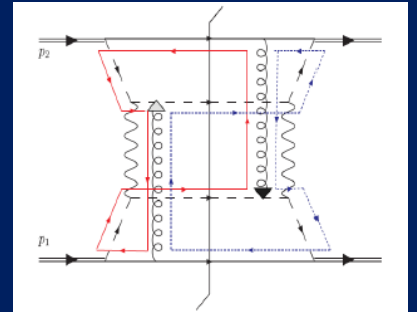
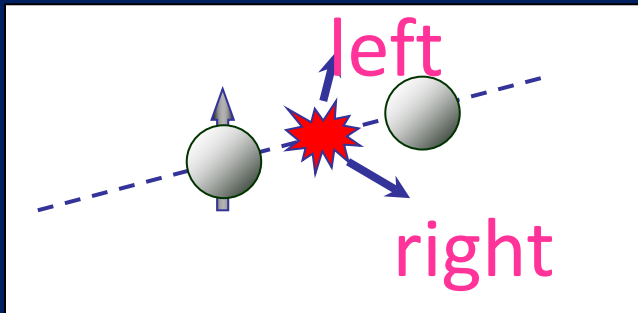


Spin-Momentum Correlations, Aharonov-Bohm, and Color Entanglement in Quantum Chromodynamics

*Christine A. Aidala
University of Michigan*



$$\psi(x)|P\rangle = e^{ig \int_x^{x'} ds_\mu A^\mu} \psi(x')|P\rangle$$

UCLA
October 17, 2016

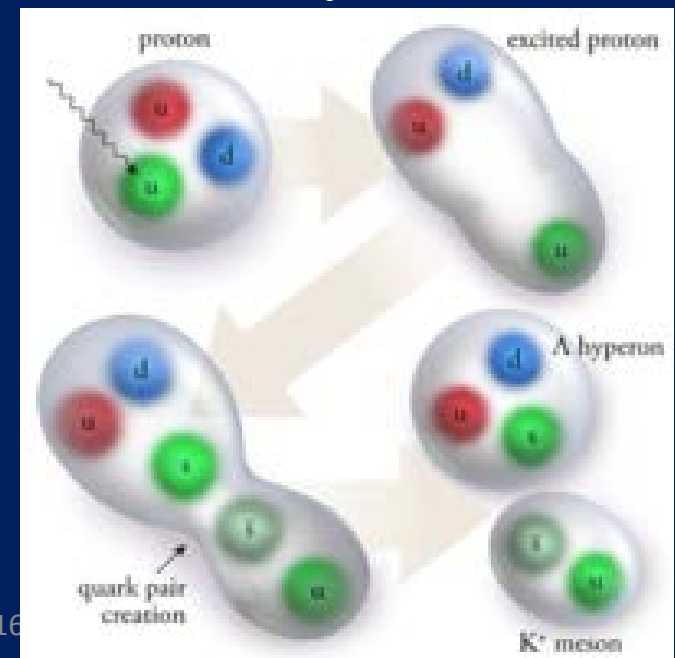


Theory of strong interactions: Quantum Chromodynamics

- Fundamental field theory in hand since the early 1970s—BUT . . .
- Quark and gluon degrees of freedom in the theory cannot be observed or manipulated directly in experiment!

Color confinement—quarks and gluons are confined to color-neutral bound states

CLAS, PRL 113, 152004 (2014)
PRL Editor's Choice Oct. 2014

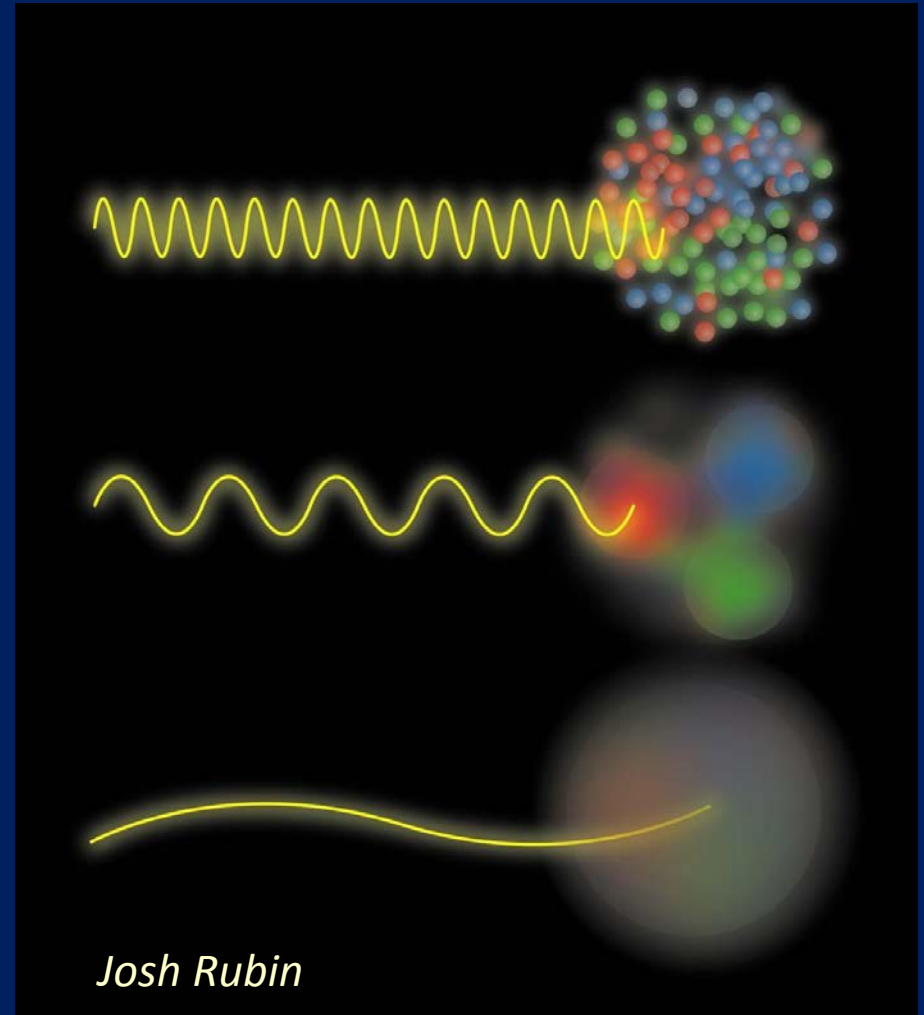


C. Aidala, U. of Michigan, Oct. 17, 2016



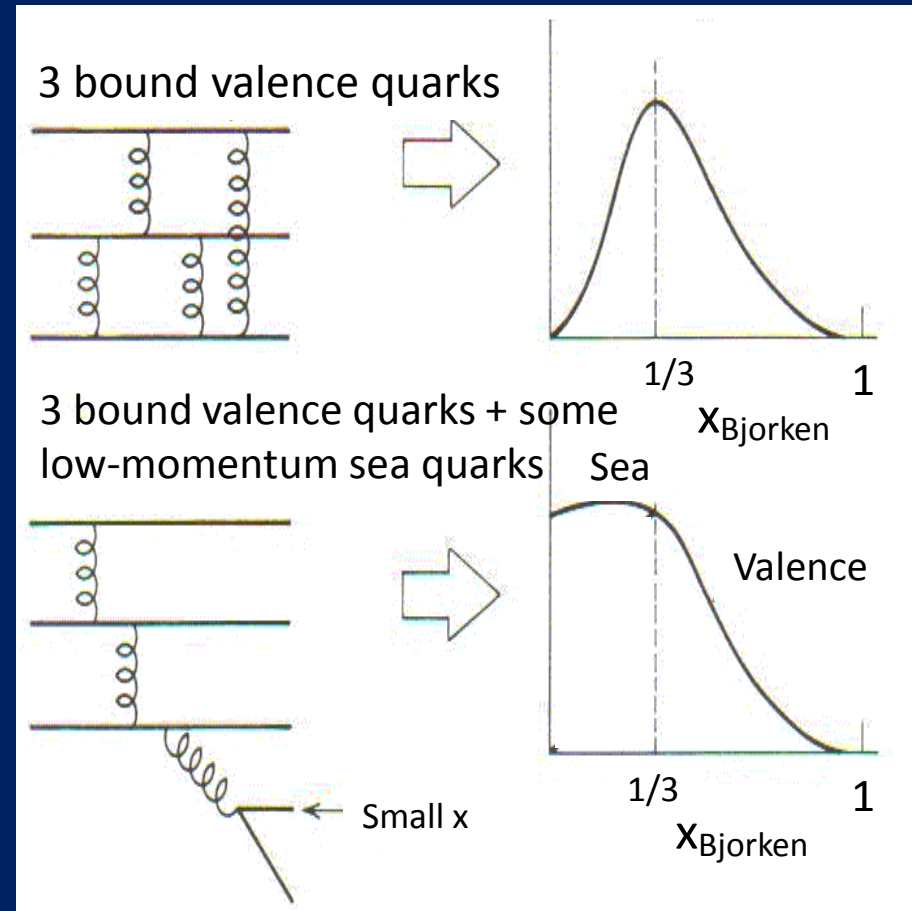
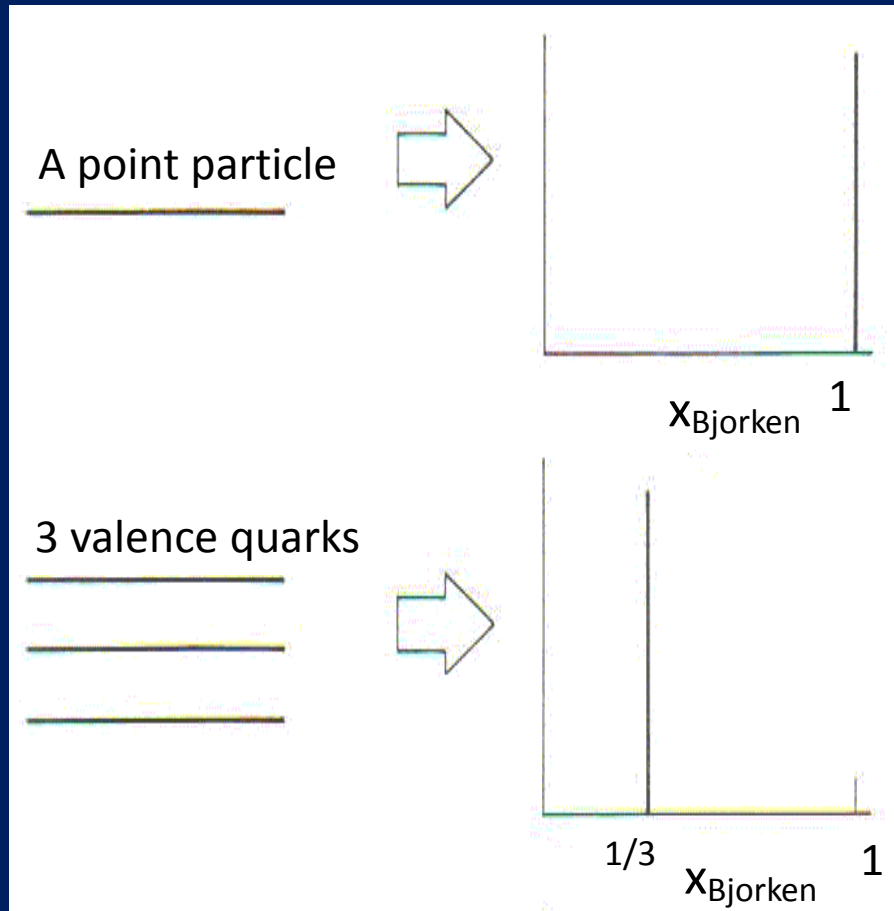
Quark and gluon structure of the nucleon

- Probing the proton at different energy scales offers information on different aspects of its internal structure



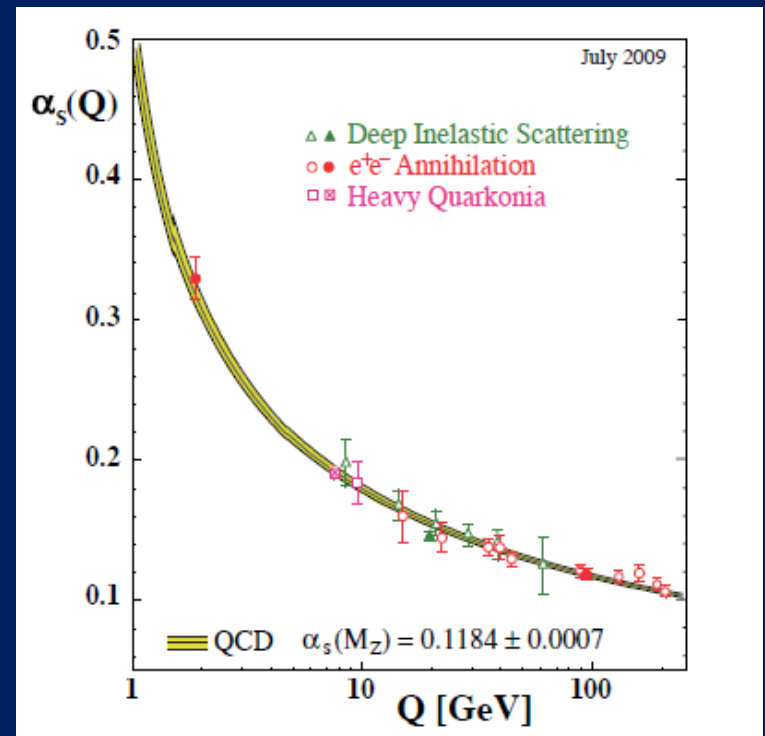
Parton distribution functions inside a nucleon: The language we've developed (so far!)

What momentum fraction would the scattering particle carry if the proton were made of ...



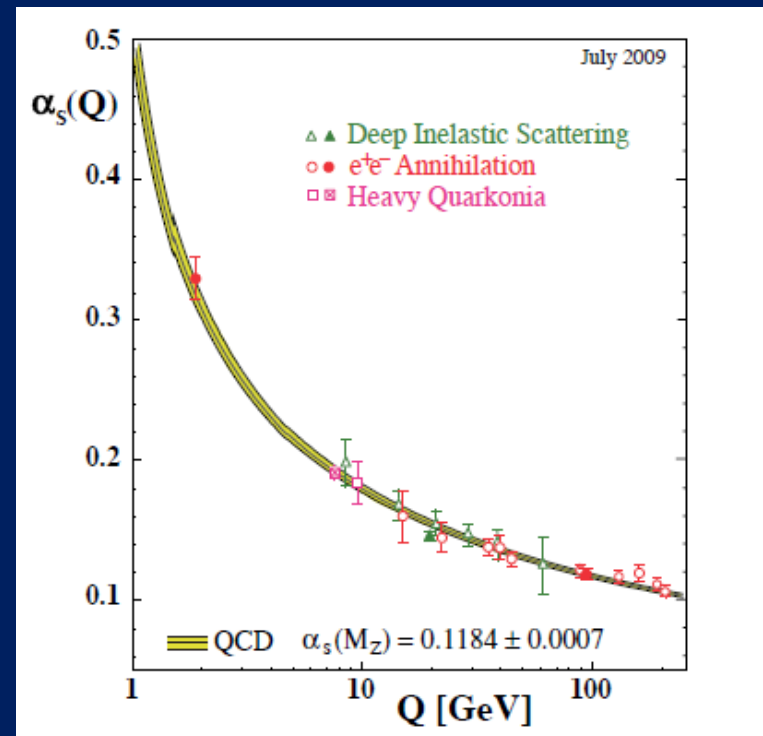
Perturbative QCD

- Take advantage of running of the strong coupling constant with energy (*asymptotic freedom*)—weak coupling at high energies (short distances)
- Perturbative expansion as in quantum electrodynamics (but many more diagrams due to gluon self-coupling!!)



Perturbative QCD

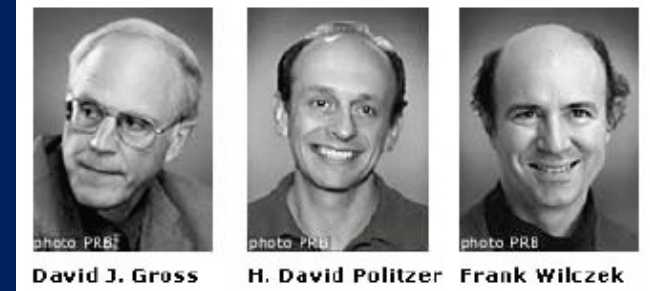
- Take advantage of running of the strong coupling constant with energy (*asymptotic freedom*)—weak coupling at high energies (short distances)
- Perturbative expansion as in quantum electrodynamics (but many more diagrams due to gluon self-coupling!!)



Provides a rigorous way of relating the fundamental field theory to a variety of physical observables!

QCD: How far have we come?

- QCD is challenging!!
 - Three-decade period after initial birth of QCD dedicated to “discovery and development”
- Symbolic closure: Nobel prize 2004 - Gross, Politzer, Wilczek for asymptotic freedom



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- QCD is challenging!!
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*Now early years of second phase:
quantitative QCD!*

Advancing into the era of quantitative QCD: Theory has been forging ahead

- In perturbative QCD, since 1990s starting to consider detailed internal *dynamics* that parts with traditional parton model ways of looking at hadrons—and perform phenomenological calculations using these new ideas/tools!

E.g.:

- Various *resummation* techniques
- *Non-linear* evolution at small momentum fractions
- *Spin-spin* and *spin-momentum* correlations in QCD bound states
- *Spatial* distributions of partons in QCD bound states



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 - *Spatial* distributions of partons in QCD bound states
- Nonperturbative methods:
 - Lattice QCD less and less limited by computing resources—since 2010 starting to perform calculations at the physical pion mass (after 36 years!). Plus recent new ideas on how to calculate previously intractable quantities.
 - AdS/CFT “gauge-string duality” an exciting recent development as first fundamentally new handle to try to tackle QCD in decades!

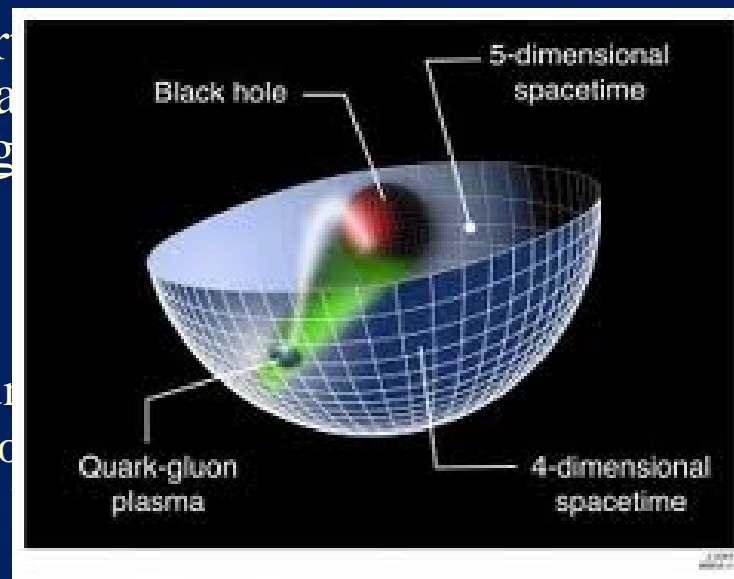


Advancing into the era of quantitative QCD: Theory has been forging ahead

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E.g.:

- Various *resummation* techniques
- *Non-linear* evolution at small momentum
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- *Spatial* distributions of partons in QCD



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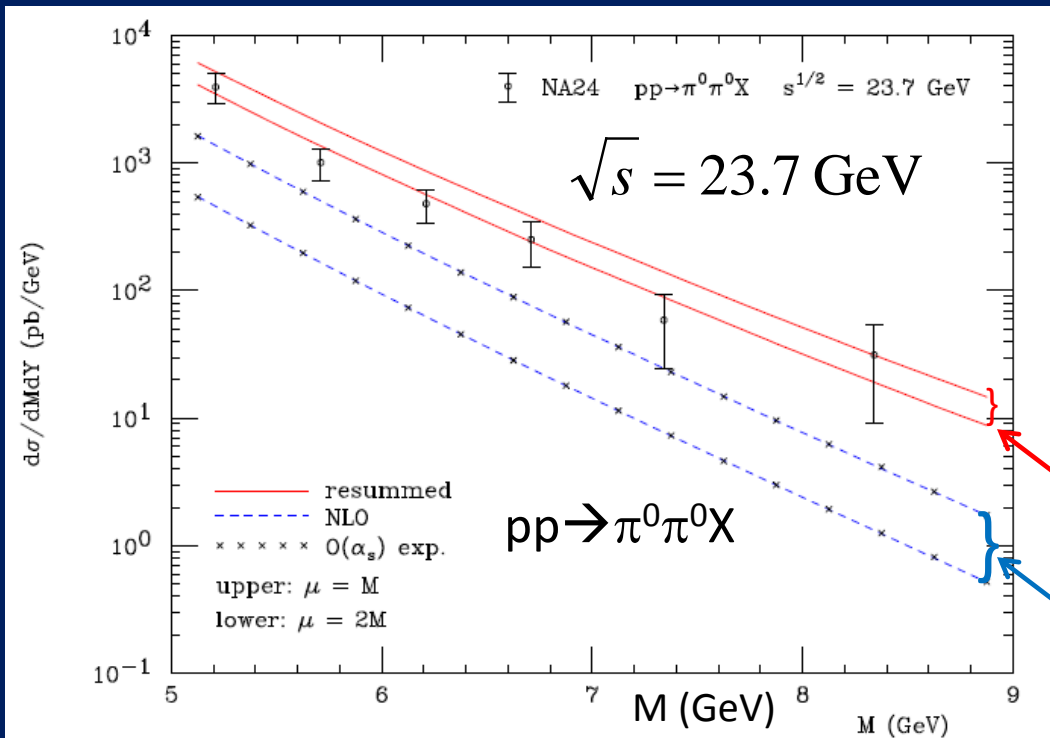
Effective field theories

- QCD exhibits different behavior at different scales—effective field theories are useful approximations within these different regimes
 - Color Glass Condensate – high energies, high densities
 - Soft-Collinear Effective Theory – new insights into performing complicated perturbative calculations very quickly
 - Heavy Quark Effective Theory, Non-Relativistic QCD, . . .
 - Many effective theories for nonperturbative QCD – chiral symmetry breaking, . . .



Example: “Threshold resummation”

Extending perturbative calculations to lower energies



For observables with two different scales, sum logs of their ratio to all orders in the strong coupling constant

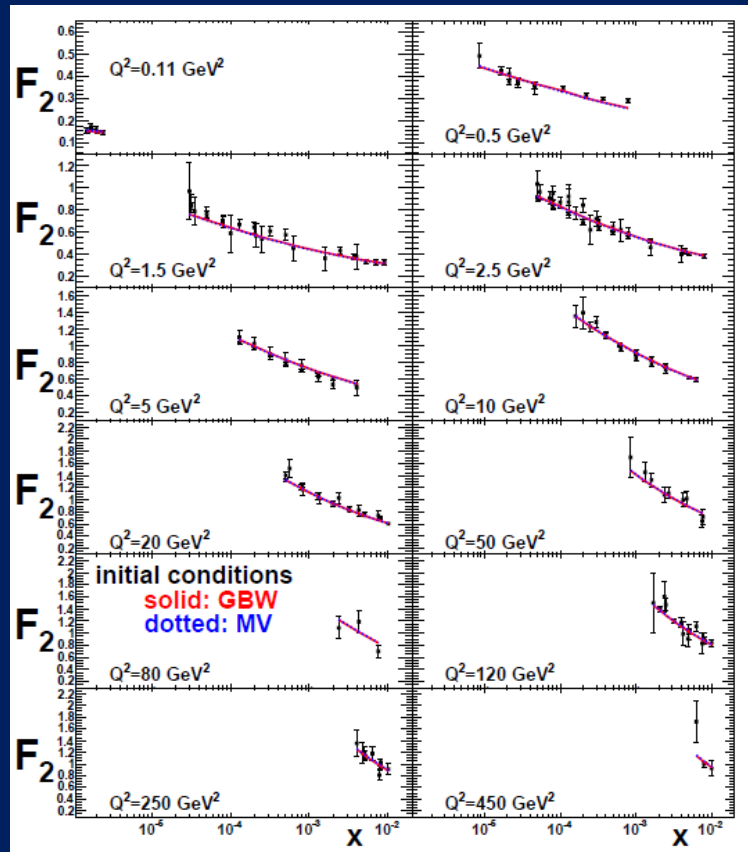
Next-to-leading-order in α_s + resum.

Next-to-leading-order in α_s

Almeida, Sterman, Vogelsang PRD80, 074016
(2009)



Example: Phenomenological applications of a non-linear gluon saturation regime



Fits to proton structure function data at low parton momentum fraction x .

Non-linear QCD meets data: A global analysis of lepton-proton scattering with running coupling BK evolution

Phys. Rev. D80, 034031 (2009)

Javier L. Albacete¹, Néstor Armesto², José Guilherme Milhano³ and Carlos A. Salgado²

Basic framework for non-linear QCD, in which gluon densities are so high that there's a non-negligible probability for two gluons to combine, developed ~1997-2001. But had to wait until “running coupling BK evolution” figured out in 2007 to compare directly to data!

Example: Spin-spin and spin-momentum correlations in QCD bound states

Unpolarized

$$f_1 = \text{circle with a dot}$$

Spin-spin correlations

$$g_{1L} = \text{circle with dot and right arrow} - \text{circle with dot and left arrow}$$

$$h_{1T} = \text{circle with dot and up arrow} - \text{circle with dot and down arrow}$$

$$g_{1T} = \text{circle with dot and up arrow} - \text{circle with dot and left arrow}$$

Spin-momentum correlations

$$S \cdot (p_1 \times p_2)$$

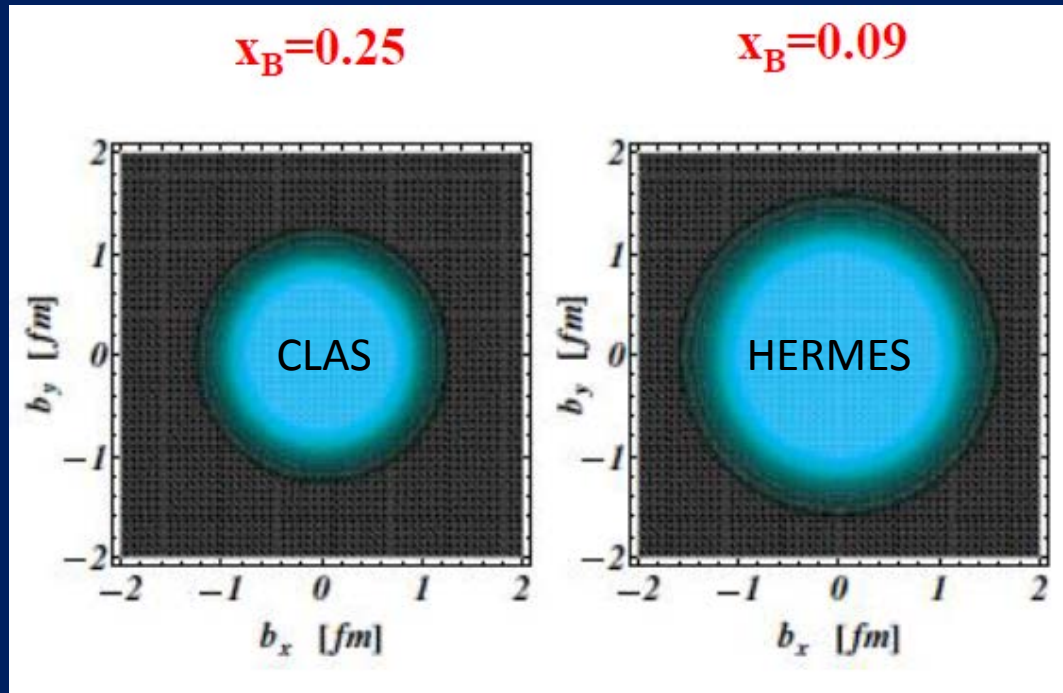
$$f_{1T}^\perp = \text{circle with dot and up arrow} - \text{circle with dot and down arrow}$$

$$h_1^\perp = \text{circle with dot and right arrow} - \text{circle with dot and left arrow}$$

$$h_{1L}^\perp = \text{circle with dot and up arrow} - \text{circle with dot and right arrow}$$

$$h_{1T}^\perp = \text{circle with dot and up arrow} - \text{circle with dot and up arrow}$$

Example: Exploring spatial distributions



Guidal, Moutarde,
Vanderhaeghen,
Rept. Prog. Phys. 76 (2013) 066202

Initial evidence that quarks carrying larger momentum fractions (25% vs. 9%) in the nucleon are distributed over a smaller volume in space

Spatial charge densities measured via deeply virtual Compton scattering

Example: Progress in lattice QCD

Recent progress in LQCD suggests the possibility to calculate the x -dependence of parton distributions

Slide from J.-C. Peng, Transversity 2014

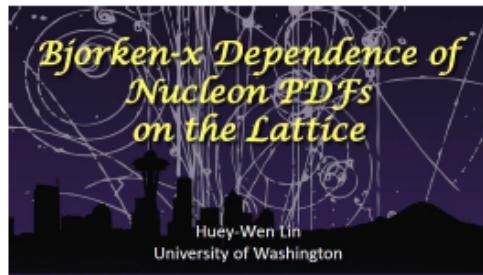
PRL 110, 262002 (2013)

PHYSICAL REVIEW LETTERS

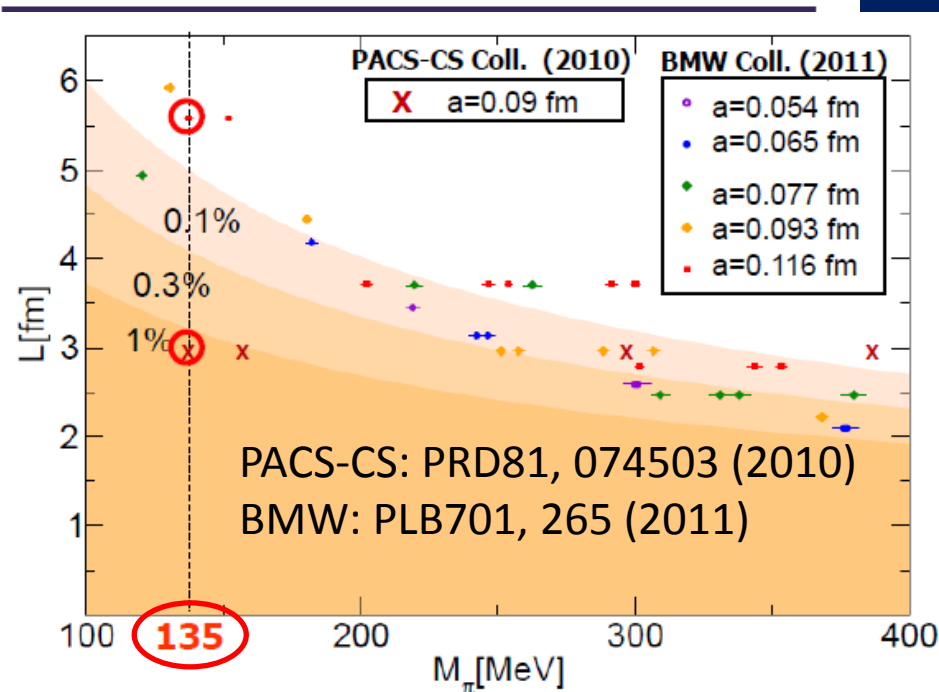
week ending
28 JUNE 2013

Parton Physics on a Euclidean Lattice

Xiangdong Ji^{1,2}



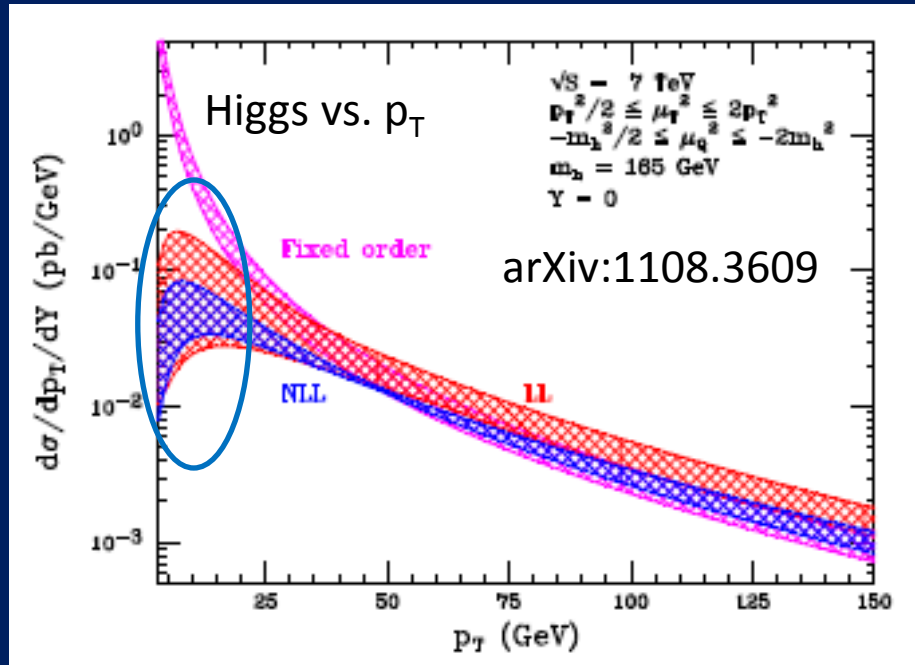
The x -dependence of the quark and antiquark transversity distributions can be calculated (not just their moments)



First calculations at physical pion mass 135 MeV

Figure from T. Hatsuda, PANIC 2011

Example: Effective field theories



TRANSVERSE MOMENTUM DISTRIBUTIONS FROM EFFECTIVE FIELD THEORY

Sonny Mantry*

University of Wisconsin at Madison
 Madison, WI 53706, USA
 mantry147@gmail.com

Frank Petriello

High Energy Physics Division, Argonne National Laboratory
 Argonne, IL 60439, USA

Department of Physics & Astronomy, Northwestern University
 Evanston, IL 60208, USA
 f-petriello@northwestern.edu

Soft Collinear Effective Theory

– Transverse momentum distribution for gluon+gluon → Higgs



Mapping out the partonic structure of the proton

What does the proton look like in terms of the quarks and gluons inside it?

- *Position*
- *Momentum*
- *Spin*



Mapping out the partonic structure of the proton

What does the proton look like in terms of the quarks and gluons inside it?

- *Position*
- *Momentum*
- *Spin*
- *Flavor*
- *Color*



Mapping out the partonic structure of the proton

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Vast majority of past four decades focused on *1-dimensional* momentum structure! Since 1990s starting to consider transverse components . . .

Mapping out the partonic structure of the proton

What does the proton look like in terms of the quarks and gluons inside it?

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- *Color*

Polarized protons first studied in 1980s. How angular momentum of quarks and gluons add up still not well understood!



Mapping out the partonic structure of the proton

What does the proton look like in terms of the quarks and gluons inside it?

- *Position*
- *Momentum*
- *Spin*
- *Flavor*
- *Color*

Good measurements of flavor distributions in valence region. Flavor structure at lower momentum fractions still yielding surprises!



Mapping out the partonic structure of the proton

What does the proton look like in terms of the quarks and gluons inside it?

- *Position*
- *Momentum*
- *Spin*
- *Flavor*
- *Color*

Theoretical and experimental concepts to describe and access position only born in mid-1990s. Pioneering measurements over past decade.



Mapping out the partonic structure of the proton

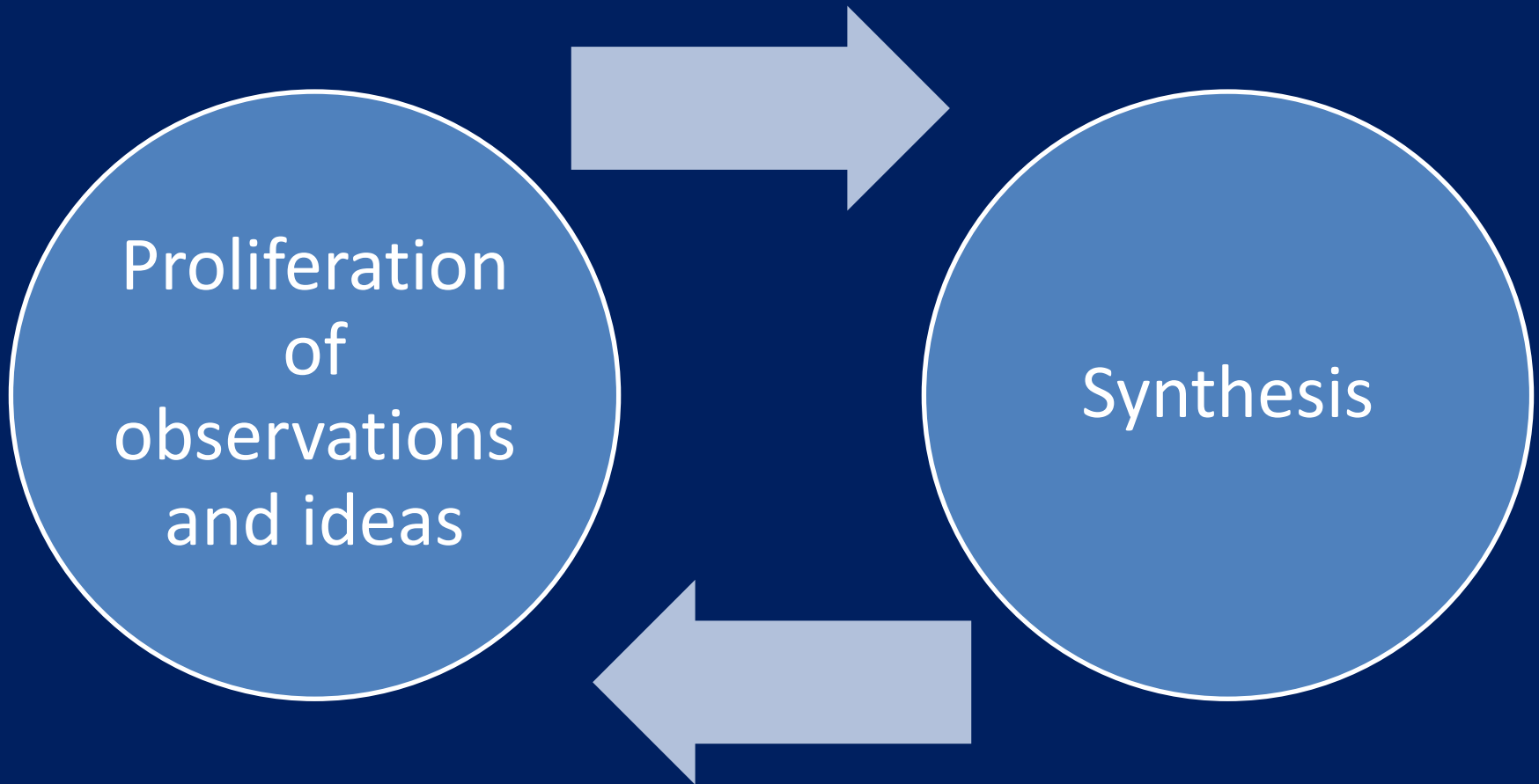
What does the proton look like in terms of the quarks and gluons inside it?

- *Position*
- *Momentum*
- *Spin*
- *Flavor*
- *Color*

Accounted for theoretically from beginning of QCD, but more detailed, potentially observable effects of color have come to forefront in last few years . . .



A cyclical process



Factorization and universality in perturbative QCD

- Need to systematically *factorize* short- and long-distance physics—observable physical QCD processes always involve at least one long-distance scale (confinement)!
- Long-distance (i.e. not perturbatively calculable) functions need to be *universal* in order to be portable across calculations for many processes

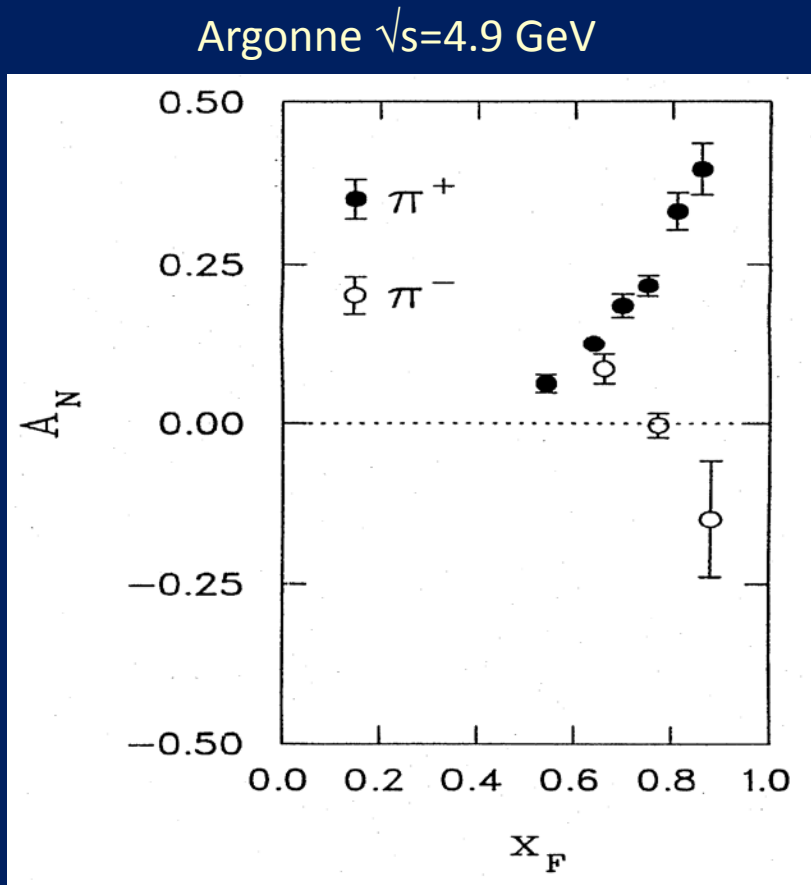


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Measure nonperturbative parton distribution functions (pdfs) and fragmentation functions in many colliding systems over a wide kinematic range → constrain by performing *simultaneous fits to world data*

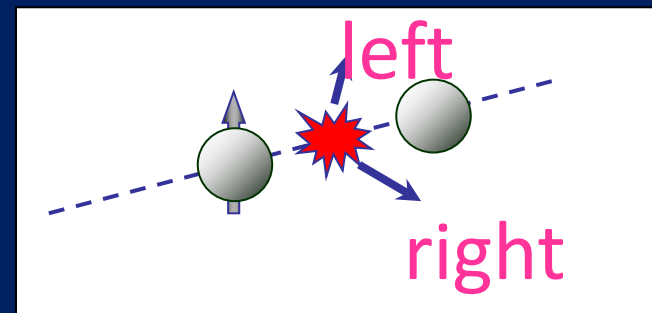
Spin-momentum correlations: 1976 discovery in $p+p$ collisions



W.H. Dragoset et al., PRL36, 929 (1976)

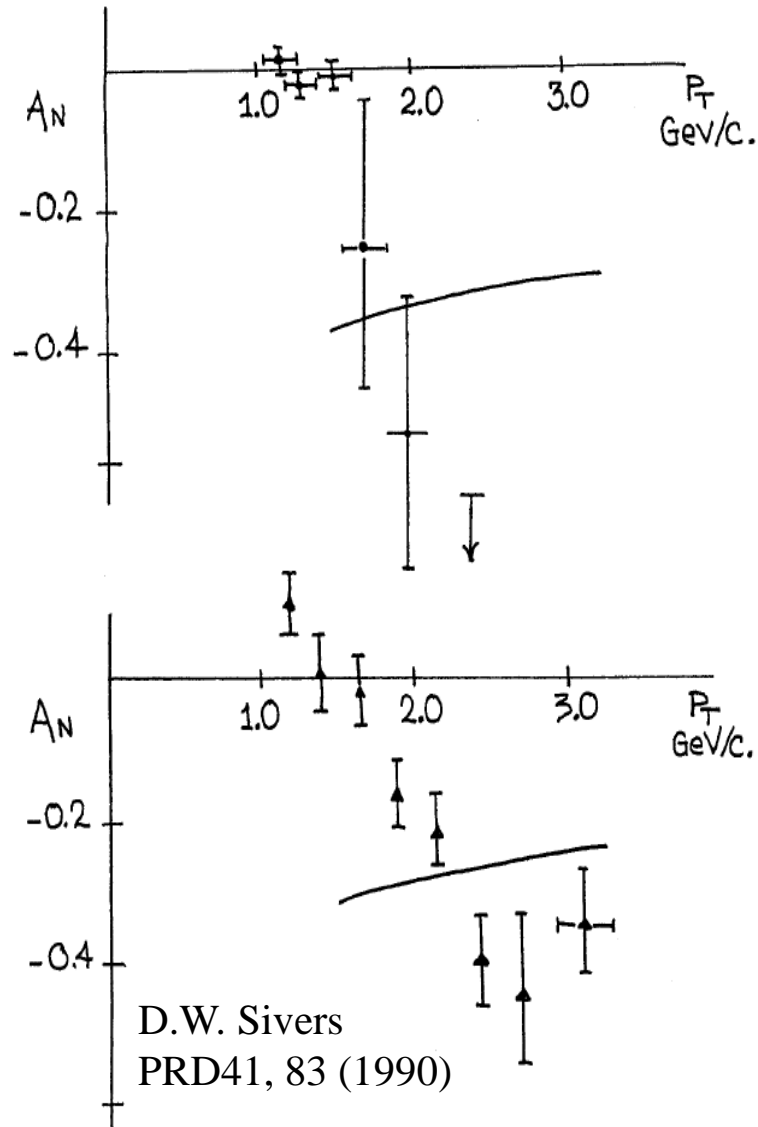
Charged pions produced preferentially on one or the other side with respect to the transversely polarized beam direction—by up to 40%!!

Had to wait more than a decade for the birth of a new subfield in order to explore the possibilities . . .



$$x_F = 2p_{\text{long}} / \sqrt{s}$$

Transverse-momentum-dependent distributions and single-spin asymmetries

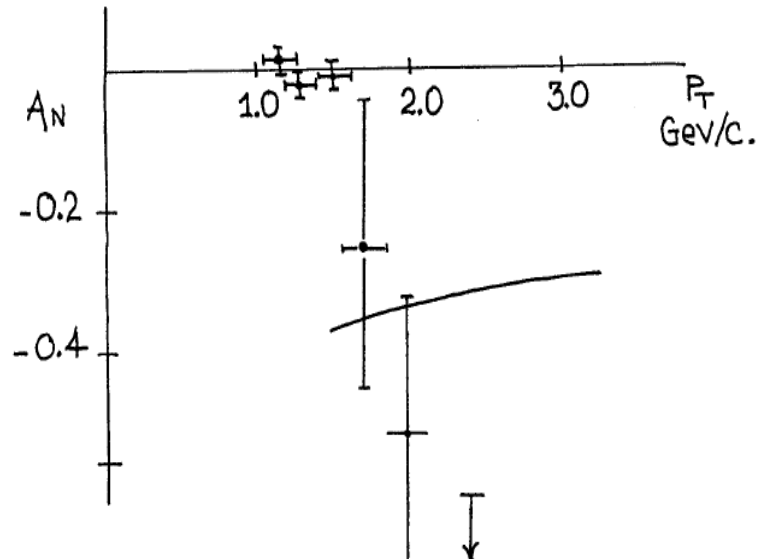


- 1990: “Sivers mechanism” proposed in attempt to understand observed asymmetries
- Departs from traditional *collinear* factorization assumption in pQCD and proposes correlation between the *intrinsic transverse motion* of the quarks and gluons and the proton’s spin

$$s \cdot (p_1 \times p_2)$$

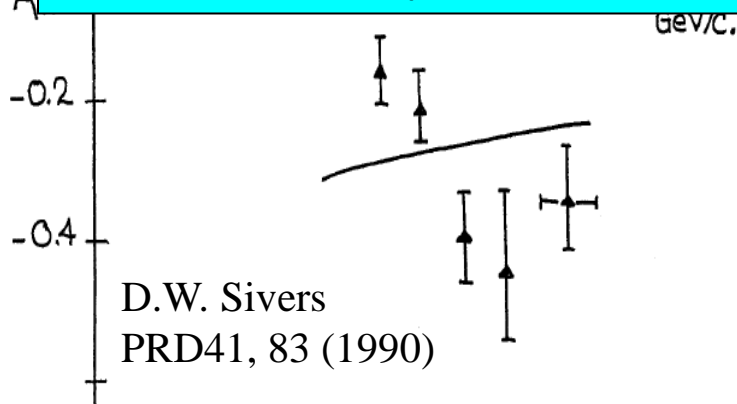
Spin and momenta of
partons and/or hadrons

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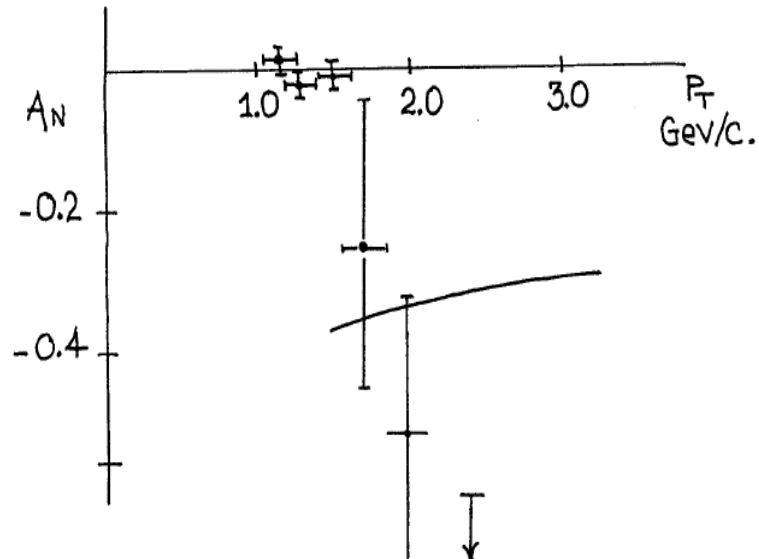
Sivers distribution: first transverse-momentum-dependent parton distribution function describing a spin-momentum correlation



$$s \cdot (p_1 \times p_2)$$

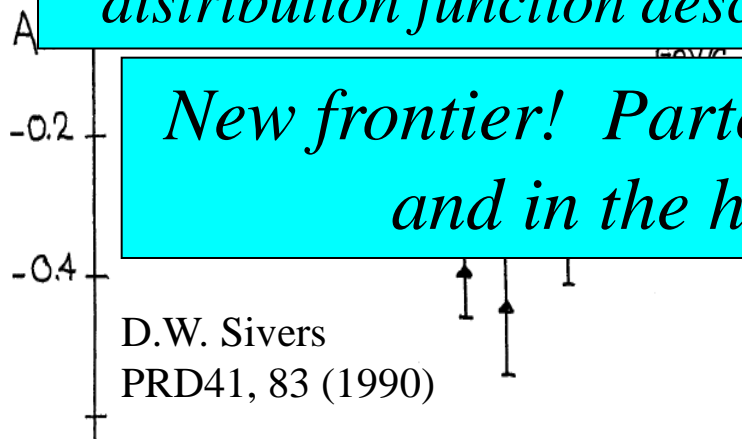
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*New frontier! Parton **dynamics** inside hadrons,
and in the hadronization process*

D.W. Sivers
PRD41, 83 (1990)

$$s \cdot (p_1 \times p_2)$$

Spin and momenta of
partons and/or hadrons

Spin-spin and spin-momentum correlations in QCD bound states

Unpolarized


$$f_1 = \text{[circle with dot]}$$

Spin-spin correlations

$$g_{1L} = \text{[circle with dot and right arrow]} - \text{[circle with dot and left arrow]} \quad \text{Helicity}$$

$$h_{1T} = \text{[circle with dot and up arrow]} - \text{[circle with dot and down arrow]} \quad \text{Transversity}$$

Worm-gear
(Kotzinian-Mulders)

$$g_{1T} = \text{[circle with dot and right arrow]} - \text{[circle with dot and left arrow]}$$


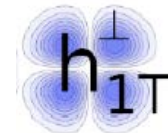
Spin-momentum correlations

$$S \cdot (p_1 \times p_2)$$

$$f_{1T}^\perp = \text{[circle with dot and up arrow]} - \text{[circle with dot and down arrow]} \quad \text{Sivers}$$

$$h_1^\perp = \text{[circle with dot and right arrow]} - \text{[circle with dot and left arrow]} \quad \text{Boer-Mulders}$$

$$h_{1L}^\perp = \text{[circle with dot and up arrow]} - \text{[circle with dot and down arrow]} \quad \text{Worm-gear} \quad h_{1T}^\perp = \text{[circle with dot and right arrow]} - \text{[circle with dot and left arrow]}$$



Pretzelosity

Spin-spin and spin-momentum correlations in QCD bound states

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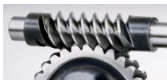
$$f_1 = \text{[circle with dot]}$$

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Helicity

Worm-gear
(Kotzinian-Mulders)

$$g_{1T} = \text{[circle with dot and right arrow]} - \text{[circle with dot and left arrow]}$$


Lots of evidence from deep-inelastic lepton-nucleon scattering experiments over past ~10 years that many of these correlations are nonzero in nature!

Spin-momentum correlations

$$S \cdot (p_1 \times p_2)$$

$$h_1^\perp = \text{[circle with dot and down arrow]} - \text{[circle with dot and up arrow]}$$

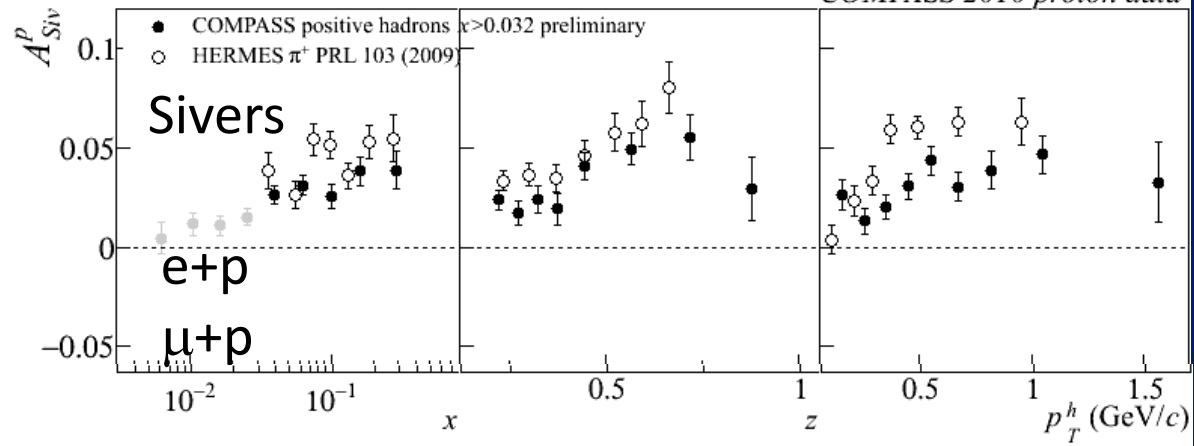
Boer-Mulders

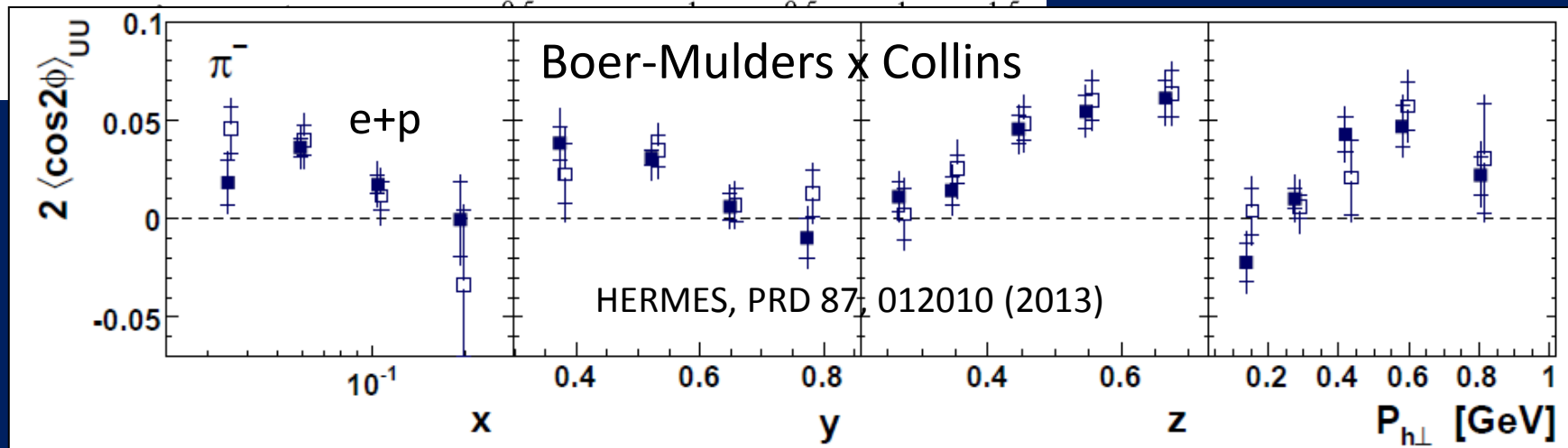
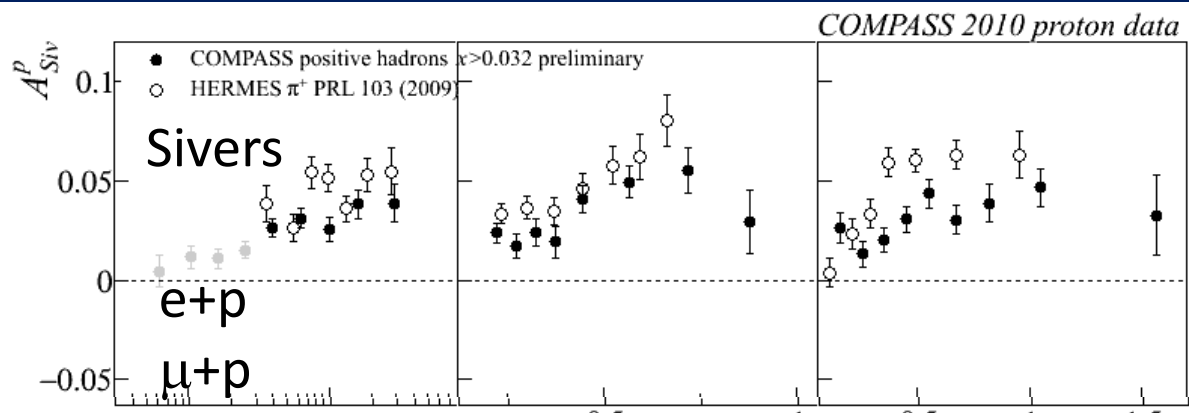
$$h_{1L}^\perp = \text{[circle with dot and right arrow]} - \text{[circle with dot and left arrow]}$$

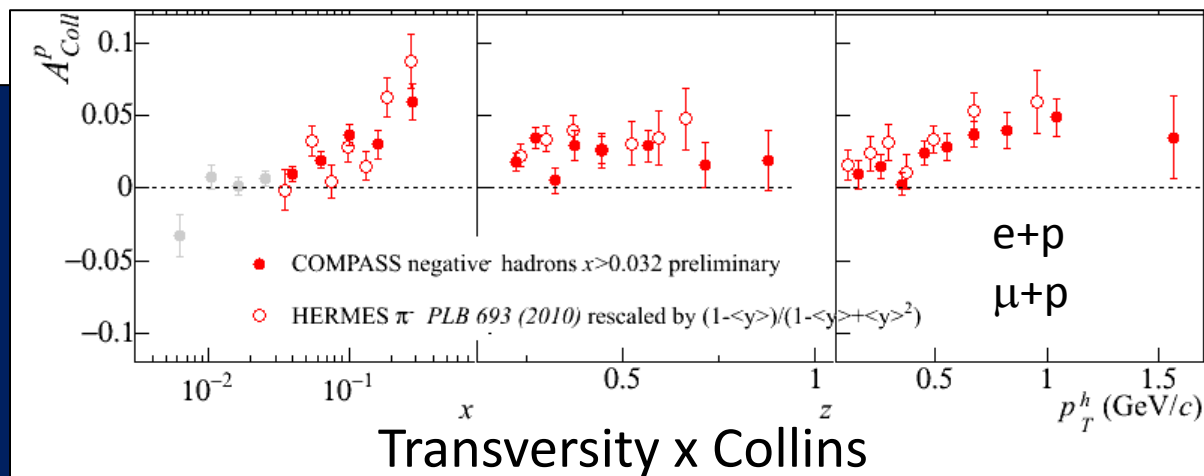
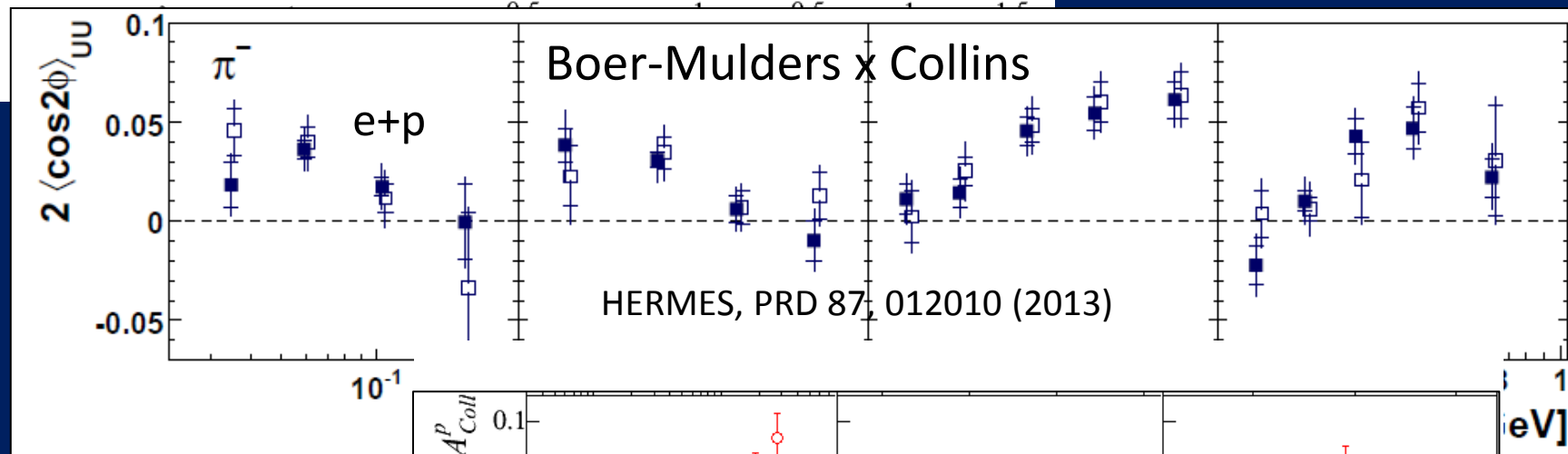
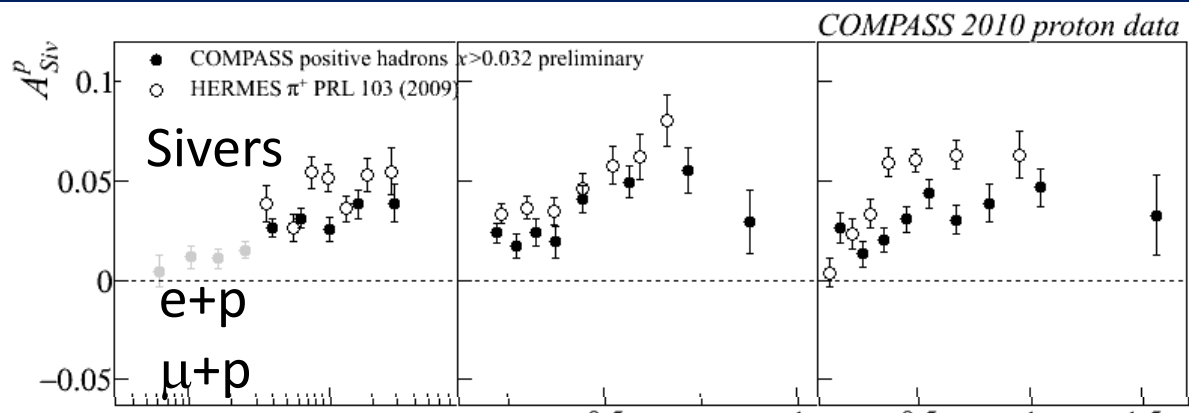
Worm-gear

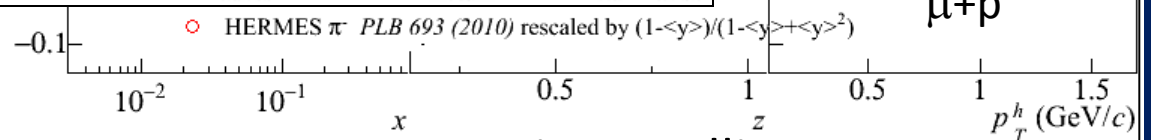
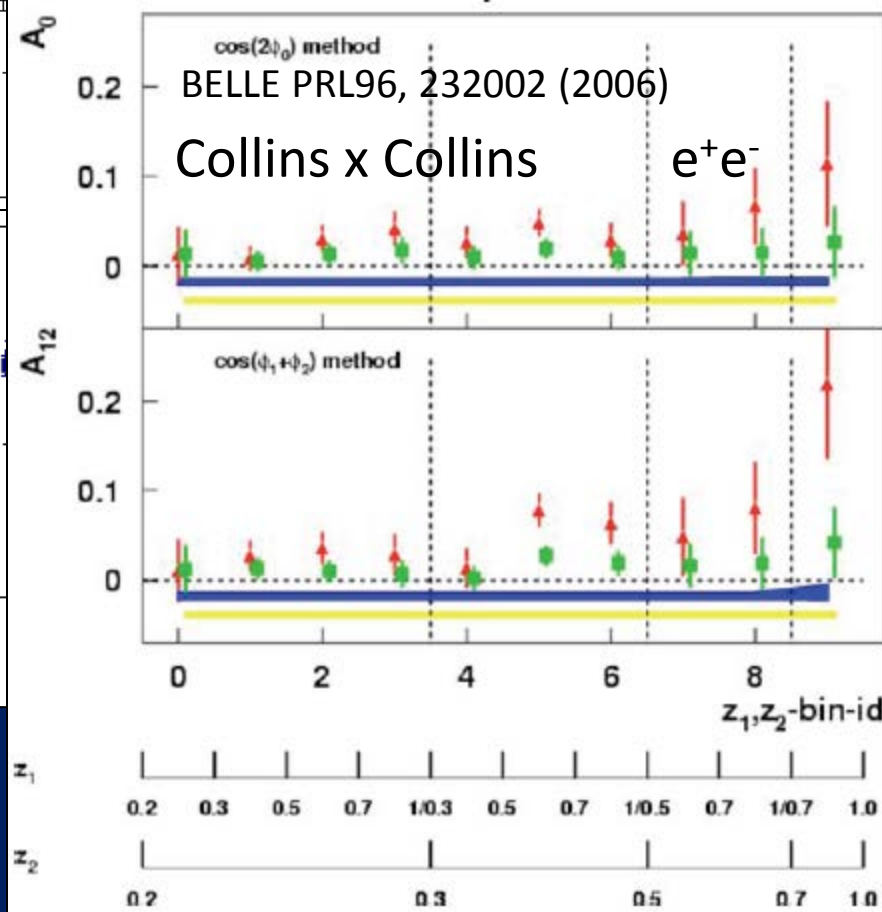
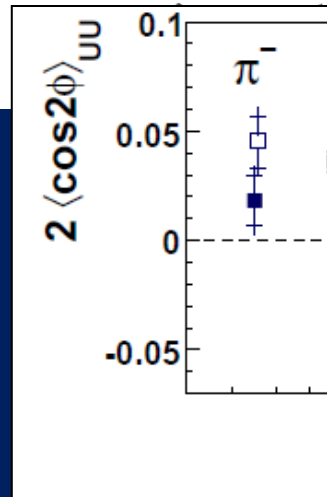
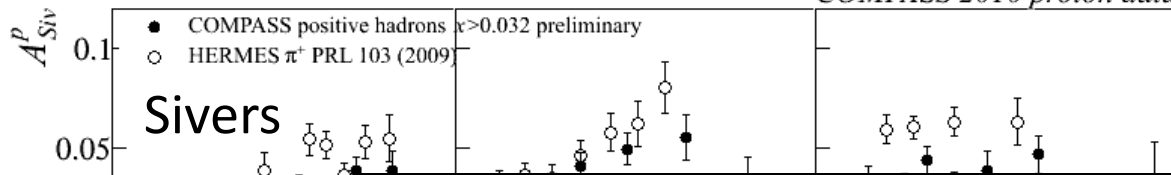
$$h_{1T}^\perp = \text{[circle with dot and up arrow]} - \text{[circle with dot and down arrow]}$$

Pretzelosity

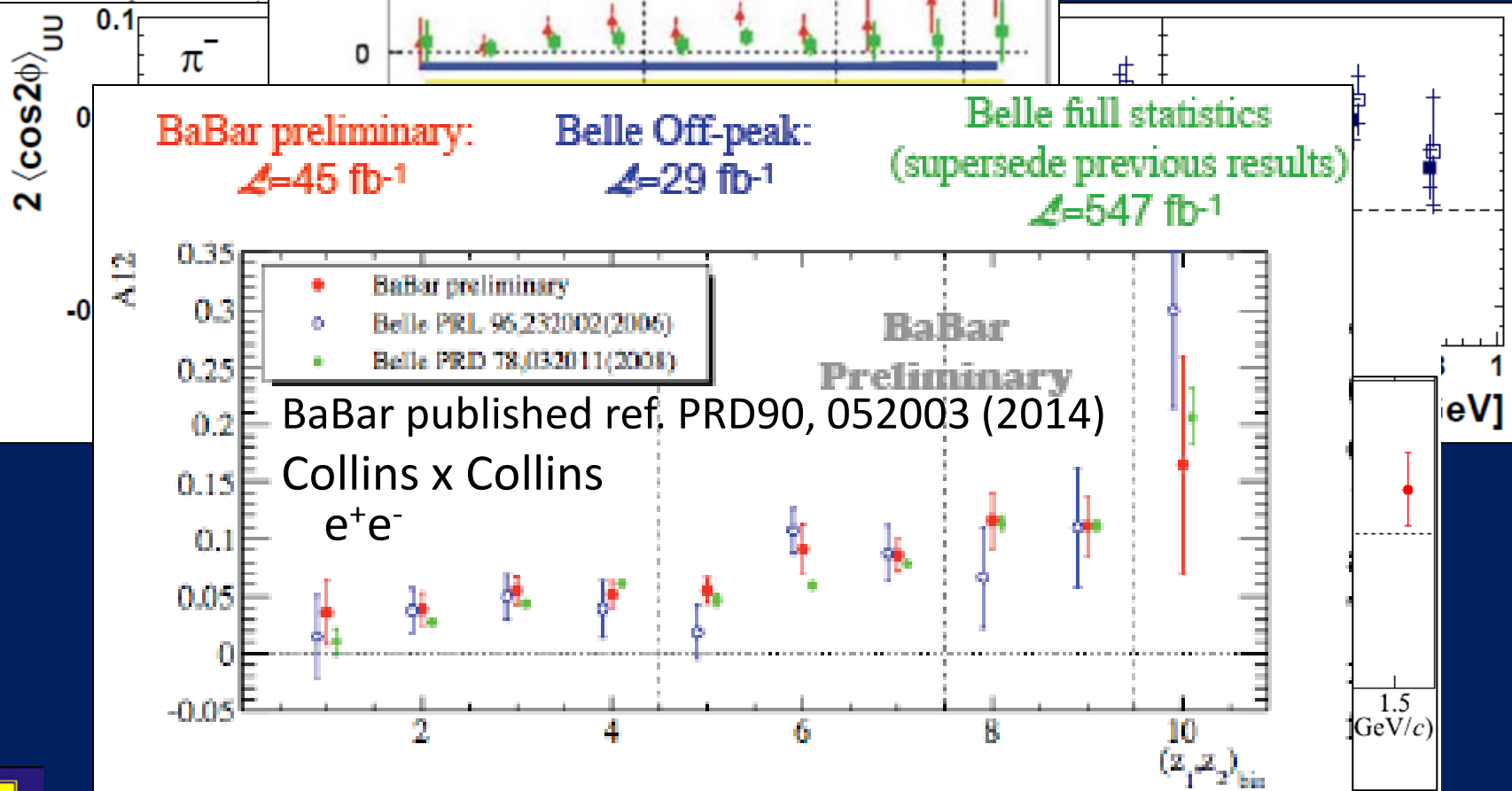
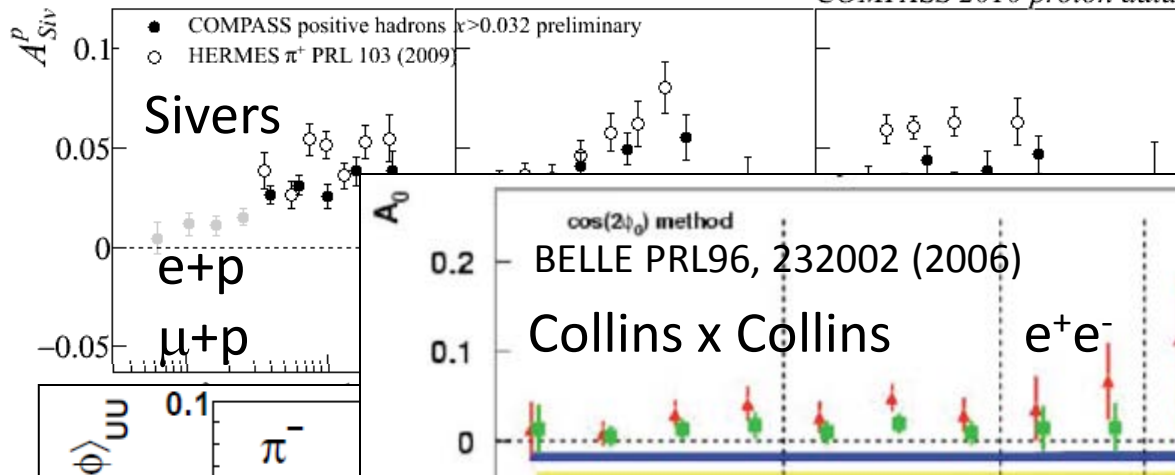








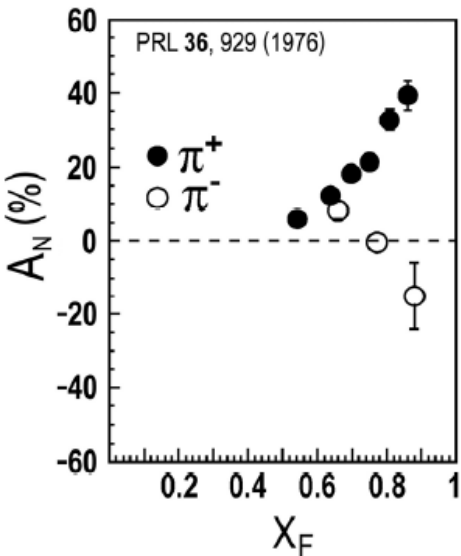
Transversity x Collins



But what about proton-proton collisions?

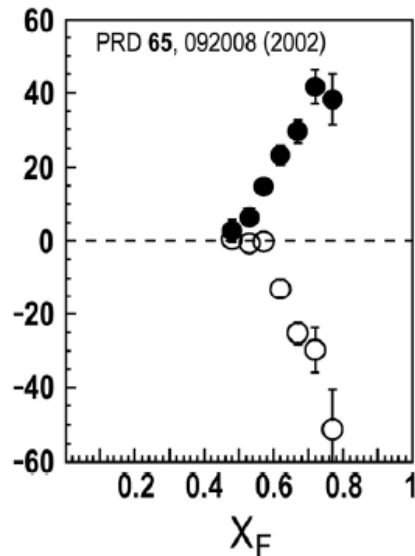
ANL

$\sqrt{s}=4.9$ GeV



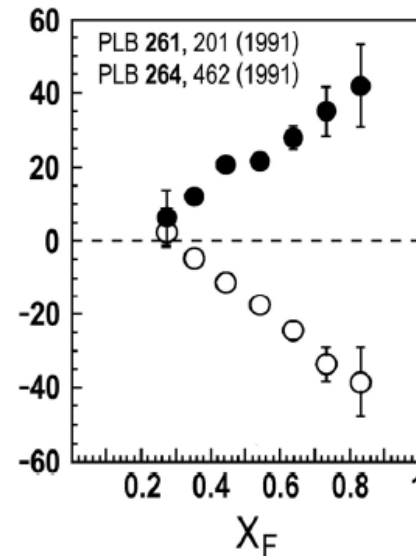
BNL

$\sqrt{s}=6.6$ GeV



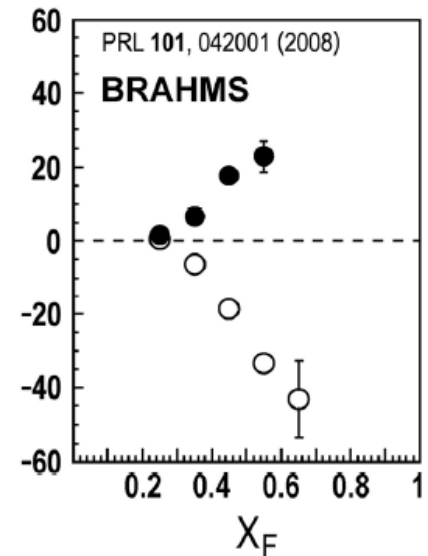
FNAL

$\sqrt{s}=19.4$ GeV

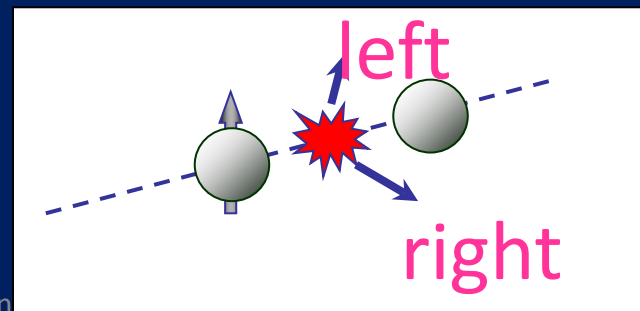


RHIC

$\sqrt{s}=62.4$ GeV



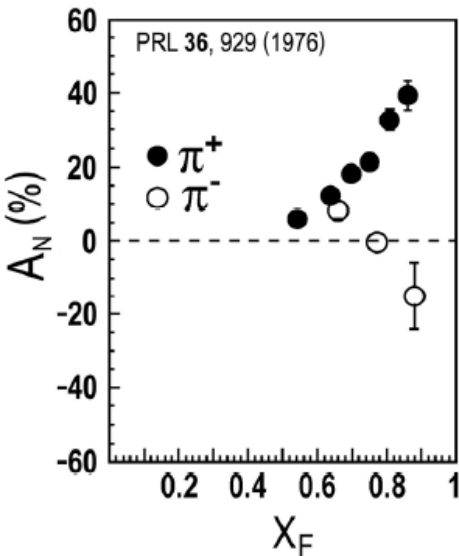
$$x_F = 2p_{long} / \sqrt{s}$$



But what about proton-proton collisions?

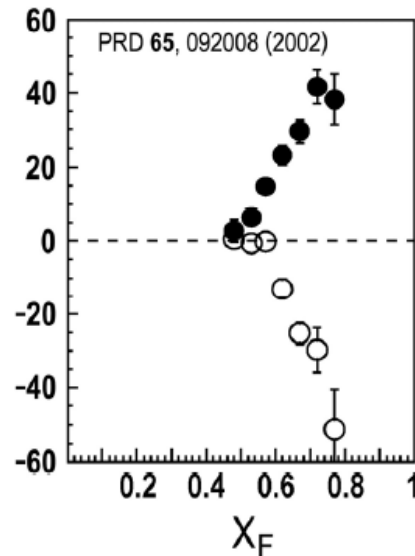
ANL

$\sqrt{s}=4.9$ GeV



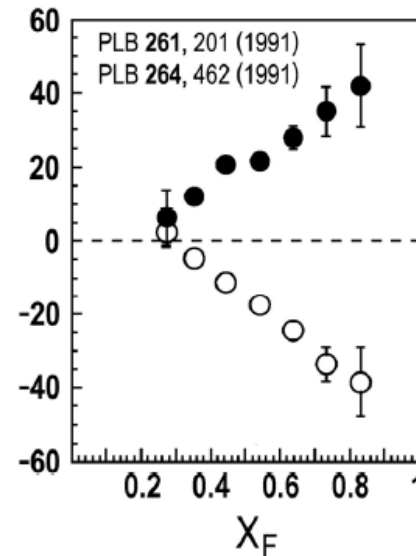
BNL

$\sqrt{s}=6.6$ GeV



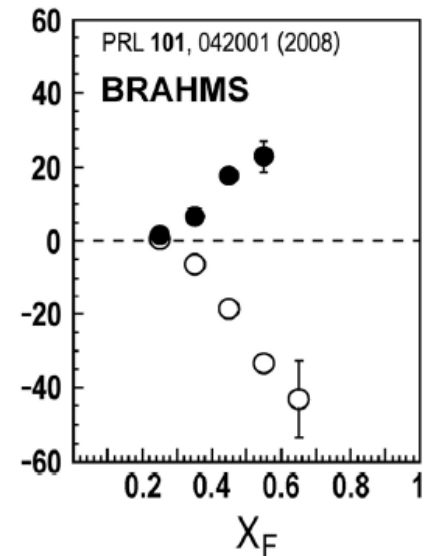
FNAL

$\sqrt{s}=19.4$ GeV



RHIC

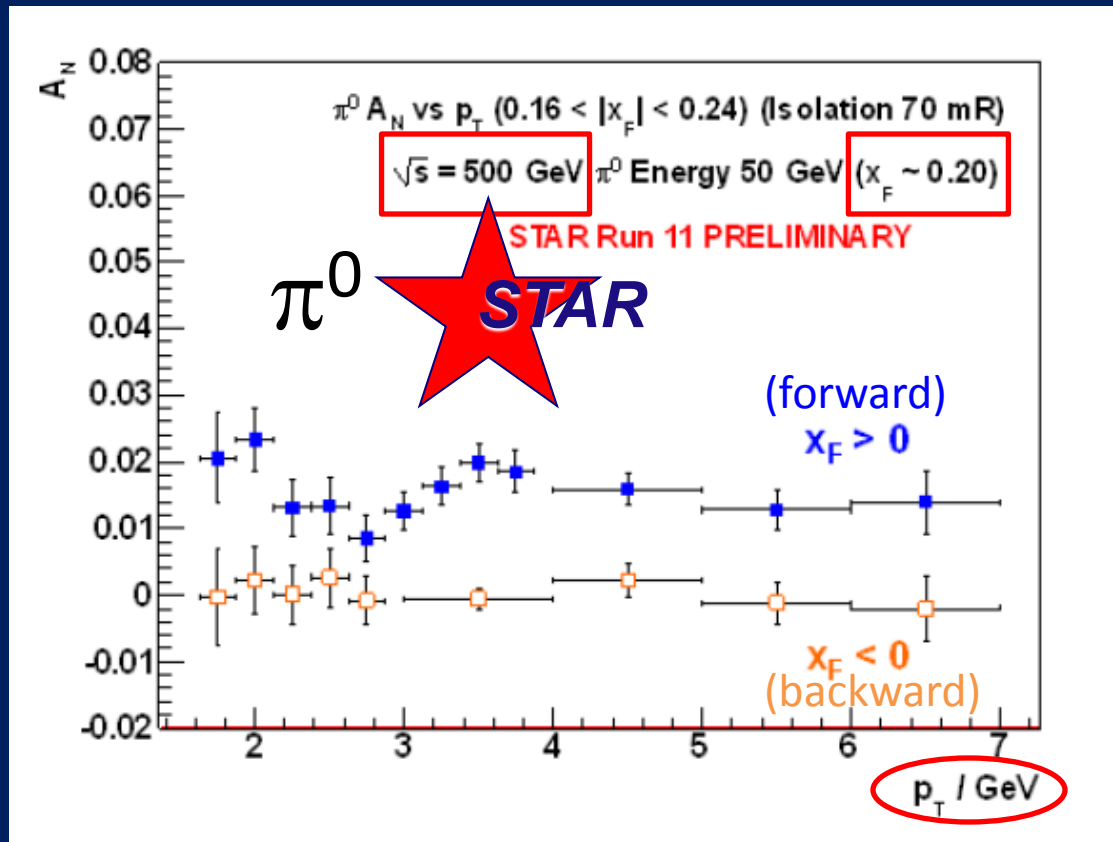
$\sqrt{s}=62.4$ GeV



Strikingly similar effects across energies!

→ Continuum between nonperturbative/nonpartonic and perturbative/partonic descriptions of this nonperturbative structure?

Single-spin asymmetries in transversely polarized $p+p$ collisions



- Effects persist to kinematic regimes where perturbative QCD techniques clearly apply
- $p_T = 7$ GeV
→ $Q^2 \sim 49$ GeV²!

$p+p \rightarrow \text{hadron} + X$:
Challenging to interpret

- Always huge effects!
- But in $p+p \rightarrow \text{pion} + X$ don't have enough information to separate initial-state (proton structure) from final-state (hadronization) effects

Properties of naive-T-odd spin-momentum correlation functions

- Sivers transverse-momentum-dependent parton distribution function is odd under “naïve-time-reversal” (actually a PT transformation)
 - As is Boer-Mulders spin-momentum correlation



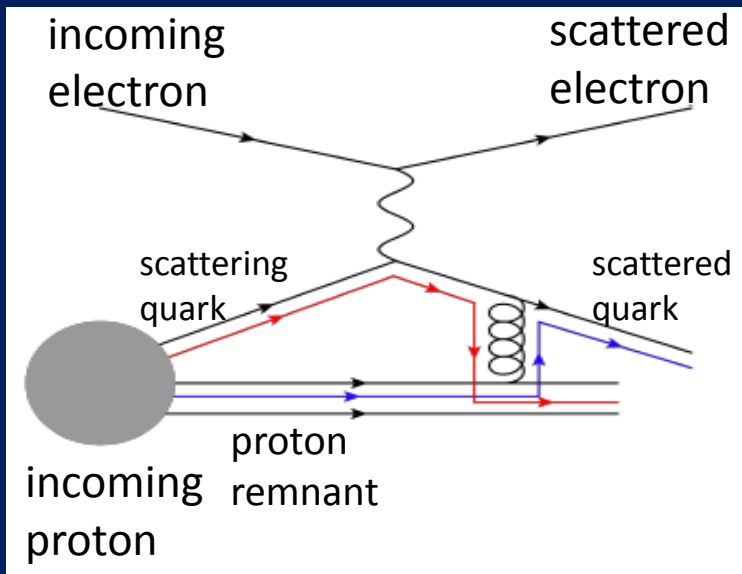
Properties of naive-T-odd spin-momentum correlation functions

- Sivers transverse-momentum-dependent parton distribution function is odd under “naïve-time-reversal” (actually a PT transformation)
 - As is Boer-Mulders spin-momentum correlation
- In 1993, after original 1990 paper by D.W. Sivers, J.C. Collins claimed such functions must vanish
- Only realized in 2002 by Brodsky, Hwang, and Schmidt that could be nonvanishing if *phase interference effects due to color interactions* present

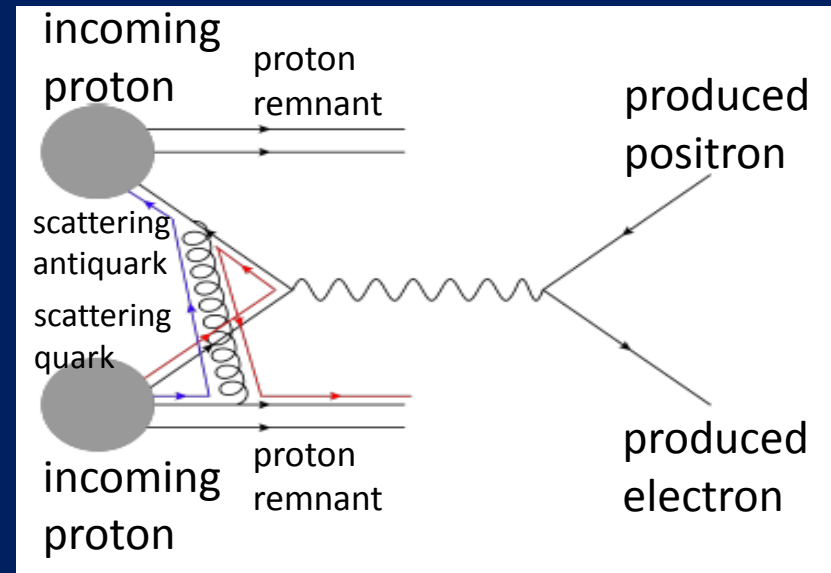


Modified universality of certain transverse-momentum-dependent distributions: *Color in action!*

Deep-inelastic lepton-nucleon scattering: Final-state color exchange

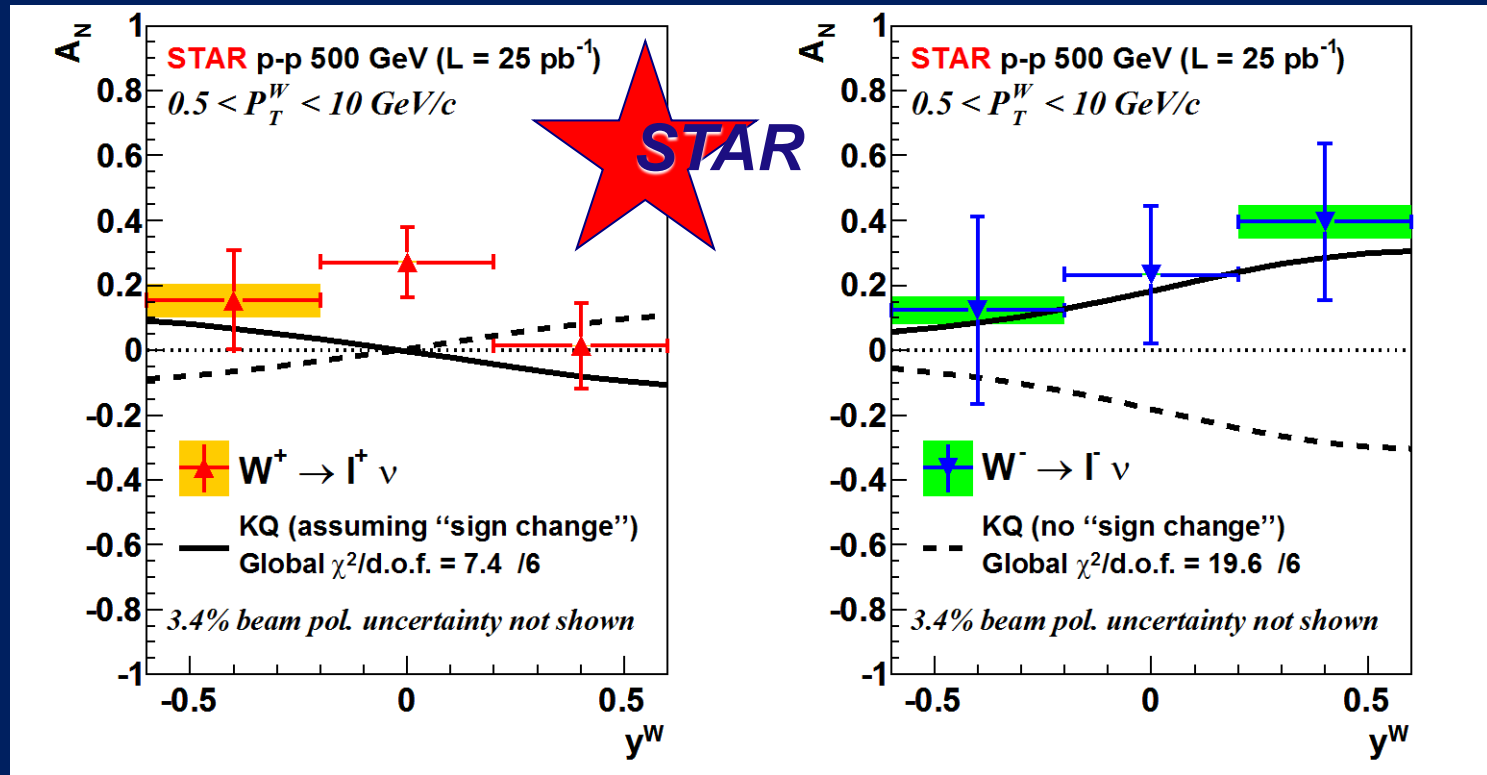


Quark-antiquark annihilation to leptons: Initial-state color exchange



As a result, get *opposite sign* for the *Sivers* transverse-momentum-dependent pdf when measure in semi-inclusive DIS versus Drell-Yan: *process-dependent* pdf! (Collins 2002)

Modified universality: Hints but not yet conclusively tested



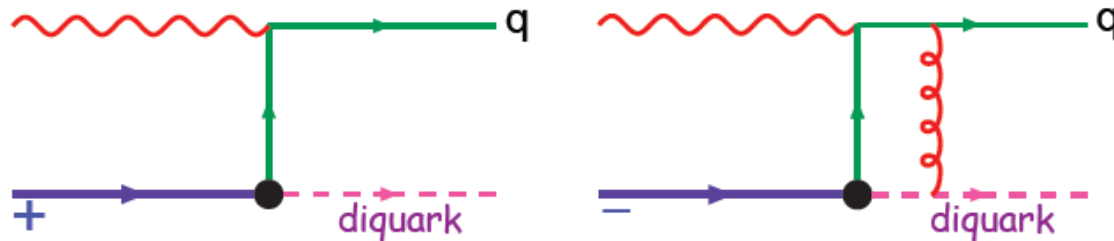
W boson measurement from STAR more consistent with sign change, but not yet conclusive. Will take more data in 2017. COMPASS experiment at CERN will also perform a measurement in next couple years.

Modified universality requires full QCD: Gauge-invariant quantum field theory

We have ignored here the subtleties needed to make this a gauge invariant definition: an appropriate path ordered exponential of the gluon field is needed [18].

From 1993 claim by J.C. Collins that such processes must vanish

gauge links have physical consequences;
quark models for non vanishing Sivvers function,
SIDIS final state interactions



Brodsky, Hwang, Schmidt, PL B530 (2002) 99 - Collins, PL B536 (2002) 43

An earlier proof that the Sivvers asymmetry vanishes because of time-reversal invariance is invalidated by the path-ordered exponential of the gluon field in the operator definition of parton densities. Instead, the time-reversal argument shows that the Sivvers asymmetry is reversed in sign in hadron-induced hard processes (e.g., Drell-Yan), thereby violating naive universality of parton densities. Previous phenomenology with time-reversal-odd parton densities is therefore validated.

$$[f_{1T}^{q\perp}]_{\text{SIDIS}} = -[f_{1T}^{q\perp}]_{\text{DY}}$$

Slide from M. Anselmino, Transversity 2014

Physical consequences of a gauge-invariant quantum theory: Aharonov-Bohm (1959)

Wikipedia:

“The Aharonov–Bohm effect is important conceptually because it bears on three issues apparent in the recasting of (Maxwell's) classical electromagnetic theory as a gauge theory, which before the advent of quantum mechanics could be argued to be a mathematical reformulation with no physical consequences. The Aharonov–Bohm thought experiments and their experimental realization imply that the issues were not just philosophical.

The three issues are:

- whether potentials are "physical" or just a convenient tool for calculating force fields;
- whether action principles are fundamental;
- the principle of locality.”



Physical consequences of a gauge-invariant quantum theory: Aharonov-Bohm (1959)

Physics Today, September 2009 :

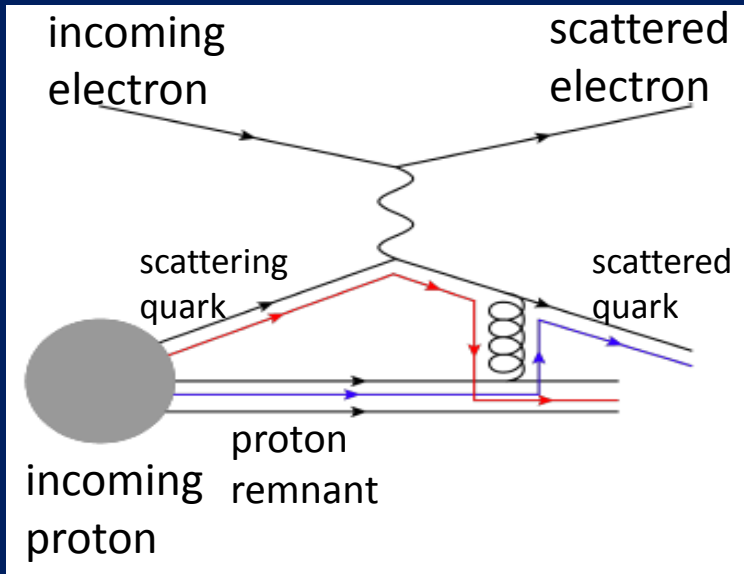
The Aharonov–Bohm effects: Variations on a subtle theme,
by Herman Batelaan and Akira Tonomura.

“Aharonov stresses that the arguments that led to the prediction of the various electromagnetic AB effects apply equally well to any other gauge-invariant quantum theory. In the standard model of particle physics, the strong and weak nuclear interactions are also described by gauge-invariant theories. So one may expect that particle-physics experimenters will be looking for new AB effects in new domains.”

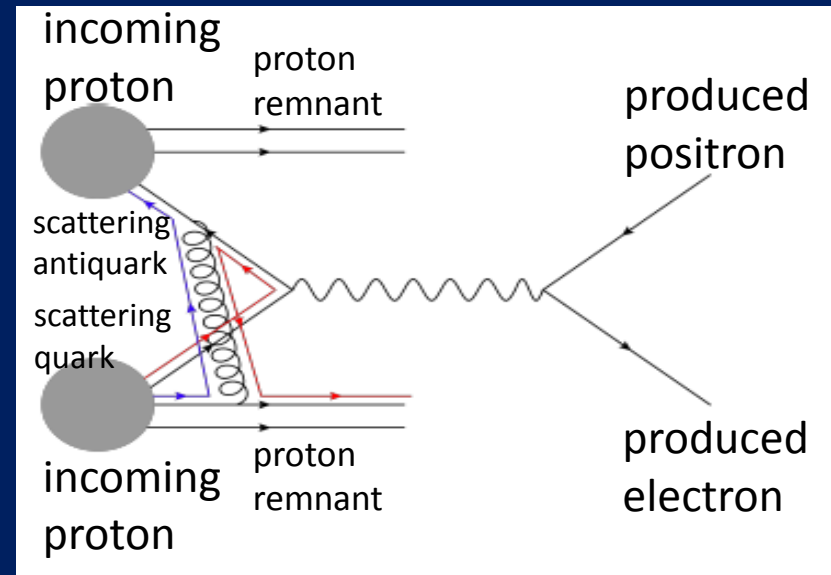


Physical consequences of a gauge-invariant quantum theory: Aharonov-Bohm effect in QCD!!

**Deep-inelastic lepton-nucleon
scattering: Final-state color exchange**



**Quark-antiquark annihilation to
leptons: Initial-state color exchange**



See e.g. Pijlman, hep-ph/0604226
or Sivers, arXiv:1109.2521

Physical consequences of a gauge-invariant quantum theory: Aharonov-Bohm effect in QCD!!

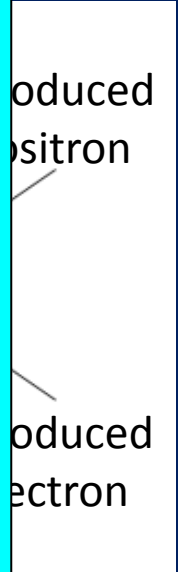
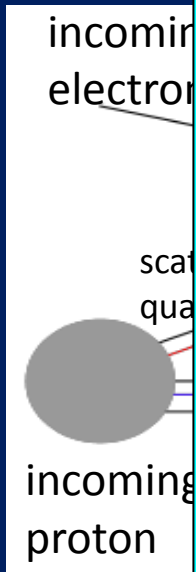
Deep-inelastic lepton-nucleon
scattering: Final-state color exchange

Quark-antiquark annihilation to
leptons: Initial-state color exchange

*Simplicity of these two processes:
Abelian vs. non-Abelian nature of the gauge
group doesn't play a major qualitative role.*

*BUT: In QCD expect additional, new effects
due to specific non-Abelian nature of the
gauge group*

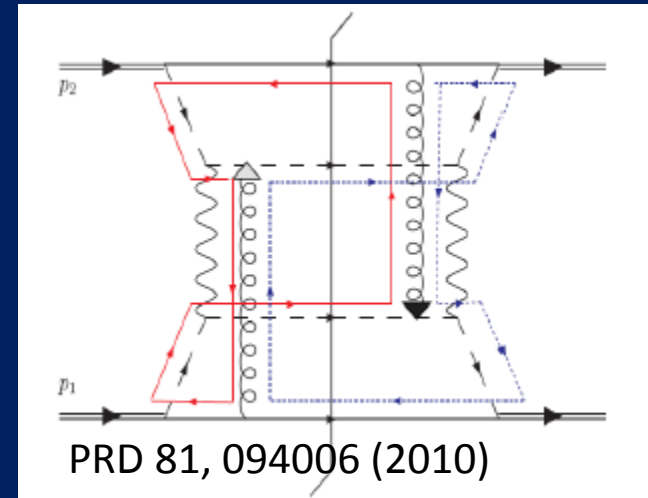
See e.g. Pijlman, hep-ph/0604226
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QCD Aharonov-Bohm effect:

Color entanglement

- 2010: Rogers and Mulders predict *color entanglement* in processes involving $p+p$ production of hadrons if quark transverse momentum taken into account
- Quarks become correlated *across* the two protons
- Consequence of QCD specifically as a *non-Abelian* gauge theory!



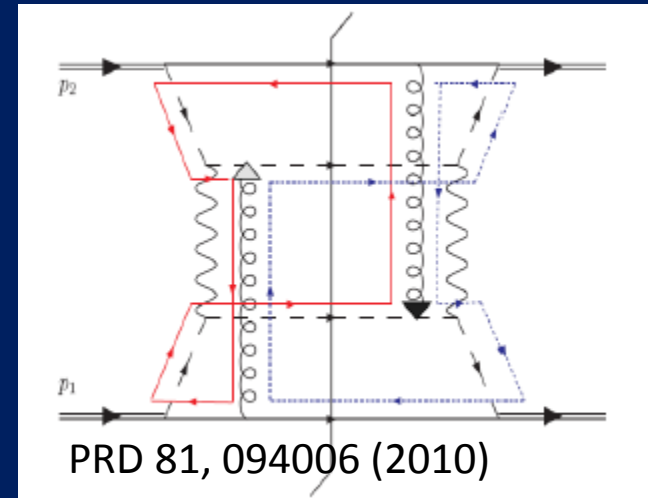
$$p + p \rightarrow h_1 + h_2 + X$$

Color flow can't be described as flow in the two gluons separately. Requires simultaneous presence of both.

QCD Aharonov-Bohm effect:

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$$p + p \rightarrow h_1 + h_2 + X$$

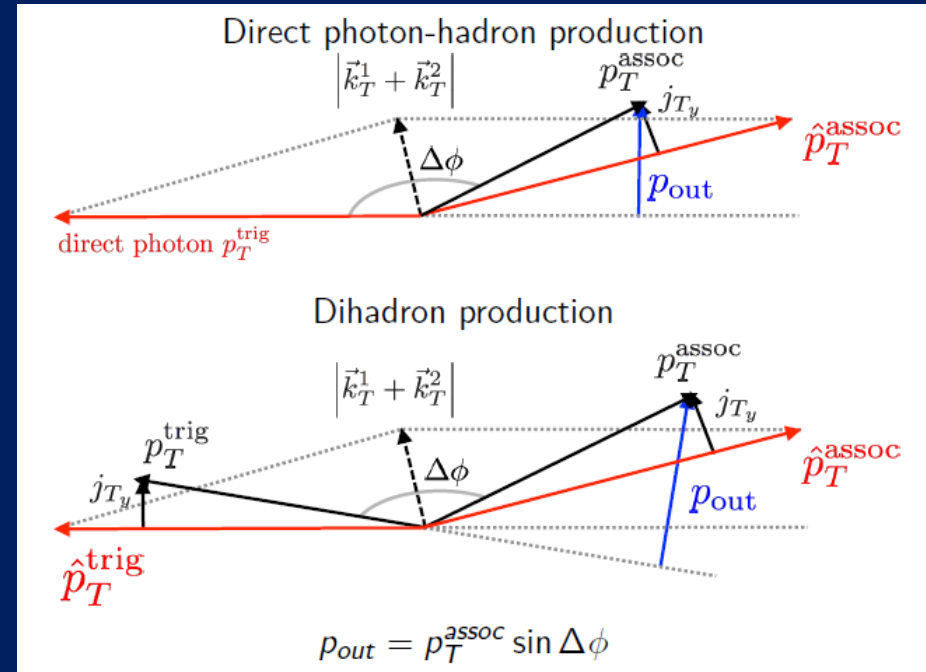
Color flow can't be described as flow in the two gluons separately. Requires simultaneous presence of both.

Huge transverse spin asymmetries in $p+p$ a color entanglement effect??

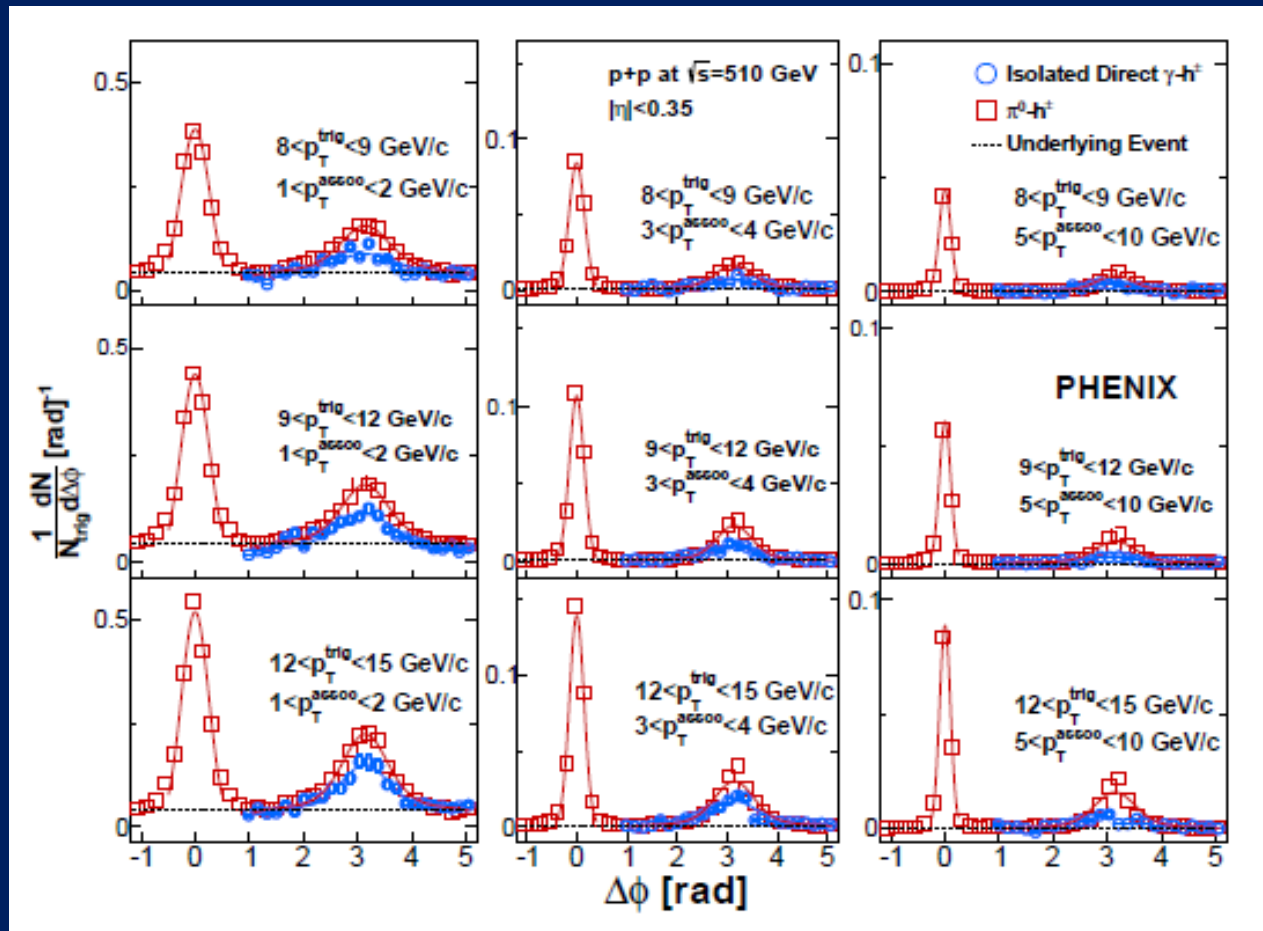
Searching for evidence of color entanglement at RHIC

- Need observable sensitive to a nonperturbative momentum scale
 - Nearly back-to-back particle production
- Need 2 initial-state hadrons
 - color exchange between a scattering parton and remnant of other proton
- And at least 1 final-state hadron
 - exchange between scattered parton and either remnant

→ In p+p collisions, measure out-of-plane momentum component in nearly back-to-back photon-hadron and hadron-hadron production

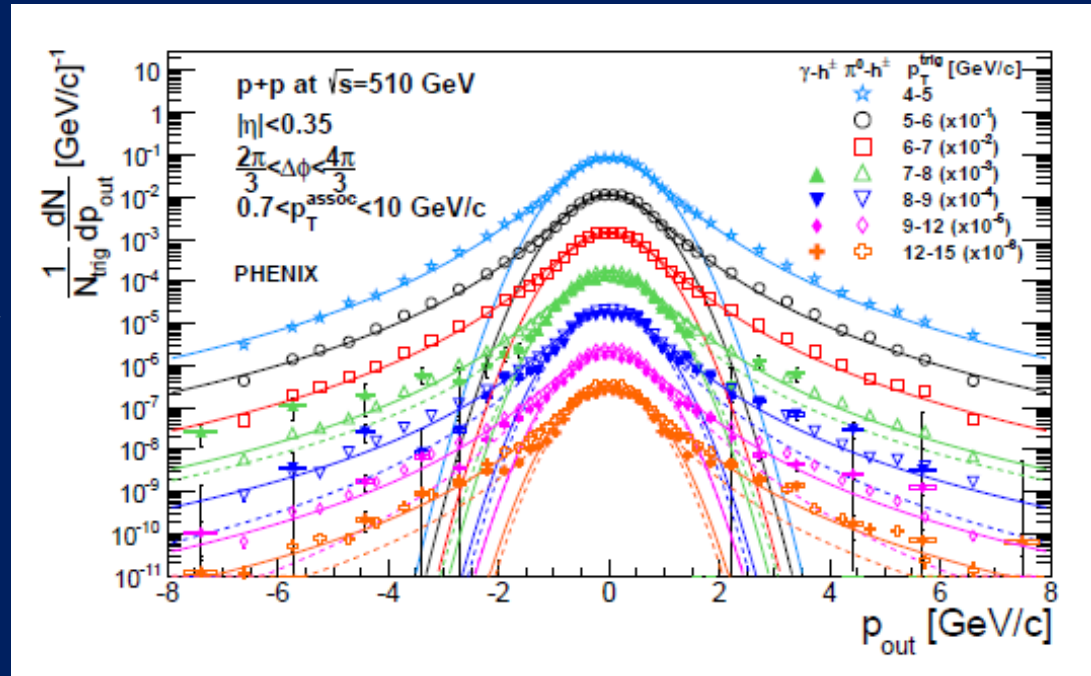


Two-particle correlation distributions show expected jet-like structure



Out-of-plane momentum component distributions

- Clear two-component distribution
 - Gaussian near 0—nonperturbative transverse momentum
 - Power-law at large p_{out} —kicks from hard (perturbative) gluon radiation
- Different colors \rightarrow different bins in hard interaction scale



Curves are fits to Gaussian and Kaplan functions, not calculations!

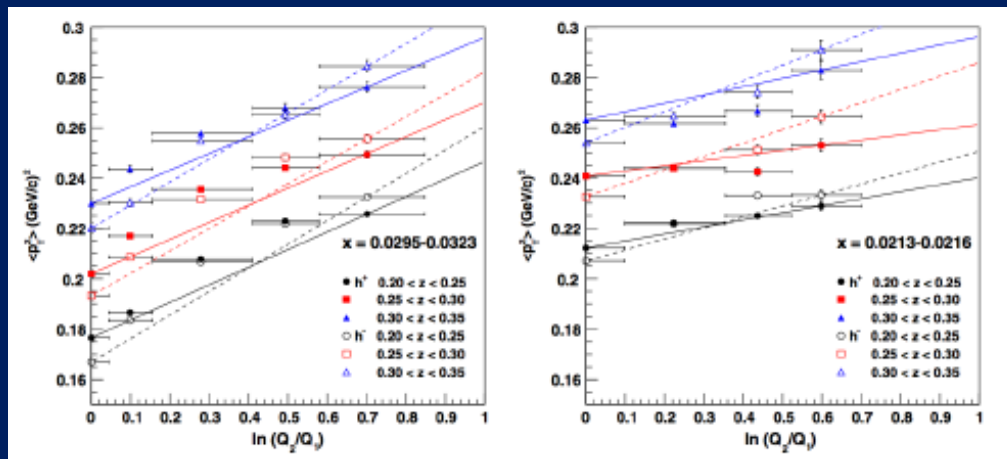
Look at evolution of nonperturbative transverse momentum widths with hard scale (Q^2)

- Theoretical proof of factorization for transverse-momentum-dependent parton distribution functions directly predicts that nonperturbative transverse momentum widths *increase* as a function of the hard scattering energy scale
 - Increased phase space for gluon radiation

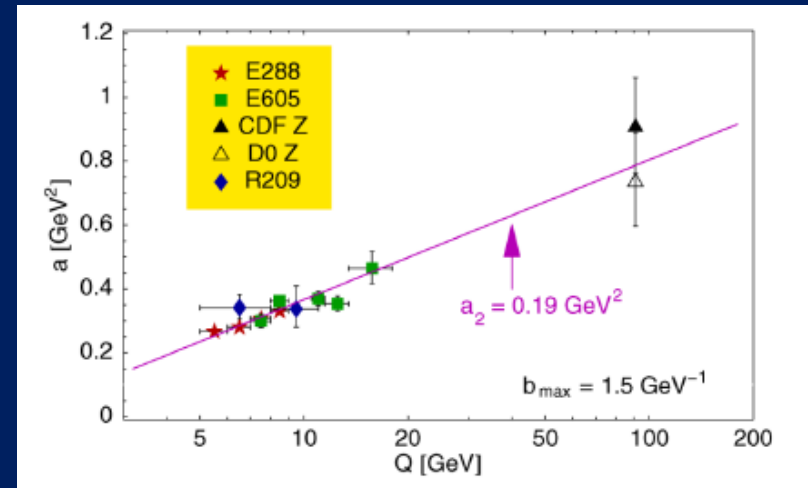


Look at *evolution* of nonperturbative transverse momentum widths with hard scale (Q^2)

- Theoretical proof of factorization for transverse-momentum-dependent parton distribution functions directly predicts that nonperturbative transverse momentum widths *increase* as a function of the hard scattering energy scale
 - Increased phase space for gluon radiation
- Observed experimentally in semi-inclusive deep-inelastic lepton-nucleon scattering (left) and quark-antiquark annihilation to leptons (right)



Aidala, Field, Gamberg, Rogers, Phys. Rev. D89, 094002 (2014)

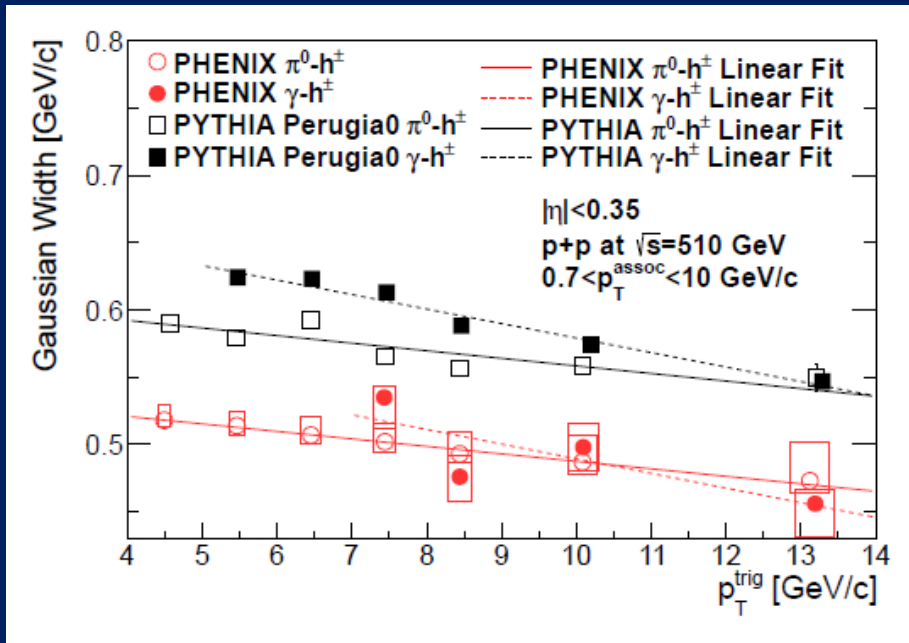


Konychev + Nadolsky, Phys. Lett. B633, 710 (2006)



Nonperturbative momentum widths decrease in processes where entanglement predicted

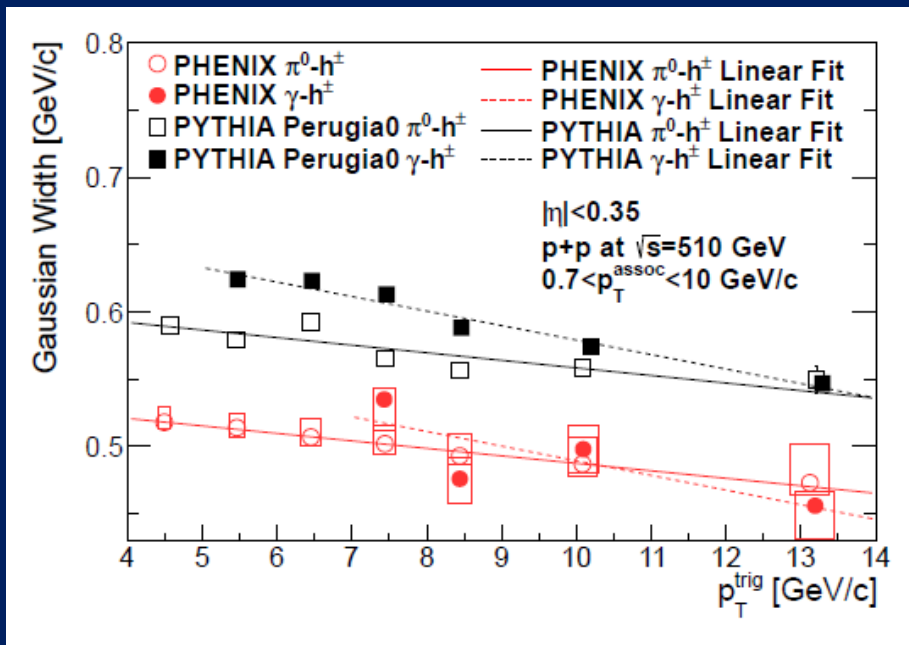
- Suggestive of quantum-correlated partons across colliding protons!
- However, have not yet completely ruled out kinematic effects



PHENIX Collab., arXiv:1609.04769,
Submitted to PRD



Nonperturbative momentum widths *decrease* in processes where entanglement predicted



PHENIX Collab., arXiv:1609.04769,
Submitted to PRD

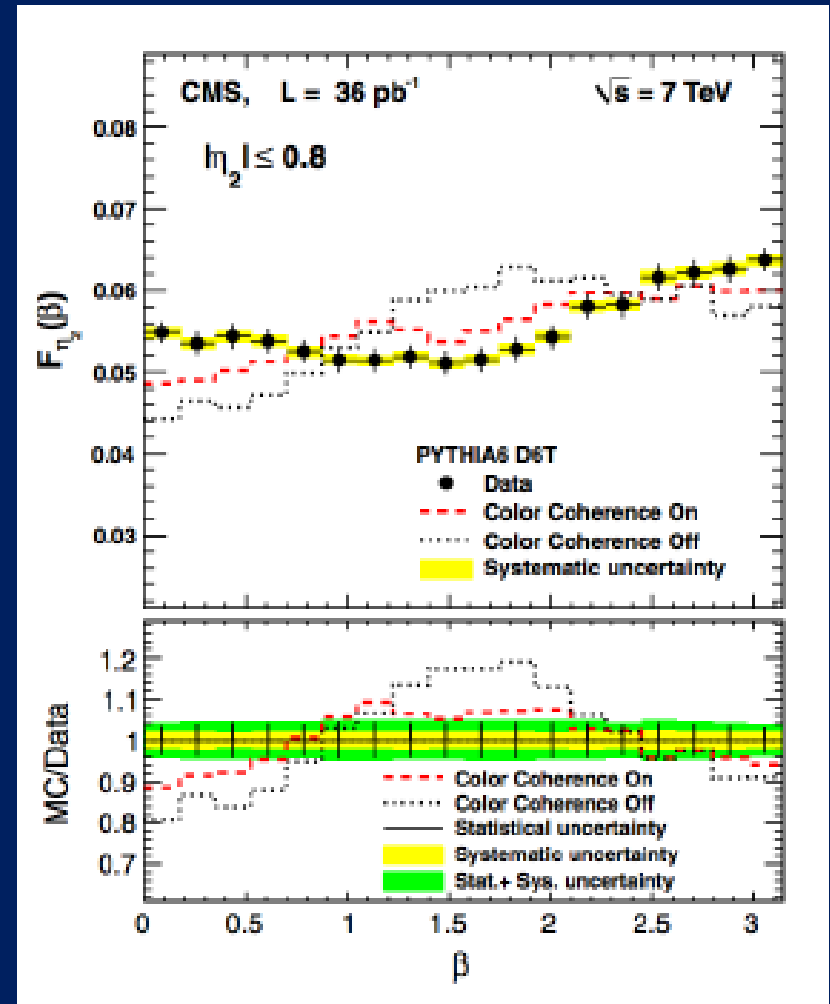
- Suggestive of quantum-correlated partons across colliding protons!
- However, have not yet completely ruled out kinematic effects
- Slope of decrease for both photon-hadron and dihadron correlations reproduced ~exactly in PYTHIA p+p event generator—could this effect be in PYTHIA??
 - Effectively yes! Unlike analytic pQCD calculations, PYTHIA forces *entire event including remnants* to color neutralize, implemented via something they call “color reconnection”

- Discussions ongoing, and follow-up studies underway . . .



Possible links to “color coherence” at Tevatron and LHC?

- D0, CDF, CMS have all published on evidence for “color coherence effects”
 - CMS: EPJ C74, 2901 (2014)
 - CDF: PRD50, 5562 (1994)
 - D0: PLB414, 419 (1997)
- Few citations—relatively little-known work thus far. Need to get different communities talking to explore detailed color effects more in upcoming years!



Summary

- Early years of rewarding new era of quantitative basic research in QCD!



Summary

- Early years of rewarding new era of quantitative basic research in QCD!
- Gradually shifting to think about QCD systems in new ways, focusing on topics/ideas/concepts that have long been familiar to the world of condensed matter physics
 - All sorts of correlations within systems
 - Quantum mechanical phase interference effects
 - Quantum entangled systems



Summary

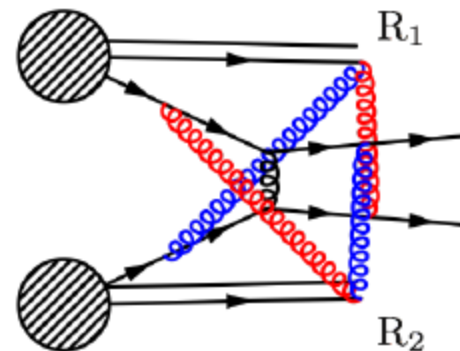
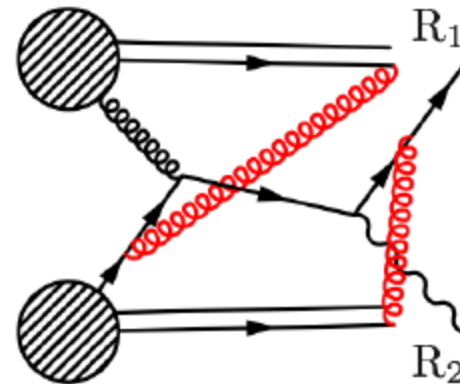
- Early years of rewarding new era of quantitative basic research in QCD!
- Gradually shifting to think about QCD systems in new ways, focusing on topics/ideas/concepts that have long been familiar to the world of condensed matter physics
 - All sorts of correlations within systems
 - Quantum mechanical phase interference effects
 - Quantum entangled systems
- Will be exciting to continue testing and exploring these ideas and phenomena in upcoming years . . .

Extra



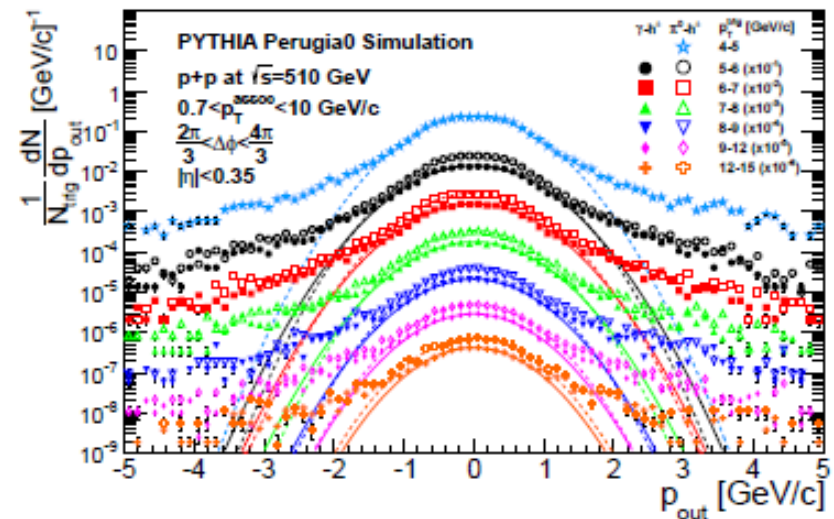
$$p+p \rightarrow \gamma+h+X \text{ vs. } p+p \rightarrow h+h+X$$

- Direct photon-hadron and dihadron correlations both predicted to be sensitive to factorization breaking effects in PHENIX
- Assuming factorization, direct photon-hadrons probe three nonperturbative functions, while dihadrons probe four
- Direct photons offer one less avenue for gluon exchange in the final-state: fewer/different effects?



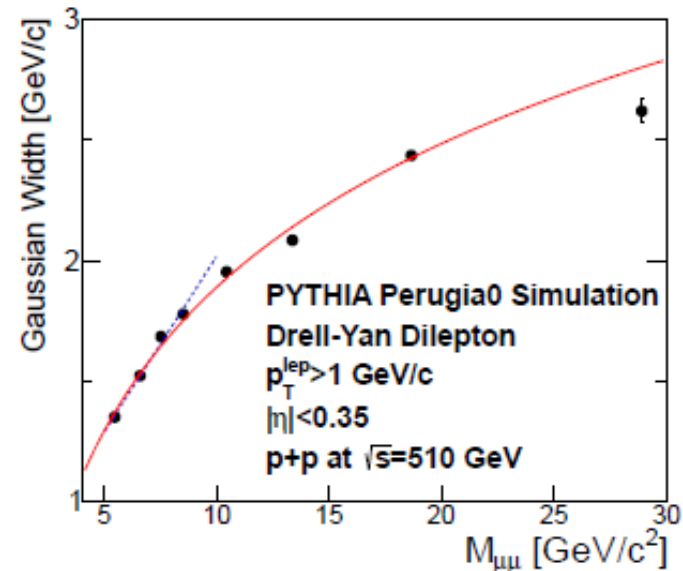
PYTHIA p_{out} distributions

- PYTHIA π^0 - h^\pm and isolated γ - h^\pm correlations analyzed similarly to data
- PYTHIA exhibits similar characteristics to data: nonperturbative transitioning to perturbative region
- Initial and final state interactions possible in PYTHIA: all particles are forced to color neutralize



PYTHIA Drell-Yan

- Can check if PYTHIA also reproduces CSS evolution with DY dimuon production
- Construct same observable
 $p_{out} = p_T^{lep} \sin \Delta\phi$
between two nearly back-to-back leptons
- PYTHIA confirms expectation from CSS evolution for same observable

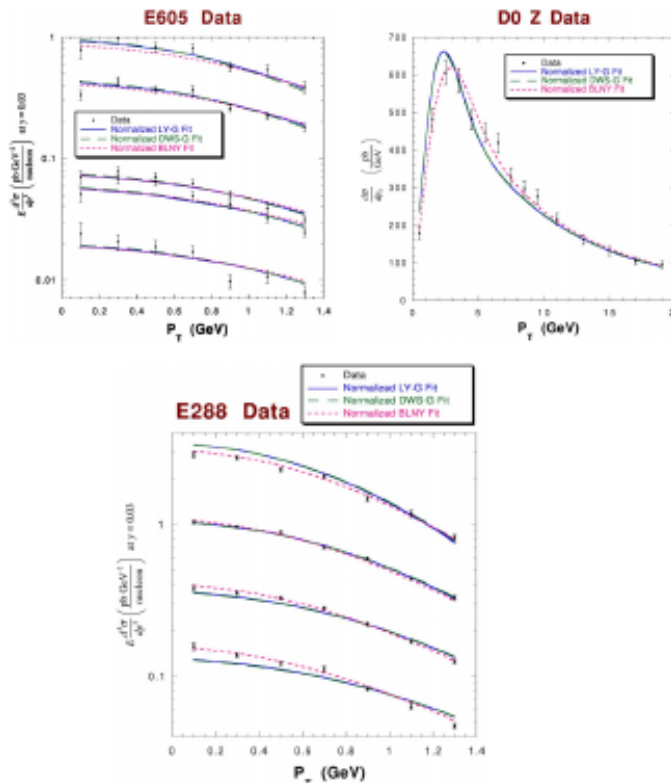


- Note rate of increase is significantly larger in magnitude also
- Red solid line shows log fit, blue dotted line shows linear fit

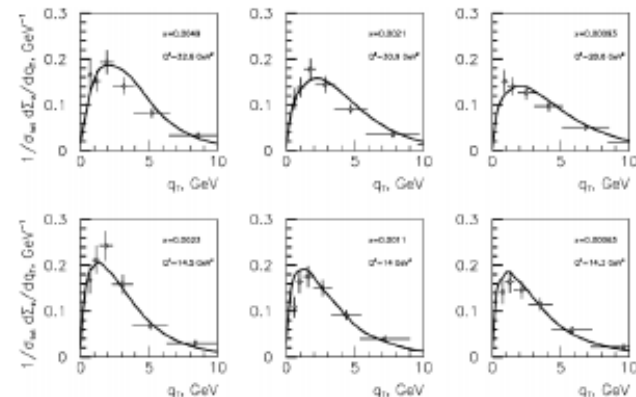
Nonperturbative momentum measurements in Drell-Yan and Z production

Other DY/Z and SIDIS Refs.

Phys. Rev. D 67, 073016 (2003)
(DY/Z)

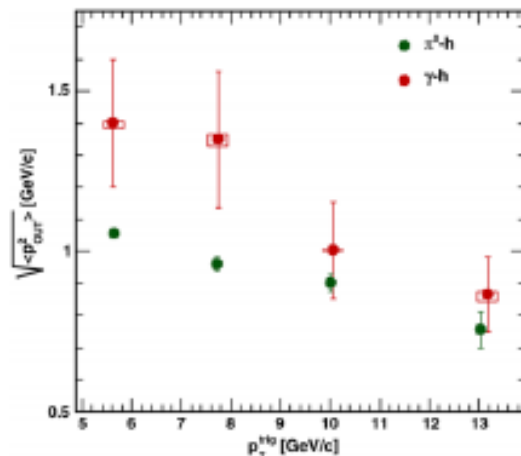


Phys. Rev. D 61, 014003 (2000)
(SIDIS)



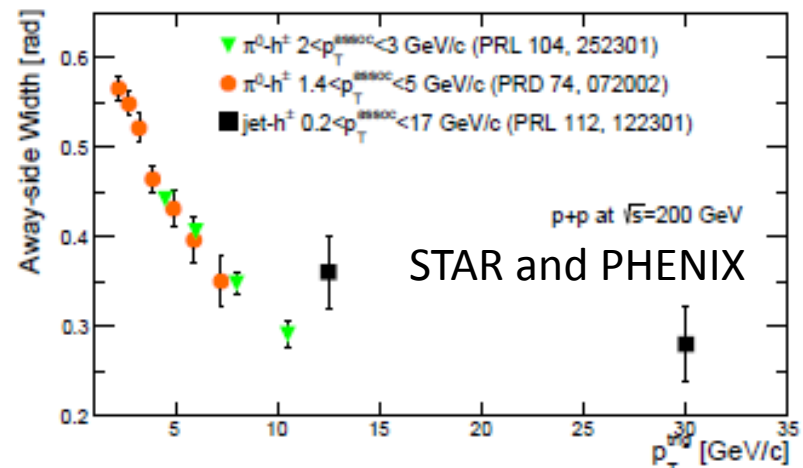
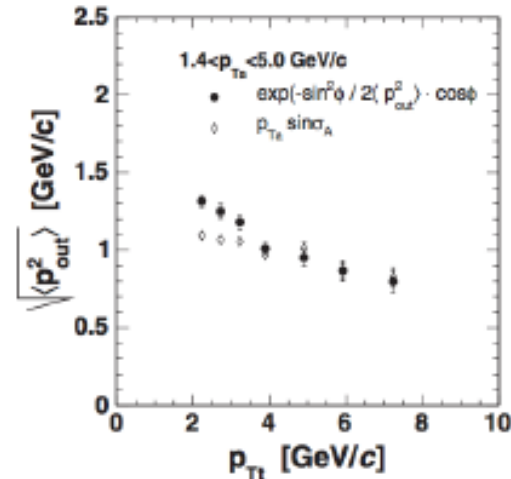
Other measurements showing decreasing nonperturbative momentum widths

- Other RHIC publications show the same effect in $\sqrt{\langle p_{out}^2 \rangle}$ and away-side width
- All previous analyses motivated by different physics goals: fragmentation functions, partonic energy loss in QGP, etc.



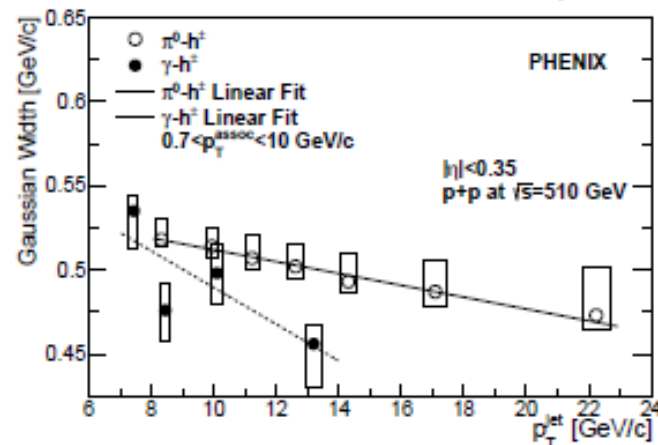
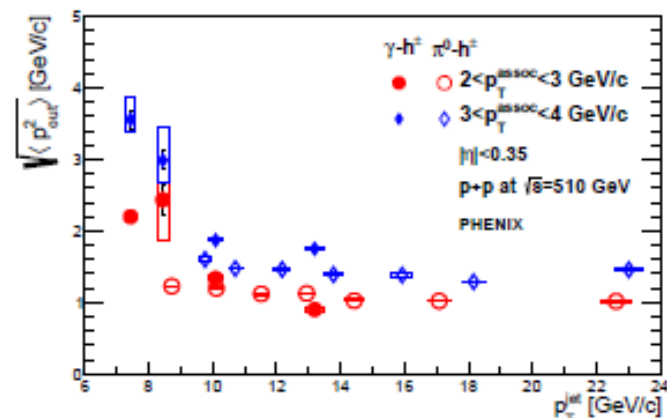
PRD 82, 072001 (2010) (PHENIX)

PRD 74, 072002 (2006) (PHENIX)

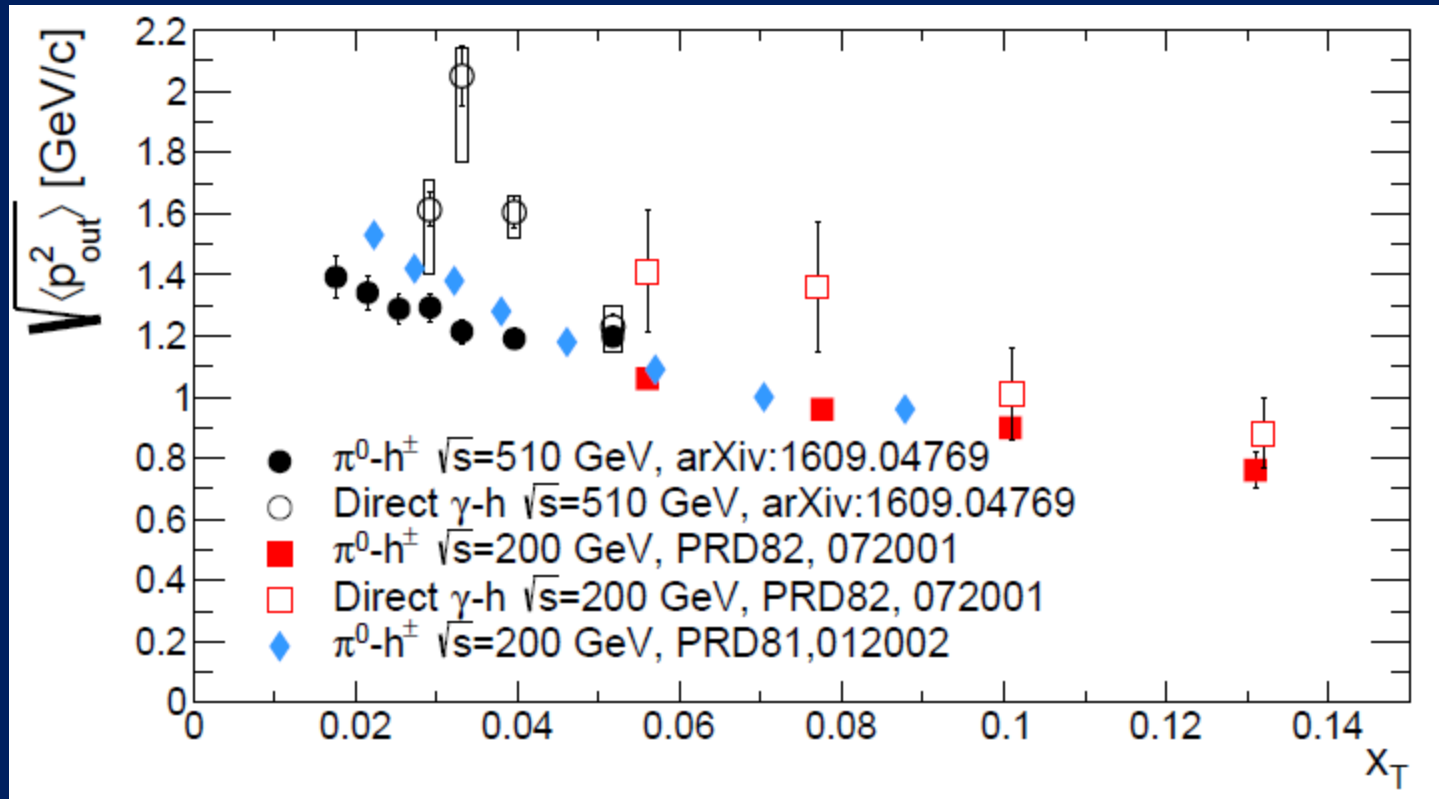


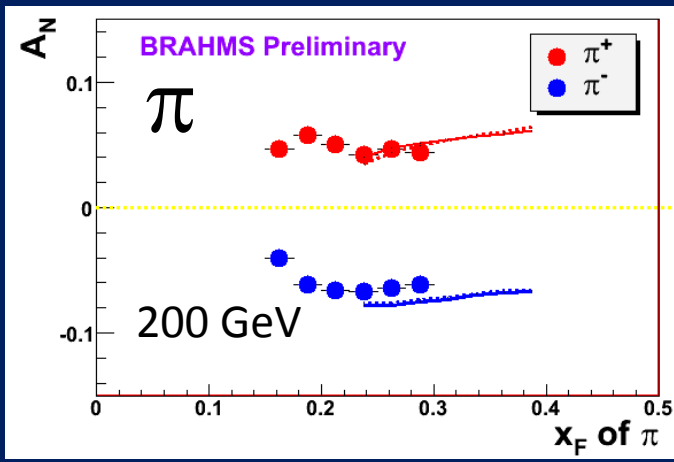
PYTHIA $\langle z_T \rangle$ Correction

- Direct photons emerge directly from hard scattering, π^0 s are a fragment
- Thus a more direct comparison is between p_T^{trig} for direct photon and jet p_T^{trig} for π^0
- Determine $\langle z_T \rangle = p_T^{\pi^0} / \hat{p}_T^{parton}$ using PYTHIA, "correct" π^0 p_T^{trig} to get $p_T^{jet} = p_T^{trig, \pi^0} / \langle z_T \rangle$

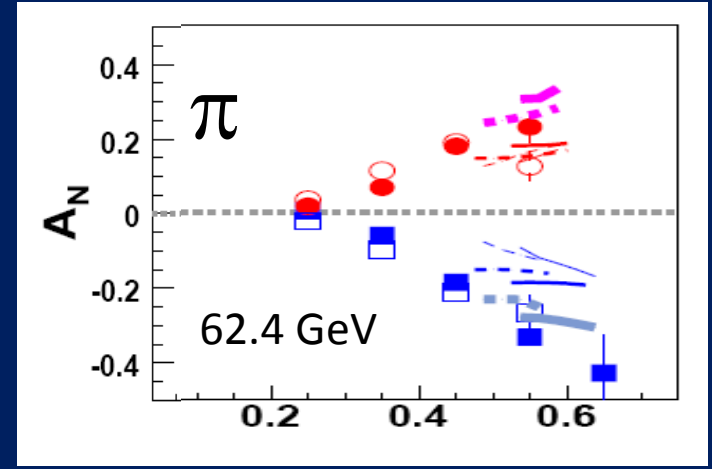


RMS p_{out} vs. x_T

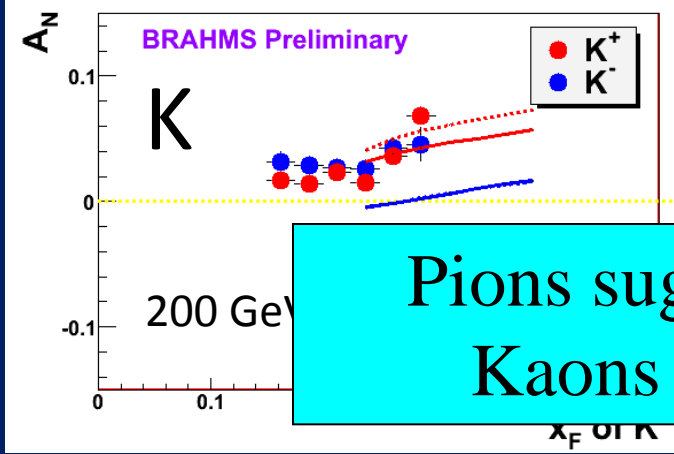




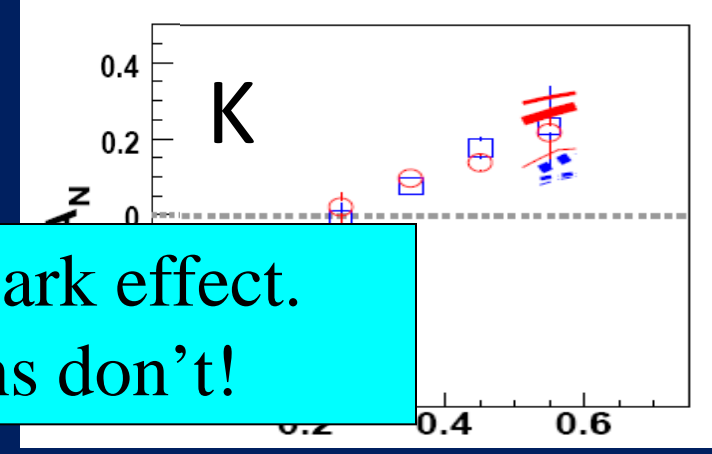
π, K, p
at 200 and
62.4 GeV



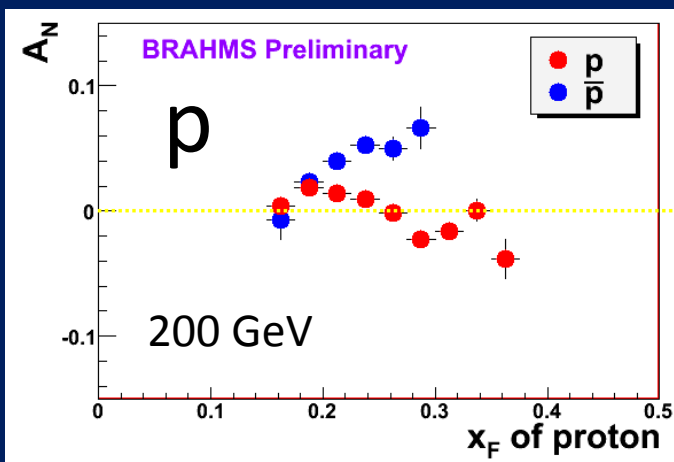
Note different scales



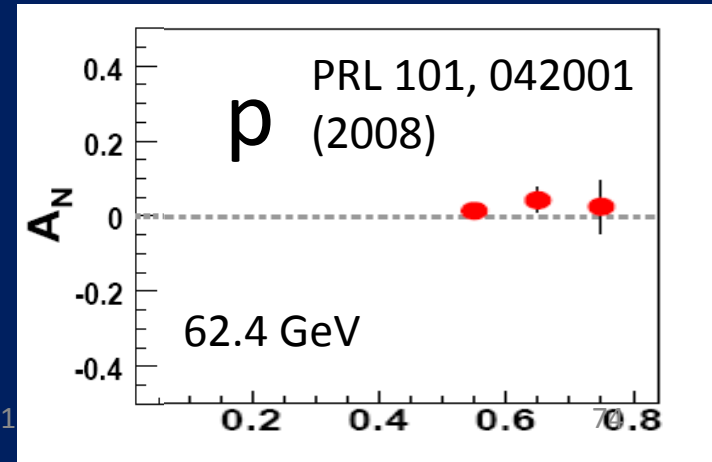
K- asymmetries
underpredicted



Pions suggest valence quark effect.
Kaons and (anti)protons don't!

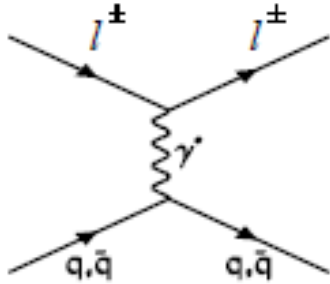


Large antiproton
asymmetry??
Unfortunately no 62.4
GeV measurement

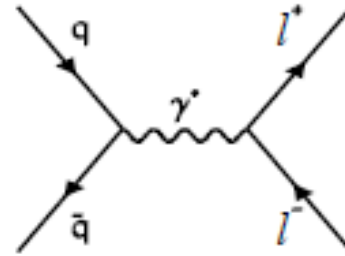


Complementarity of Drell-Yan and DIS

DIS

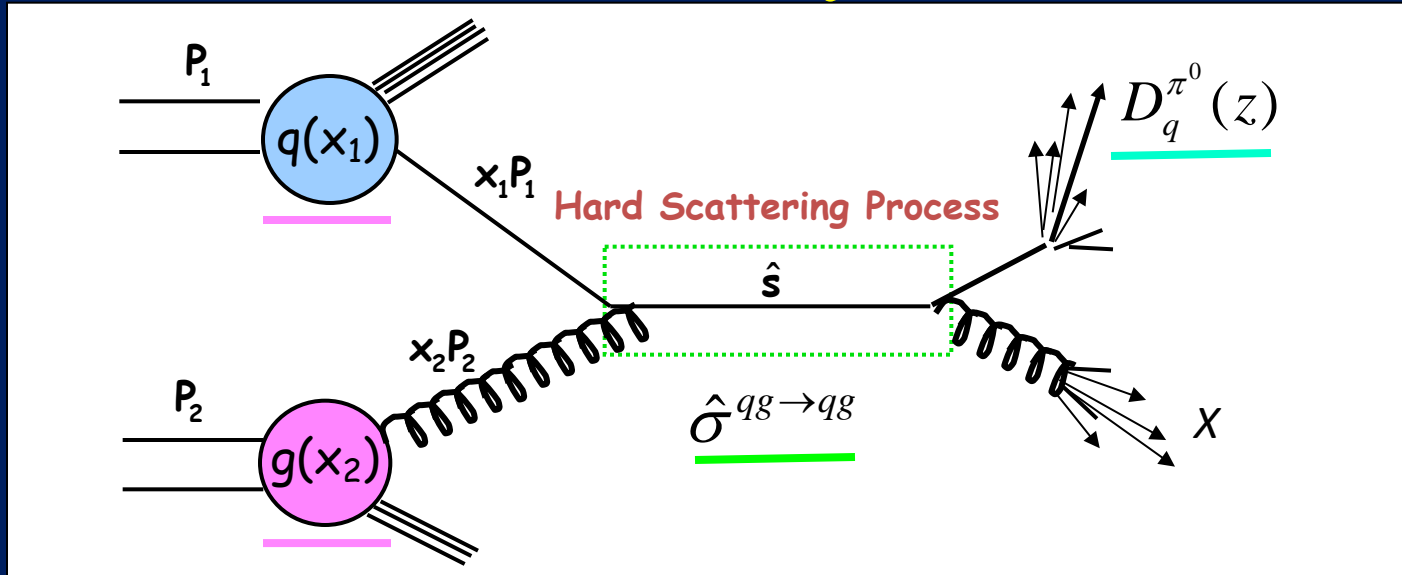


Drell-Yan



- Both deep-inelastic lepton-nucleon scattering (DIS) and quark-antiquark annihilation to leptons (Drell-Yan process) are tools to probe the quark and antiquark structure of hadrons

Parton distribution functions in perturbative QCD calculations of observables



$$\sigma(pp \rightarrow \pi^0 X) \propto \underline{q(x_1)} \otimes \underline{g(x_2)} \otimes \underline{\hat{\sigma}^{qg \rightarrow qg}(\hat{s})} \otimes \underline{D_q^{\pi^0}(z)}$$

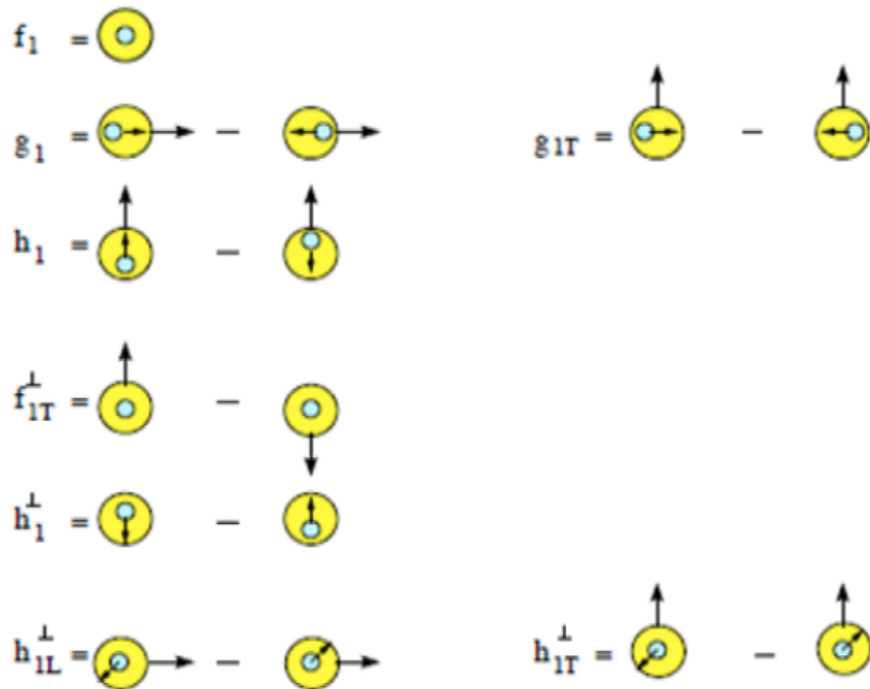
High-energy processes have predictable rates given:

- Partonic hard scattering rates (calculable in pQCD)
- Parton distribution functions (experiment or lattice)
- Fragmentation functions (experiment or lattice)

} Universal non-perturbative factors

Spin-spin and spin-momentum correlations in QCD bound states

U = unpolarized L = longitudinally polarized T = transversely polarized
N = nucleon q = quark



| N \ q | q | | |
|-------|----------------|----------------|--------------------|
| | U | L | T |
| U | f_1 | | h_1^\perp |
| L | | g_1 | h_{1L}^\perp |
| T | f_{1T}^\perp | g_{1T}^\perp | $h_1 h_{1T}^\perp$ |

Magnetic and electric A-B effects; Type-I and Type-II A-B effects

Box 1. Types and duals

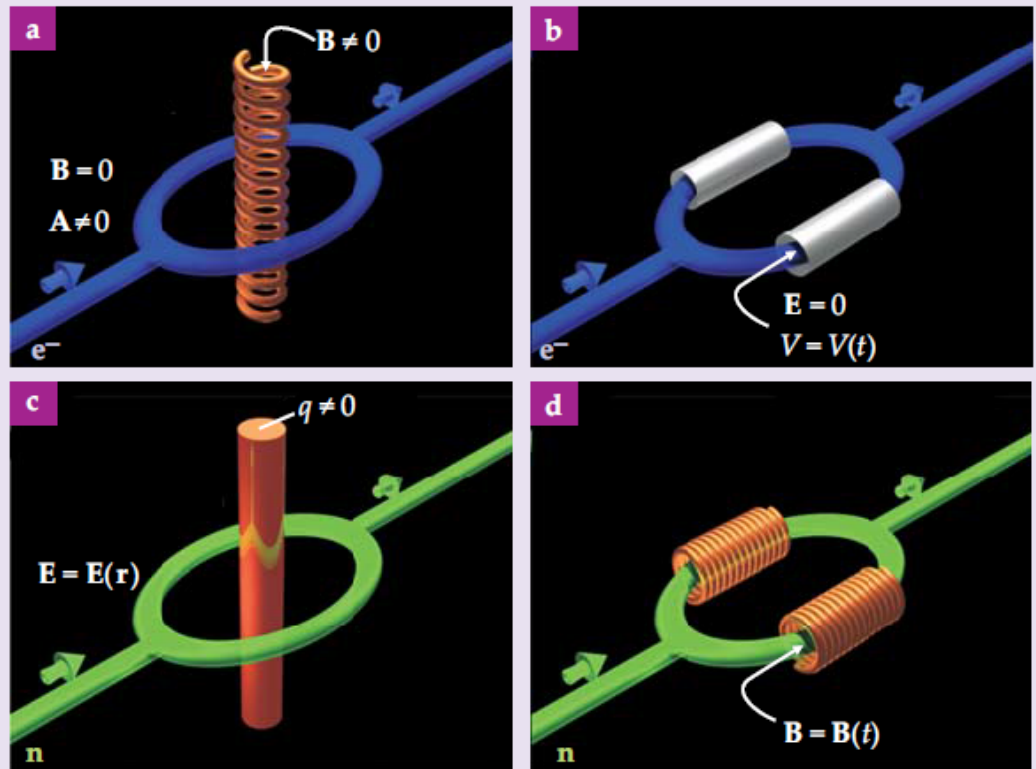
Physics Today, September 2009

The original magnetic and electric Aharonov–Bohm effects (panels a and b) are type I effects in the sense that in an ideal experiment, the electron sees no \mathbf{B} or \mathbf{E} fields, though it does traverse different potentials \mathbf{A} and V . In their respective dual effects—the Aharonov–Casher effect (panel c) and the so-called neutron-scalar AB effect (panel d)—polarized neutrons (neutral particles with magnetic dipole moments) replace unpolarized electrons, and electrostatic configurations change places with solenoids.¹⁰ In panel c, a neutron interferometer encloses a line of charge, and in panel d, neutrons pass through pulsed solenoids. These duals are classified as type II effects because the neutron must traverse a nonvanishing \mathbf{E} or \mathbf{B} field.

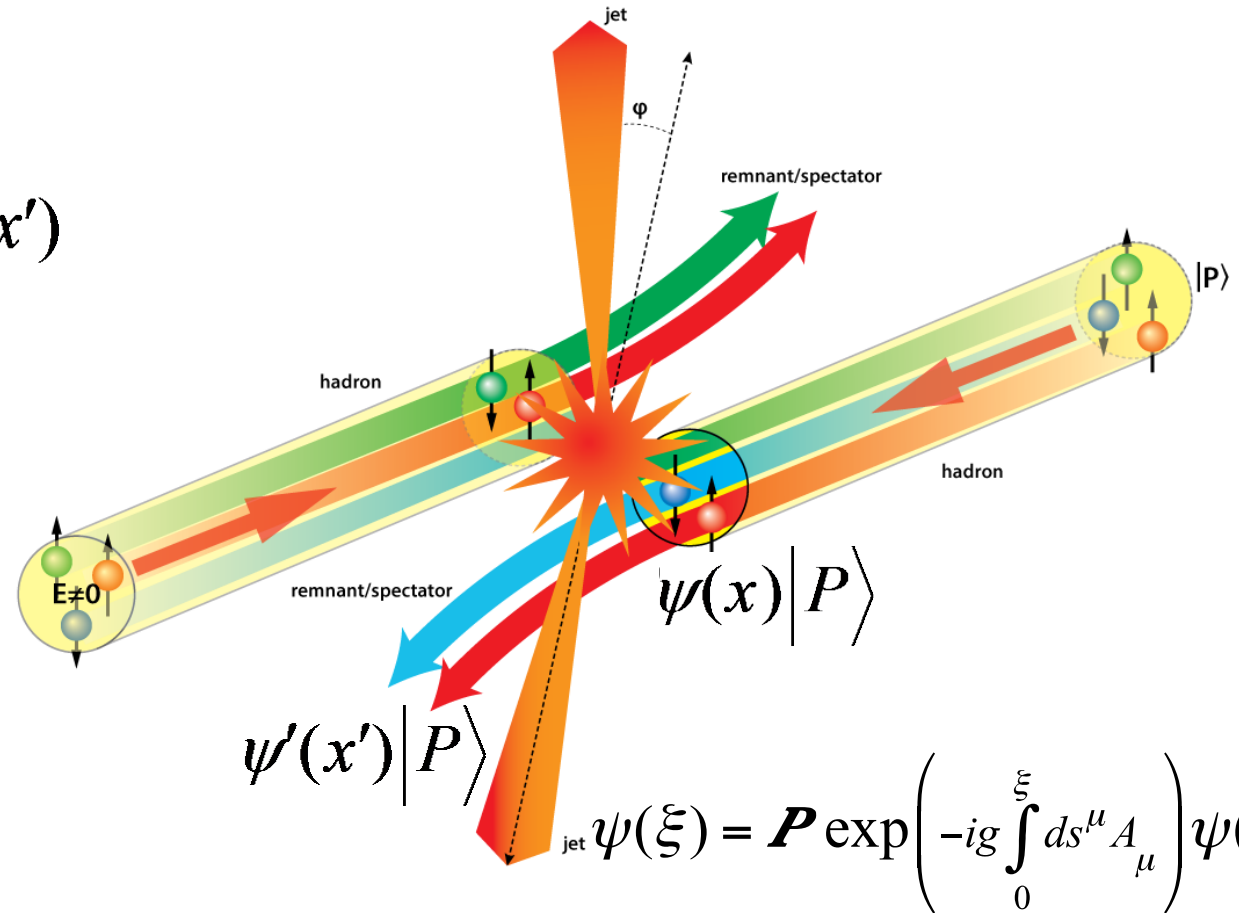
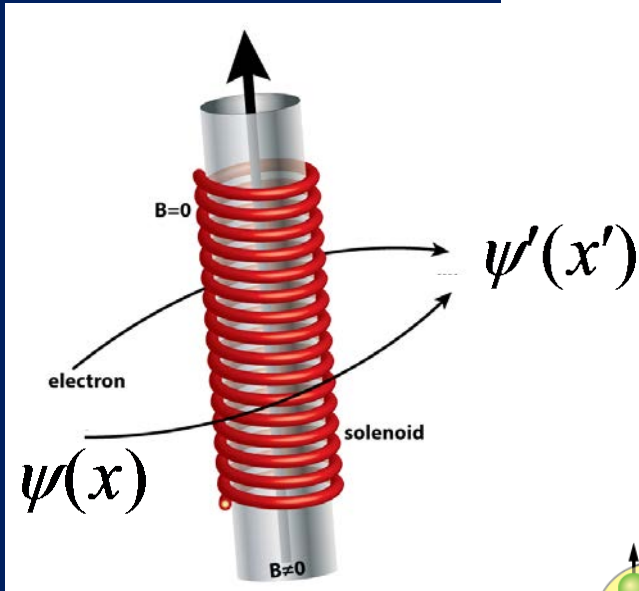
In either case, to acquire an AB phase shift, the electron or neutron must pass through a region of nonzero electromagnetic potential. That quantum mechanical result seems to elevate the status of the potentials to a physical reality absent from classical electromagnetism. Yakir Aharonov has pointed out that the potentials do overdetermine the experimental outcome; the phase shift need only be known modulo 2π . An alternative view is that the original magnetic AB effect shows electromagnetic fields acting nonlocally.¹

For type II effects, the wavepackets can plow straight through force fields, and forces are allowed in the interaction. But the AB interpretation requires

that the emerging wavepackets not be deflected or delayed in any way. Quantum mechanical descriptions generally circumvent the notion of forces. But one can use here an operational definition of forces that might be mimicking an AB effect: If the interaction has produced no deflection or delay, there were no forces.



Opportunities to see color-induced phases in QCD



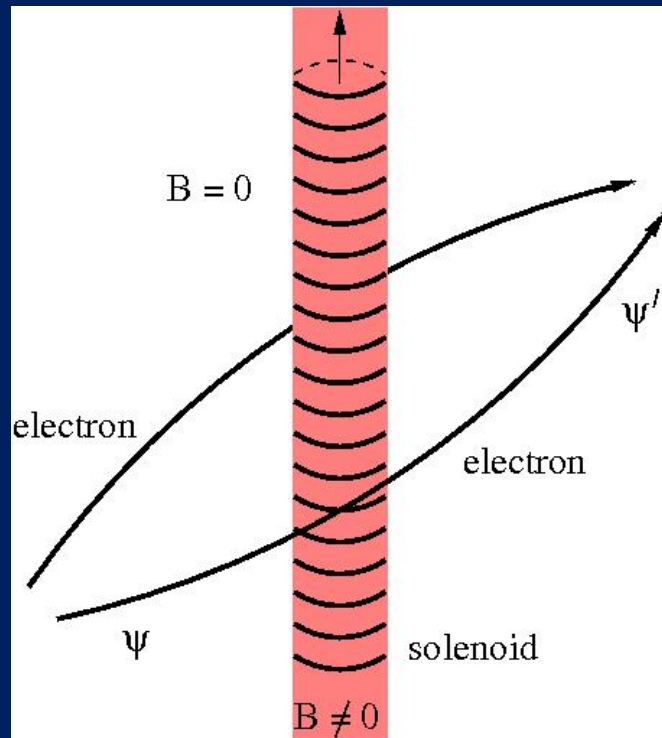
Slide from P. Mulders

Figures by Kees Huyser

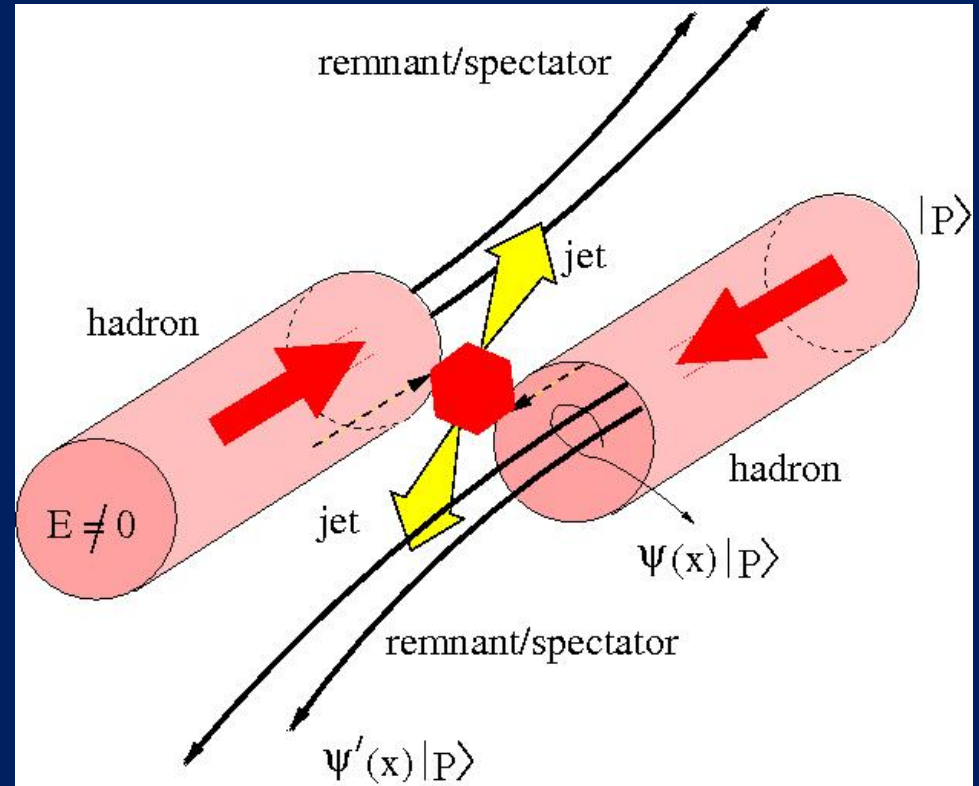
$$\psi(\xi) = \mathbf{P} \exp \left(-ig \int_0^\xi ds^\mu A_\mu \right) \psi(0)$$

Featuring: phases in gauge theories

Slide from P. Mulders



$$\psi' = P e^{ie \int ds \cdot A} \psi$$



$$\psi_i(x) |P\rangle = P e^{-ig \int_x^{x'} ds_\mu A^\mu} \psi_i(x') |P\rangle$$

Example: Fits to quark and gluon distributions including much wider range of data

- Incorporate corrections for target mass, “higher-twist,” and nuclear effects
- Can in turn make predictions for future measurements in extended kinematic regions

