QCD and Baryon Polarization Lecture 2: Spin-momentum correlations in the nucleon

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Outline of lectures

- 1. Introduction; collinear and TMD nucleon structure
- 2. Spin-momentum correlations in the nucleon in terms of TMD PDFs and collinear twist-3 multiparton correlations
- 3. Hadronization: collinear and TMD fragmentation functions, collinear twist-3 correlations; dihadron FFs; different hadronization mechanisms/pictures
- 4. Hadron structure and hadronization: Sea quarks/non-valence quarks
- 5. Hyperon and heavy flavor baryon polarization I
- 6. Hyperon and heavy flavor baryon polarization II



Spin-momentum correlations: 1976 discovery in p+p collisions

Argonne \sqrt{s} =4.9 GeV



W.H. Dragoset et al., PRL36, 929 (1976)

Charged pions produced preferentially on one or the other side with respect to the transversely polarized beam direction—by up to 40%!!

Had to wait more than a decade for the birth of a new subfield in order to explore the possibilities . . .



$$x_F = 2 p_{long} / \sqrt{s}$$



Transverse-momentum-dependent distributions and single-spin asymmetries



1990: D.W. Sivers departs from traditional collinear factorization assumption in pQCD and proposes correlation between the *intrinsic transverse motion* of the quarks and gluons and the proton's spin



Spin and momenta of quarks and/or bound states

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First quark distribution function describing a spin-momentum correlation in the proton

 $s \cdot (p_1 \times p_2)$



Fig. 1

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New frontier! Quark **dynamics** inside QCD bound states, and in their formation process

$$s \cdot (p_1 \times p_2)$$

Spin and momenta of quarks and/or bound states

-0.2

-0.4

D.W. Sivers

PRD41, 83 (1990)

Transverse-momentum-dependent PDFs



Transverse-momentum-dependent PDFs





TMD PDFs predict specific angular dependences of cross sections

But pQCD can't make ab initio predictions of *magnitude* of effects—TMD PDFs are nonperturbative and need to be constrained from experimental data



General expression for angular dependence of *unpolarized* Drell-Yan (qqbar \rightarrow dileptons): $\left(\frac{1}{\sigma}\right)\left(\frac{d\sigma}{d\Omega}\right) = \left[\frac{3}{4\pi}\right]\left[1 + \lambda\cos^2\theta + \mu\sin 2\theta\cos\phi + \frac{\nu}{2}\sin^2\theta\cos 2\phi\right]$

Boer-Mulders TMD PDF leads to $\cos 2\phi$ dependence. Correlation between transverse spin of (anti-)quark and its own transverse momentum (spin-orbit) $(h_{\perp}^{\perp})(\bar{h}_{\perp}^{\perp})$





Drell-Yan angular-dependent cross section in terms of structure functions

+ Schlegel,

5 (2009)

$$\begin{aligned} \frac{d\sigma}{d^{4}qd\Omega} &= \frac{\alpha_{em}^{2}}{Fq^{2}} \times \\ &\left\{ \left((1 + \cos^{2}\theta) F_{UU}^{1} + (1 - \cos^{2}\theta) F_{UU}^{2} + \sin 2\theta \cos \phi F_{UU}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{UU}^{\cos 2\phi} \right) \\ &+ S_{aL} \left(\sin 2\theta \sin \phi F_{UU}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{UU}^{\sin 2\phi} \right) \\ &+ S_{bL} \left(\sin 2\theta \sin \phi F_{UU}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{UU}^{\sin 2\phi} \right) \\ &+ S_{bL} \left(\sin 2\theta \sin \phi F_{UU}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{UU}^{\sin 2\phi} \right) \\ &+ \left| \vec{S}_{aT} \right| \left[\sin \phi_{a} \left((1 + \cos^{2}\theta) F_{IU}^{1} + (1 - \cos^{2}\theta) F_{U}^{2} + \sin 2\theta \cos \phi F_{TU}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{TU}^{\cos 2\phi} \right) \\ &+ \cos \phi_{a} \left(\sin 2\theta \sin \phi F_{UT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{UT}^{\sin 2\phi} \right) \right] \\ &+ \left| \vec{S}_{aT} \right| \left[\sin \phi_{b} \left((1 + \cos^{2}\theta) F_{IU}^{1} + (1 - \cos^{2}\theta) F_{UT}^{2} + \sin 2\theta \cos \phi F_{UT}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{UT}^{\cos 2\phi} \right) \\ &+ \cos \phi_{b} \left(\sin 2\theta \sin \phi F_{UT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{UT}^{\sin 2\phi} \right) \right] \\ &+ S_{aL} S_{bL} \left((1 + \cos^{2}\theta) F_{LL}^{1} + (1 - \cos^{2}\theta) F_{LL}^{2} + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{LL}^{\cos 2\phi} \right) \\ &+ \sin \phi_{b} \left(\sin 2\theta \sin \phi F_{LT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{LT}^{\sin 2\phi} \right) \right] \\ &+ \left| \vec{S}_{aT} \right| \left[\cos \phi_{b} \left((1 + \cos^{2}\theta) F_{LL}^{1} + (1 - \cos^{2}\theta) F_{LL}^{2} + \sin 2\theta \cos \phi F_{LL}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{LL}^{\cos 2\phi} \right) \\ &+ \sin \phi_{a} \left(\sin 2\theta \sin \phi F_{LT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{LT}^{\sin 2\phi} \right) \right] \\ &+ \left| \vec{S}_{aT} \right| \left| \vec{S}_{bL} \right| \left[\cos \phi_{a} \left((1 + \cos^{2}\theta) F_{LL}^{1} + (1 - \cos^{2}\theta) F_{LL}^{2} + \sin 2\theta \cos \phi F_{TL}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{TL}^{\cos 2\phi} \right) \\ &+ \sin \phi_{a} \left(\sin 2\theta \sin \phi F_{LT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{TL}^{\sin 2\phi} \right) \right] \\ &+ \left| \vec{S}_{aT} \right| \left| \vec{S}_{bL} \right| \left[\cos \phi_{a} \left((1 + \cos^{2}\theta) F_{TL}^{1} + (1 - \cos^{2}\theta) F_{TL}^{2} + \sin 2\theta \cos \phi F_{TL}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{TL}^{\cos 2\phi} \right) \\ &+ \sin \phi_{a} \left(\sin 2\theta \sin \phi F_{TT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{TL}^{\sin 2\phi} \right) \right] \\ &+ \cos \phi_{a} \left(\sin 2\theta \sin \phi F_{TT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{TT}^{\sin 2\phi} \right) \\ &+ \sin (\phi_{a} - \phi_{b}) \left((1 + \cos^{2}\theta) F_{TT}^{1} + (1 - \cos^{2}\theta) F_{TT}^{2} + \sin 2\theta \cos \phi F_{TT}^{\cos \phi} + \sin^{2}\theta \cos 2\phi F_{TT}^{\cos 2\phi} \right) \\ &+ \sin (\phi_{a} - \phi_{b}) \left(\sin 2\theta \sin \phi F_{TT}^{\sin \phi} + \sin^{2}\theta \sin 2\phi F_{TT}^{\sin 2\phi} \right) \right] \right\} .$$

TMD PDFs: Some properties

- PDFs involving transversely polarized quarks are chiralodd—can only be observed experimentally in conjunction with a second chiral-odd function
 - Another PDF or a FF
 - Transversity, Boer-Mulders, pretzelosity, one of the worm gears (all the 'h' PDFs)
- TMD PDFs involving a single spin vector are "naïve-time-reversal-odd," i.e. PT-odd
 - "Sivers" and "Boer-Mulders"



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• What about TMD PDFs for gluons?

- Gluons massless—can't have helicity-flip states, i.e. transversely polarized gluons
- However, can have TMD PDFs for linearly polarized gluons, similar to linearly polarized photons—relatively less explored than TMD PDFs for quarks
- Unpolarized and longitudinally polarized TMD PDFs no problem for gluons



Transverse single-spin asymmetries

- General form for transverse single-spin asymmetries: $S \cdot (p_1 \times p_2)$
 - Collinear momenta would produce no effect
 - Thus importance of transverse momentum dependence



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 - Thus importance of transverse momentum dependence
- Spin could be of initial proton, struck quark, fragmenting quark, produced hadron
- Possible momentum vectors include initial proton momentum, final-state particle or jet momentum, k_T of parton within proton, j_T of final-state particle with respect to jet axis
- Lots of combinations possible!



Spin-spin and spin-momentum correlations in QCD bound states





Boer-Mulders TMD PDF ×Collins TMD FF asymmetry from semi-inclusive DIS



- Boer-Mulders TMD PDF correlation between quark transverse spin and its own transverse momentum. Chiral-odd. Zero if orbital angular momentum zero.
- Chiral-odd \rightarrow need another chiral-odd function to measure it. Here the Collins TMD FF
- Clearly nonzero for positive and negative hadrons
- Also measured by HERMES PRD87, 012010 (2013) (see backup)



Sivers TMD PDF asymmetry from semi-inclusive DIS

Charged pions, HERMES (e+p) and COMPASS (μ +p)



- Sivers TMD PDF correlation between proton transverse spin and quark transverse momentum. Zero if orbital angular momentum zero.
- Clearly nonzero for positive pions. Cancellations between up and down quarks lead to smaller negative pion asymmetries.



Example: Fit of Sivers TMD PDF

Boglione, D'Alesio, Flore, Gonzalez-Hernandez, JHEP 07, 148 (2018)



- Recall: Sivers TMD PDF describes correlation between tranverse spin of proton and intrinsic motion of quark.
- Fit based on semi-inclusive DIS data on transversely polarized targets from COMPASS and HERMES.
- 220 points (compare to thousands for unpolarized, collinear PDF extraction).
- <u>Clear</u> opposite spin-momentum correlation for up versus down quarks.



Different symmetry properties for different spin-momentum correlations

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Different symmetry properties for different spin-momentum correlations

- Some transverse-momentum-dependent quark distribution functions odd under a parity- and time-reversal (PT) transformation
- In 1993, after original 1990 paper by D.W. Sivers, J.C. Collins claimed such functions must vanish
- Only realized in 2002 by Brodsky, Hwang, and Schmidt that could be nonvanishing if *phase interference effects due to color interactions* present



Modified universality of PT-odd correlations: Color in action!

Deep-inelastic lepton-nucleon scattering: Final-state color exchange



Quark-antiquark annihilation to leptons: Initial-state color exchange





Figures by J.D. Osborn

Modified universality of PT-odd correlations: Color in action!

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Figures by J.D. Osborn

Deep-inelastic lepton-nucleon scattering: Final-state color exchange

incoming scattered incoming proton electron proton electron produced remnant positron scattering scattering scattered antiquark quark guark scattering quark produced proton proton incoming incoming electron remnant remnant proton proton

Opposite sign for PT-odd spin-momentum correlations in the proton measured in these two processes: process-dependent! (Collins 2002)



Modified universality: Initial experimental hints



First measurements by STAR at the Relativistic Heavy Ion Collider and COMPASS at CERN suggestive of predicted sign change in colorannihilation processes compared to quark knock-out by a lepton. More statistics forthcoming . . .

High-energy QCD: Thinking in terms of individual partons

- PDFs are *single-parton* functions in *single* nucleons
 - Or in nuclei, but typically still think of partons in individual nucleons within nucleus

• Can we go beyond this single-parton picture while staying in the hard (short-distance) limit of perturbative QCD?



An alternative approach to describing the large single-spin asymmetries: Higher-twist multiparton correlations

• Extend our ideas about (single-parton) PDFs to correlation functions that can't be associated with a single parton



An alternative approach to describing the large single-spin asymmetries: Higher-twist multiparton correlations

- Extend our ideas about (single-parton) PDFs to correlation functions that can't be associated with a single parton
- Non-perturbative structure → matrix elements involving the quantum mechanical *interference* between scattering off of a (quark+gluon) and scattering off of a single quark (of the same flavor and at the same x)
 - Can also have interference between (gluon+gluon) and single gluon
 - No explicit dependence on partonic transverse momentum
 - Efremov+Teryaev 1981, 84; Qiu+Sterman 1991, 98



Beware: Two common usages of the term "twist"

- Formal definition of twist: "mass dimension minus spin" of the operator in a matrix element within the Operator Product Expansion
 - "Leading twist" is twist-2
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 - Twist-n matrix element carries a factor of $1/Q^{(n-2)}$
- But *observables* with measurable contributions from terms suppressed by a factor of $1/Q^{(n-2)}$ often referred to as sensitive to "twist-n" contributions
 - Never measure a matrix element, only matrix elements squared!
 - To get 1/Q term describing an *observable*, need interference term in the square modulus:
 - $A = order 1 + order 1/Q + order 1/Q^2 + \dots$
 - $|A|^2 = |order 1|^2 + |order 1/Q|^2 + (order 1)(order 1/Q)^* + (order 1)^*(order 1/Q) + \dots$
 - So twist-3 term in matrix element times *twist-2* term gives 1/Q
 - Square modulus of *twist-3* term gives $1/Q^2$, sometimes referred to as "twist-4"



Transverse single-spin asymmetries provide <u>new</u> information on hadron structure

- <u>Leading</u> contribution to transverse single-spin asymmetries comes from *either*:
 - Convolution of two twist-2 *transverse-momentumdependent* parton distribution functions and/or fragmentation functions, or . . .
 - Convolution of one twist-2 collinear PDF or fragmentation function and one twist-3 (collinear) *multiparton correlation* matrix element



Transverse-momentum-dependent functions and twist-3 multiparton correlators

- Twist-3 (collinear) multiparton correlators are related to k_Tmoments of (twist-2) TMD PDFs and fragmentation functions
 - NPB667, 201 (2003); PRL97, 082002 (2006)



Transverse-momentum-dependent functions and twist-3 multiparton correlators

- Twist-3 (collinear) multiparton correlators are related to k_Tmoments of (twist-2) TMD pdfs and fragmentation functions
 - NPB667, 201 (2003); PRL97, 082002 (2006)
- To directly constrain TMD functions with experimental data, need *two* scales
 - Hard momentum
 - Observable sensitive to parton intrinsic momentum
 - In semi-inclusive DIS, have (large) Q^2 from scattered lepton and (small) p_T of measured hadron
 - In qqbar \rightarrow dileptons, have (large) invariant mass and (small) p_T of lepton pair
 - Recall: Original p+p→pion+X asymmetries only measured a single scale



Transverse single-spin asymmetries in p+p across energies BRAHMS **FNAL** RHIC **BNL** ANL √s=19.4 GeV √s=62.4 GeV $\sqrt{s}=6.6 \text{ GeV}$ √s=4.9 GeV 60 60 60 60 PRL 36, 929 (1976) PRL 101, 042001 (2008) PRD 65, 092008 (2002) PLB 261, 201 (1991) PLB 264, 462 (1991) BRAHMS 40 40 40 40 π 20 20 20 20 A_N (%) $\circ \pi$ 0 0 0 0 Ō 0 Ο \circ þ Ο -20 -20 -20 Ο -20 0 φĄ -40 -40 -40 -40 -60 -60 -60 -60 0.2 0.2 0.4 0.6 0.8 0.2 0.4 0.6 0.8 0.2 0.4 0.6 0.8 0.4 0.6 0.8 X_{F} X_F X_{F} X_F

$$x_F = 2 p_{long} / \sqrt{s}$$





Transverse single-spin asymmetries in p+p across energies



$$x_F = 2p_{long} / \sqrt{s}$$



right



Transverse single-spin asymmetry in $p+p \rightarrow hadron + X$: Only measure one momentum scale

- For high enough p_T of produced hadron (>1-2 GeV) have hard scale, so can apply perturbative calculations
 - Clear nonzero asymmetries out to 8 GeV \rightarrow Q² ~ 64 GeV²
- Can have contributions from initialstate and final-state effects (spinmomentum correlations in the proton and in hadronization)
- Inclusive measurement—don't measure the combination of a hard plus a nonperturbative momentum scale required to (directly) apply TMD framework in pQCD calculations





First global fit of transverse-single spin asymmetries using both TMD functions and twist-3 correlators

- Fit
 - semi-inclusive DIS,
 - Drell-Yan,
 - e+e- annihilation, and
 - − $p+p \rightarrow h+X$
- *Simultaneously* extract nonperturbative spin-momentum correlations as TMD PDFs and twist-3 correlators in the proton and in hadronization



Jefferson Lab Angular Momentum Collaboration, arXiv:2002.08384





Summary: Lecture 2

- Striking spin-momentum correlations (up to tens of percent!) in collisions involving transversely polarized protons were first observed in the 1970s
 - Opened up new field of parton dynamics within the proton





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 - Opened up new field of parton *dynamics* within the proton
- Since the 1990s two theoretical frameworks within perturbative QCD have been developed to characterize these nonperturbative spin-momentum correlations
 - Transverse-momentum-dependent PDFs
 - Collinear twist-3 multiparton correlators in the proton
 - Both of these encode additional information beyond traditional collinear PDFs
 - Multiparton correlators are related to k_T moments of TMD PDFs





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- Clear experimental evidence from semi-inclusive DIS, Drell-Yan, and p+p that two spin-momentum correlations described by TMD PDFs are nonzero in nature, analogous to the two spin-orbit couplings in the hydrogen atom
 - Sivers, describing "spin-orbit" coupling between proton spin and quark orbital motion
 - Boer-Mulders, describing "spin-orbit" coupling between *quark* spin and quark orbital motion
 - Others remain to be measured with higher precision





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 - Others remain to be measured with higher precision
- Food for thought Should we expect isospin symmetry between the proton and neutron when investigating parton *dynamics*??







Drell-Yan vector diagram for TMD PDF measurements

- One plane formed by produced lepton pair; other by incoming quark and antiquark
 - Incoming quark and antiquark cannot be collinear, otherwise can't define a plane!



Dilepton rest frame



Semi-inclusive DIS vector diagram for TMD PDF measurements





Effects persist up to transverse momenta of 7(!) GeV/c at $\sqrt{s}=500$ GeV



- Can try to interpret these non-perturbative effects within the framework of perturbative QCD.
- Haven't yet
 disentangled all the
 possible contributing
 effects to the (messy)
 process of p+p to
 pions



HERMES Sivers for pions





Sivers asymmetry in SIDIS for kaons and protons



Nonzero for positive kaons and protons, hints for negative kaons and antiprotons. Identified particles help give *flavor separation* for Sivers TMD PDF



Boer-Mulders ×Collins asymmetry from SIDIS





Boer-Mulders ×Boer-Mulders asymmetry from Drell-Yan E866, PRL 99, 082301 (2007);

PRL 102, 182001 (2009)



Huge $\cos 2\phi$ dependence in pion-induced Drell-Yan

- Significantly reduced in proton-induced Drell-Yan
- Suggests sea quark transverse spinmomentum correlations small?

Boer - Mulders function h_1^{\perp}

 $v(\pi W \rightarrow \mu^* \mu^* X) \sim [valence h_1^{\perp}(\pi)] * [valence h_1^{\perp}(p)]$ $v(pd \rightarrow \mu + \mu - X) \sim [valence h_1^{\perp}(p)] * [sea h_1^{\perp}(p)]$

Other TMD PDF measurements in SIDIS



Other TMD PDF measurements in SIDIS





35

<u>Inclusive</u> hadron transverse single-spin asymmetries in e+p

 Striking enhancement if measure scattered electron



HERMES, PLB728, 183 (2014)



Twist-3 multiparton correlations to interpret inclusive A_N *data from RHIC*



Find dominant contribution from twist-3 correlation in hadronization



Midrapidity direct photon transverse single-spin asymmetry in p+p

Will improve constraints on twist-3 trigluon spinmomentum correlator



