

WHAT IS THE PROTON SPIN PROBLEM ALL ABOUT,

STEFANO FORTE
UNIVERSITÀ DI MILANO

**WHAT IS THE PROTON SPIN
PROBLEM ALL ABOUT,
AND WHEN ARE WE GOING TO SORT IT OUT?**

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SUMMARY

- **WHERE DO WE COME FROM?: THE NUCLEON SPIN PROBLEM**
 - POLARIZED STRANGENESS IN THE NUCLEON: Δ_s
 - THE AXIAL ANOMALY AND POLARIZED GLUE: Δ_g
 - THE QCD VACUUM AND INSTANTONS
 - WHAT WE'D LIKE TO KNOW
- **WHERE DO WE STAND NOW?: UNPOL. EXPERIENCE VS POL. RESULTS**
 - THE GLUON: SCALING VIOLATIONS
 - STRANGENESS: NEUTRINO DATA
 - POLARIZED TROUBLE AND SEMI-INCLUSIVE DATA
 - POLARIZED GLOBAL FITS
- **WHERE ARE WE GOING?: HADRONIC PROCESSES AND GLOBAL FITS**
 - THE ROLE OF HADRONIC DATA IN UNPOLARIZED FITS
 - HADRONIC DATA: POLARIZED EXPECTATIONS
 - CAN WE TRUST PARTON FITS?
 - NEW IDEAS: NEURAL NETWORKS

PAST

THE NUCLEON SPIN: WHAT IS THE PROBLEM?

AXIAL CHARGES (for flavor i): $\langle N; p, s | J_{5,i}^\mu | N; p, s \rangle = a_i M_N s^\mu$

SINGLET AXIAL CHARGE \Rightarrow **QUARK SPIN FRACTION (??)**

$$a_0 \equiv \sum_{i=1}^{n_f} a_i = \int_0^1 dx \sum_{i=1}^{n_f} (\Delta q_i + \Delta \bar{q}_i) + O(\alpha_s)$$

MEASURED TO BE SMALL: $a_0 = 0.18_{-0.11}^{+0.17}$ ($Q^2 = 10$) GeV^2

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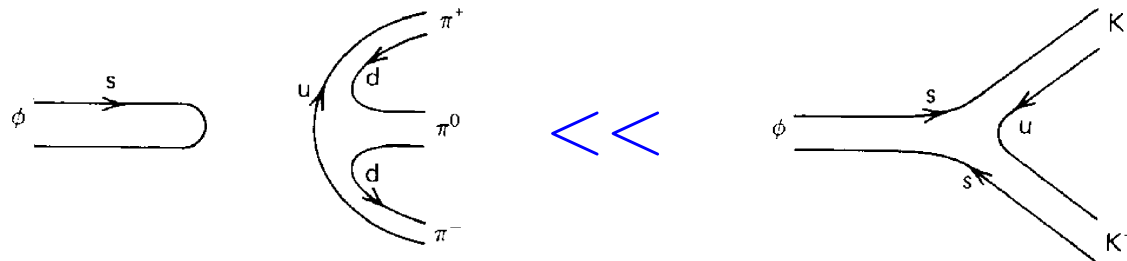
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WHY IS IT A PROBLEM?

OZI VIOLATION!

OZI RULE:



EXPECT $a_0 \equiv \Delta u + \Delta d + \Delta s \approx a_8 \equiv \Delta u + \Delta d - 2\Delta s$ (OZI \rightarrow Ellis-Jaffe S. R.)

GET $a_8 = 0.58 \pm 0.03 \gg a_0 = 0.18_{-0.11}^{+0.17}$

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SOLUTIONS

- THE OCTET AXIAL CHANNEL IS SPECIAL BECAUSE OF
SU(3) SPIN STRUCTURE \Rightarrow **POLARIZED STRANGENESS** Δ_s
(SEA POLARIZATION, SKYRMIONS, . . .)

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- THE SINGLET AXIAL CHANNEL IS SPECIAL BECAUSE OF THE
ANOMALY \Rightarrow POLARIZED GLUONS Δg
(SCHEME AND SCALE DEPENDENCE, INSTANTONS, . . .)

STRANGENESS IN THE NUCLEON

INDICATIONS OF SIZABLE NUCLEON STRANGENESS:

- THE SIGMA TERM:

$$2 \frac{\langle p | \bar{s}s | p \rangle}{\langle p | \bar{u}u + \bar{d}d | p \rangle} \sim 0.5 \quad \text{BUT LARGE UNCERTAINTY (100%?)}$$

- UNPOLARIZED PDFS MOMENTUM FRACTIONS $M_i \equiv \int_0^1 x q_i(x) dx$:

$$\frac{M_s + M_{\bar{s}}}{M_{\bar{u}} + M_{\bar{d}}} \sim 0.5 \quad \text{BUT} \quad \frac{M_s + M_{\bar{s}}}{M_u + M_d + M_{\bar{u}} + M_{\bar{d}}} \sim 0.03$$

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MODELS OF THE NUCLEON

- LARGE INTRINSIC STRANGENESS FROM MESON CLOUDS

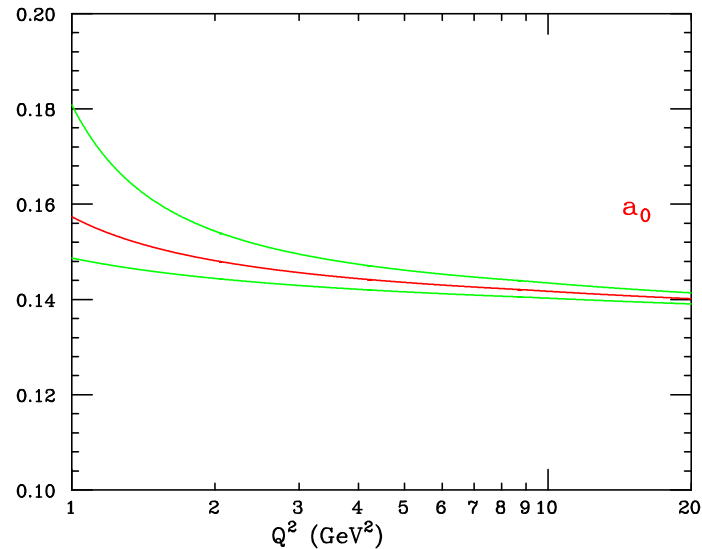
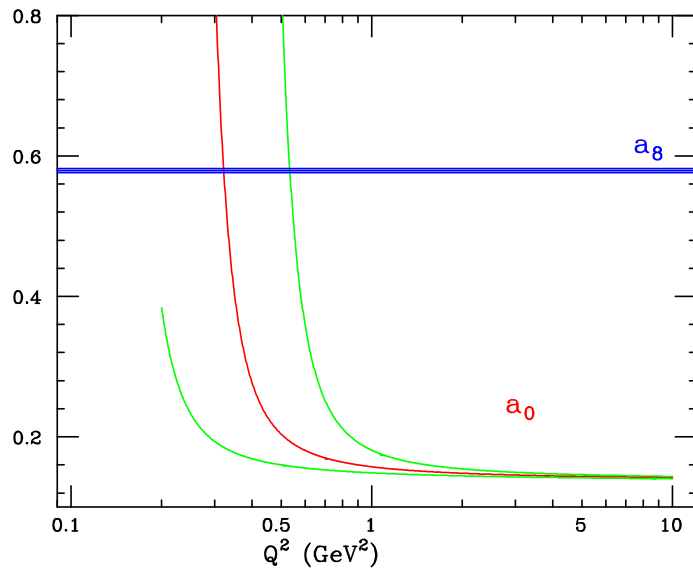
$$\Leftrightarrow \text{OZI EVASION } p\bar{p} \rightarrow \gamma\phi$$

- SKYRME MODEL: NUCLEON FROM MESON OCTET (π , K , η_8) BUT NO COUPLING TO $\eta_0 \Leftrightarrow \Delta s = -(\Delta u = \Delta d)$
(Brodsky, Ellis, Karliner, 1988)

THE AXIAL ANOMALY

SCALE DEPENDENCE OF THE AXIAL CHARGE

THE SINGLET a_0 , UNLIKE THE OCTET a_8 , IS SCALE DEPENDENT



SCALE DEPENDENCE OF $a_0 \iff$ NON-CONSERVATION OF j_5^μ (Axial Anomaly)

$$\partial_\mu j_5^\mu = n_f \frac{\alpha_s}{2\pi} \epsilon^{\mu\nu\rho\sigma} \text{tr} F_{\mu\nu} F_{\rho\sigma}$$

PERTURBATIVE RESULTS

SCALE DEPENDENCE OF $a_0 \Rightarrow$ SCALE DEPENDENCE OF $\int_0^1 dx \Delta g$

$$\begin{aligned} \frac{d}{d \ln Q^2} a_0(Q^2) &= \frac{d}{d \ln Q^2} \left(-2n_f \frac{\alpha_s}{2\pi} \int_0^1 dx \Delta g \right) \\ &= \gamma(\alpha_s) a_0(Q^2); \quad \gamma(\alpha_s) = \gamma^{(2)} \alpha_s^2 + \gamma^{(3)} \alpha_s^3 + \dots \end{aligned}$$

- $\int_0^1 dx \Delta g \sim \frac{1}{\alpha_s} \Leftrightarrow$ THE GLUON DOES NOT DECOUPLE AS $Q^2 \rightarrow \infty$
- CAN DEFINE A SCALE-INDEPENDENT QUARK (choice of factn. scheme)

$$a_0 = \int_0^1 dx \sum_{i=1}^{n_f} (\Delta q_i + \Delta \bar{q}_i) - \frac{\alpha_s}{2\pi} \int_0^1 dx \Delta g$$

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IF Δg IS LARGE $\frac{\alpha_s}{2\pi} \int_0^1 dx \Delta g \sim 0.5$

- SCHEME DEPENDENCE OF Δq AS LARGE AS OZI VIOLATION
- NO OZI VIOLATION IN SCHEME WHERE ALL QUARKS $\int_0^1 (\Delta q_i + \Delta \bar{q}_i)$ SCALE INDEP. (Altarelli and Ross, 1988)

INSTANTONS?

- QCD VACUUM CAN CONTRIBUTE TO $\langle j_5^\mu \rangle$ DUE TO ANOMALY

$$\langle j_5^\mu \rangle = \frac{\alpha}{\pi} \epsilon^{\mu\nu\rho\sigma} \text{tr} \int d^3x (A_\nu(\vec{x}) \partial_\rho A_\sigma(\vec{x}) + A_\mu(\vec{x}) A_\nu(\vec{x}) A_\rho(\vec{x})) + \text{index}$$

nonlocal index term required for gauge invariance

- TUNNELING BETWEEN VACUA \Leftrightarrow CREATION OF AXIAL CHARGE $\langle j_5^\mu \rangle$

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- IN $n_f = 1, m_q = 0$ CASE

FULL SCREENING OF QUARK AXIAL CHARGE IN INSTANTON VACUUM

$$\langle q | j_5^\mu | q \rangle = 0$$

(S.F., 1989, S.F. and Shuryak 1991)

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IF SO, Δg SMALL, a_0 SMALL (ZERO?), $\Delta s = \Delta \bar{s}$

can view instanton contribution as polarized quark sea

DESIDERATA

WHAT WOULD WE LIKE TO KNOW?

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IF NOT, NO INSTANTONS
⇒ NEED DIRECT MEASUREMENT OF $\Delta s + \Delta \bar{s}$ w.o. using SU(3)
- IS THE STRANGENESS “VALENCELIKE”?
IF NOT, NO SKYRMIONS
⇒ NEED MEASUREMENT OF $\Delta s - \Delta \bar{s}$

PRESENT

UNPOLARIZED EXPERIENCE

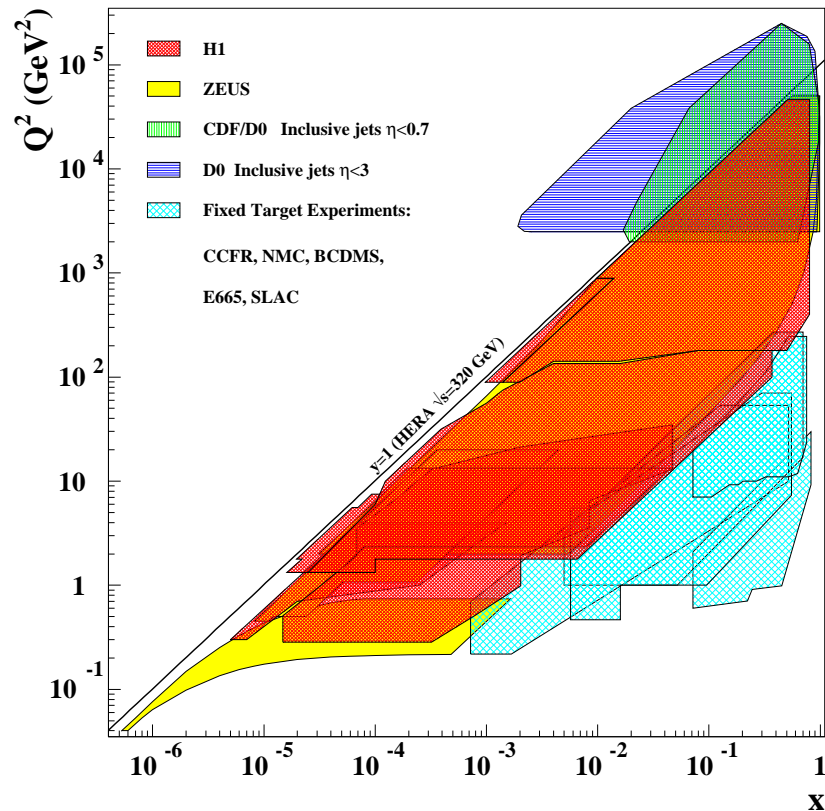
HOW ARE PDFS DETERMINED?

GLOBAL FITS

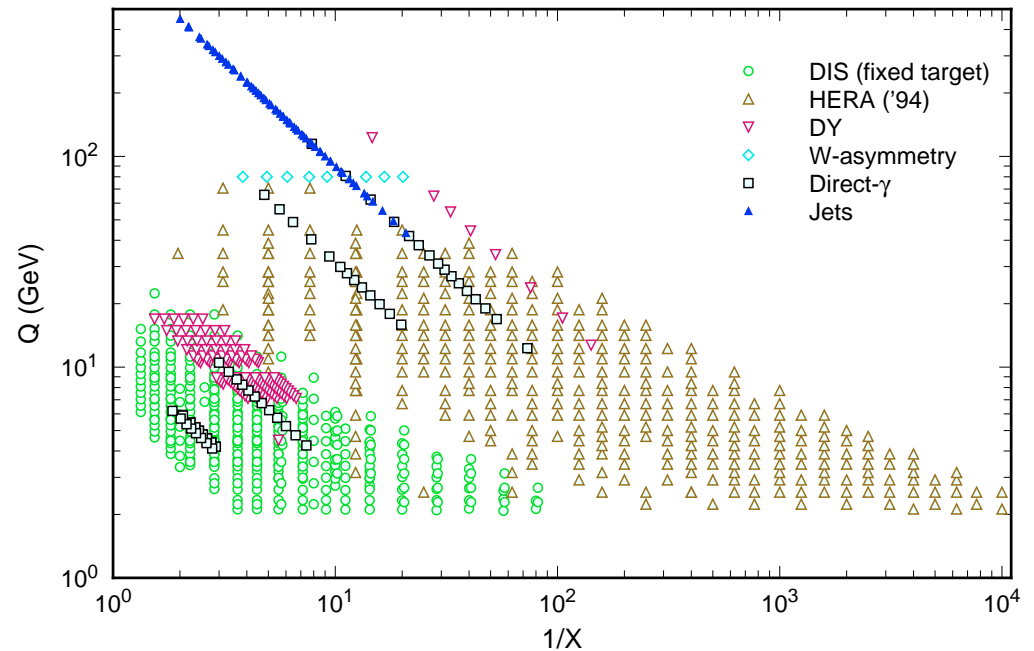
KINEMATIC COVERAGE

OF UNPOLARIZED EXPERIMENTS

ν, μ FIXED TARG.; $e^\pm p$ COLL.



DATA IN CTEQ5 PARTON FIT



THE GLUON

DETERMINED BY SINGLET SCALING VIOLATIONS

$$\frac{d}{d \ln Q^2} F_2^s(N, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} [\gamma_{qq}(N) F_2^s + 2 n_f \gamma_{qg}(N) g(N, Q^2)] + O(\alpha_s^2)$$

$$F_2(N, Q^2) \equiv \int_0^1 dx x^{N-1} F_2(x, Q^2); \quad \gamma_{ij}(N) \equiv \int_0^1 dx x^{N-1} P_{ij}(x, Q^2)$$

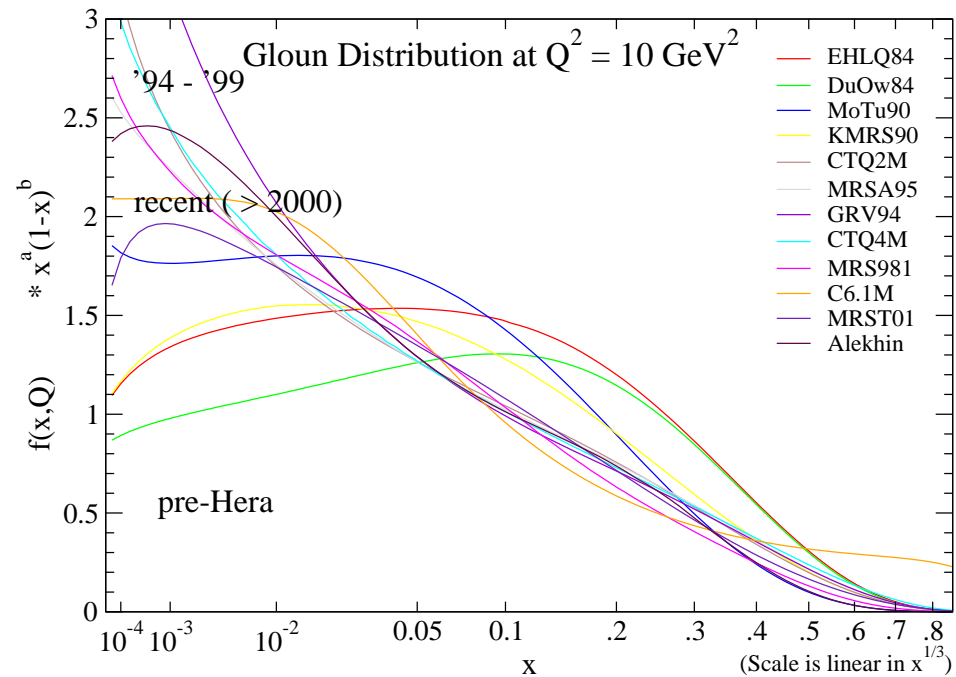
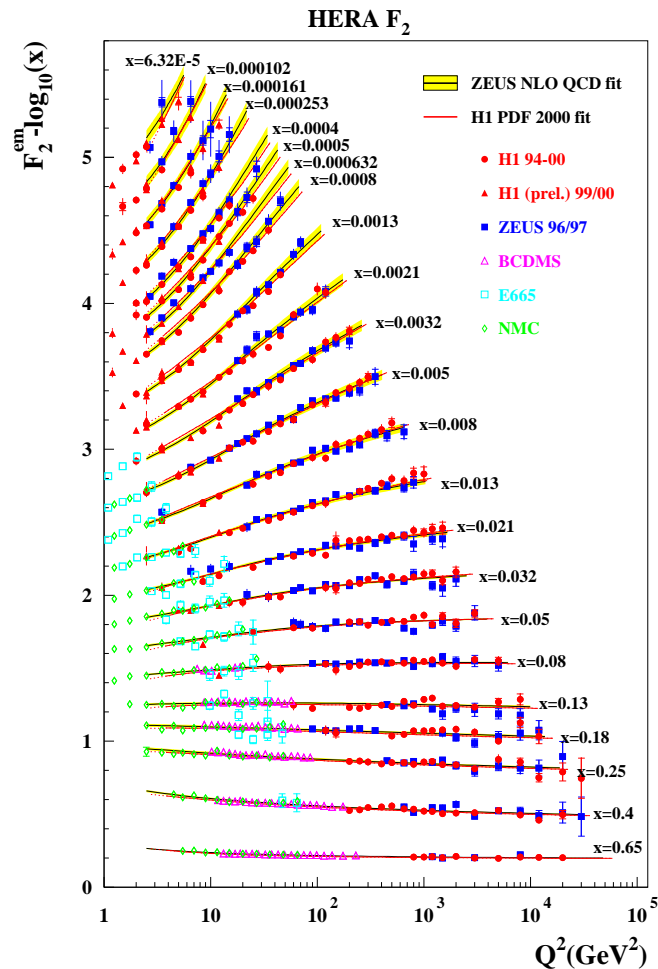
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NEED LARGE LEVER ARM \Rightarrow ONLY REALLY KNOWN AFTER HERA



W.K. Tung, 2004

STRANGENESS

γ^* SCATTERING vs. W^\pm SCATTERING:

IN NC, CHARGED LEPTON DIS, ONLY MEASURE COMBINATION $\sum_i e_i^2 (q_i + \bar{q}_i)$

- CANNOT DETERMINE STRANGENESS
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IN NEUTRINO DIS, CAN DISENTANGLE INDIVIDUAL PDFS BY LINEAR COMBINATION: AT LO

$$\begin{aligned} \frac{1}{2} \left(F_1^{W^-} + \frac{1}{2} F_3^{W^-} \right) &= u + c; & \frac{1}{2} \left(F_1^{W^+} - \frac{1}{2} F_3^{W^+} \right) &= \bar{u} + \bar{c} \\ \frac{1}{2} \left(F_1^{W^+} + \frac{1}{2} F_3^{W^+} \right) &= d + s; & \frac{1}{2} \left(F_1^{W^-} - \frac{1}{2} F_3^{W^-} \right) &= \bar{d} + \bar{s} \end{aligned}$$

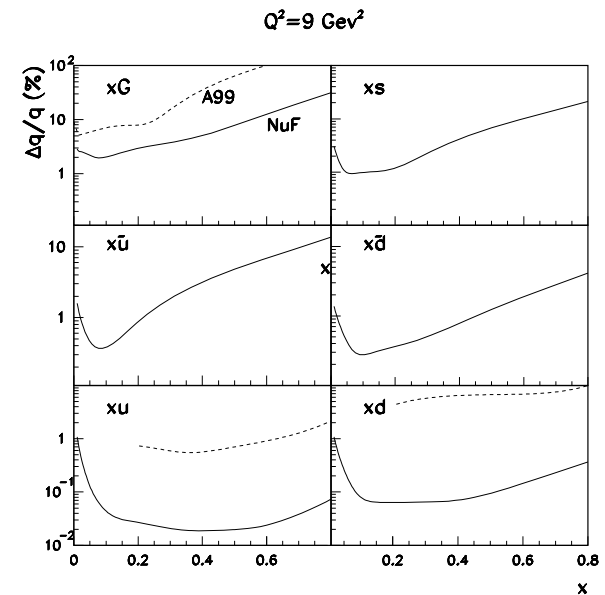
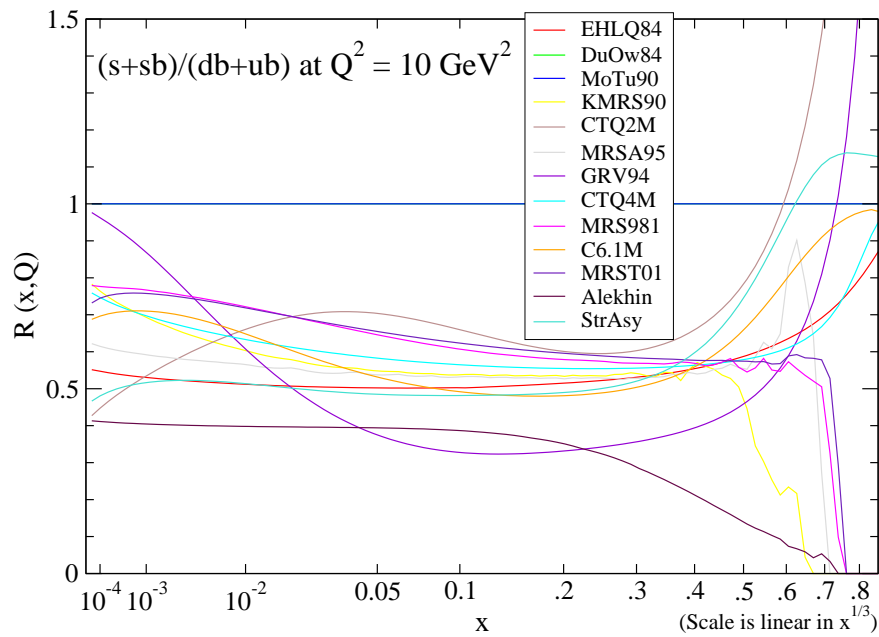
c, \bar{c}, s, \bar{s} only present above charm threshold

STRANGENESS: TOTAL

NEUTRINO DATA

PRESENT: ONLY $\frac{M_s + M_{\bar{s}}}{M_{\bar{u}} + M_{\bar{d}}} \sim 0.5$
REALLY KNOWN

FUTURE: UNCERTAINTIES AT A
NEUTRINO FACTORY
(50 GEV MUONS)



STRANGENESS: VALENCE

COMBINING INCLUSIVE AND EXCLUSIVE INFORMATION

STRANGENESS TAGGED FROM CHARM PRODUCTION: $W^+ + s \rightarrow c$

easily tagged through dimuon signal, 2nd muon from subsequent c decay

W^\pm ASYMMETRY $\Leftrightarrow s - \bar{s}$ ASYMMETRY RELEVANT FOR NUTEV ANOMALY

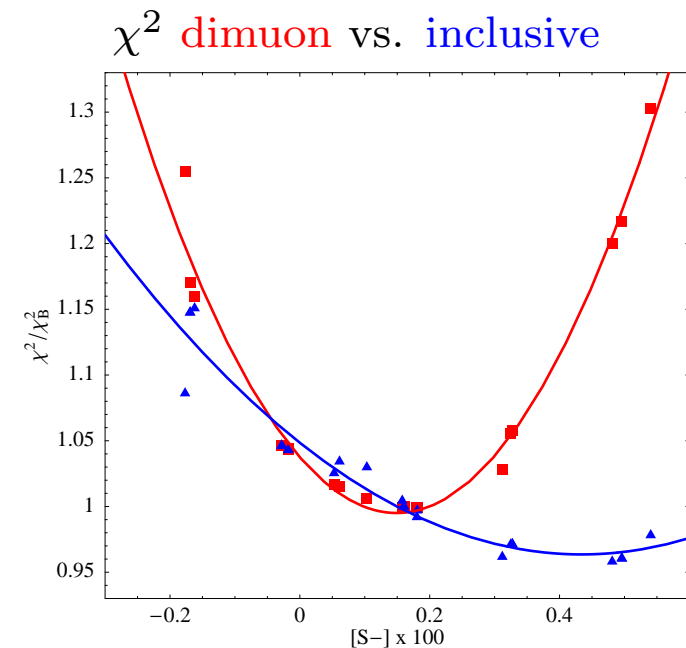
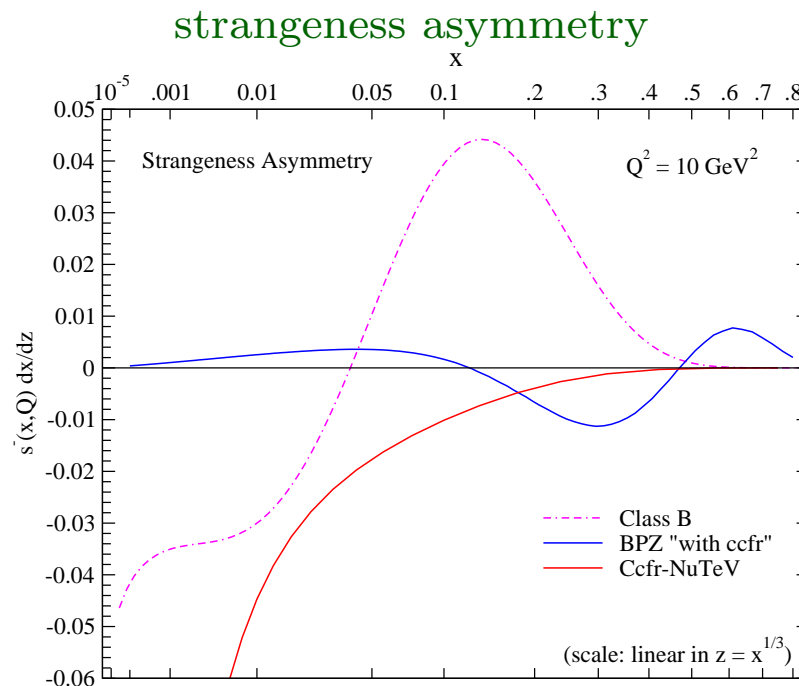
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W^\pm ASYMMETRY $\Leftrightarrow s - \bar{s}$ ASYMMETRY RELEVANT FOR NU ν TEV ANOMALY

- $\int_0^1 (s(x) - \bar{s}(x)) dx = 0$ IN PROTON
 \Rightarrow EITHER $s(x) - \bar{s}(x)$ ZERO, OR MUST CHANGE SIGN
- DIMUON DATA $\Rightarrow [s(x) - \bar{s}(x)] < 0$ FOR $x \lesssim 0.05$
- POSITIVE MOM. FRACTION $s - \bar{s} \approx 0.02$ (CTEQ 2004)



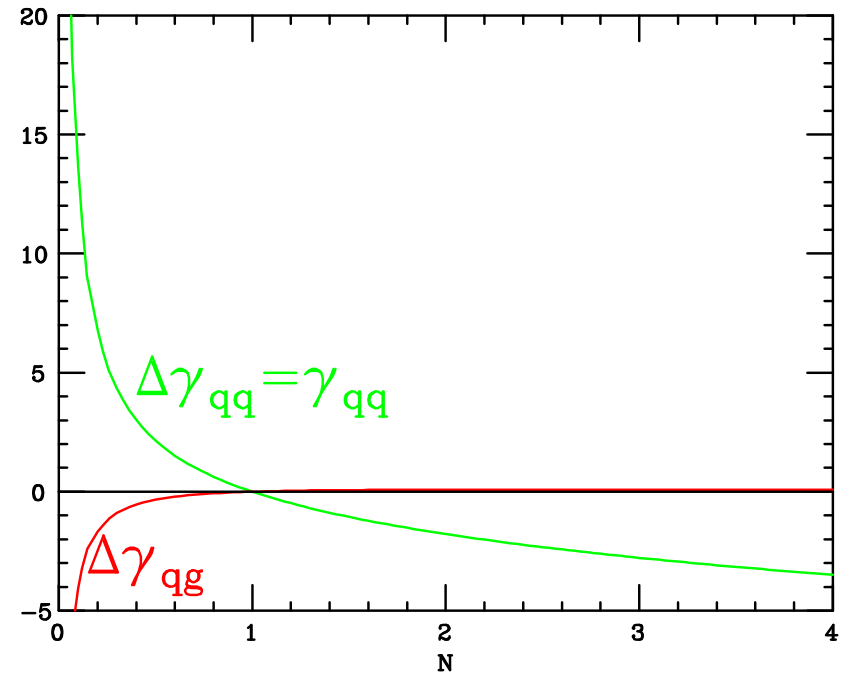
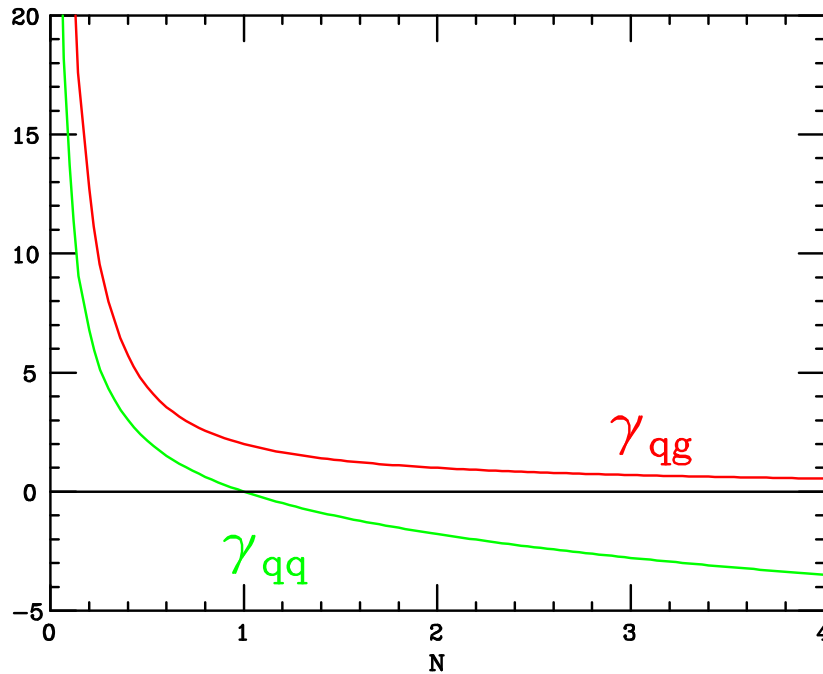
POLARIZED GLUE?

PROBLEMS:

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POLARIZED GLUE? PROBLEMS:

- NO POLARIZED COLLIDER! LIMITED KINEMATIC COVERAGE IN Q^2
- GLUON DRIVES SCALING VIOLATIONS ONLY AT LOW x !
WHERE THERE ARE NO DATA...



$$\frac{d}{d \ln Q^2} F_2^s(N, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} [\gamma_{qq}(N) F_2^s + 2 n_f \gamma_{qg}(N) g(N, Q^2)] + O(\alpha_s^2)$$

$$\frac{d}{d \ln Q^2} g_1^s(N, Q^2) = \frac{\alpha_s(Q^2)}{2\pi} [\Delta\gamma_{qq}(N) g_1^s + 2 n_f \Delta\gamma_{qg}(N) \Delta g(N, Q^2)] + O(\alpha_s^2)$$

POLARIZED STRANGENESS?

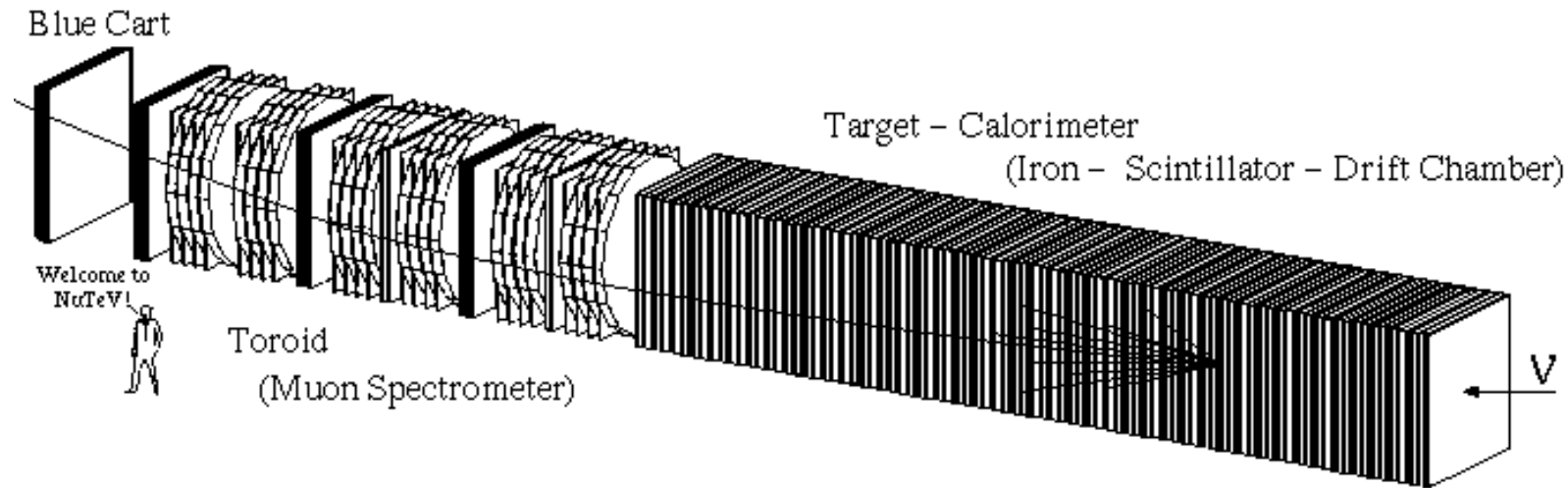
PROBLEM:

NO POLARIZED NEUTRINO DATA!

POLARIZED STRANGENESS?

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The NUTEV target/spectrometer

CAN'T FIT TARGET INTO MAGNET!

SEMI-INCLUSIVE CHANNELS?

CAN DETERMINE PARTON CONTENT BY TAGGING FINAL STATE HADRONS

E.G. IN SIDIS: $\gamma^* p \rightarrow \pi^\pm + X \Rightarrow u/d, \bar{d}/\bar{u}$; $\gamma^* p \rightarrow K^\pm + X \Rightarrow \bar{s}/s$, ETC.

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PROBLEMS:

set	χ^2	χ^2_{DIS}	χ^2_{SIDIS}	δu_v	δd_v	$\delta \bar{u}$	$\delta \bar{d}$	$\delta \bar{s}$	δg	$\delta \Sigma$	
NLO	KRE	430.91	206.01	224.90	0.936	-0.344	-0.0487	-0.0545	-0.0508	0.680	0.284
	KKP	436.17	205.66	230.51	0.700	-0.255	0.0866	-0.107	-0.0454	0.574	0.311
LO	KRE	457.54	213.48	244.06	0.697	-0.248	-0.0136	-0.0432	-0.0415	0.121	0.252
	KKP	448.71	219.72	228.99	0.555	-0.188	0.0497	-0.0608	-0.0365	0.187	0.271

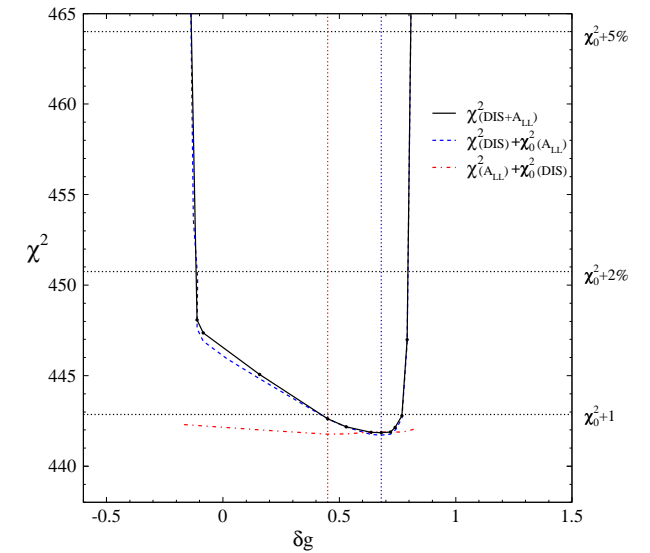
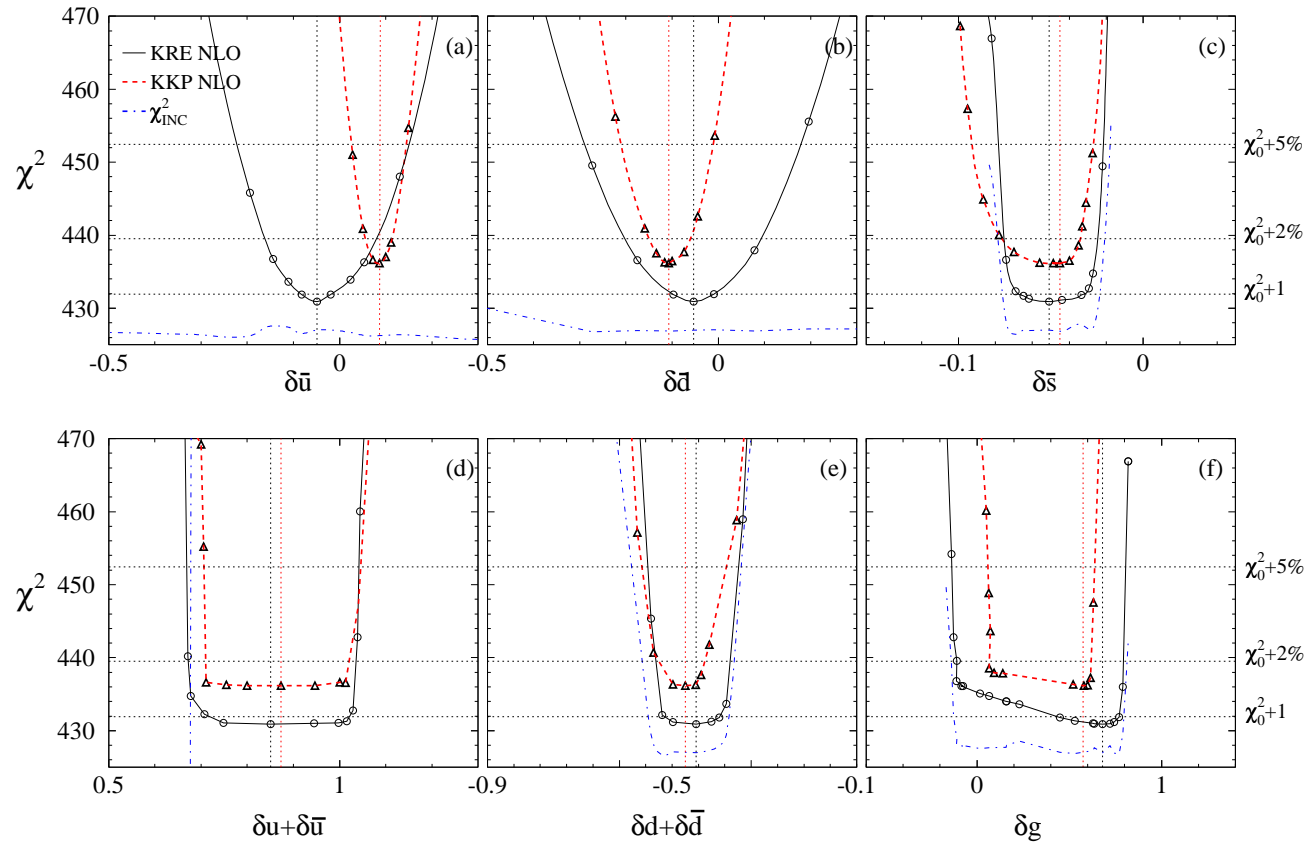
(De Florian, Navarro, Sassot, 2005)

- SENSITIVITY TO POORLY KNOWN FRAGMENTATION FCTN
 $\Delta \bar{u}$ CHANGES SIGN WITH DIFFERENT FRAGMENTATION
- LARGE NLO CORRECTIONS
 $\Delta \bar{u}$ CHANGES BY FACTOR TWO FROM LO TO NLO

THE POLARIZED CASE: ARE WE THERE YET?

GLOBAL FIT WITH SIDIS DATA (De Florian, Navarro, Sassot 2005)

THE IMPACT OF THE PHENIX
 $A_{LL}^{\pi^0}$ (Navarro, Sassot 2006)

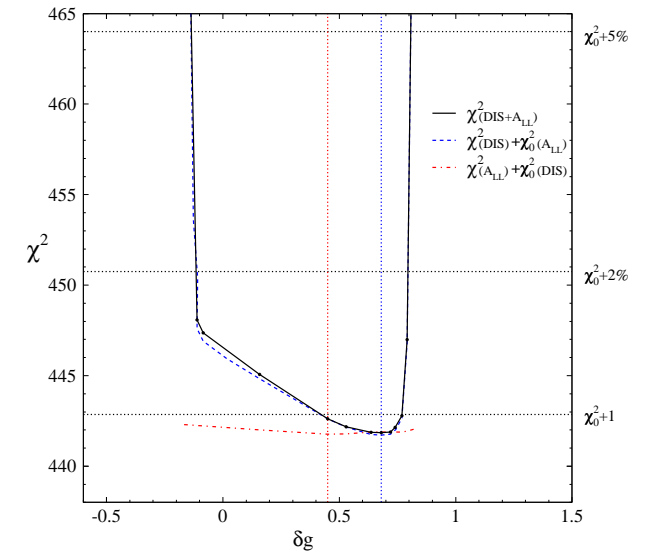
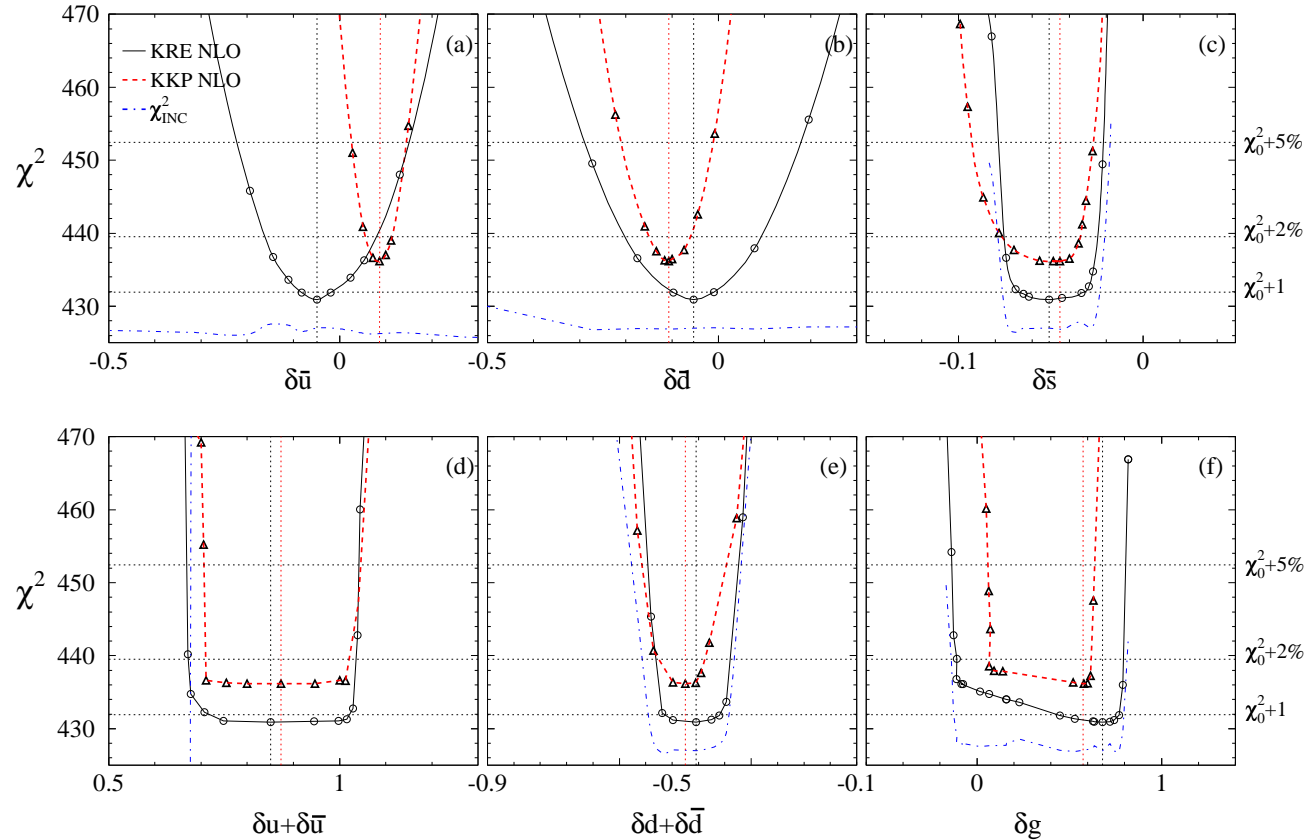


- LARGE TH. UNCERTAINTIES (FRAGMENTATION)

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- LARGE TH. UNCERTAINTIES (FRAGMENTATION)
- UNKNOWN PARAMETRIZATION BIAS (E.G. POSITIVITY)
- ERRORS NOT PARABOLIC \Rightarrow ONE SIGMA ERROR BANDS FROM $\Delta\chi^2 \sim 20$

THE NUCLEON SPIN

WHAT HAVE WE LEARNT FROM GLOBAL FITS?

FIRST MOMENTS

- SINGLET AXIAL CHARGE:

$$a_0(10 \text{ GeV}^2) = 0.18 \pm 0.03^{+0.17}_{-0.11} \text{ [ABF]};$$

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$$a_0(10 \text{ GeV}^2) = 0.18 \pm 0.03^{+0.17}_{-0.11} \text{ [ABF];}$$

- **GLUON:**

$$\Delta g(1 \text{ GeV}^2) = 0.8 \pm 0.2 \pm 0.8; \text{ [ABF]; } = 0.4^{+0.1}_{-0.5} \text{ [DNS];}$$

- **STRANGENESS:**

$$\Delta s^{\overline{MS}}(1 \text{ GeV}^2) = -0.13 \pm 0.02 \pm 0.09; \text{ [ABF]; } -0.10 \pm 0.06 \text{ [DNS]}$$

[ABF]: ALTARELLI, BALL, S.F., RIDOLFI 98-01:

- STAT. ERRORS OBTAINED FROM ONE-SIGMA CONTOURS
- MAIN **SYSTEMATICS: SMALL x EXTRAP., FUNCTIONAL FORM**
- SU(3) USED TO DETERMINE OCTET, 30% VIOLATION ALLOWED;

[DNS]: DE FLORIAN, NAVARRO, SASSOT 05

- ERRORS ESTIMATED FROM $\Delta\chi^2 = 0.02\chi_{min}^2 \approx 20$
- **STRANGENESS DETERMINED DIRECTLY (SU(3) NOT USED)**
- **GLUON INCLUDES PHENIX DATA**

THE NUCLEON SPIN

WHAT HAVE WE LEARNT FROM GLOBAL FITS?

FIRST MOMENTS

- **SINGLET AXIAL CHARGE:**

$$a_0(10 \text{ GeV}^2) = 0.18 \pm 0.03^{+0.17}_{-0.11} \text{ [ABF];}$$

- **GLUON:**

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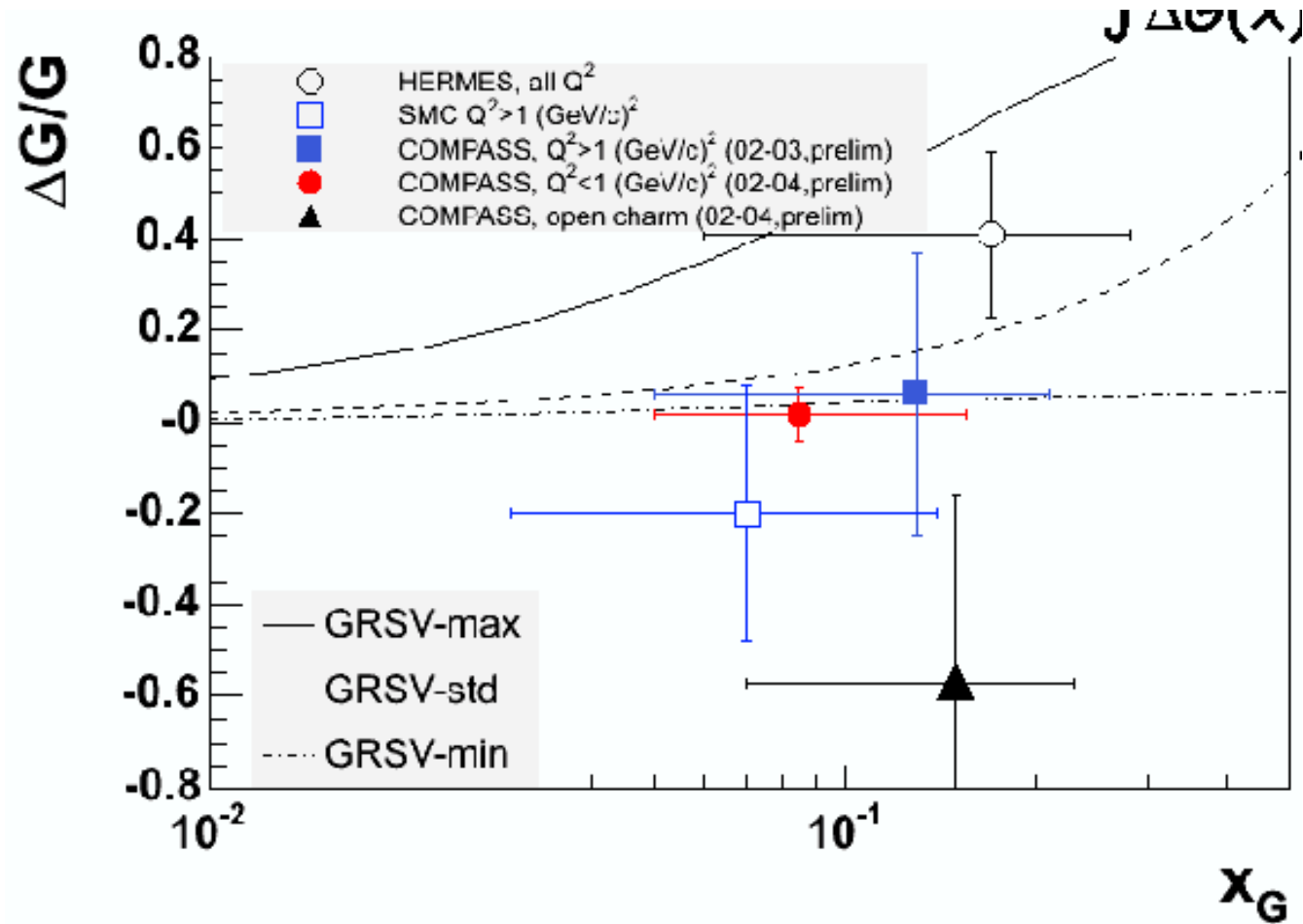
- **THEORETICAL ERRORS DOMINATE**

- SEMI-INCLUSIVE DATA HAVE MODERATE IMPACT ON $\Delta g, \Delta s$

- SU(3) NOT TESTED YET

- NO INFORMATION OF $\Delta s - \Delta \bar{s}$

INFORMATION FROM HIGH p_T HADRONS (AND CHARM)



- NLO CALCULATIONS NOT AVAILABLE
- DEPENDENCE ON MONTE CARLO
- DECONVOLUTION OF GLUON USING MEAN VALUES

(NO) CONCLUSION ON NUCLEON SPIN YET
EXISTING RESULTS VERY HARD TO COMPARE:

(NO) CONCLUSION ON NUCLEON SPIN YET

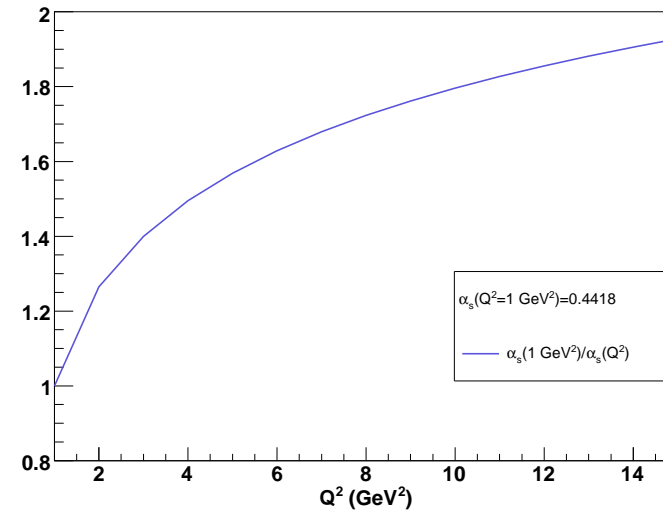
EXISTING RESULTS VERY HARD TO COMPARE:

- **SOME LO, SOME NLO:** DIFFERENCE TYPICALLY 20%

- ΔG **DEPENDS STRONGLY ON SCALE:**

$$\Delta G \sim \frac{1}{\alpha_s(Q^2)}, \text{ SO}$$

$$\Delta G(2m_c) \sim 2\Delta G(1 \text{ GEV})$$



(NO) CONCLUSION ON NUCLEON SPIN YET

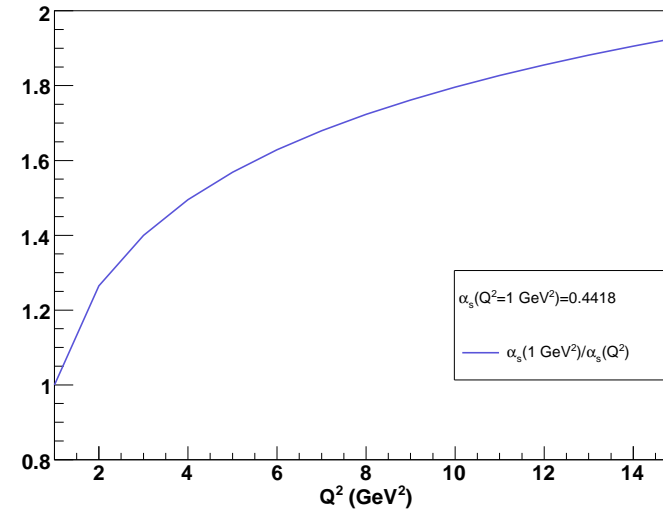
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- **UNCERTAINTIES** DETERMINED WITH $\Delta\chi^2 = 1$ UP TO $\Delta\chi^2 = 12$
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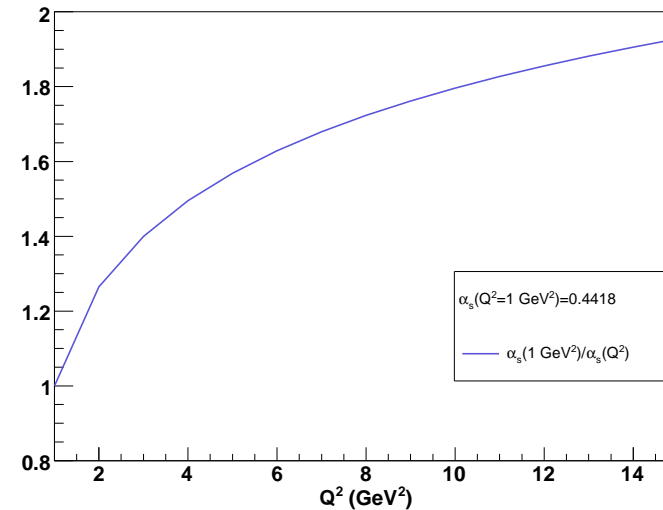
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ΔG **PROBABLY POSITIVE** BUT PERHAPS **SMALL** $\left(\frac{N_f}{2\pi}\right) \alpha_s(Q^2) \Delta G \lesssim \frac{1}{2} \Delta\Sigma$,

PERHAPS ZERO

Δs **PROBABLY NEGATIVE**, **MAYBE SMALL**

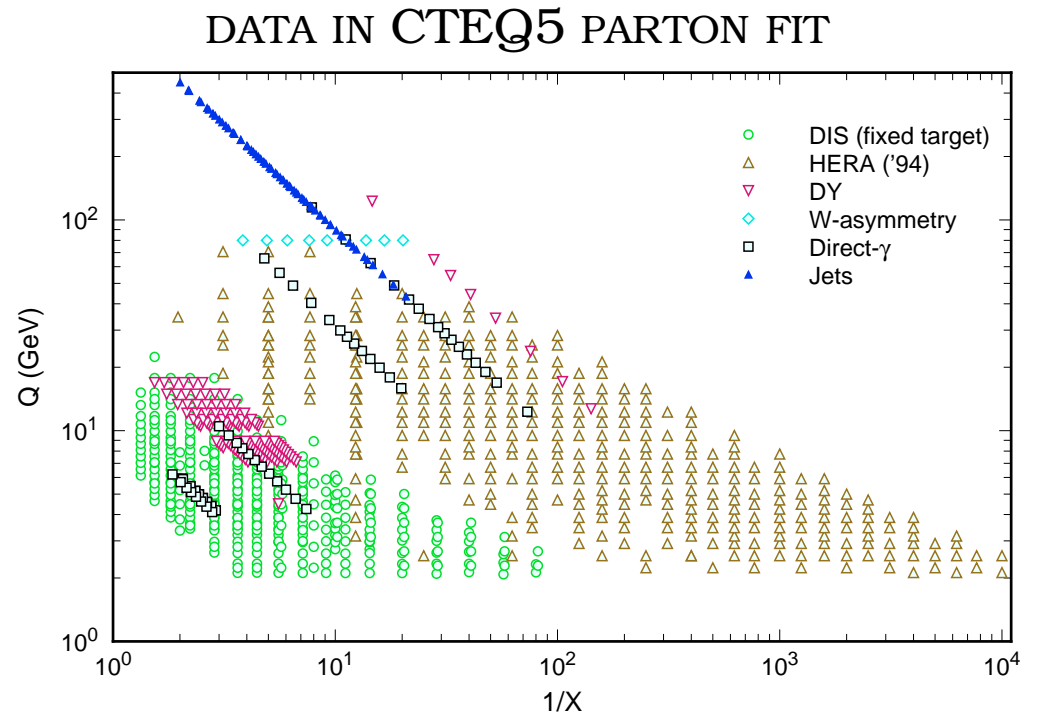
FUTURE

HADRONIC CHANNELS

HADRONIC CHANNELS

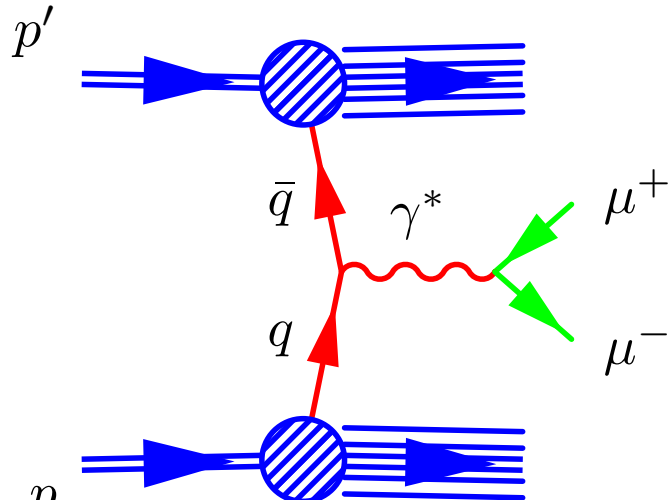
UNPOLARIZED GLOBAL FITS

- DRELL-YAN
⇒ \bar{u}/\bar{d} QUARK ASYMMETRY
- W^\pm PRODUCTION
⇒ u/d QUARK ASYMMETRY
- DIRECT γ ⇒ GLUON
(IMPACT NEGLIGIBLE W.R. TO DIS)
- LARGE E_T JETS
⇒ LARGE x GLUON



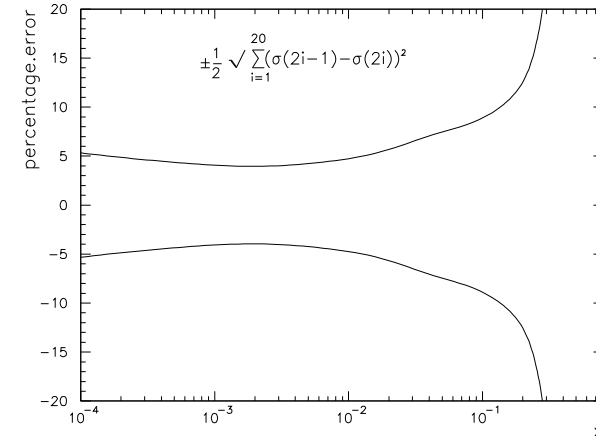
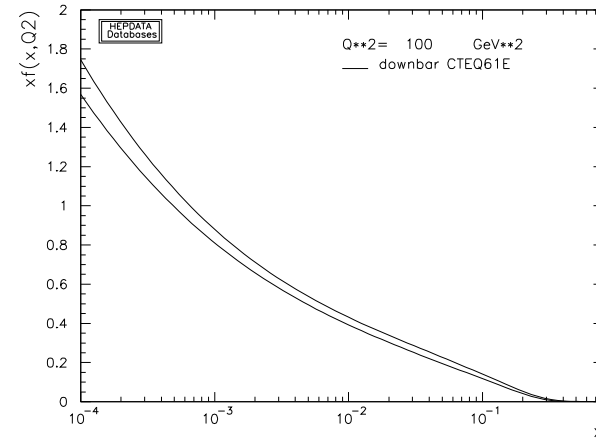
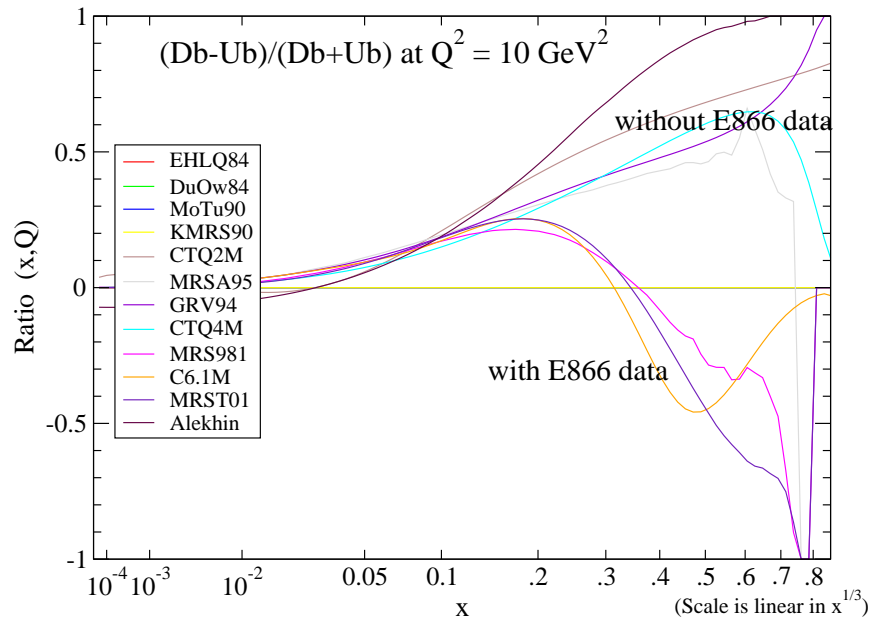
HOW WELL DOES IT WORK?

DRELL-YAN p/d ASYMMETRY



$$\left. \frac{\sigma^{pd}}{\sigma^{pp}} \right|_{x_1 \gg x_2} \approx \frac{1}{2} \left(1 + \frac{\bar{d}(x_2)}{\bar{u}(x_2)} \right)$$

5 – 10% ACCURACY ON \bar{u} , \bar{d}



POLARIZED CASE
WHAT DO WE EXPECT?

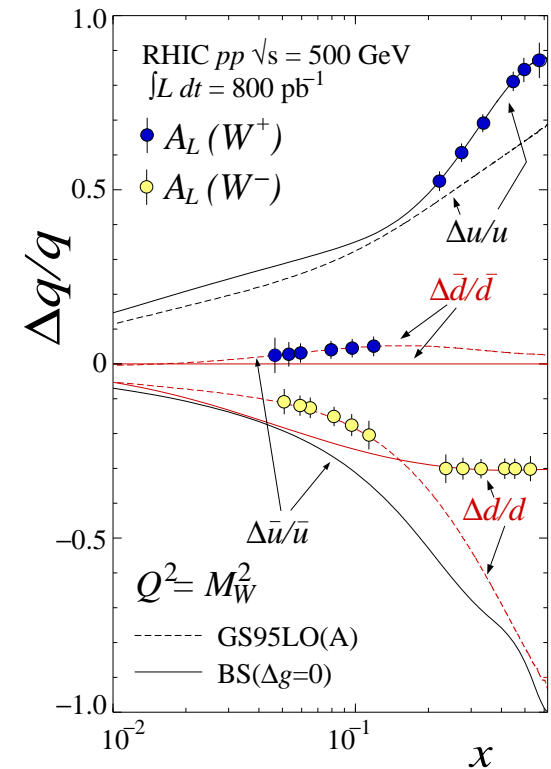
POLARIZED CASE WHAT DO WE EXPECT?

GOOD NEWS

- CAN DETERMINE Δu , Δd , $\Delta \bar{u}$, $\Delta \bar{d}$ TO A FEW PERCENT ACCURACY from W^\pm production

RESULTS

- $\int (\Delta s + \Delta \bar{s}) dx$
COMBINING LIGHT QUARKS WITH a_0 FROM DIS



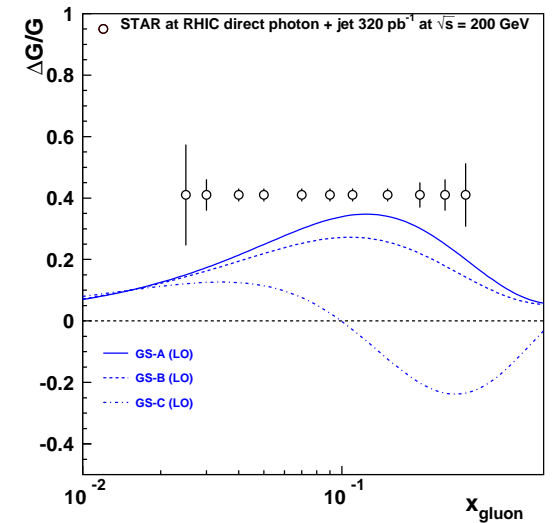
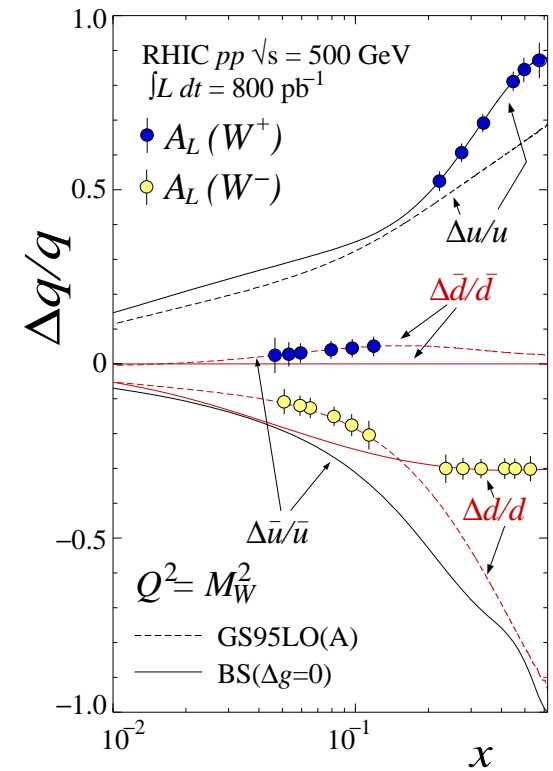
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- CAN DETERMINE Δg AT MANY INDIVIDUAL x VALUES from direct γ , jets, single inclusive hadrons

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- $\int (\Delta s + \Delta \bar{s}) dx$ COMBINING LIGHT QUARKS WITH a_0 FROM DIS
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POLARIZED CASE WHAT DO WE EXPECT?

GOOD NEWS

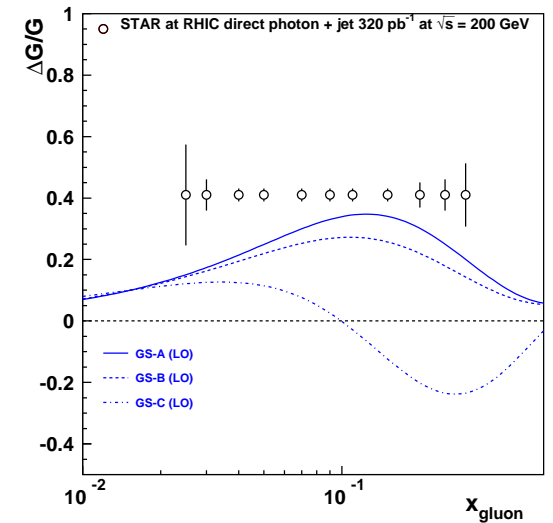
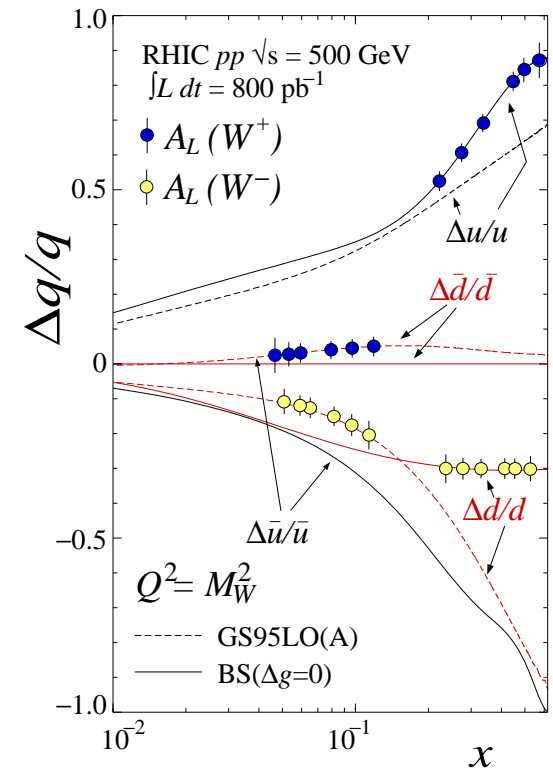
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RESULTS

- $\int (\Delta s + \Delta \bar{s}) dx$ COMBINING LIGHT QUARKS WITH a_0 FROM DIS
- $\int \Delta g dx$ COMBINING $\Delta g(x)$ WITH SCAL VLNS.

BAD NEWS

- NO WAY TO DETERMINE $\Delta s - \Delta \bar{s}$ WITHOUT NEUTRINO DIS
- NEED A GLOBAL FIT!



CAN WE TRUST GLOBAL FITS?

PARTON SETS DO NOT AGREE WITHIN RESPECTIVE ERRORS!

W PRODUCTION CROSS-SECTION TEVATRON

PDF SET	XSEC [NB]	PDF UNCERTAINTY
ALEKHIN	2.73	± 0.05 (TOT)
MRST2002	2.59	± 0.03 (EXPT)
CTEQ6	2.54	± 0.10 (EXPT)

THORNE 2003

- ALEKHIN VS. MRST/CTEQ
→ W PRODUCTION XSECT AT
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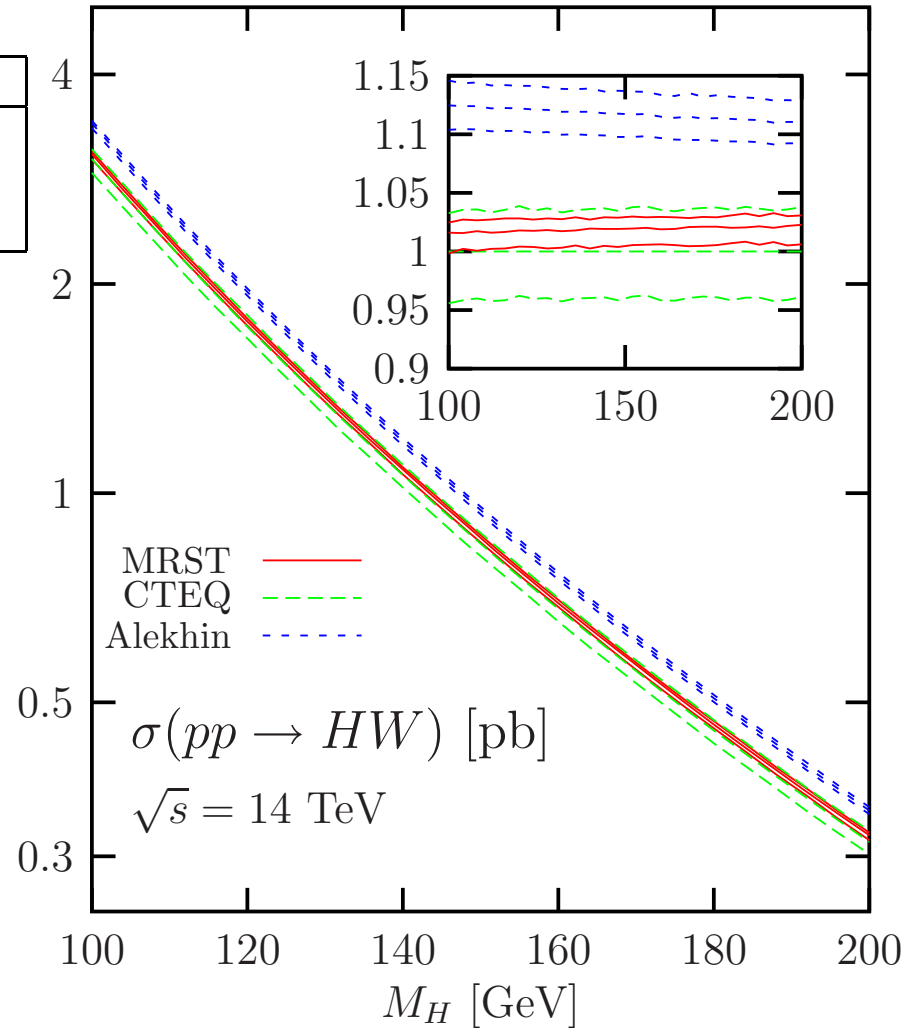
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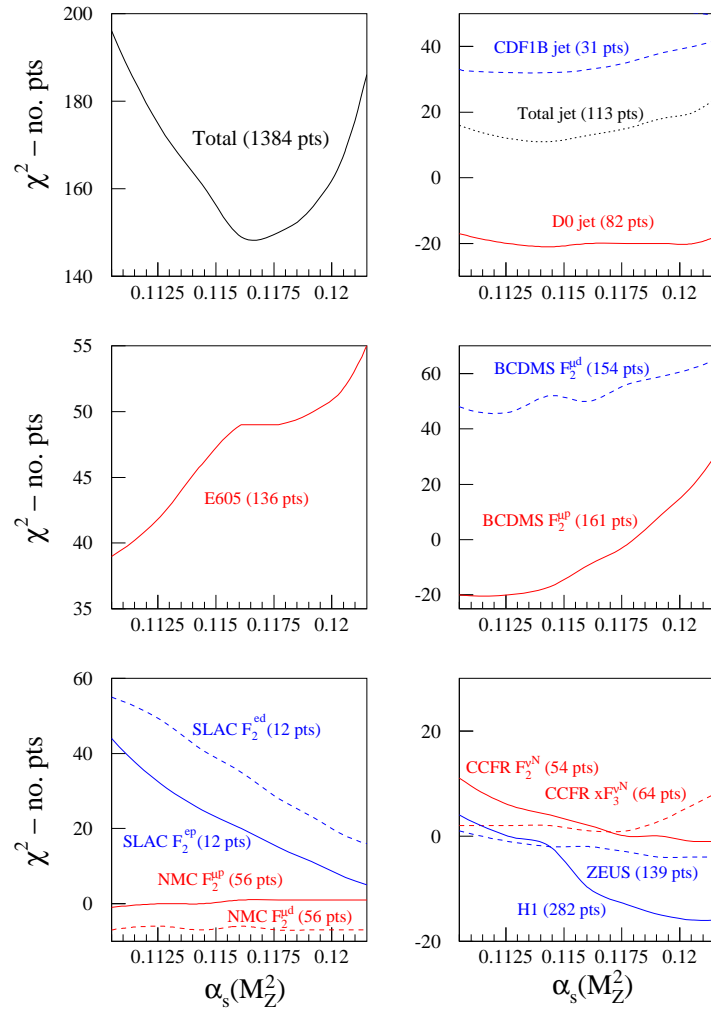
- ALEKHIN VS. MRST/CTEQ
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- ALEKHIN VS. MRST/CTEQ
→ PREDICTIONS FOR ASSOCIATE HIGGS W PRODUCTION LHC DO NOT AGREE WITHIN RESPECTIVE ERRORS

HIGGS PRODUCTION AT LHC



INCOMPATIBLE DATA?

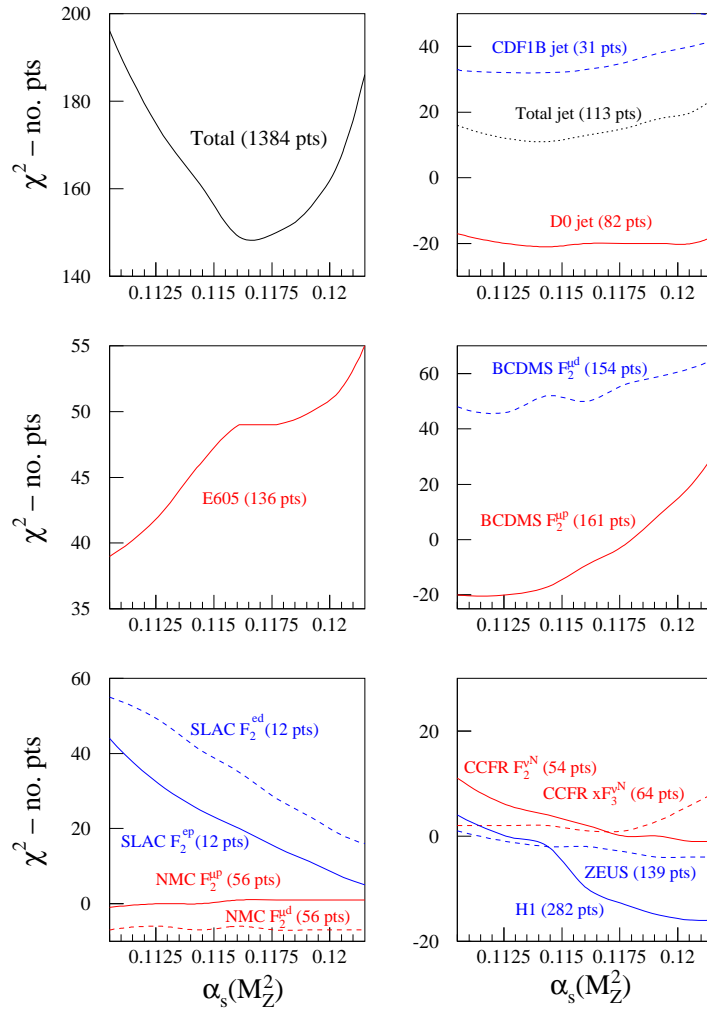
GLOBAL χ^2 MINIMUM MAY NOT
CORRESPOND TO LOCAL MINIMA



MRST 2003

INCOMPATIBLE DATA?

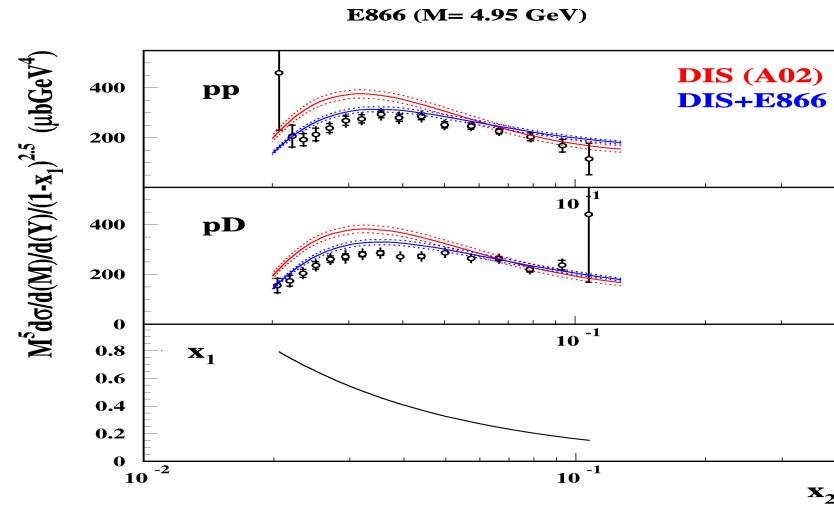
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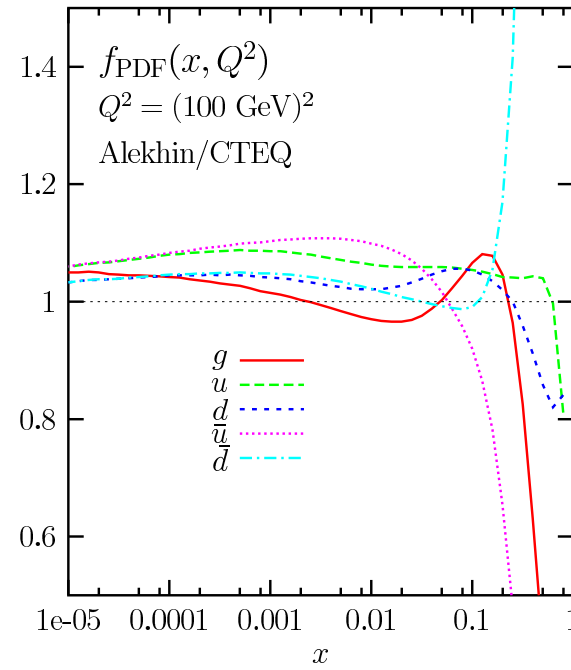
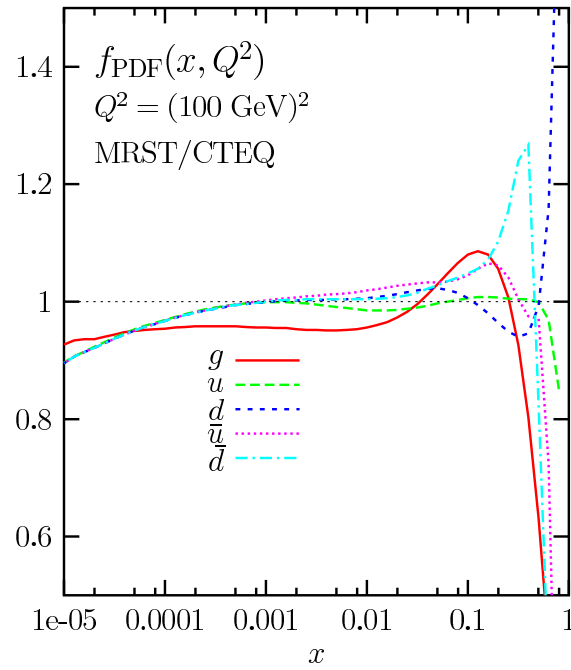
E866 DY DATA DISAGREE WITH DIS DATA:

$\sigma_{DY} \sim q(x_1)q(x_2)$ DISAGREES WITH DIS QUARK AT SAME x AND Q^2



ALEKHIN 2005

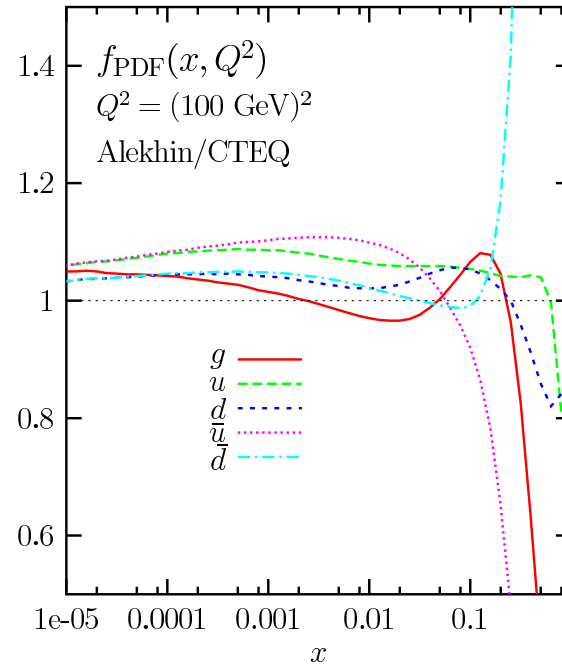
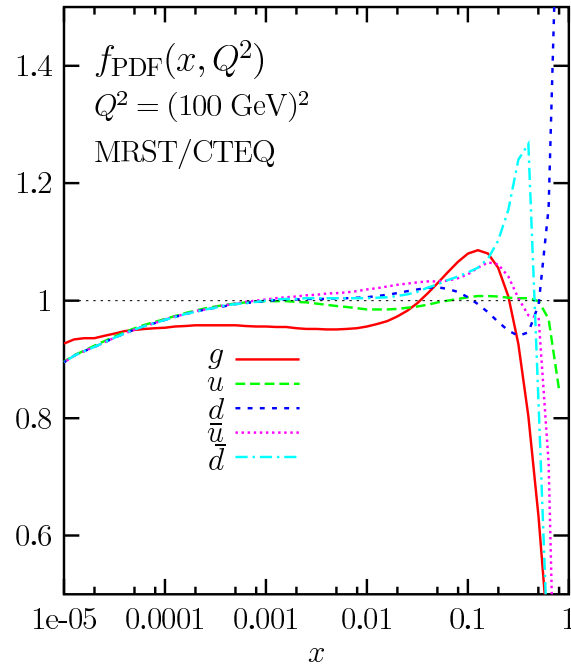
PARAMETRIZATION BIAS?



MRST & CTEQ
→ SIMILAR PARTONS

Djouadi and Ferrag 2003

PARAMETRIZATION BIAS?



MRST & CTEQ
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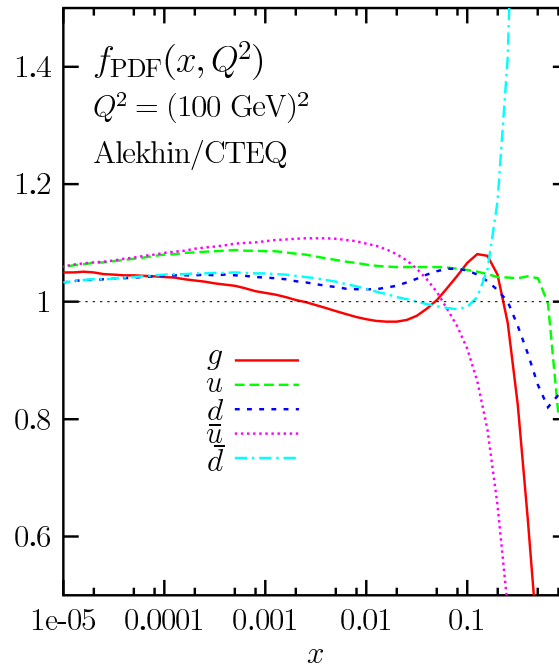
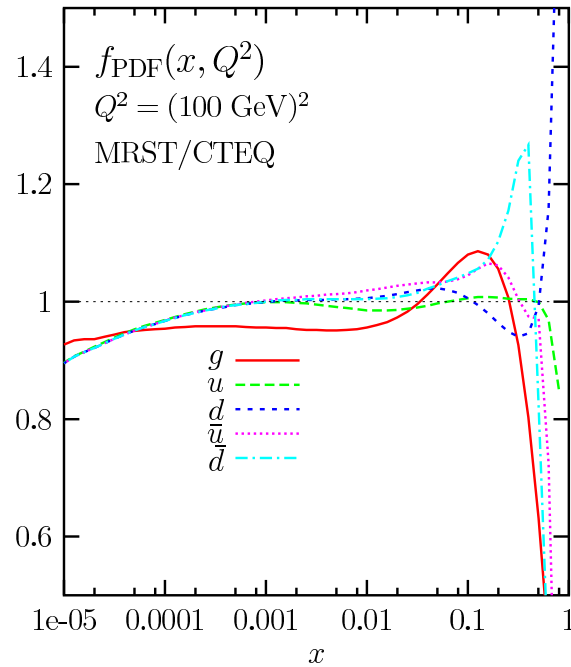
Djouadi and Ferrag 2003

THE W XSECT. AGAIN...

SIMILAR PARTONS
 ⇒ SIMILAR RESULTS

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ALEKHIN	LHC	215	± 6 (TOT)
MRST2002	LHC	204	± 4 (EXPT)
CTEQ6	LHC	205	± 8 (EXPT)

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We do not seem to have the optimum parameterization for both finding the best fit and also investigating fluctuations about this best fit (...) This might then influence our error analysis...(MRST 2004)

SOLUTIONS: CTEQ TOLERANCE CRITERION

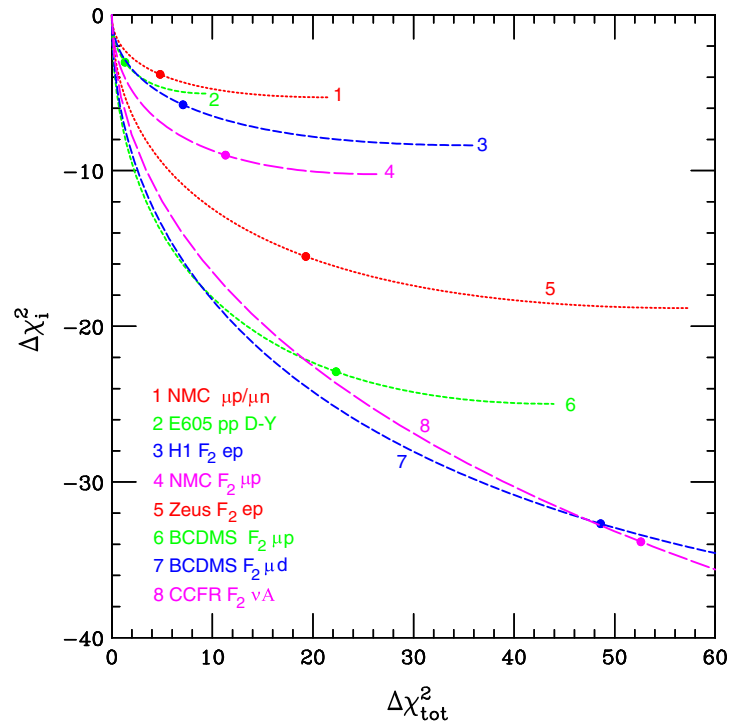
SINGLE OUT INCONSISTENT DATA

- HOW MANY PARAMETERS ARE SIGNIFICANTLY DETERMINED BY EACH DATASET?
- HOW CONSISTENT ARE THE DATA FROM ONE SET WITH THE REST?

STUDY MINIMUM ALLOWED χ_i^2

FOR i -TH EXP. AS

GLOBAL χ^2 ALLOWED TO INCREASE



Collins, Pumplin 2001

CCFR, BCDMS

INCOMPATIBLE WITH THE REST

SOLUTIONS: CTEQ TOLERANCE CRITERION

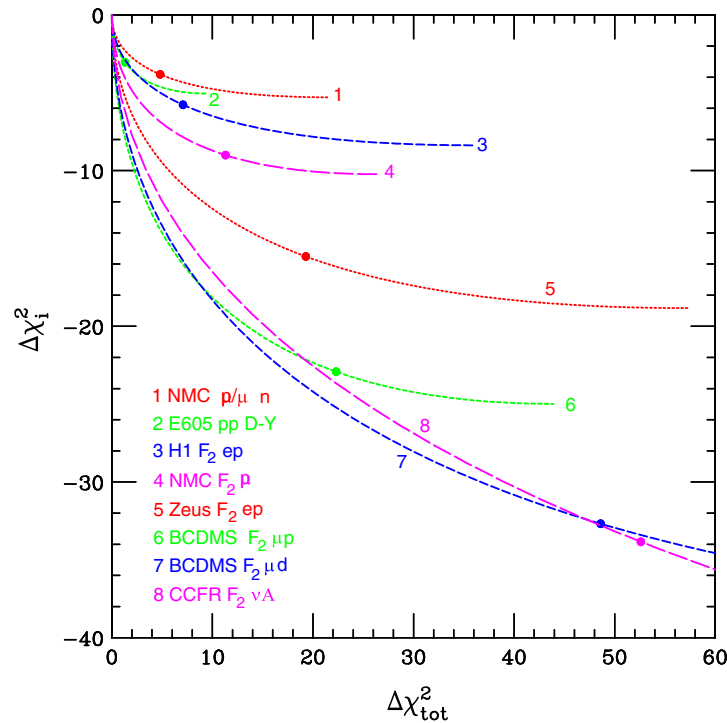
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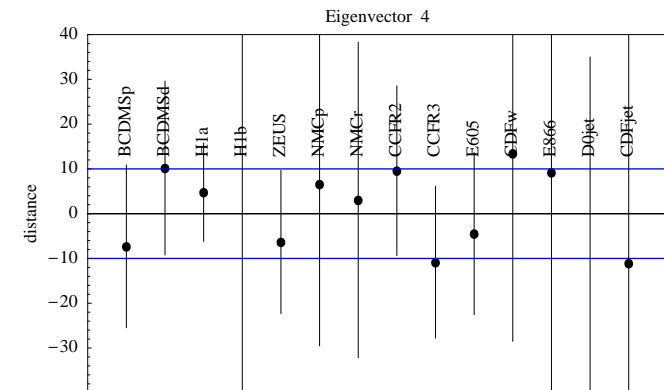
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INCOMPATIBLE WITH THE REST

OPTIONS

- discard incompatible experiments
- reweight individual contributions
- **INCORPORATE IN ERROR, TOLERATING FIXED MAX DEVIATION FOR EACH EXPERIMENT & EACH FIT PARAMETER**



SOLUTIONS: CTEQ TOLERANCE CRITERION

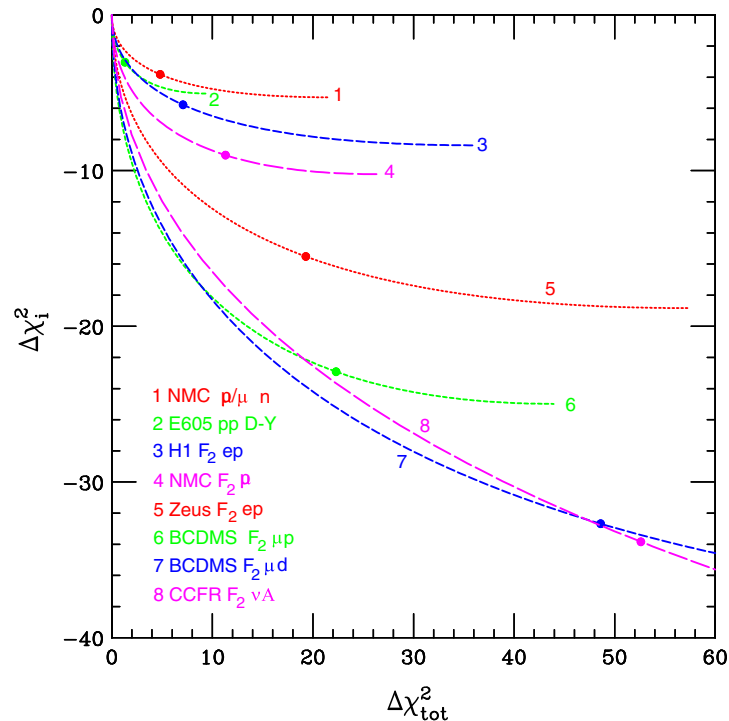
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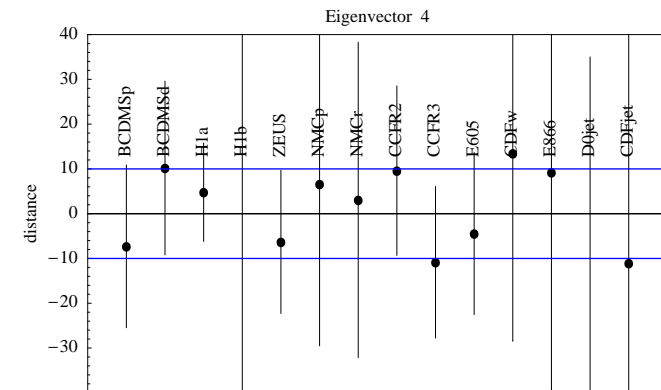
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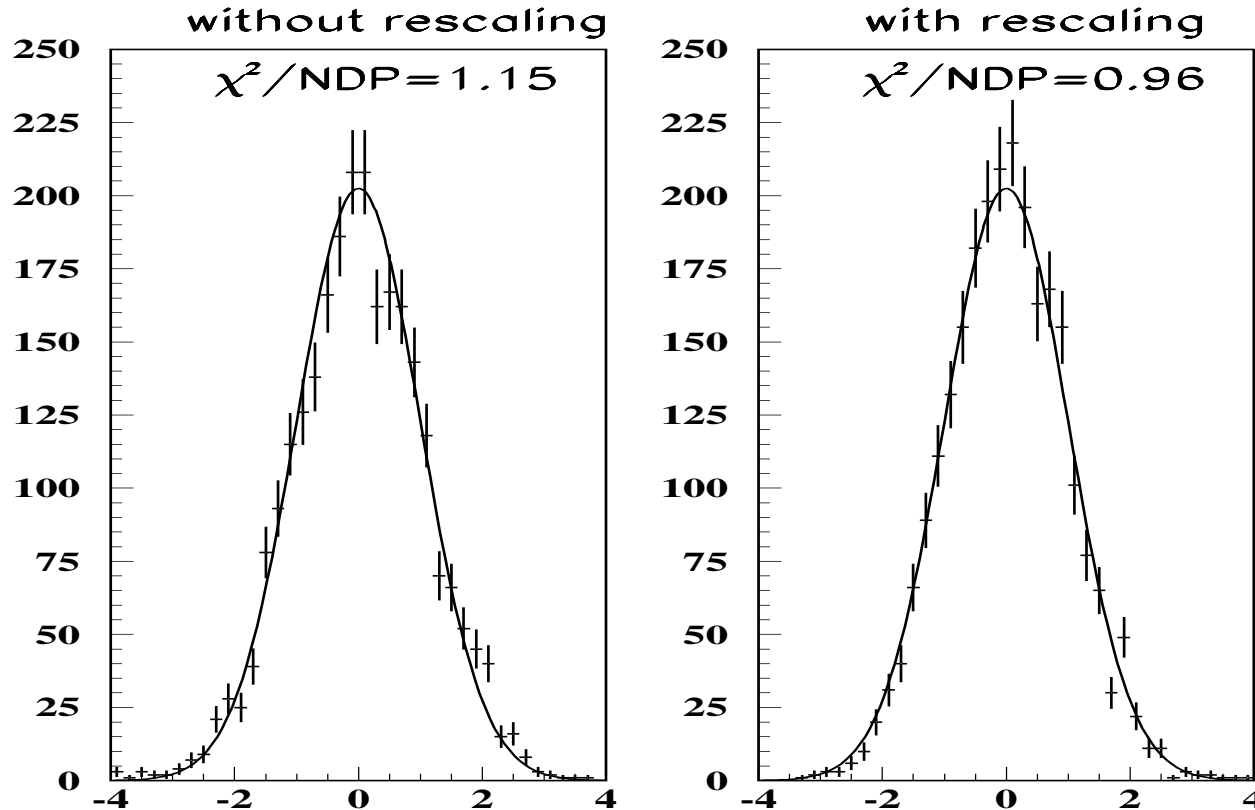
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$$\Delta\chi^2 = 100 \text{ (CTEQ6)}$$

SOLUTIONS: ERROR RESCALING

HOW CAN DATA FROM INCONSISTENT SETS BE INCLUDED?
ASSUME INCONSISTENCY DUE TO UNDERESTIMATED (SYST.) ERROR:

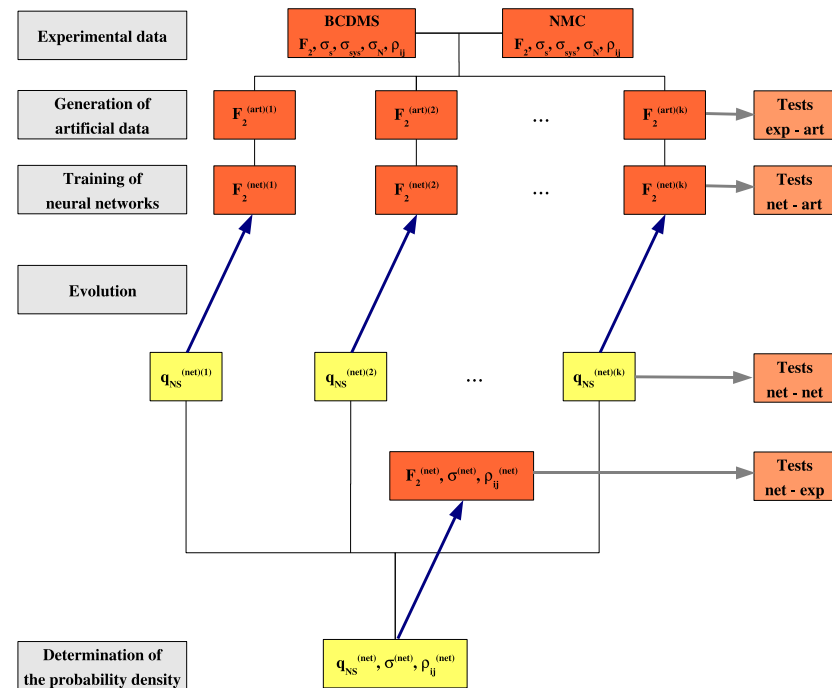


For the experiments with $\chi^2 > 1$ the statistical errors in data were rescaled in order to get $\chi^2 = 1$

THE FUTURE?

THE NEURAL MONTE CARLO APPROACH (S.F., GARRIDO, LATORRE, PICCIONE 2002)

BASIC IDEA: USE NEURAL NETWORKS AS UNIVERSAL UNBIASED INTERPOLANTS

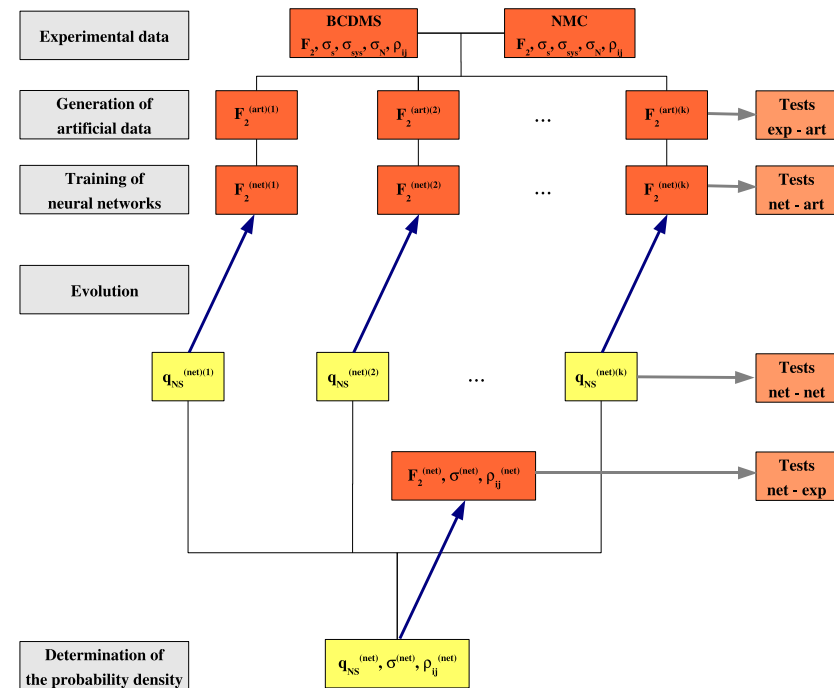


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- GENERATE A SET OF MONTE CARLO REPLICAS $\sigma^{(k)}(p_i)$ OF THE ORIGINAL DATASET $\sigma^{(\text{data})}(p_i)$
 \Rightarrow REPRESENTATION OF $\mathcal{P}[\sigma(p_i)]$ AT DISCRETE SET OF POINTS p_i

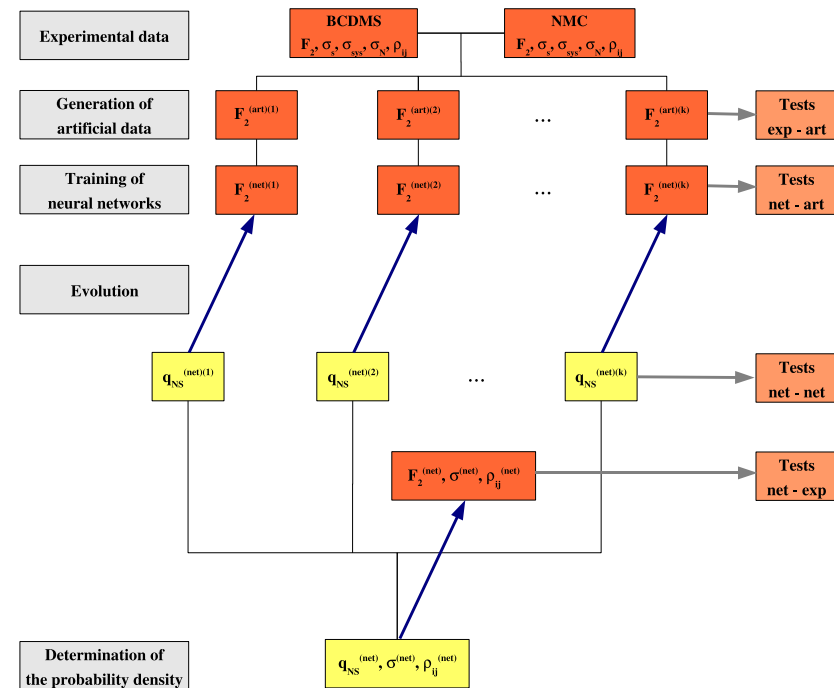


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- TRAIN A NEURAL NET FOR EACH PDF ON EACH REPLICA
 \Rightarrow NEURAL REPRESENTATION OF THE PDFS $f_i^{(\text{net}), (k)}$



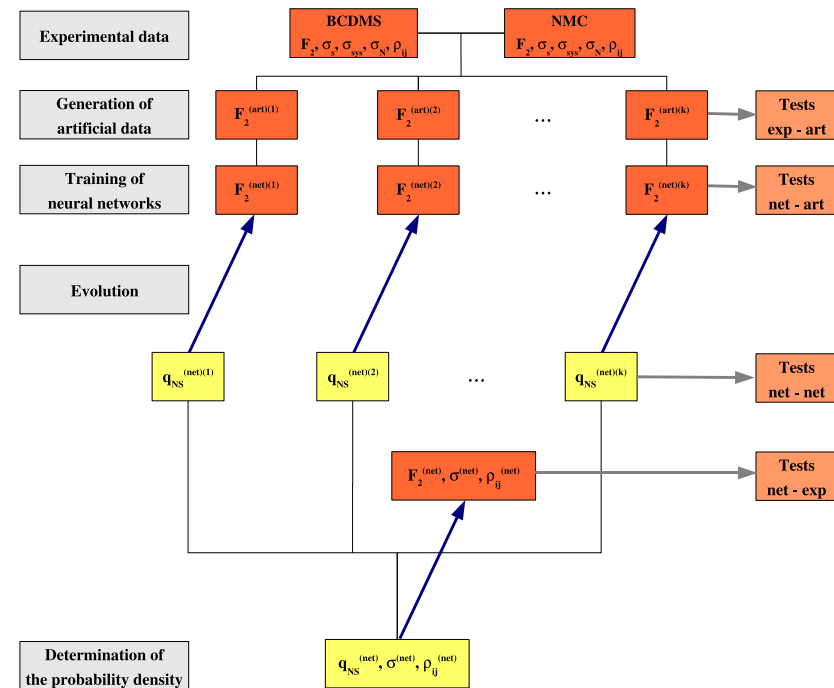
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- TRAIN A NEURAL NET FOR EACH PDF ON EACH REPLICA
 \Rightarrow NEURAL REPRESENTATION OF THE PDFS $f_i^{(\text{net}), (k)}$
- THE SET OF NEURAL NETS IS A REPRESENTATION OF THE PROBABILITY DENSITY

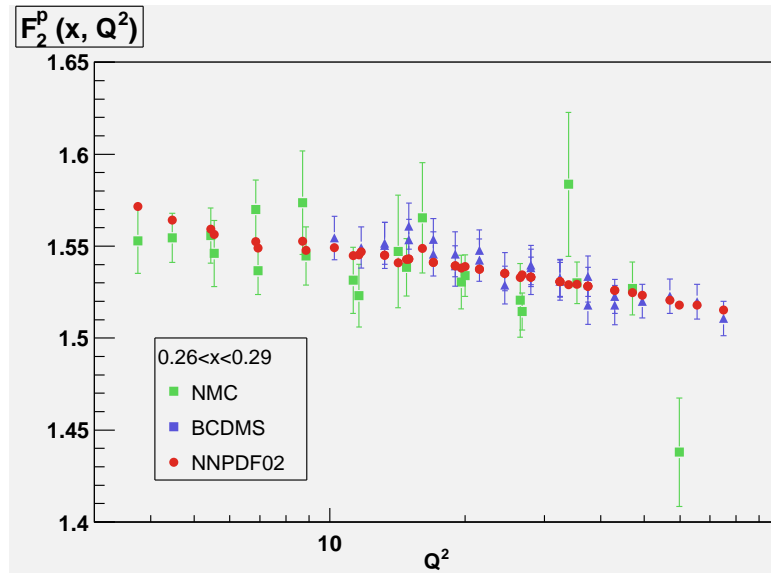
$$\langle \sigma [f_i] \rangle = \frac{1}{N_{rep}} \sum_{k=1}^{N_{rep}} \sigma [f_i^{(\text{net}) (k)}]$$



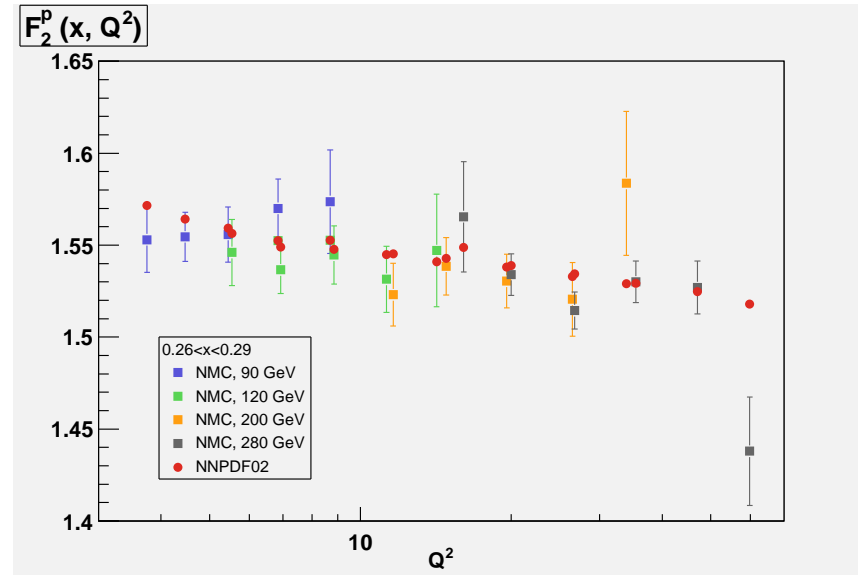
NEURAL HANDLING OF INCOMPATIBLE DATA

SOME NMC DATA ARE INCOMPATIBLE WITH OTHER DATA

Blow-up of proton data/nets



NMC proton data/nets

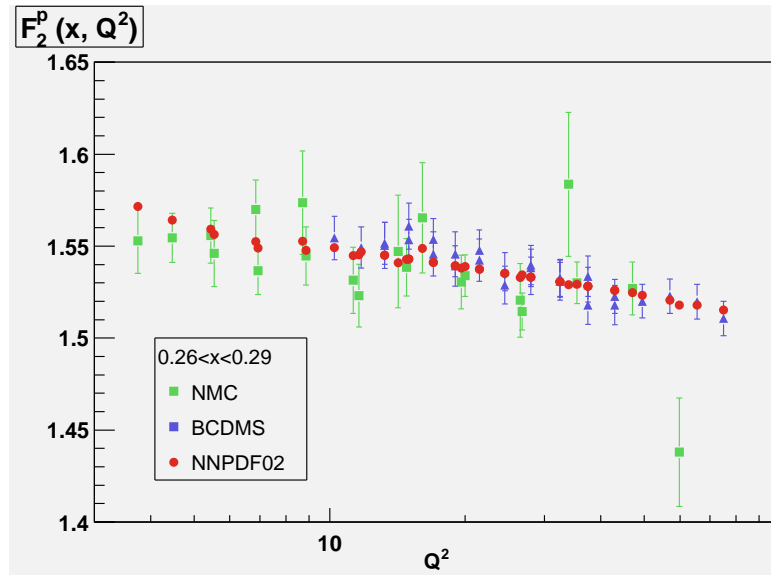


NEURAL NET DISCARDS INCONSISTENT DATA & PROVIDES GOOD FIT TO THE REST

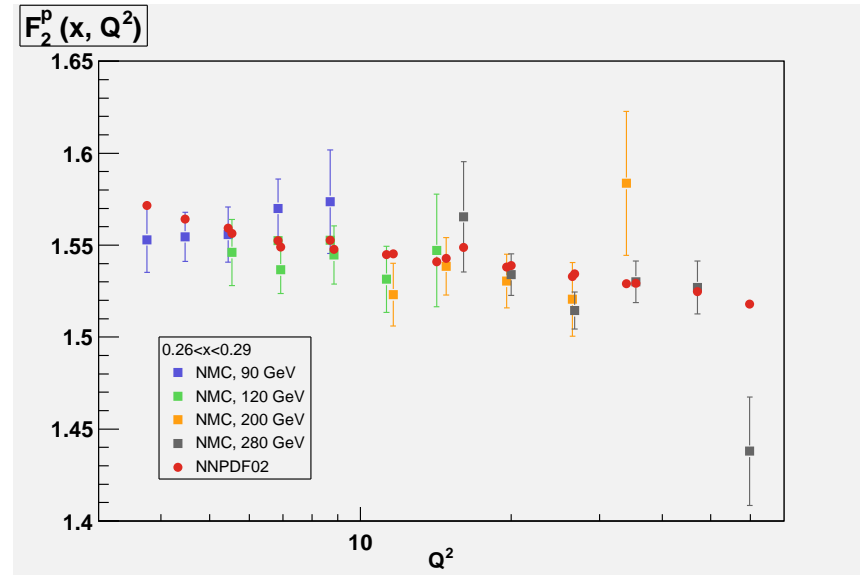
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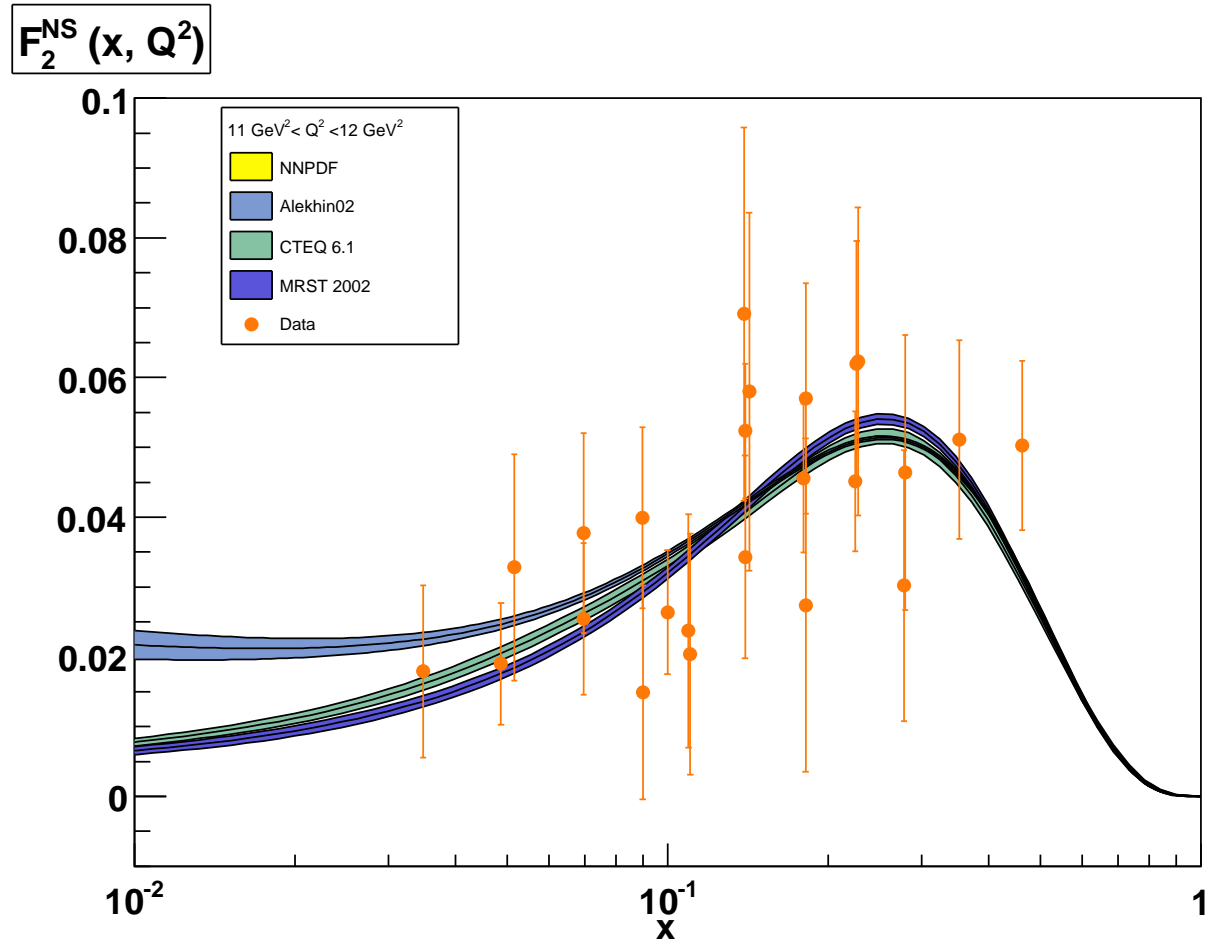
- NEURAL NET IS REDUNDANT
- SMOOTHNESS DECREASES DURING THE FIT
- STOPPING CRITERION: (NOT MINIMUM χ^2)
DIVIDE DATA IN TWO SETS, USE ONE SET FOR FITTING, STOP FIT WHEN FIT TO OTHER SET STARTS DETERIORATING

NEURAL PARTONS:

THE NONSINGLET CASE

the NNPDF collaboration: L. Del Debbio, S.F. J. Latorre, A. Piccione, J. Rojo

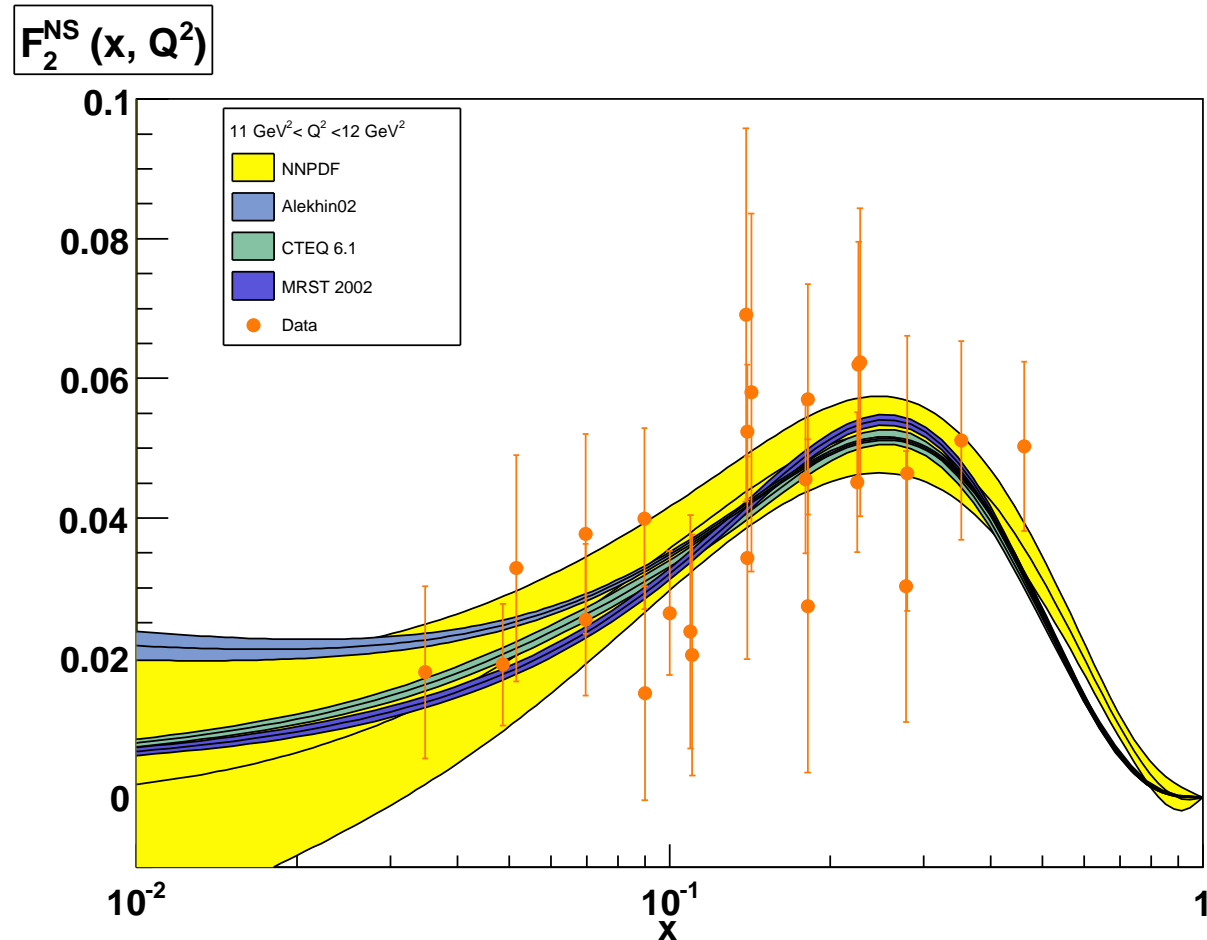
FULL NLO, NNLO FITS TO ALL AVAILABLE DATA



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the NNPDF collaboration: L. Del Debbio, S.F. J. Latorre, A. Piccione, J. Rojo
FULL NLO, NNLO FITS TO ALL AVAILABLE DATA



- FITS WITH FIXED FUNCTIONAL FORM SUBSTANTIALLY UNDERESTIMATE ERROR (ESPECIALLY ON EXTRAPOLATION)
- NO EVIDENCE FOR “WELL-KNOWN” SMALL x RISE OF NONSINGLET

CONCLUSION

IN SPIN PHYSICS,
ONE HAS TO BE VERY CLEVER
TO GET AN ANSWER