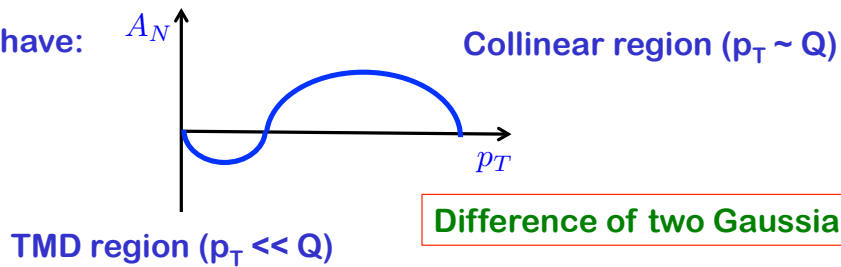
Sign change of $\Delta g(x)$:

$$A(s) \propto \sigma(s) - \sigma(-s)$$



□ **A_N of Drell-Yan p_T distribution:**

We could have:



Important measurement for understanding A_N of hadronic pion

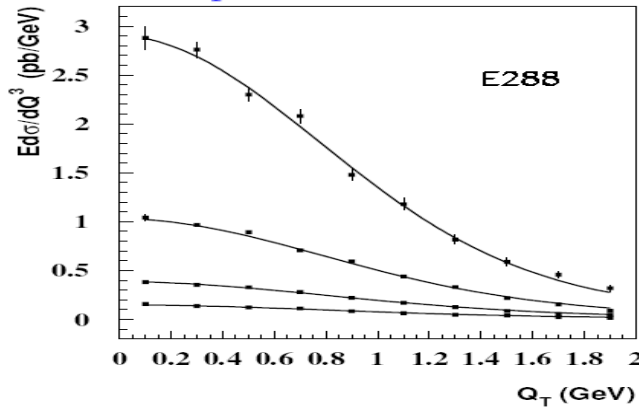
Unpolarized Drell-Yan cross section

□ The denominator of the Asymmetry:

$$\frac{d\sigma}{d^4q} \quad \frac{d\sigma}{d^4q d\Omega}$$

□ Angular integrated Drell-Yan is under control:

- Fermilab E288 data at $p_{\text{beam}} = 400 \text{ GeV}$



But, Drell-Yan lepton angular distribution needs work!

Lam-Tung relation

□ Normalized Drell-Yan lepton angular distribution:

$$\frac{dN}{d\Omega} \equiv \left(\frac{d\sigma}{d^4q}\right)^{-1} \frac{d\sigma}{d^4q d\Omega} = \frac{3}{4\pi} \left(\frac{1}{\lambda + 3}\right) \left[1 + \lambda \cos^2\theta + \mu \sin(2\theta) \cos\phi + \frac{\nu}{2} \sin^2\theta \cos(2\phi) \right]$$

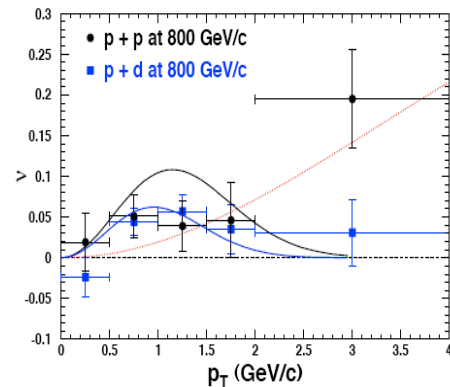
□ Lam-Tung relation:

$$1 - \lambda - 2\nu = 0$$

□ Collinear factorization:

$$\lambda = \frac{W_T - W_L}{W_T + W_L} \approx \frac{W_T^{\text{Resum}} - W_L^{\text{Resum}}}{W_T^{\text{Resum}} + W_L^{\text{Resum}}} = \frac{1 - \frac{1}{2} Q_{\perp}^2 / Q^2}{1 + \frac{3}{2} Q_{\perp}^2 / Q^2}$$

$$\nu = \frac{2W_{\Delta\Delta}}{W_T + W_L} \approx \frac{2W_{\Delta\Delta}^{\text{Resum}}}{W_T^{\text{Resum}} + W_L^{\text{Resum}}} = \frac{Q_{\perp}^2 / Q^2}{1 + \frac{3}{2} Q_{\perp}^2 / Q^2}$$



Still valid after resummation of large logs

Berger, Qiu, Rodriguez, 2007

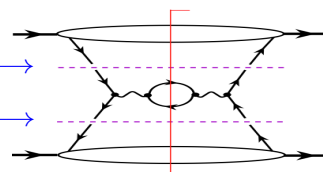
□ TMD factorization:

– Boer – Mulder function – $\cos(2\Phi)$:

$$h_1^{\perp \text{DY}}(x) = -h_1^{\perp \text{SIDIS}}(x)$$

$$\sigma^{+\alpha} k_{\perp\alpha}$$

$$\sigma^{+\beta} k'_{\perp\beta}$$



Theory challenge: Q-dependence of BM function, $\cos(\Phi)$, ...

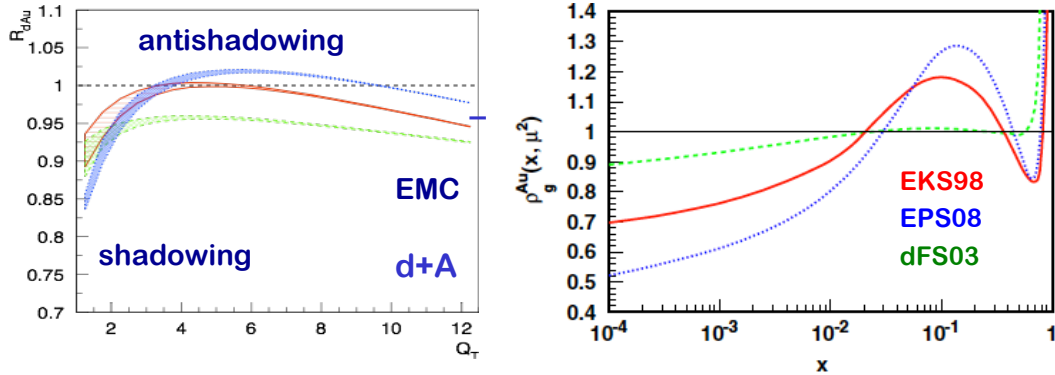
Excellent probe of gluon distribution

□ Nuclear modification factor:

Kang, Qiu, Vogelsang, PRD 2009

$$R_{dAu} \equiv \frac{1}{\langle N_{coll} \rangle} \frac{d^2 N^{dAu}/dQ_T dy}{d^2 N^{pp}/dQ_T dy} \stackrel{\text{min.bias}}{=} \frac{\frac{1}{2A} d^2 \sigma^{dAu}/dQ_T dy}{d^2 \sigma^{pp}/dQ_T dy}$$

□ RHIC kinematics – if dominated by single scattering:

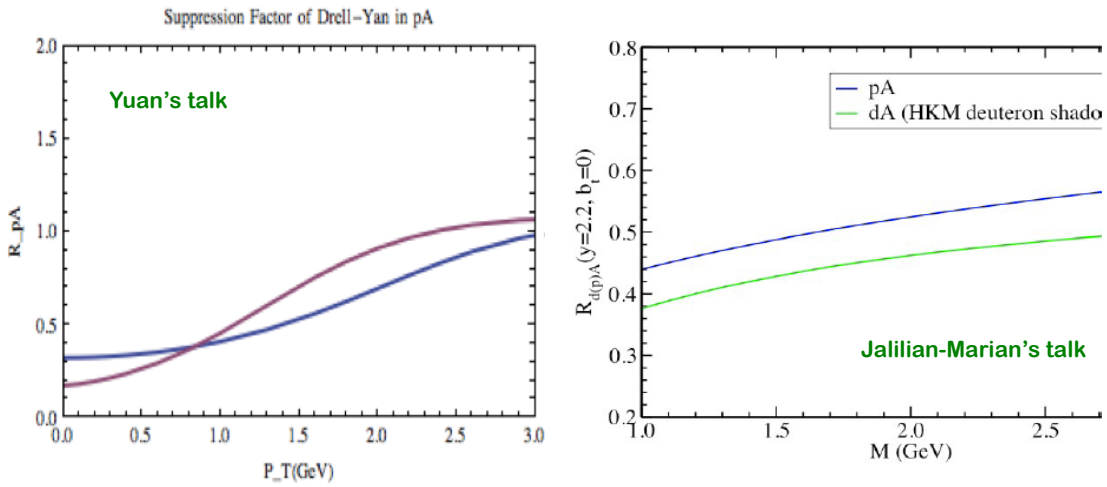
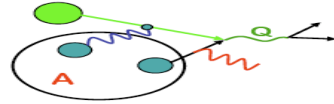


- The band is given by $\kappa=1$ (top lines) and $\kappa=0$ (bottom lines)
- Ratio follows the feature of gluon distribution if turns off isospin
- No suppression if removing isospin effect

Saturation and CGC physics

□ Forward rapidity ($y \gg 0$):

If $Q_T \sim Q_s$, collinear factorization fails

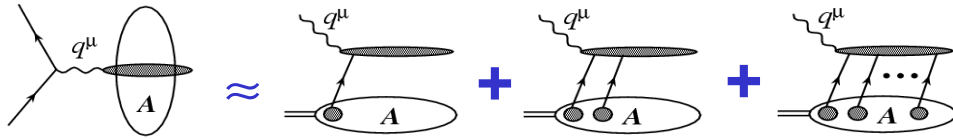


Nuclear shadowing cannot produce such suppression!

Theory challenge: Role of p_T ?

Another sign change

- Power correction to DIS – single scale:

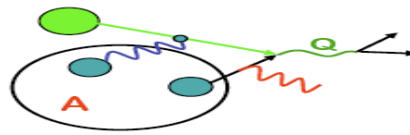


$$F_{eA}(x, Q^2) = F_{eA}^{\text{LP}}(x, Q^2) + \frac{1}{Q^2} F_{eA}^{\text{NLP}}(x, Q^2) + \dots$$

Negative - suppression

- Power correction to inclusive DY – single scale:

$$\frac{d\sigma_{pA}}{dQ^2} = \frac{d\sigma_{pA}^{\text{LP}}}{dQ^2} + \frac{1}{Q^2} \frac{d\sigma_{pA}^{\text{NLP}}}{dQ^2} + \dots$$



Positive - enhancement

Compton gives negative contribution in CO factorization

Summary and outlook

- Drell-Yan process is one of the oldest hard process proposed to test QCD – it still a very good one!
- The proof of QCD factorization for Drell-Yan is solid (LP + NLP for collinear, LP for TMD)
- The test of the sign change of the Sivers function is a critical test of TMD factorization!
- Drell-Yan could provide much more than the sign change of Sivers function

Thank you!