Heavy Flavor Working Group Report

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Abstract
We present the findings of the RHIC II heavy flavor working group. The introduction briefly summarizes the main points, followed by sections with more detailed discussions of the detector and accelerator capabilities expected for RHIC II, and discussions of the physics issues for open heavy flavor and quarkonium in the RHIC and RHIC II context.

1. Introduction

1.1 Motivation

Because of the large charm and bottom quark masses, they are produced almost exclusively in the initial parton-parton interactions in heavy ion collisions at RHIC energies. In the absence of any nuclear effects, the heavy flavor cross sections in A+A collisions at RHIC would simply scale with the number of binary collisions. Thus departures from binary scaling for heavy flavor production in A+A collisions provide information about nuclear effects. These can be divided into two categories: effects due to embedding of the colliding partons in a nucleus (cold matter effects) and effects due to the large energy density in the final state. The main focus of the heavy flavor program at RHIC is to investigate the properties of the dense matter produced in A+A collisions by studying its effects on open heavy flavor and quarkonium production. This in turn requires a detailed understanding of cold matter effects so that they can be unfolded from the dense matter effects.

The program thus requires detailed measurements and calculations of p+p and p+A heavy flavor cross sections to characterize the cold matter effects if we are to quantify the differences between QGP and non-QGP effects. Up-to-date benchmark calculations of the total open heavy flavor (charm and bottom hadrons) and quarkonium (J/ψ and Υ families) yields and spectra are imperative. Cold matter effects that need to be included are nuclear shadowing, for both open heavy flavor and quarkonium production, and nuclear absorption of quarkonium. Recent calculations of charm and bottom production to FONLL in p+p collisions [1] have been published, along with a discussion of the theoretical uncertainties inherent in these calculations and reference calculations of heavy quark, heavy flavor meson and decay lepton spectra. Similar calculations have been done for quarkonium production, including studies of shadowing and absorption effects as a function of rapidity and centrality in d+Au [2] and A+A [3] collisions at RHIC.

A number of dense matter effects on heavy flavor production have been predicted. Some of these effects do not change the total cross section but, instead, modify the p_T spectra of heavy flavor hadrons and their decay products. Heavy quark energy loss [4, 5, 6, 7, 8] by collisional and radiative processes, steepens the p_T distribution relative to that in p+p collisions. On the other hand, random p_T kicks result in transverse momentum broadening, increasing the average p_T in both cold nuclear matter [9] and in passage through hadron bubbles in the mixed phase of a QGP [10]. If the medium surrounding the heavy quarks after production exhibits collective motion, such as transverse flow [11, 12], the low p_T heavy quarks ($p_T < m$) may be caught in this flow. Some of the effects of energy loss [13, 14] on heavy flavor decays to electrons and charm flow [14] have already been seen in Au+Au collisions at RHIC. Studying heavy flavor energy loss through the single electron channel requires being able to separate electrons

from c and b decays, since the large mass difference suggests that energy loss is weaker for bottom than for charm [4]. Some QGP studies require an accurate baseline for the total heavy flavor cross sections to interpret other effects. For example, if more than one $c\overline{c}$ pair is produced in an A+A event, uncorrelated c and $\overline{c}$ quarks might come together to form a J/ψ in a QGP [16, 17, 18, 19]. The total $c\overline{c}$ yield is needed to normalize the J/ψ production rate from this process.

Suppression of J/ψ production was one of the most exciting proposed QGP signatures at the CERN SPS [20]. J/ψ suppression was predicted to occur due to the shielding of the $c\overline{c}$ binding potential by color screening, leading to the break up of the quarkonium states, first the χ_c and ψ', and finally the J/ψ itself as the temperature increased [21, 22]. The QGP suppression may not be so simple, as lattice gauge theory studies of the J/ψ spectral function above the critical temperature for deconfinement, T_c, attest. The J/ψ may exist as a bound state for temperatures considerably larger than T_c [23]. However, the J/ψ may instead be dissociated by hot thermal gluons in medium [24] before it could be suppressed by color screening. Secondary quarkonium production from uncorrelated $Q\overline{Q}$ pairs, either in the plasma phase [17, 19, 26, 27, 28] or in the hadron phase [29, 30], could counter the effects of suppression, ultimately leading to enhanced quarkonium production. Such secondary J/ψ production would lead to different kinematic distributions than the initial production. Because the underlying $c\bar{c}$ distribution falls rapidly with p_T, the p_T distribution produced by coalescence will be softer. If the underlying $c\bar{c}$ distribution peaks at mid rapidity, coalescence will produce a narrower rapidity distribution also. The effect on the rapidity distribution should be calculated with shadowing effects on the underlying $c\bar{c}$ distribution taken into account, however, since these can cause the $c\bar{c}$ distribution to flatten in more central A+A collisions [3]. Elliptic flow effects are also expected on quarkonium production, as well as on open heavy flavors [11, 12].

With higher luminosity at RHIC, the Υ states could also be measured. Since the Υ radius is smaller than that of the J/ψ [22], direct color screening in the QGP would not occur until much higher temperatures. The higher bottomonium states, however, would likely be suppressed at RHIC, as are the χ_c and ψ' in the charmonium family. The feed down structure is more complicated for the Υ since there are three S states (Υ, Υ' and Υ'') and two sets of P states (χ_{b1} and χ_{b2}) below the $B\overline{B}$ threshold. The Υ family suppression should be measurable over a large p_T range, with QGP suppression possible on the Υ' and Υ'' up to $p_T \sim 40$ GeV/c [25]. Because of the small number of $b\bar{b}$ pairs at RHIC energy, bottomonium formation by coalescence of unrelated pairs should be negligible.

1.2 Overview of results from the heavy flavor program at RHIC

Heavy flavor measurements capable of discriminating between theoretical models need large integrated luminosity. In RHIC runs so far we have acquired useful data sets at 200 GeV for p+p, d+Au, Cu+Cu and Au+Au collisions. These data sets are not yet fully analyzed for Runs 4 (Au+Au) and 5 (Cu+Cu), where the highest luminosities were captured for heavy ions, but preliminary heavy flavor results are already available for all runs and species.

The data collected to date for p+p collisions provide an essential reference for the heavy ion program in the form of the underlying heavy flavor production rates as a function of rapidity and p_T. Equally essential, the data from d+Au collisions provide baseline information about cold nuclear matter effects which must also contribute to heavy flavor production in heavy ion collisions. The existing d+Au data provide useful tests of models that include the effects of shadowing on heavy flavor production and of absorption of J/ψ in cold nuclear matter [2].

Two very striking and unexpected results have already been seen for open heavy flavor in heavy ion collisions at RHIC. The first of these is the observation that the nuclear modification factor for electrons from open heavy flavor, R_{AA}, shows very strong suppression in central Au+Au collisions [31, 32], similar to that seen for pions. The second striking result is that the elliptic flow parameter (v_2) of electrons from open heavy flavor appears to favor charm quark flow at low p_T [31]. Until recently, it had been expected that heavy quark energy loss would be considerably smaller than that for light quarks

due to interference effects [4]. Generating the necessary energy loss for charm and bottom quarks with realistic gluon densities in the material is a major challenge for models [4, 33]. The relatively large v_2 values at low p_T imply at least some degree of charm quark equilibration with the medium. This also implies very strong interactions of charm quarks with the medium at lower p_T [11, 12].

The first high statistics charmonium results for heavy ion collisions at RHIC were presented at Quark Matter 2005 [34]. Preliminary results were shown for the J/ψ nuclear modification factors, R_{AA}, as a function of the number of participant nucleons for Cu+Cu and for Au+Au in the rapidity intervals $|y| < 0.35$ and $1.2 < |y| < 2.2$. Comparison with existing models shows that cold nuclear matter baseline calculations [3] which approximately reproduce the PHENIX d+Au J/ψ rapidity distributions [2] somewhat underpredict the suppression observed in Cu+Cu and Au+Au collisions. On the other hand, several suppression models [19, 35, 36] which were successful in describing J/ψ suppression at the SPS are found to strongly overpredict the suppression at RHIC. Models which incorporate strong suppression combined with J/ψ regeneration from uncorrelated $c\bar{c}$ pairs seem to agree best with the data, although the existing models slightly underpredict the suppression.

In the last few years, theorists have begun exploring the consequences of J/ψ regeneration by coalescence on observables other than the centrality dependence of the nuclear modification factor [16]. This work has led to the prediction that J/ψ's formed by coalescence of uncorrelated $c\bar{c}$ pairs will have narrower rapidity and p_T distributions due to the presumed shape of the underlying charm quark distributions. The coalescence contribution to J/ψ production will cause many observables to change with centrality, including the rapidity and p_T dependence of R_{AA}, the shape of the p_T distribution (quantified by the average p_T^2, $\langle p_T^2 \rangle$), and the J/ψ elliptic flow parameter, v_2. Quantitative predictions have been made for $\langle p_T^2 \rangle$ as a function of centrality for Au+Au and Cu+Cu collisions, with and without coalescence [16]. These predictions were compared to the preliminary Au+Au and Cu+Cu $\langle p_T^2 \rangle$ data for the first time at Quark Matter 2005. The data favor the calculations that include coalescence. On the other hand, the large coalescence contributions predicted for central Au+Au (and even central Cu+Cu) collisions are qualitatively expected to narrow the J/ψ rapidity distributions if the underlying charm distributions are peaked at mid rapidity. The preliminary data presented at Quark Matter 2005 show no evidence of this narrowing in central collisions. There is still work to do to quantify both the theoretical predictions and the experimental observables. The existing data sets will not provide a useful measurement of the J/ψ v_2 due to insufficient yield.

The running schedule for RHIC over the next five years is not settled. Based on the beam use proposal discussions prior to Run 6, it seems likely that more Au+Au data will be collected, providing up to an order of magnitude increase in integrated luminosity (if there are two more Au+Au runs). There will also be a very large increase in the integrated luminosity for p+p collisions due to the requirements of the Spin program. Such luminosity increases will quantitatively improve the measurements of many heavy flavor observables, especially as a function of centrality. The J/ψ R_{AA} and $\langle p_T^2 \rangle$ as well as R_{AA} and v_2 measurements of charm and bottom semileptonic decays to single electrons will all improve significantly, allowing more definitive tests of models. Measurements of other observables will be qualitatively improved. Examples are definitive semileptonic decay v_2 measurements at intermediate to high p_T, where we might hope to see the transition from charm to bottom dominance and flow to non-flow, a possible first J/ψ v_2 measurement, definitive measurements of J/ψ R_{AA} with rapidity, to quantify the coalescence contribution, and measurements of J/ψ R_{AA} to higher p_T, invaluable for understanding coalescence and formation time effects. Finally, it seems likely that a first, low statistics, Υ suppression measurement would be possible.

However, it is clear that the RHIC heavy flavor program will be limited by the capabilities of the accelerator after about 5 more years. The luminosity increase brought by RHIC II, combined with the detector upgrades in place by that time, will be required for the heavy flavor program at RHIC to move to the next level, as described below.

1.3 Overview of the proposed heavy flavor program at RHIC II

The order of magnitude increase in luminosity, combined with the increased power of the upgraded PHENIX and STAR detectors, will make it possible to add many powerful new probes to the heavy flavor program at RHIC.

One of the most powerful additions from the luminosity upgrade will be the ability to measure yields of the excited states of charmonium - the $\psi\prime$ and the χ_c states. Lattice calculations predict much smaller melting temperatures for the $\psi\prime$ and the χ_c than for the more tightly bound J/ψ, and these excited states should not be able to exist at all in the QGP at RHIC. Therefore comparison of the $\psi\prime$ and χ_c yields with the J/ψ yield as a function of centrality is considered to be a direct test for deconfinement.

Testing models in which the observed J/ψ yield in heavy ion collisions is due to competition between gluon dissociation of J/ψ and coalescence formation of J/ψ in the GQP requires very large luminosity. Tests of charm coalescence models include measuring J/ψ elliptic flow as a function of p_T, Measuring $J\psi$ R_{AA} to much higher p_T to see if the suppression goes away as the J/ψ formation time approaches the QGP crossing time, and measuring the J/ψ polarization as a function of collision centrality. Also of great interest is the rapidity and p_T dependence of the R_{AA} as a function of collision energy and centrality, a measurement that requires enough luminosity to make precision measurements at multiple energies - something that is not possible in a reasonable number of years at the present RHIC luminosity.

The study of bottomonium states, the Υ family, is possible only at RHIC II luminosities. Like the charmonium states, the bottomonium states are predicted by lattice calculations to melt at different temperatures due to their different binding energies. There are two important differences from charmonium. First, the binding energies are higher for the bottomonium states than for charmonium, so they are predicted to become unbound at a higher temperature. Only the higher bottomonium states are likely to be unbound at RHIC energy. Second, The $b\bar{b}$ production rate in a central Au+Au collision is only $\sim$ 0.05 pairs per collision, so that coalescence production of bottomonium is much less likely than it is for charmonium. The result is that bottomonium production at RHIC II will provide a very different window on color screening effects than charmonium production. The bottomonium yields at RHIC II should be sufficient to allow a measurement of R_{AA} vs centrality in heavy ion collisions for the three Υ states. Υ yields at RHIC II and at the LHC will not be sufficient for measurements of v_2 or polarization.

As mentioned earlier, measurements of semi-leptonic decay products from open heavy flavor at RHIC have already produced results that were strikingly different from expectations. The strongly suppressed R_{AA} and large v_2 suggest very large energy loss of heavy quarks in the medium. But these semi-leptonic decay spectra suffer from the complication that they contain both charm and bottom contributions. The separation of open charm and bottom can be done in several ways. Charm can be observed via the $D \rightarrow K\pi$ hadronic decays, and this has been done by STAR. But a precise measurement of R_{AA} and v_2 by this method will be difficult. These events can not be triggered on, and so they must be taken from a minimum bias data set that samples only a small fraction of the available luminosity. Also the combinatorial background is very large, making statistical precision difficult. The addition of a displaced vertex measurement in STAR will dramatically reduce combinatorial background, but the problem of not being able to trigger on these events has yet to be solved. At RHIC II luminosity, bottom can be observed very cleanly in both PHENIX and STAR via $B \rightarrow J/\psi$ decays with a displaced vertex measurement. This should provide good measurements of the open bottom cross section and R_{AA}. However the yields will likely be too small to allow v_2 measurements at RHIC II or the LHC. Finally, the combination of RHIC II luminosity and a displaced vertex measurement should allow the statistical separation of charm and beauty contributions to the semi-leptonic decay spectra, using the different decay lengths of c and b quarks. Such semi-leptonic decay measurements, while less clean than the direct D and B decay measurements, have the advantage of much larger yields, and separate v_2 measurements for charm and bottom should be possible.

The independent measurement of open charm and bottom R_{AA} and v_2 to high p_T at RHIC II will

be very important. At low p_T the v_2 and R_{AA} reflect the degree of thermalization of the heavy quarks in the medium. At high p_T, the heavy quark R_{AA} and v_2 must reflect the energy loss of heavy quarks in the medium, providing an independent measurement of the energy density from that given by light quark energy loss measurements. The thermalization mechanism for low p_T and energy loss mechanism for high p_T may be quite different, due to the possible importance of resonance scattering at low p_T. RHIC II will provide the opportunity to measure separately the v_2 and R_{AA} of b quark semileptonic decay products, and to extend measurements of the v_2 and R_{AA} of c quark semileptonic decay products to much higher p_T.

1.4 Overview of the relationship of RHIC II to the LHC program

The heavy flavor production cross sections are significantly higher at the LHC than at RHIC since the Pb+Pb energy at the LHC is a factor of 27.5 higher than the maximum Au+Au energy at RHIC. The $c\overline{c}$ and $b\overline{b}$ cross sections increase by factors of 15 and 100 respectively [9] while the J/ψ and Υ cross sections increase by factors of 13 and 55 respectively [37]. But, because of the higher luminosity and the longer heavy ion runs, the integrated luminosity for Au+Au at RHIC II is projected to be 36 times higher than for Pb+Pb at LHC. Therefore the heavy flavor yields per year are expected to be similar.

At $\sqrt{s} = 200$ GeV, bottom decays to leptons begin to dominate the single electron spectrum at $p_T \sim 4$ GeV. As the collision energy increases, the lepton spectra from B and D decays move closer together rather than further apart [9]. Thus, the large increase in the $b\overline{b}$ cross section relative to $c\overline{c}$ does not make single leptons from B and D decays easier to separate. Preliminary calculations show that the $B \to e$ decay does become larger than that of $D \to e$, but at $p_T > 10$ GeV. The two lepton sources differ by less than a factor of two to $p_T \sim 50$ GeV in the range $|y| \leq 1$. Thus interpretation of single lepton results on heavy flavors will be more difficult at the LHC. Other means of separating charm and bottom must be found. ALICE can reconstruct hadronic D^0 decays from $p_T \sim 0$ to $p_T \sim 25$ GeV [38] but, like STAR, will have to rely on minimum bias data for these measurements because of the lack of a trigger. While it is not yet clear what CMS and ATLAS will do to reconstruct charm, they should be able to do b jets well, similar to the Tevatron measurements. One way that B mesons can be measured at the LHC is through their decays to J/ψ, as discussed further below. It has also been suggested that the $B\overline{B}$ contribution to the dimuon continuum, the dominant contribution above the Υ mass, can be used to measure energy loss [39]. That channel would be fairly clean at the LHC but more difficult at RHIC.

The RHIC II upgrades and the high LHC energies make detailed studies of Υ production and suppression possible. At the LHC, higher initial temperatures make Υ suppression more likely than at RHIC II. But the higher $b\bar{b}$ production rate (~ 5 per central Pb+Pb collision) means that, unlike RHIC, significant coalescence contributions to Υ production are expected at the LHC. Thus measurements at the two energies complement each other. At RHIC II, it is likely that PHENIX will be able to measure and resolve the three Υ S states. STAR will see Υ yields similar to those in ALICE, but the mass resolution will require fitting to extract yields. At the LHC, all 3 S states will also be measurable. CMS has the mass resolution to separate all three. The Υ states can be measured to $p_T \sim 0$ at all LHC detectors. Only ALICE will be able to measure J/ψ production to $p_T \sim 0$ without a special trigger [37] since CMS and ATLAS require high single muon p_T so that typically only J/ψ with $p_T > 5$ GeV are accepted. (However, CMS is working on a higher-level trigger to measure lower p_T J/ψ's [40].) The larger $b\overline{b}$ cross section at the LHC means that J/ψ production from $B \to J/\psi X$ cannot be neglected. These decay J/ψ's should be separable from the initial production using displaced vertices [37].

2. Detector upgrade program at RHIC

Both PHENIX and STAR have extensive upgrade programs underway that are extremely important for the heavy flavor program. The upgrades that are most relevant to heavy flavor measurements are described here. The impact on the heavy flavor program of these detector upgrades, in combination with the RHIC II luminosity increase, will be discussed in sections 4 and 5.

PHENIX upgrades:

Several PHENIX detector upgrades that greatly enhance the heavy flavor capability of the experiment are expected to be available in the RHIC II time frame.

The Silicon Vertex Detector (SVTX) [41] consists of a central arm barrel and two endcap detectors covering the muon arm acceptance. The SVTX will provide a displaced vertex resolution of 50 μm, and will provide inner tracking with full azimuthal coverage for $|\eta| < 2.4$. By connecting to tracks in the central and muon arms it will be able to tag heavy flavor decays using displaced vertices, improve the quarkonium invariant mass resolution, eliminate decays of light hadrons, and reduce high p_T background. In addition to identifying semileptonic heavy flavor decays via displaced vertices, the same technique can be used to identify J/ψ from B meson decays, since all other J/ψ's are prompt. The displaced vertex measurement will also help greatly with $D \to K\pi$ measurements, presently very difficult in PHENIX, by reducing the background dramatically.

The muon trigger upgrade [42] is required for PHENIX to be able to take complete advantage of the RHIC II luminosity upgrade. Consisting of three layers of Resistive Plate Chamber (RPC) detectors in the muon arms, it provides an online momentum measurement to improve the level 1 trigger rejection for both single muons and muon pairs. It will also provide improved high multiplicity background rejection.

The PHENIX Nose Cone Calorimeters [43], tungsten-silicon calorimeters that will replace the central arm magnet nosecones, will provide coverage for $0.9 < |\eta| < 3.5$. They should provide good acceptance for $\chi_c \to J/\psi + \gamma$ decays, as well as allow $\gamma\pi^0$ separation for $p_T < 30$ GeV/c and jet identification.

STAR upgrades:

While work on answering the questions discussed in this document are underway in STAR, to complete many of the challenging measurements, upgrades to the STAR detector are needed. The collaboration has planned a series of upgrades for the near and intermediate term to overcome the current shortcomings and enhance its heavy flavor capabilities. Implementation of these upgrades will also allow optimum utilization of the increased luminosity expected from RHIC II.

The current layout of the STAR detector is depicted in Fig. 1. The medium term upgrades to the detector relevant for heavy flavor physics include: a full barrel Time of Flight detector (TOF) replacing the current ToF patch and the Central Trigger Barrel (CTB); new front end electronics for the large Time Projection Chamber (TPC); an upgrade to the data acquisition system (DAQ-1000), and a tracking upgrade including a barrel section with two inner layers of silicon pixel sensors (HFT) and three layers of silicon strip detectors (IST).

The new time of flight system covering the full outer barrel of the TPC is planned for construction and installation in STAR over the next three years. The system uses the Multi-gap Resistive Plate Chamber (MRPC) technology developed at CERN and will consist of 3840 MRPC modules with 23,000 channels of readout. The modules will cover the TPC outer barrel ($-1 < \eta < 1$, $0 < \phi < 2\pi$) and will be mounted in 120 trays which will replace the existing CTB (Central Trigger Barrel scintillation counter) trays.

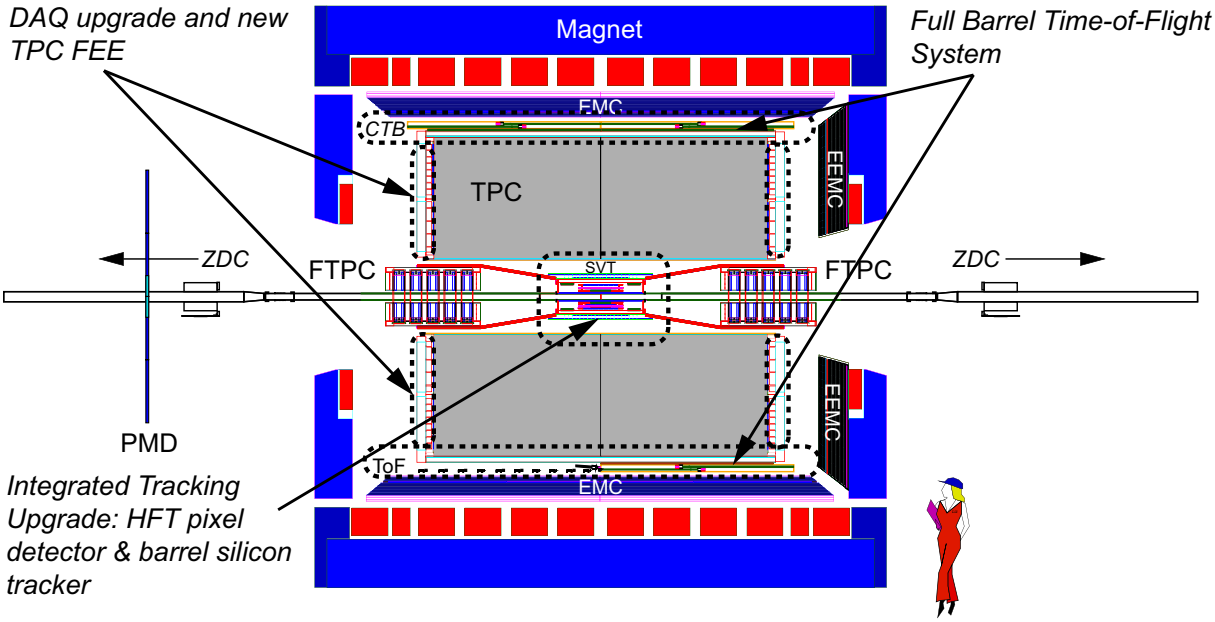

Fig. 1: Layout of the STAR experiment 2005/2006 [44]. The locations of the planned upgrades are circled by dashed lines. See text for details.

The ToF doubles the current momentum range over which π ,K, and p can be identified and thus improves the reconstruction of charmed mesons and baryons considerably. When the TOF measurement is combined with the TPC dE/dx measurement, electrons can be cleanly identified from the lowest momentum measured ($\sim 200\ \mathrm{MeV}/c$) up to a few GeV/c. This capability complements the electromagnetic calorimeter which works well for momenta above $\sim 2\ \mathrm{GeV}/c$. It will allow STAR to reconstruct soft to medium momentum electrons with high efficiency and purity thus providing the capability to conduct a comprehensive J/ψ measurement. The TOF in conjunction with the EMC also allows STAR to implement a level-2 trigger scheme to select $J/\psi \rightarrow e^+e^-$decays in A+A collisions.

A series of improvements to the STAR data acquisition system over the past several years has brought the capability from the original design rate of 1 Hz recorded events to 50-100 Hz. To acquire the very large data samples and high data rates needed for heavy flavor measurements an upgrade has been initiated with the goal of achieving a recorded event rate of at least 1 kHz. This rate could produce data volumes which would significantly exceed the capacity for analysis and storage. The rare-trigger data sets will especially benefit from the upgrade since the pipelined architecture being implemented will virtually eliminate the front end dead time allowing to make full use of rate triggers such as the one for the Υ.

This increase in read out speed can be achieved by replacing the TPC frontend electronics (making use of circuits developed for the ALICE experiment at CERN) in conjunction with an upgrade of the STAR DAQ. In addition to the increased physics capabilities from the DAQ upgrade, the replacement of the TPC front end electronics, specifically the RDO boards that collect data from the FEE boards, will make space for a future precision tracking chamber between the TPC end planes and the end-cap calorimeter. Replacing the TPC front end electronics also assures that this system will be maintainable for the next decade or more. The readout for the other existing detectors which will remain in place for the RHIC II era can be adapted to the new high speed DAQ with minor changes.

In order to address questions of heavy quark energy loss and thermalization it will be necessary to cleanly identify open charm. The recent results from both STAR and PHENIX on the suppression and flow of non-photonic electrons are intriguing; however, without an identified sample of charm the questions of contribution from semi-leptonic open-beauty decays and systematic errors on background

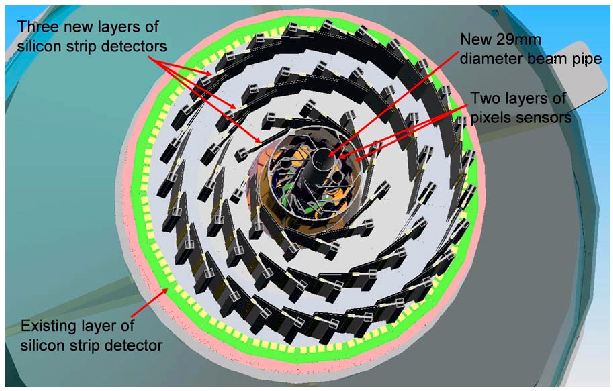

Fig. 2: Layout of the new inner tracker for STAR. End view of small radius beam pipe (1.4 cm radius), two layers of pixel sensor, three new layers of silicon strip sensors, and the existing layer of silicon strip sensor.

subtraction make a clear interpretation of these results difficult. Measurement of the yields of various charm species will also allow a study of the charm hadro-chemistry.

Efficient topological reconstruction of open charm decays requires tracking ”point-back” resolution to the primary collision vertex of $\sim 50\mu$m or better. Further, the beam pipe and innermost layers of detector must be very thin to allow measurement of particles at low transverse momentum which comprise the bulk of the cross section. This will minimize the systematic errors in extrapolating the measured yield to the total yield. A thin beam pipe and inner detector layers are also key elements in efficiently vetoing photon conversion electrons which in combination with the electron identification from the TPC, TOF and electromagnetic calorimeter will enable measurement of the soft lepton and di-lepton spectra. To achieve this STAR is developing a tracking upgrade for the central rapidity region ($-1 < \eta < 1, 0 < \phi < 2\pi$). The essential elements under consideration for this upgrade are a new thin, small-radius beam pipe (0.5 mm thick, 14 mm radius), two layers of thinned (50 μm) CMOS pixel detectors at average radii of 1.5 and 4.5 cm (HFT) and three layers of conventional silicon strip detectors at average radii of 10, 15 and 20 cm (IST) (see Fig. 2). The existing layer of double-sided silicon strip sensors at a radius of 25 cm (SSD) will be kept. The three new layers of conventional silicon strip sensors (IST) will allow connecting tracks from the TPC and SSD to hits in the pixel layers. These layers will replace the existing three layers of silicon drift detector (SVT). It will be necessary to replace the SVT since when RHIC II becomes operational, the SVT will be over 10 years old, its readout is too slow to be compatible with the upgraded DAQ, and it has a large amount of infrastructure (cables and cooling) in the region $1 < \eta < 2$.

3. Projected RHIC II yields

In this section we present some estimates of the quality of the heavy flavor measurements that can be achieved at RHIC II luminosities with the upgraded detectors.

Table 1: The anticipated luminosity per week delivered by RHIC. The delivered luminosity has to be reduced by a factor that accounts for detector up time and collision vertex cuts imposed by the detectors. The RHIC I projected luminosities are maximum values achieved by 2008 taken from the RHIC Collider Accelerator Division projections. They represent the performance of a mature RHIC I accelerator. Because the length of the collision diamond is smaller for RHIC II, the gain in usable luminosity is larger than the ratio of delivered luminosities when going to RHIC II. There are no projections available for d+Au and Cu+Cu for RHIC I. The numbers in the "obtained" column are the best weekly luminosities from previous runs.

Species	Energy	Units	RHIC I Obtained/week	RHIC I Projected/week	RHIC II Projected/week
Au+Au	200	μb^{-1}	160	327	2500
Cu+Cu	200	nb^{-1}	2.4	–	25
d+Au	200	nb^{-1}	4.5	–	62
p+p	200	pb^{-1}	0.9	26	33
p+p	500	pb^{-1}	-	50	166

Table 1 is a summary of the weekly integrated luminosity estimates for RHIC and RHIC II. The weekly luminosity expectations are taken from RHIC Collider Accelerator Division guidance. The projected weekly luminosities for RHIC II and for RHIC I in 2008 and beyond (projected RHIC I) are used to estimate the yields shown in later tables.

Table 2 summarizes the projected PHENIX yields for critical heavy flavor signals for the mature RHIC acceleratot (in 2008 and beyond) and for RHIC II for a 12 week physics run. Table 2 also includes the yields observed in recent RHIC runs. The yields are based on the following:

- The cross sections are taken from Ref. [45].
- The ratio of ψ' to J/ψ yields is 0.14.
- The detector acceptances are from PHENIX simulations.
- For charmonia, the cross sections were reduced in Au+Au by a factor of 0.43, which approximately accounts for the suppression measured by PHENIX in Au+Au collisions. for the Υ family, no suppression is assumed.
- The minimum bias trigger efficiency for hard processes is assumed to be 0.75 for p+p and 0.92 for Au+Au.
- An additional realistic level 1 trigger efficiency of 0.8 is used, where appropriate.
- It is assumed that 80% of the RHIC beam is in the central bucket and thus usable by experiments.
- It is assumed that the root-mean square (RMS) of the collision diamond is 20 cm at RHIC and 10 cm at RHIC II.
- Use of the PHENIX vertex detector is assumed for the projected yields, requiring a collision vertex cut of ± 10 cm.
- A realistic pair-reconstruction efficiency is used (0.8 in p+p and 0.4 in Au+Au collisions).
- An efficiency of 0.4 is assumed when using a displaced vertex cut of 1 mm to identify $B \to J/\psi$ decays.

Table 3 contains a summary of the projected yields from STAR for various critical heavy flavor signals for RHIC and for RHIC II. The detector acceptances are from STAR simulations. Otherwise, the assumptions are identical to those used for the PHENIX yields presented in table 2.

Table 2: The projected yields of several heavy flavor signals in PHENIX for 12 week physics runs at RHIC II and a mature RHIC I. The approximate yields obtained at RHIC to date are also shown. These reflect the fact that RHIC had not yet achieved the full luminosity development for Au+Au by Run 4, or for p+p by Run 5. The yields are shown for both p+p and Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV. The projected RHIC I and RHIC II values assume that the PHENIX SVTX detector is in place, limiting the usable collision vertex range to ± 10 cm. The SVTX detector has a much larger impact at RHIC I, where the collision diamond RMS is 20 cm, than at RHIC II where the collision diamond RMS is 10 cm.

Species	signal	$\lvert\eta\rvert$	To Date	Projected RHIC I	Projected RHIC II
p+p	$J/\psi \to e^+e^-$	< 0.35	$\sim$ 1,000	25,000	55,000
	$J/\psi \to \mu^+\mu^-$	$1.2 - 2.4$	$\sim$ 10,000	208,000	470,000
	$\psi' \to e^+e^-$	< 0.35	–	440	990
	$\psi' \to \mu^+\mu^-$	$1.2 - 2.4$	–	3,700	8,500
	$\chi_c \to e^+e^-\gamma$	< 0.35	–	1,600	3,600
	$\chi_c \to \mu^+\mu^-\gamma$	$1.2 - 2.4$	–	62,000	139,000
	$\Upsilon \to e^+e^-$	< 0.35	–	90	200
	$\Upsilon \to \mu^+\mu^-$	$1.2 - 2.4$	$\sim$ 27	230	500
	$B \to J/\psi \to e^+e^-$	< 0.35	–	130	300
	$B \to J/\psi \to \mu^+\mu^-$	$1.2 - 2.4$	–	1,300	3,000
Au+Au	$J/\psi \to e^+e^-$	< 0.35	$\sim$ 800	3,300	45,00
	$J/\psi \to \mu^+\mu^-$	$1.2 - 2.4$	$\sim$ 7,000	29,000	395,00
	$\psi' \to e^+e^-$	< 0.35	–	60	800
	$\psi' \to \mu^+\mu^-$	$1.2 - 2.4$	–	520	7,100
	$\chi_c \to e^+e^-\gamma$	< 0.35	–	220	2,900
	$\chi_c \to \mu^+\mu^-\gamma$	$1.2 - 2.4$	–	8,600	117,000
	$\Upsilon \to e^+e^-$	< 0.35	–	30	400
	$\Upsilon \to \mu^+\mu^-$	$1.2 - 2.4$	–	80	1,040
	$B \to J/\psi \to e^+e^-$	< 0.35	–	40	570
	$B \to J/\psi \to \mu^+\mu^-$	$1.2 - 2.4$	–	420	5,700

Table 3: The projected yields of several heavy flavor signals in STAR for 12 week physics runs at RHIC II and a mature RHIC I. The approximate yields obtained at RHIC to date are also shown. The projected RHIC I and RHIC II values assume that the STAR Heavy Flavor Tracker is in place, limiting the usable collision vertex range to ± 10 cm. The $D \to K\pi$ yields assume 100 Hz of recorded minimum bias data.

Species	signal	$\lvert\eta\rvert$	To Date	Projected RHIC I	Projected RHIC II
p+p	$J/\psi \to e^+e^-$	< 1.0	–	1,260,000	1,600,000
	$\psi' \to e^+e^-$		–	23,000	29,000
	$\Upsilon \to e^+e^-$		–	6,600	8,300
	$B \to J/\psi \to e^+e^-$		–	15,000	19,000
Au+Au	$J/\psi \to e^+e^-$	< 1.0	?	16,000	220,000
	$\psi' \to e^+e^-$		–	300	4,000
	$\Upsilon \to e^+e^-$		?	830	11,200
	$B \to J/\psi \to e^+e^-$		–	190	2,500
	$D \to K\pi$		–	30,000	30,000

Table 4: Projected heavy flavor yields in PHENIX and STAR for 12 weeks of $\sqrt{s_{NN}} = 500$ GeV p+p running at RHIC II and a mature RHIC I. The projected RHIC I and RHIC II values assume that both the the PHENIX SVTX detector and STAR HF tracker are in place, limiting the usable collision vertex range to ± 10 cm. These detectors have a much larger impact at RHIC I, where the collision diamond RMS is 20 cm, than at RHIC II where the collision diamond RMS is 10 cm.

Experiment	signal	$\|\eta\|$	Projected RHIC I	Projected RHIC II
PHENIX	$J/\psi \to e^+e^-$	< 0.35	183,000	600,000
	$J/\psi \to \mu^+\mu^-$	$1.2 - 2.4$	1,650,000	5,500,000
	$\psi' \to e^+e^-$	< 0.35	3,300	11,000
	$\psi' \to \mu^+\mu^-$	$1.2 - 2.4$	30,000	100,000
	$\chi_c \to e^+e^-\gamma$	< 0.35	31,000	100,000
	$\chi_c \to \mu^+\mu^-\gamma$	$1.2 - 2.4$	1,200,000	4,800,000
	$\Upsilon \to e^+e^-$	< 0.35	900	3000
	$\Upsilon \to \mu^+\mu^-$	$1.2 - 2.4$	2,300	7,700
	$B \to J/\psi \to e^+e^-$	< 0.35	2,300	7,700
	$B \to J/\psi \to \mu^+\mu^-$	$1.2 - 2.4$	23,000	77,000
STAR	$J/\psi \to e^+e^-$	< 1.0	3,700,000	12,000,000
	$\psi' \to e^+e^-$		76,000	220,000
	$\Upsilon \to e^+e^-$		25,000	84,000
	$B \to J/\psi \to e^+e^-$		346,000	1,100,000

Table 5: The estimated heavy flavor yields in the LHC experiments for a 10^6 s Pb+Pb run with 500 μb^{-1} integrated luminosity (the planning numbers for one year). The estimates were all reported by the LHC experiments. As for the RHIC experiment tables, the Υ rates are for the three states combined. The range of values for ATLAS corresponds to different assumptions about trigger thresholds.

Species	signal	ALICE	$\|\eta\|$	CMS	$\|\eta\|$	ATLAS	$\|\eta\|$
Pb+Pb	$J/\psi \to \mu^+\mu^-$	740,000	2.5-4	24,000	< 2.4	8K-100K	< 2.5
	$J/\psi \to e^+e^-$	9,500	< 0.9				
	$\psi' \to \mu^+\mu^-$	14,000	2.5-4	440	< 2.4	1,400-1,800	< 2.5
	$\psi' \to e^+e^-$	190	< 0.9				
	$\Upsilon \to \mu^+\mu^-$	8,400	2.5-4	26,000	< 2.4	15,000	< 2.5
	$\Upsilon \to e^+e^-$	2,600	< 0.9				
	$D \to K\pi$	8,000	< 0.9				

Table 4 contains a summary of expected heavy flavor yields at $\sqrt{s_{NN}} = 500$ GeV for PHENIX and STAR. Although not directly comparable with heavy ion yields from collisions at $\sqrt{s_{NN}} = 200$, the order of magnitude larger heavy flavor yields at 500 GeV should help greatly with understanding the reaction mechanisms in p+p collisions.

Table 5 contains a summary of the projected yields from the LHC detector collaborations for various critical heavy flavor signals in a $\sqrt{s_{NN}} = 5.5$ TeV Pb+Pb run of 10^6 s, which is the standard planning number for a year of running [46, 47, 48]. Note that the estimates by the LHC collaborations generally assume more optimistic reconstruction efficiencies than those used for the RHIC detectors.

Comparison of tables 2, 3 and 5 reveal that the projected heavy flavor yields for one year of running are similar at the LHC and at RHIC II. The much larger heavy flavor cross sections at the higher LHC energy are largely compensated at RHIC II by the integrated luminosities that result from three times longer runs and order of magnitude higher luminosity.

4. Open heavy flavor

In this section we present a more detailed discussion of the theoretical motivation for studying open heavy flavor in HI collisions, of the present experimental and theoretical status, and of the proposed experimental program of open heavy flavor measurements at RHIC II.

As described in the Introduction, dense matter effects in nuclear collisions may change the kinematic distributions and the total cross sections of open heavy flavor production. Effects such as energy loss and flow can significantly modify the heavy flavor p_T distributions but do not, in fact, change the total yields. In a finite acceptance detector, however, the measured yields may appear to be enhanced or suppressed, depending on the acceptance. Energy loss steepens the slope of the heavy flavor p_T distribution because the heavy quark p_T is reduced. If the momentum is reduced sufficiently for the quarks to be stopped within the medium, the heavy quarks can take the same velocity as the surrounding medium and 'go with the flow'. The first RHIC results on R_{AA} and v_2 for heavy flavor decays to leptons show that these effects are indeed important for charm quarks. However, higher p_T measurements and reconstructed charm hadrons are needed to solidify and quantify the results. In addition, reconstructed bottom measurements are necessary to cleanly separate leptons from charm and bottom decays to determine the bottom quark's importance in the measured electron R_{AA}.

Effects that may modify the total heavy flavor yields are the initial parton distributions in the nuclei and secondary charm production in the medium. The parton distribution functions needed for perturbative QCD calculations of heavy flavor production are modified in the nucleus, as was observed in nuclear deep-inelastic scattering [49]. At very small momentum fractions, x, the gluon fields may be treated as classical color fields. The modifications of the parton distributions in nuclei relative to free protons would affect the total yields. The effect is expected to be small at mid rapidity and moderate p_T at RHIC but is likely to be more important at large rapidity, where lower x values are probed. Although thermal charm production from the medium is likely to be small at RHIC energies, it could moderately enhance the total yields.

Since the J/ψ yields may be enhanced in nuclear collisions due to coalescence of uncorrelated c and $\overline{c}$ quarks in the medium, it is important for charmonium production in heavy ion collisions to be properly normalized. The ratio of J/ψ to open charm production in p+p collisions is not a strong function of energy. Thus the total charm yield sets the scale against which J/ψ suppression relative to enhancement can be quantified. Production of J/ψ through $c\overline{c}$ coalescence is discussed in more detail in the section on quarkonium.

4.1 Open Heavy Flavor Theoretical Results

4.11 Theoretical Baseline Results

We now discuss the most recent theoretical baseline calculations of the transverse momentum (p_T) distributions of charm and bottom quarks, the charm and bottom hadron distributions resulting from fragmentation and, finally, the electrons produced in semileptonic decays of the hadrons [1].

The theoretical prediction of the electron spectrum includes three main components: the p_T and rapidity distributions of the heavy quark Q in p+p collisions at $\sqrt{S} = 200$ GeV, calculated in perturbative QCD; fragmentation of the heavy quarks into heavy hadrons, H_Q, described by phenomenological input extracted from e^+e^- data; and the decay of H_Q into electrons according to spectra available from other measurements. This cross section is schematically written as

$$\frac{Ed^3\sigma(e)}{dp^3} = \frac{E_Q d^3\sigma(Q)}{dp_Q^3} \otimes D(Q \to H_Q) \otimes f(H_Q \to e) \quad (1)$$

where the symbol $\otimes$ denotes a generic convolution. The electron decay spectrum, $f(H_Q \to e)$, accounts for the branching ratios.

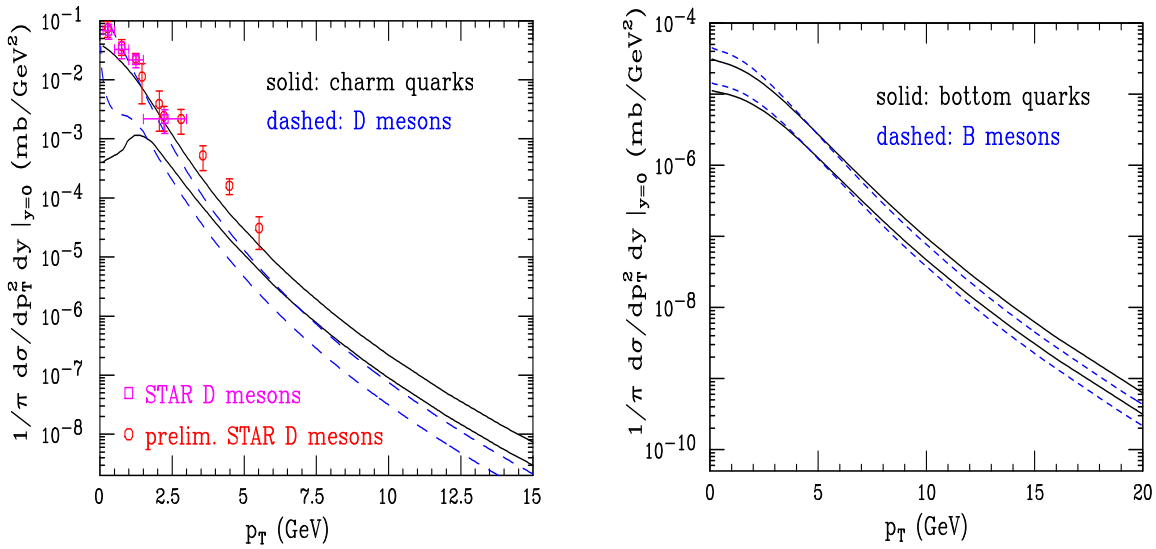

Fig. 3: Left-hand side: The theoretical uncertainty bands for c quark and D meson p_T distributions in p+p collisions at $\sqrt{S} = 200$ GeV, using BR($c \to D$) = 1. The fi nal [50] and preliminary [51] STAR d+Au data (scaled to p+p using $N_{\rm bin}$ = 7.5) are also shown. Right-hand side: The same for b quarks and B mesons.

The distribution $Ed^3\sigma(Q)/dp_Q^3$ is evaluated at Fixed-Order plus Next-to-Leading-Log (FONLL) level, implemented in Ref. [52]. In addition to including the full fixed-order NLO result [53, 54], the FONLL calculation also resums [55] large perturbative terms proportional to $\alpha_s^n \log^k(p_T/m)$ to all orders with next-to-leading logarithmic (NLL) accuracy (i.e. $k = n,\, n-1$) where m is the heavy quark mass. The perturbative parameters are m and the value of the strong coupling, α_s. We take $m_c = 1.5$ GeV and $m_b = 4.75$ GeV as central values and vary the masses in the range $1.3 < m_c < 1.7$ GeV for charm and $4.5 < m_b < 5$ GeV for bottom to estimate the mass uncertainties. The five-flavor QCD scale is the CTEQ6M value, $\Lambda^{(5)} = 0.226$ GeV. The perturbative calculation also depends on the factorization (μ_F) and renormalization (μ_R) scales. The scale sensitivity is a measure of the perturbative uncertainty. We take $\mu_{R,F}^2 = \mu_0^2 = p_T^2 + m^2$ as the central value and vary the two scales independently within a 'fiducial' region defined by $\mu_{R,F} = \xi_{R,F}\mu_0$ with $0.5 \le \xi_{R,F} \le 2$ and $0.5 \le \xi_R/\xi_F \le 2$ so that $\{(\xi_R, \xi_F)\}$ = {(1,1), (2,2), (0.5,0.5), (1,0.5), (2,1), (0.5,1), (1,2)}. The envelope containing the resulting curves defines the uncertainty. The mass and scale uncertainties are added in quadrature.

These inputs lead to a FONLL total $c\bar{c}$ cross section in p+p collisions of $\sigma_{c\bar{c}}^{\rm FONLL} = 256^{+400}_{-146}$ μb at $\sqrt{S} = 200$ GeV. The theoretical uncertainty is evaluated as described above. The corresponding NLO prediction is 244^{+381}_{-134} μb. The predictions in Ref. [9], using $m_c = 1.2$ GeV and $\mu_R = \mu_F = 2\mu_0$ gives $\sigma_{c\bar{c}}^{\rm NLO} = 427$ μb, within the uncertainties. Since the FONLL and NLO calculations tend to coincide at small p_T, which dominates the total cross section, the two results are very similar. Thus the two calculations are equivalent at the total cross section level, within the large perturbative uncertainties. The total cross section for bottom production is $\sigma_{b\bar{b}}^{\rm FONLL} = 1.87^{+0.99}_{-0.67}$ μb.

The fragmentation functions, $D(c \to D)$ and $D(b \to B)$, where D and B indicate a generic admixture of charm and bottom hadrons, are consistently extracted from e^+e^- data in the context of FONLL [56].

The measured spectra for primary $B \to e$ and $D \to e$ decays are assumed to be equal for all bottom and charm hadrons, respectively. The contribution of electrons from secondary B decays, $B \to D \to e$, was obtained by convoluting the $D \to e$ spectrum with a parton-model prediction of $b \to c$ decay. The resulting electron spectrum is very soft, giving a negligible contribution to the total. The decay spectra are normalized using the branching ratios for bottom and charm hadron mixtures [57]: BR$(B \to e) = 10.86 \pm 0.35\%$, BR$(D \to e) = 10.3 \pm 1.2\%$, and BR$(B \to D \to e) = 9.6 \pm 0.6\%$.

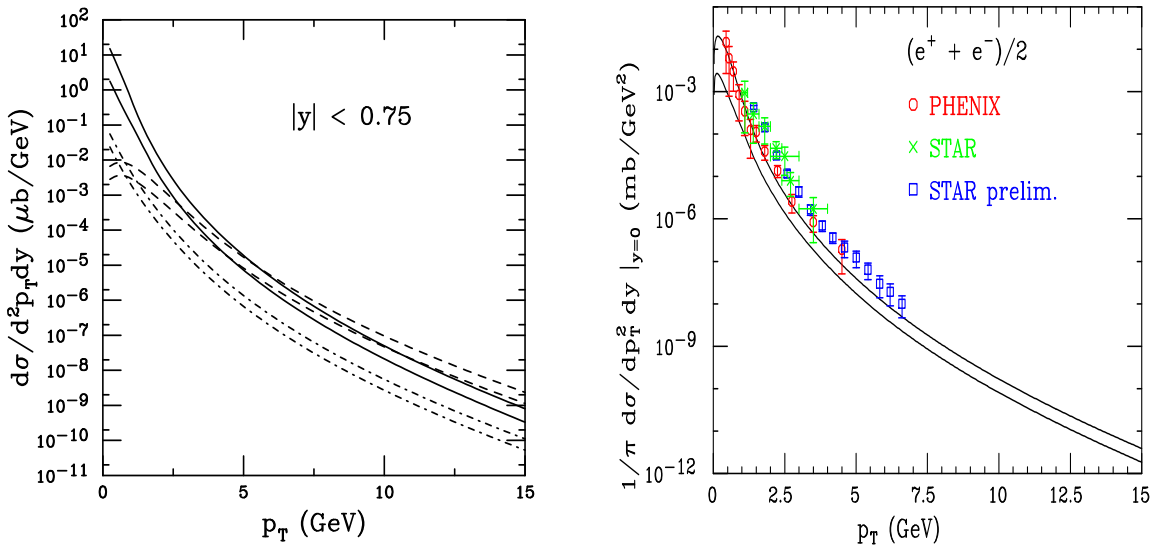

Fig. 4: Left-hand side: The theoretical uncertainty bands for $D \to e$ (solid), $B \to e$ (dashed) and $B \to D \to e$ (dot-dashed) as a function of p_T in $\sqrt{S} = 200$ GeV p+p collisions for $|y| < 0.75$. Right-hand side: The final electron uncertainty band in p+p collisions is compared to the PHENIX [58] and STAR (final [50] and preliminary [51]) data.

The left-hand side of Fig. 3 shows the theoretical uncertainty bands for c quarks and D mesons, obtained by summing the mass and scale uncertainties in quadrature. The band is broader at low p_T due to the large value of α_s and the behavior of the CTEQ6M parton densities at low scales as well as the increased sensitivity of the cross section to the charm quark mass. The rather hard fragmentation function causes the D meson and c quark bands to separate only at $p_T > 9$ GeV. The right-hand side of Fig. 3 shows the same results for b quarks and B mesons. The harder $b \to B$ fragmentation function causes the two bands to partially overlap until $p_T \simeq 20$ GeV/c.

Fig. 4 shows the individual uncertainty bands for the $D \to e$, $B \to e$ and $B \to D \to e$ decays to electrons on the left-hand side and compares the RHIC data to the total band on the right-hand side. The upper and lower limits of the total band are obtained by summing the upper and lower limits for each component. The secondary $B \to D \to e$ spectrum is extremely soft, only exceeding the primary $B \to e$ decays at $p_T < 1$ GeV. It is always negligible with respect to the total yield. While, for the central parameter sets, the $B \to e$ decays begin to dominate the $D \to e$ decays at $p_T \simeq 4$ GeV/c, a comparison of the bands shows that the crossover may occur over a rather broad range of electron p_T.

4.2 Models of Heavy Quark Energy Loss

While the heavy quarks are in the medium, they can undergo energy loss by two means: elastic collisions with light partons in the system (collisional) and gluon bremsstrahlung (radiative). We will briefly review some of the predicted results for $-dE/dx$ of heavy quarks for both collisional and radiative loss. We then show the predicted effect on the charm and bottom contributions to single electrons at RHIC [4].

The collisional energy loss of heavy quarks through processes such as $Qg \to Qg$ and $Qq \to Qq$ depends logarithmically on the heavy quark momentum, $-dE/dx \propto \ln(q_{\max}/q_{\min})$. Treatments of the collisional loss vary with the values assumed or calculated for the cutoffs. These cutoffs are sensitive to the energy of the heavy quark and the temperature and strong coupling constant in the medium. Thus the quoted value of the energy loss is usually for a certain energy and temperature. The calculation was first done by Bjorken [59] who found $-dE/dx \approx 0.2$ GeV/fm for a 20 GeV quark at $T = 250$ MeV. Further work refined the calculations of the cutoffs [60, 61, 62], with similar results. Braaten and Thoma calculated the collisional loss in the limits $E \ll m_Q^2/T$ and $E \gg m_Q^2/T$ in the hard thermal loop approximation, removing the cutoff ambiguities. They obtained $-dE/dx \approx 0.3$ GeV/fm for a 20 GeV

charm quark and 0.15 GeV/fm for a 20 GeV bottom quark at $T = 250$ MeV [63].

Other models of heavy quark energy loss were presented in the context of J/ψ suppression: Could a produced $c\overline{c}$ pair stay together in the medium long enough to form a J/ψ? Svetitsky [64] calculated the effects of diffusion and drag on the $c\overline{c}$ pair in the Boltzmann approach and found a strong effect. The drag stopped the $c\overline{c}$ pair after traveling about 1 fm but Brownian diffusion drove them apart quickly. The diffusion effect increased at later times. Essentially he predicted that the heavy quarks would be stopped and then go with the flow. His later calculations of D meson breakup and rehadronization [10] while moving through plasma droplets reached a similar conclusion. Koike and Matsui calculated energy loss of a color dipole moving through a plasma using kinetic theory and found $-dE/dx \sim 0.4 - 1.0$ GeV/fm for a 10 GeV $Q\overline{Q}$ [65]. The collisional loss was thus predicted to be small, less than 1 GeV/fm for reasonable assumptions of the temperature. The loss increases with energy and temperature. Using the hard thermal loop approach, Mustafa *et al.* found $-dE/dx \approx 1 - 2$ GeV/fm for a 20 GeV quark at $T = 500$ MeV [66].

The first application of radiative loss to heavy quarks was perhaps by Mustafa *et al.* [66]. They included the effects of only a single scattering/gluon emission, $Qq \to Qqg$ or $Qg \to Qgg$. In this case, the loss grows as the square of the logarithm $\ln^2(q_{\max}/q_{\min})$, one power more than the collisional loss, but is of the same order in the strong coupling constant [63]. Thus the radiative loss is guaranteed to be higher than the collisional in this approximation. The heavy quark mass enters their expressions only in the definition of $q_{\max}$ so that the mass dependence of the energy loss is rather weak. They found, for a 20 GeV quark at $T = 500$ MeV, $-dE/dx \approx 12$ GeV/fm for charm and 10 GeV/fm for bottom.

These large values suggested that energy loss could be quite important for heavy quarks. If true, there would be a strong effect on the $Q\overline{Q}$ contribution to the dilepton continuum. Shuryak [67] was the first to consider this possibility for A+A collisions. He assumed that low mass $Q\overline{Q}$ pairs would be stopped in the medium, suppressing the dilepton contribution from these decays substantially. However, the stopped heavy quarks should at least expand with the medium rather than coming to rest, as discussed by Svetitsky [64]. Lin *et al.* then calculated the effects of energy loss at RHIC, including thermal fluctuations, for a constant $-dE/dx = 0.5 - 2$ GeV/fm [8]. These results showed that the heavy quark contributions to the dilepton continuum would be reduced albeit not completely suppressed. In any case, the energy loss does not affect the total cross section. The heavy quarks are thus piled up at low p_T and at mid rapidity if stopped completely.

Dokshitzer and Kharzeev pointed out that soft gluon radiation from heavy quarks is suppressed at angles smaller than $\theta_0 = m_Q/E$ [7]. Thus bremsstrahlung is suppressed for heavy quarks relative to light quarks by the factor $(1 + \theta_0^2/\theta^2)^{-2}$, the 'dead cone' phenomenon. The radiative energy loss of heavy quarks could then be quite small. However, Armesto *et al.* [6] later showed that medium-induced gluon radiation could 'fill the dead cone', leading to non-negligible energy loss for heavy flavors. They also found that the energy loss would be larger for charm than bottom quarks.

So far the RHIC measurements are not for heavy flavored hadrons but for the electrons from their semileptonic decays. If the effects of energy loss are substantially different for charm and bottom quarks, then the results in Fig. 4 which show that at high p_T, the single electron spectrum is dominated by b decays, would suggest that, if charm quarks lose more energy than bottom quarks, this b-quark dominance of the electron spectra would begin at smaller values of electron p_T in A+A collisions. This would, in turn, limit the electron suppression factor, R_{AA}, at moderate p_T since the large bottom contribution would make R_{AA} larger than expected if the spectrum arose primarily from charm quark decays.

The left-hand side of Fig. 5 compares the c and b distributions at mid rapidity, as well as their contributions to single electrons. Single electrons from bottom dominate the single electron spectra at $p_T \sim 5$ GeV for all gluon rapidity densities. This conclusion is further supported by the right-hand side of Fig. 5, where the ratio of charm relative to bottom decays to electrons is shown. In all cases,

[0] His drag coefficient $A(p^2)$ is related to the energy loss per unit length through $A(p^2) = (-dE/dx)/p^2$.

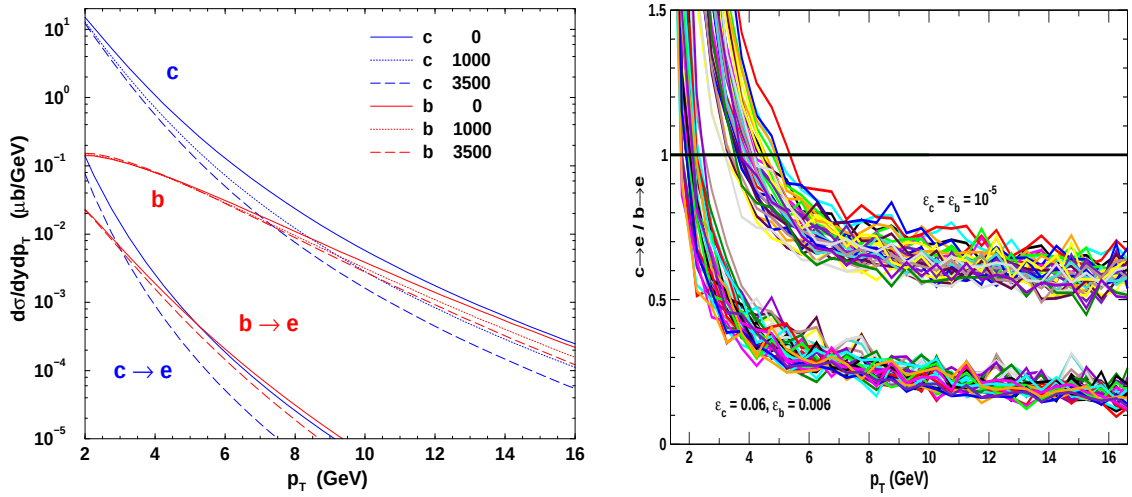

Fig. 5: Left-hand side: The differential cross section (per nucleon pair) of charm and bottom quarks calculated to NLO in QCD [1] compared to single electron distributions calculated with the fragmentation and decay scheme of Ref. [1]. The solid, dotted and long dashed curves show the effects of heavy quark energy loss with initial gluon rapidity densities of $dN_g/dy = 0, 1000,$ and 3500, respectively. Right-hand side: The ratio of charm to bottom decays to electrons obtained by varying the quark masses and scale factors. The effect of changing the Peterson function parameters from $\epsilon_c = 0.06$, $\epsilon_b = 0.006$ (lower band) to $\epsilon_c = \epsilon_b = 10^{-5}$ (upper band) is also illustrated. From Ref. [4].

the bottom contribution to single electrons is large and cannot be neglected in the computation of single electron suppression, shown in Fig. 6. Since bottom energy loss is greatly reduced relative to charm [4], the possible effect on the electron spectrum is reduced, leading to $R_{AA}(p_T < 6 \text{ GeV}; e) > 0.5 \pm 0.1$. A calculation by Armesto *et al.*, with a somewhat different model of energy loss, showed similar results to those in Fig. 6.

Recently two groups, Moore and Teaney [68] and Rapp *et al.* [69, 70] have calculated R_{AA} and the non-photonic electron elliptic flow, v_2, in a Langevin model of the time evolution of heavy quarks in the medium. Both these groups emphasize that elastic (collisional) energy loss should be important at low p_T relative to radiative loss since the boost for heavy flavor hadrons in the medium should not be large. Both also find a strong correlation between R_{AA} and v_2. The approaches differ somewhat but the trends are similar in the two calculations.

Moore and Teaney [68] calculate the diffusion and drag coefficients for charm quarks in perturbative QCD. The diffusion coefficient is proportional to the inverse square of the strong coupling constant, α_s, *e.g.* $D(2\pi T) \propto \alpha_s^{-2}$. They present the effects of a range of values for $D(2\pi T)$ on R_{AA} and v_2, finding the largest effects at high p_T for small $D(2\pi T)$, corresponding to large α_s or strong coupling in the plasma.

Rapp *et al.* [69, 70] calculated the diffusion and drag coefficients assuming that resonant D and B states in the QGP elastically scatter in the medium. Resonance scattering reduces the thermalization times for heavy flavors relative to those calculated with perturbative QCD matrix elements for fixed $\alpha_s = 0.4$. The effect is larger for charm than for the more massive bottom quarks. Including these states thus reduces the electron R_{AA} at high p_T relative to the results in Ref. [4] while increasing the electron v_2 to $\sim 10\%$ at $p_T \sim 2$ GeV, in relative agreement with the data.

Thus, given sufficiently strong coupling and/or resonant states, both R_{AA} and v_2 can be described within transport approaches using elastic scattering. More and better data is necessary to distinguish the two approaches.

4.3 Open heavy flavor measurements to date at RHIC

Open heavy flavor production cross sections can be measured by reconstructing the invariant mass of the heavy quark hadron from its hadronic decay products or by detecting leptons from semileptonic

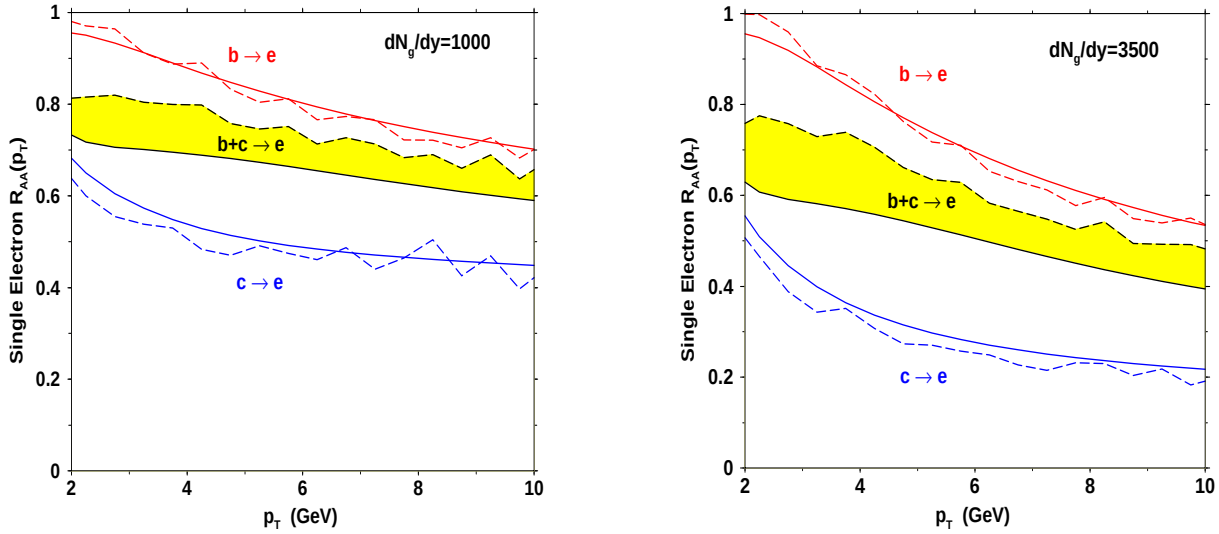

Fig. 6: Single electron attenuation pattern for $dN_g/dy = 1000$, left, and $dN_g/dy = 3500$, right. The solid curves employ the fragmentation scheme and lepton decay parameterizations of Ref. [1] while the dashed curves use the Peterson function with $\epsilon_c = 0.06$ and $\epsilon_b = 0.006$ and the decay to leptons employed by the PYTHIA Monte Carlo. Even for the extreme case on the right, the less quenched b quarks dilute R_{AA} so much that the modification of the combined electron yield from both c and b decays does not fall below $\sim 0.5 - 0.6$ near $p_T \sim 5$ GeV. From Ref. [4].

decays of those hadrons. While both PHENIX and STAR can measure heavy flavor cross sections by either technique, PHENIX has some advantages for semileptonic decay measurements and STAR has advantages for the hadronic decay measurements.

In both cases, the signal to background rate can be greatly improved if a precise measurement of the decay vertex position is available since hadrons containing c or b quarks typically travel several hundred microns from the collision point before decaying. Both PHENIX and STAR have plans to add secondary vertex detectors capable of the necessary precision but they will not be implemented for several years. In addition to reducing the background rates for open heavy flavor decays to leptons and hadrons, the secondary vertex detectors open up the possibility of a clean bottom cross section measurement using displaced vertex J/ψ decays.

Open heavy flavor cross section measurements based on semileptonic decays of charm and bottom mesons are feasible because a small lepton signal can be identified in a very large hadron background. In addition, background lepton sources are both small and well enough understood that they can be subtracted to get the open heavy flavor signal. However a disadvantage of semileptonic decay measurements is the loss of information about the decaying heavy meson due to the recoil kinematics. As a consequence, charm and bottom decays cannot easily be distinguished. Open charm measurements using hadronic decay products have two advantages: the D meson kinematic properties are reconstructed and separation of charm from bottom is far easier because only a small fraction of D mesons arise from bottom decays [1]. A disadvantage of hadronic decay measurements is the huge combinatorial background in heavy ion collisions.

PHENIX has measured open heavy flavor yields via semileptonic decays to electrons at mid rapidity ($|\eta| < 0.35$) using the Ring Imaging Cerenkov detector and electromagnetic calorimeter for electron identification. At forward and backward rapidity ($1.2 < |\eta| < 2.2$) the two muon spectrometers are used. PHENIX results are available for p+p at mid rapidity [71] and forward rapidity [72] as well as for d+Au [73] and Au+Au at mid rapidity [31]. No open charm results from hadronic decays have yet been reported by PHENIX since the small central arm acceptance is a disadvantage for such measurements.

STAR has measured open heavy flavor yields at mid rapidity ($|\eta| < 1.0$) via semileptonic decays

using either a combination of the time projection chamber (TPC) and time of flight (TOF) for electron identification or a combination of the TPC and the electromagnetic calorimeter. The backgrounds that must be subtracted are much larger than they are for PHENIX because of the larger photon conversion rates in STAR and the lack of of a hadron blind electron identifier, but this is compensated somewhat by the larger acceptance. STAR electron results are available for p+p [50], d+Au [50] and Au+Au collisions [32]. STAR has also measured open charm yields in the range $|\eta| < 1.0$ through hadronic D meson decays [50] for d+Au collisions.

Because the charm cross section is much larger than the bottom cross section at RHIC and dominates the semileptonic decay spectrum for $p_T < 2.5$ GeV$/c$, the integrated non-photonic lepton spectrum is usually assumed to be equal to the charm cross section.

4.31 Baseline measurements

Before any conclusions can be drawn about the hot, dense final state from the results for heavy ion collisions, some baseline information is needed. Data from p+p collisions are needed to establish the underlying cross sections and kinematic distributions for open heavy flavor. Also, p+A data are needed to study effects due to gluon saturation and the intrinsic k_T distributions in the colliding nuclei.

Both PHENIX and STAR have measured the charm production cross section at mid rapidity. These measurements have been extrapolated to all rapidities to yield total cross sections. These total cross sections are compared to each other, to results at other energies, and to pQCD calculations in Fig. 7. Results for d+Au and Au+Au collisions are scaled by the number of binary collisions ($N_{\rm coll}$) for direct comparison to the p+p results. The STAR data points are from combined fits to hadronic decay and semileptonic decay data. The PHENIX points are from semileptonic decay measurements only.

Although the STAR values are somewhat higher than those for PHENIX, the total charm cross sections are in acceptable agreement within the fairly large systematic uncertainties, as can be seen more clearly in Fig. 8, discussed later. The p_T-dependent slopes at mid rapidity are in good agreement for p+p collisions [74].

Since charm and bottom quarks are expected to be produced only in the initial nucleon-nucleon interactions, their yield should scale as the number of binary collisions, $N_{\rm coll}$. Fig. 8 shows the PHENIX measurement of the charm invariant yield in Au+Au collisions, scaled by $N_{\rm coll}$, at mid rapidity as a function of $N_{\rm coll}$ [31]. The p+p invariant yields from PHENIX [71] and the STAR invariant d+Au [50] and minimum bias Au+Au yields [32] are also shown. The PHENIX data are consistent with no $N_{\rm coll}$ dependence, as expected. As mentioned previously, the STAR yields are somewhat higher but, within the systematic uncertainties, the two experiments are in acceptable agreement. The linear fits shown in the figure are attempts to quantify deviations from a result independent of $N_{\rm coll}$. The slope fit to the PHENIX data is not significant within errors.

Note that because these yields are integrated over all p_T, no modification of the charm p_T distributions in the final state medium can be determined from these data. The p_T-dependent effects are discussed in the following section.

A NLO pQCD theoretical uncertainty band [1] is also shown in Fig. 8. While the data tend to lie above the calculation, they agree within the limits of their systematic uncertainties.

Unlike J/ψ measurements, the current d+Au and p+p open heavy flavor results are not precise enough for any conclusions to be drawn about either shadowing or k_T broadening. Obtaining meaningful open heavy flavor baseline results is an important priority for the RHIC program over the next few years.

4.32 Heavy ion measurements

Both PHENIX and STAR presented striking preliminary results on suppression of single electrons from open heavy flavor decays in central Au+Au collisions at Quark Matter 2005. PHENIX also showed

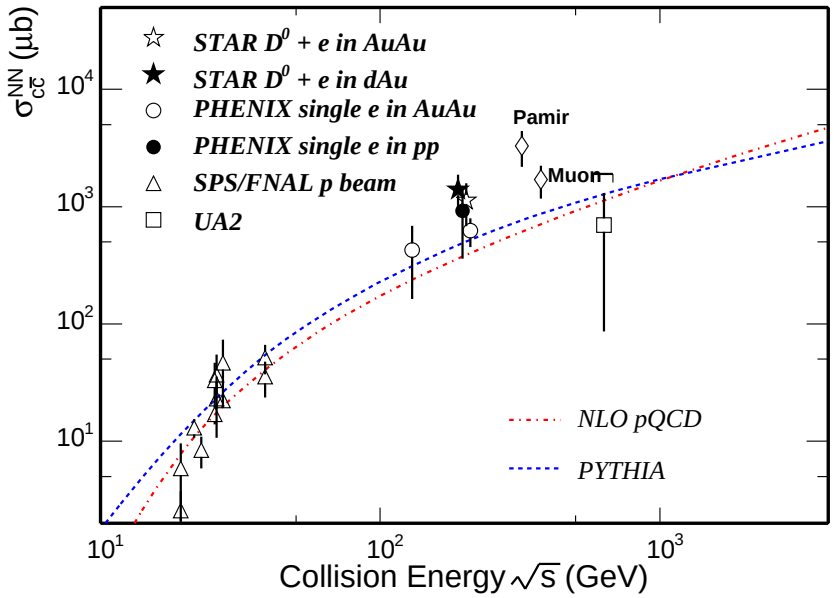

Fig. 7: Comparison of total cross section measurements. The STAR and PHENIX results are given as cross section per binary collisions [74]. The theory curves are from Ref. [50]. The NLO pQCD calculation uses parameters optimized using the low energy data [9].

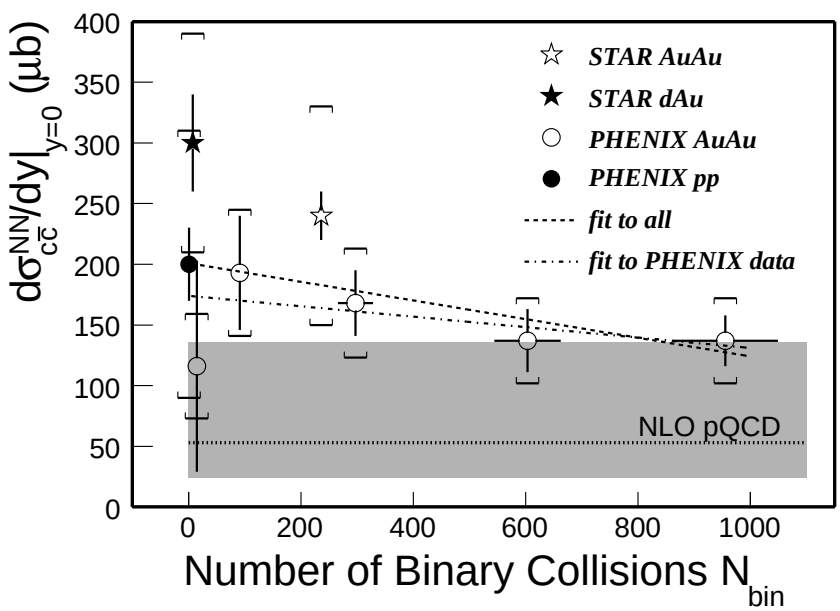

Fig. 8: Comparison of STAR and PHENIX measurements of the mid rapidity invariant charm yields scaled by $N_{\rm coll}$ as a function of $N_{\rm coll}$ [74].

preliminary results for the electron v_2 (elliptic flow parameter) from open heavy flavor decays

Nuclear modification factors for electrons from semileptonic decays of open heavy flavor in Au+Au central collisions from PHENIX [14] and STAR [32] are compared to theory calculations in Fig. 9 [74]. The data from the two experiments are in reasonable agreement. Both show strong suppression in central collisions at high p_T. The suppression factor, R_{AA}, is 0.2-0.3, similar to that seen for light quark hadrons [75].

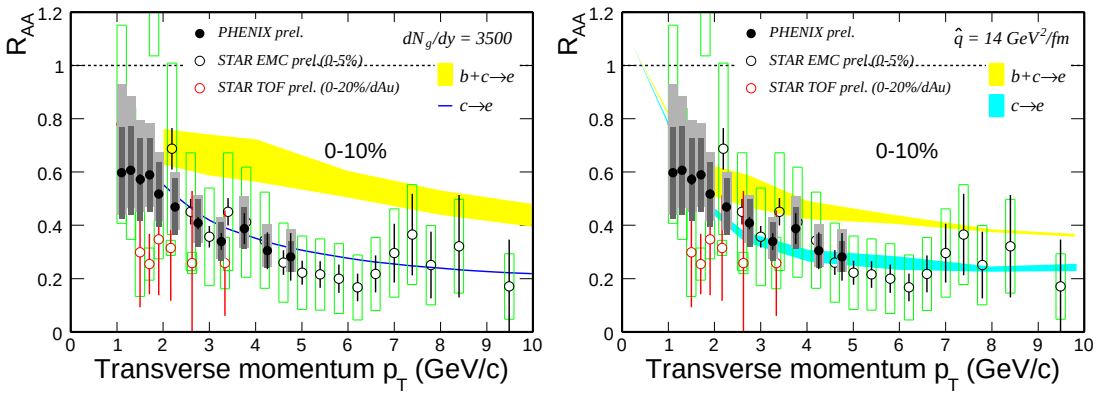

Fig. 9: Comparison of non-photonic electron R_{AA} measurements in central Au+Au collisions by PHENIX and STAR [74]. The data are identical on the two sides. Theory predictions from Ref. [4] (left-hand side) and [33] (right-hand side) are shown.

When comparing the non-photonic electron R_{AA} data to theory, recall that while the electron data contain contributions from both charm and bottom decays, the bottom contribution is expected to dominate for $p_T \sim 4$ GeV/c [1]. Figure 9 shows pQCD calculations of radiative energy loss (discussed earlier) from Djordevic *et al.* [4] (left) and Armesto *et al.* [33] (right). In the two cases, the calculations are shown for both charm alone and for combined charm and bottom decays. A range of parameters (gluon multiplicity, dN_g/dy, [4] and transport coefficient, $\hat{q}$, [33]) were determined from comparisons with central PHENIX π^0 data. Since the upper limits were used here, these calculations represent the maximum suppression consistent with the π^0 data. The predicted suppression for charm decay electrons is similar to the suppression seen in the data. However, when bottom quark decays are included, as they must be, the calculations predict much less suppression than observed. Thus the non-photonic electron R_{AA} data suggest the need for other energy loss mechanisms. Since light hadron suppression appears to be well described by radiative energy loss, this is a considerable challenge for theorists.

The elliptic flow parameter, v_2, has been measured for non-photonic electrons in minimum bias Au+Au collisions by PHENIX [15]. The preliminary results are shown in Fig. 10.

While there is no model independent way to extract the charm quark v_2 from the semileptonic decay electron v_2, there are some theoretical calculations available for the expected behavior of the electron v_2 [11, 77]. So far, the main interest has been for $p_T < 2$ GeV/c, where the primary contribution to the lepton spectrum is expected to be from charm quark decays and where the light hadron v_2 results are in good agreement with hydrodynamic models [75].

Figure 10 shows calculations of the electron v_2 from D meson decays in minimum bias Au+Au collisions by Greco *et al.*, [11] in a quark coalescence model. The dotted curve shows the semileptonic decay electron v_2 if the charm quark v_2 is zero and the D thus acquires all its v_2 from the light quark at coalescence. This represents the extreme case of no reinteraction of the charm quark in the medium. The solid curve shows the electron v_2 if the charm quark is assumed to be completely thermalized in the medium, the opposite extreme. At or below $p_T \sim 1.5$ GeV/c the data favor the assumption of complete charm thermalization. This result has been widely interpreted as providing evidence of at least partial

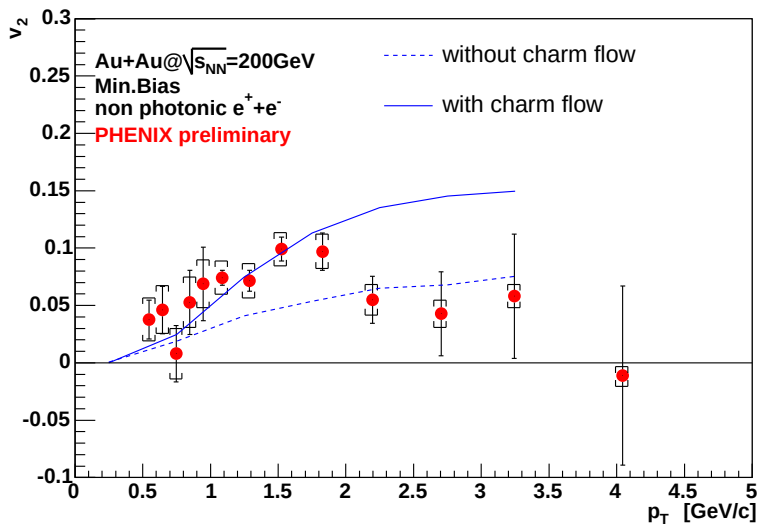

Fig. 10: Comparison of measurement of non-photonic electron v_2 in minimum bias Au+Au collisions from PHENIX [15]. The theory results are from Ref. [11].

charm thermalization.

Figure 10 also shows calculations by Zhang *et al.* [77] using the AMPT dynamical transport model. This model uses the charm quark scattering cross section in the medium is a parameter. The dot-dashed curve shows the predicted v_2 for a 3 mb charm rescattering cross section while the dashed curve assumes 10 mb. They conclude that the large elliptic flow observed for non-photonic electrons at $p_T \sim 2$ GeV$/c$ requires a charm rescattering cross section too large to be consistent with pQCD.

There is considerable interest in the behavior of the non-photonic electron v_2 for $p_T > 2$ GeV$/c$ where the bottom contribution is expected to become important [1]. The PHENIX v_2 results fall toward zero here, consistent with both a larger bottom contribution and a smaller degree of thermalization for higher p_T charm quarks. But, while the PHENIX data have small systematic errors at high p_T, their statistical precision is poor, keeping them from being definitive.

4.4 Proposed open heavy flavor experimental program at RHIC II

Here we focus on the new open heavy flavor physics that becomes available with the combination of the detector upgrades and the RHIC II luminosity upgrade.

With a displaced vertex measurement and RHIC II luminosity, the $B \to J/\psi$ decay channel can provide a very clean measurement of open bottom production by both PHENIX and STAR (see tables 2 and 3). The displaced vertex distributions for prompt J/ψ and for $B \to J/\psi$ decays into the PHENIX muon arms are compared in the left panel of Figure 11. The yields in tables 2 and 3 assume a displaced vertex cut of 1 mm. A good measurement of the cross section and R_{AA} vs p_T and y for open bottom production will be possible using $B \to J/\psi$. Even at RHIC II luminosity, however, the yields are not expected to be large enough to permit a v_2 measurement.

With a displaced vertex measurement and RHIC II luminosity, the separation of charm and bottom contributions to the semileptonic decay spectra can be done statistically by using the different decay lengths for open charm and open bottom mesons (see the right panel of Figure 11). By analyzing data samples with different decay length cuts, the fraction of the signal due to b quarks can be varied. The

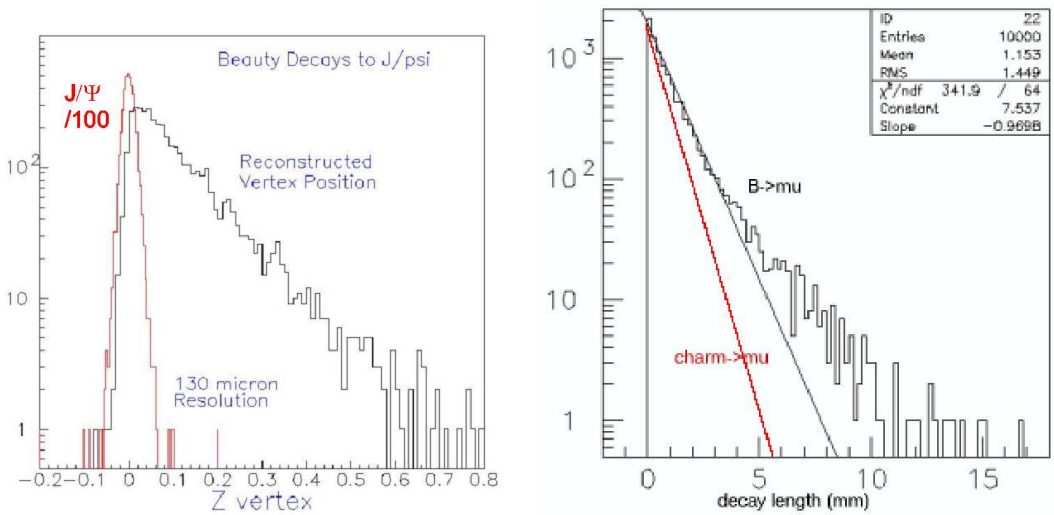

Fig. 11: Left: Comparison of prompt J/ψ displaced vertex distribution (in cm) with that from $B \rightarrow J/\psi$ decays. Note that the prompt J/ψ distribution is scaled down by a factor of 100. Right: Decay length distributions (in mm) from simulations for open charm and open beauty.

addition of a displaced vertex measurement will also help by reducing background in the measurement of single muons from open heavy flavor, since a displaced vertex cut of 1 cm reduces the yield of muons from light hadron decays by about one order of magnitude. As a result, measurements of R_{AA} and v_2 vs p_T and y should be possible for open charm and open bottom separately at RHIC II.

Tight displaced vertex cuts will also greatly reduce the background for the $D \rightarrow K\pi$ measurement of open charm yields by eliminating most of the prompt hadron tracks from the combinatorial background. The background reduction will result in a much improved cross section and R_{AA} vs p_T measurement. Without a trigger for $D \rightarrow K\pi$ decays, however, this measurement will not benefit greatly from the increased luminosity at RHIC II. Therefore it is not clear if a useful v_2 measurement can be expected. The situation will be similar at the LHC.

5. Hidden heavy flavor: quarkonium

In this section we present a more detailed discussion of the theoretical motivation for studying heavy quarkonia in heavy ion collisions. We also summarize the present experimental and theoretical status and describe the proposed RHIC II experimental quarkonia program.

5.1 Theoretical results

5.11 Cross sections in p+p collisions

We discuss quarkonium production in the color evaporation model (CEM) which can be used to calculate the total quarkonium cross sections. The CEM was first discussed some time ago [78, 79] and has enjoyed considerable phenomenological success. In the CEM, the quarkonium production cross section is some fraction F_C of all $Q\overline{Q}$ pairs below the $H\overline{H}$ threshold where H is the lowest mass heavy flavor hadron. Thus the CEM cross section is simply the $Q\overline{Q}$ production cross section with a cut on the pair mass but without any constraints on the color or spin of the final state. The produced $Q\overline{Q}$ pair then neutralizes its color by interaction with the collision-induced color field—"color evaporation". The Q and the $\overline{Q}$ either combine with light quarks to produce heavy-flavored hadrons or bind with each other to form quarkonium. The additional energy needed to produce heavy-flavored hadrons when the partonic center of mass energy, $\sqrt{\hat{s}}$, is less than $2m_H$, the heavy hadron threshold, is obtained nonperturbatively from the color field in the interaction region. Thus the yield of all quarkonium states may be only a small fraction of the total $Q\overline{Q}$ cross section below $2m_H$. At leading order, the production cross section of quarkonium state C in an AB collision is

$$\sigma_C^{\rm CEM} = F_C \sum_{i,j} \int_{4m^2}^{4m_H^2} d\hat{s} \int dx_1 dx_2 \; f_i^A(x_1,\mu^2) \; f_j^B(x_2,\mu^2) \; \hat{\sigma}_{ij}(\hat{s}) \; \delta(\hat{s} - x_1 x_2 s) \; , \tag{2}$$

where $ij = q\overline{q}$ or gg and $\hat{\sigma}_{ij}(\hat{s})$ is the $ij \to Q\overline{Q}$ subprocess cross section. The total $Q\overline{Q}$ cross section takes $\hat{s} \to s$ in the upper limit of the integral over $\hat{s}$ in Eq. (2).

The fraction F_C must be universal so that, once it is fixed by data, the quarkonium production ratios should be constant as a function of $\sqrt{s}$, y and p_T. The actual value of F_C depends on the heavy quark mass, m, the scale, μ^2, the parton densities, $f_i^A(x,\mu^2)$ and the order of the calculation. It was shown in Ref. [45] that the quarkonium production ratios were indeed constant, as expected by the model.

Of course the leading order calculation in Eq. (2) is insufficient to describe high p_T quarkonium production since the $Q\overline{Q}$ pair p_T is zero at LO. Therefore, the CEM was taken to NLO [45, 80] using the exclusive $Q\overline{Q}$ hadroproduction code of Ref. [81]. At NLO in the CEM, the process $gg \to gQ\overline{Q}$ is included, providing a good description of the quarkonium p_T distributions at the Tevatron [80]. In the exclusive NLO calculation [81], both the Q and $\overline{Q}$ variables are integrated to obtain the pair distributions. Thus, although $\mu \propto m$ in analytic LO calculations, at NLO, $\mu^2 \propto m_T^2 = m_Q^2 + p_T^2$ where p_T is that of the $Q\overline{Q}$ pair, $p_T^2 = 0.5(p_{T_Q}^2 + p_{T_{\overline{Q}}}^2)$.

We use the same parton densities and parameters that agree with the $Q\overline{Q}$ total cross section data, given in Table 6, to determine F_C for J/ψ and Υ production. The fit parameters [82, 83] for the parton densities [84, 85, 86], quark masses and scales are given in Table 6 while the $Q\overline{Q}$ cross sections calculated with these parameters are compared to $pp \to Q\overline{Q}$ and $\pi^- p \to Q\overline{Q}$ data in Fig. 12.

We now describe the extraction of F_C for the individual quarkonium states. The J/ψ has been measured in p+p and p+A interactions up to $\sqrt{s} = 63$ GeV. The data are of two types: the forward cross section, $\sigma(x_F > 0)$, and the cross section at zero rapidity, $d\sigma/dy|_{y=0}$. All the cross sections are inclusive with feed down from χ_c and ψ' decays. To obtain $F_{J/\psi}$ for inclusive J/ψ production, the normalization of Eq. (2) is obtained from a fit using the $c\overline{c}$ parameters in Table 6. The comparison of $\sigma_{J/\psi}^{\rm CEM}$ to the $x_F > 0$ data for all four fits is shown on the left-hand side of Fig. 13. The ratios of the direct production

Table 6: Parameters used to obtain the 'best' agreement to the $Q\overline{Q}$ cross sections. The quark mass is given in GeV. The inclusive J/ψ production fraction, $F_{J/\psi}$, and the inclusive Υ production fraction, F_Υ, obtained from the data are also given.

$c\overline{c}$					$b\overline{b}$				
Label	PDF	m	μ/m_T	$F_{J/\psi}$	Label	PDF	m	μ/m_T	F_Υ
$\psi 1$	MRST HO	1.2	2	0.0144	$\Upsilon 1$	MRST HO	4.75	1	0.0276
$\psi 2$	MRST HO	1.4	1	0.0248	$\Upsilon 2$	MRST HO	4.5	2	0.0201
$\psi 3$	CTEQ 5M	1.2	2	0.0155	$\Upsilon 3$	MRST HO	5.0	0.5	0.0508
$\psi 4$	GRV 98 HO	1.3	1	0.0229	$\Upsilon 4$	GRV 98 HO	4.75	1	0.0225

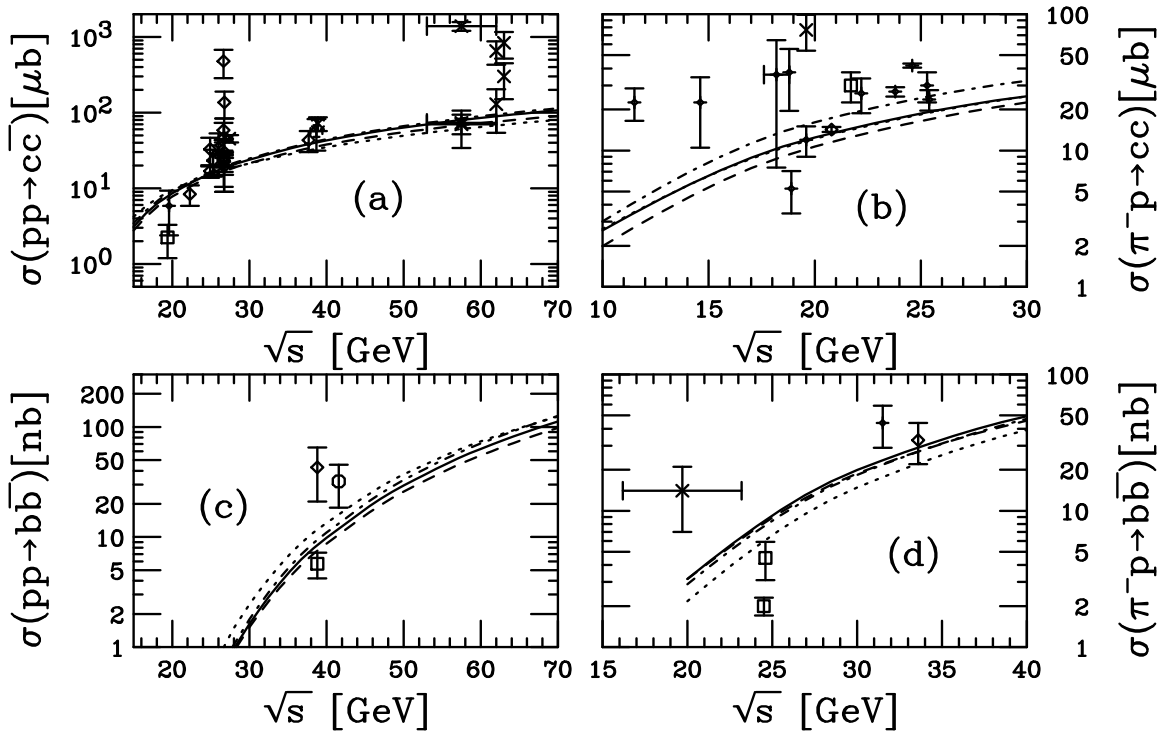

Fig. 12: The $c\overline{c}$, (a) and (b), and $b\overline{b}$, (c) and (d), total cross section data in p+p and $\pi^- p$ interactions compared to NLO calculations. In (a) and (b), we show $\psi 1$ (solid), $\psi 2$ (dashed), $\psi 3$ (dot-dashed) and $\psi 4$ (dotted). In (c) and (d), we show $\Upsilon 1$ (solid), $\Upsilon 2$ (dashed), $\Upsilon 3$ (dot-dashed) and $\Upsilon 4$ (dotted).

cross sections to the inclusive J/ψ cross section can be determined from data on inclusive cross section ratios and branching fractions. These direct ratios, R_C, given in Table 7, are multiplied by the inclusive fitted $F_{J/\psi}$, also shown in Table 6 to obtain the direct production fractions, $F_C^{\rm dir} = F_{J/\psi} R_C$.

Table 7: Direct quarkonium production ratios, $R_C = \sigma_C^{\rm dir}/\sigma_{C'}^{\rm inc}$ where $C' = J/\psi$ and Υ. From Ref. [87].

	J/ψ	ψ'	χ_{c1}	χ_{c2}	Υ	Υ'	Υ''	$\chi_b(1P)$	$\chi_b(2P)$
R_C	0.62	0.14	0.60	0.99	0.52	0.33	0.20	1.08	0.84

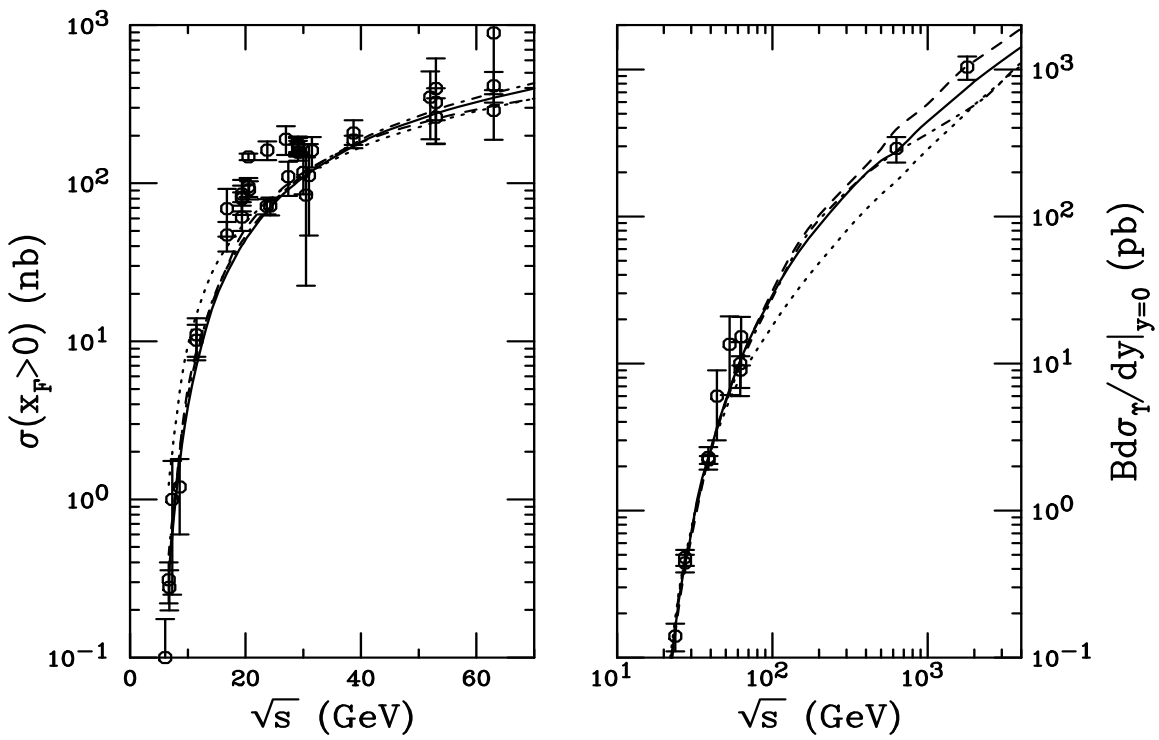

Fig. 13: Forward J/ψ (left) and combined $\Upsilon + \Upsilon' + \Upsilon''$ inclusive (right) cross sections calculated to NLO in the CEM. On the left-hand side, we show $\psi 1$ (solid), $\psi 2$ (dashed), $\psi 3$ (dot-dashed) and $\psi 4$ (dotted). On the right-hand side, we show $\Upsilon 1$ (solid), $\Upsilon 2$ (dashed), $\Upsilon 3$ (dot-dashed) and $\Upsilon 4$ (dotted).

The same procedure, albeit somewhat more complicated due to the larger number of bottomonium states below the $B\overline{B}$ threshold, is followed for the bottomonium. For most data below $\sqrt{s} = 100$ GeV, the three bottomonium S states were either not separated or their sum was reported. No x_F-integrated cross sections were available so that we fit the CEM Υ cross section to the effective lepton pair cross section at $y = 0$ for the three $\Upsilon(nS)$ states. The extracted fit fraction, $F_{\sum\Upsilon}$, combined with $\sigma_\Upsilon^{\rm CEM}$ and compared to the data for all parameter sets in Table 6, is shown on the right-hand side of Fig. 13. Using the individual branching ratios of the Υ, Υ' and Υ'' to lepton pairs and the total cross sections reported by CDF [88], it is possible to extract the inclusive Υ fit fraction, F_Υ, given in Table 6. The direct production ratios obtained in Ref. [25] have been updated in Ref. [87] using recent CDF χ_b data. The resulting direct to inclusive Υ ratios, R_C, are also given in Table 7. The sub threshold $b\bar{b}$ cross section is then multiplied by $F_C^{\rm dir} = F_\Upsilon R_C$ to obtain the direct bottomonium cross sections.

The total cross sections for the charmonium and bottomonium states in p+p collisions at $\sqrt{s} = 200$ GeV are shown in Tables 8 and 9 respectively.

The energy dependence shown in Fig. 13 for both states is well reproduced by the NLO CEM. All the fits are equivalent for $\sqrt{s} = 100$ GeV but differ by up to a factor of two at 2 TeV. The high energy

Case	$\sigma^{\rm inc}_{J/\psi}$	$\sigma^{\rm dir}_{J/\psi}$	$\sigma_{\chi_{c1}}$	$\sigma_{\chi_{c2}}$	$\sigma_{\psi'}$
$\psi 1$	2.35	1.46	1.41	2.33	0.33
$\psi 2$	1.76	1.09	1.06	1.74	0.25
$\psi 3$	2.84	1.76	1.70	2.81	0.40
$\psi 4$	2.10	1.31	1.26	2.08	0.29

Table 8: The charmonium cross sections (in μb) for 200 GeV p+p collisions. The inclusive and direct J/ψ cross sections are both given.

Case	$\sigma^{\rm inc}_{\Upsilon}$	$\sigma^{\rm dir}_{\Upsilon}$	$\sigma_{\Upsilon'}$	$\sigma_{\Upsilon''}$	$\sigma_{\chi_b(1P)}$	$\sigma_{\chi_b(2P)}$
$\Upsilon 1$	6.60	3.43	2.18	1.32	7.13	5.54
$\Upsilon 2$	7.54	3.92	2.49	1.51	8.15	6.34
$\Upsilon 3$	5.75	2.99	1.90	1.15	6.21	4.83
$\Upsilon 4$	4.31	2.24	1.42	0.86	4.66	3.62

Table 9: The direct bottomonium cross sections (in nb) for p+p collisions at 200 GeV. The production fractions for the total Υ are multiplied by the appropriate ratios determined from data.

Υ data seem to agree best with the energy dependence of $\Upsilon 1$ and $\Upsilon 2$ although $\Upsilon 1$ underestimates the Tevatron result by a factor of ≈ 1.4. A similar check cannot be made for the J/ψ because the high lepton p_T cut excludes J/ψ acceptance for $p_T = 0$ at the Tevatron in Run I.

5.12 Cold nuclear matter effects on quarkonium production at RHIC

It is essential that the A dependence be understood in cold nuclear matter to set a proper baseline for quarkonium suppression in A+A collisions. The NA50 collaboration has studied the J/ψ A dependence and attributed its behavior to J/ψ break up by nucleons in the final state, referred to as nuclear absorption. However, the parton distributions are modified in the nucleus relative to free protons. This modification, referred to here as shadowing, is increasingly important at higher energies, as emphasized in Ref. [89]. In this section, we discuss the interplay of shadowing and absorption in d+Au and A+A collisions at RHIC.

Shadowing, the modification of the parton densities in the nucleus with respect to the free nucleon, is taken into account by replacing f^p_j in Eq. (2) by $F^A_j(x,\mu^2,\vec{b},z) = \rho_A(\vec{b},z)S^j(A,x,\mu^2,\vec{b},z)f^p_j(x,\mu^2)$ and adding integrals over the spatial coordinates. Here S^j is the shadowing parameterization. The density distribution of the deuteron is also included in these calculations but the small effects of shadowing in deuterium are ignored. The PHENIX J/ψ d+Au data as a function of rapidity show a dependence consistent with nuclear shadowing plus a small absorption cross section of 1-3 mb. We have calculated J/ψ production in the CEM using Eq. (2) with the same mass and scale as in $c\overline{c}$ production. The calculations of the d+Au/p+p and A+A/p+p ratios are done at LO to simplify the calculations since the LO and NLO ratios are equivalent [2].

To implement nuclear absorption of the J/ψ in d+Au collisions, the dN production cross section is weighted by the survival probability, $S^{\rm abs}$ [90]

$$S^{\rm abs}(\vec{b},z) = \exp\left\{-\int_z^\infty dz'\rho_A(\vec{b},z')\sigma_{\rm abs}(z'-z)\right\} \tag{3}$$

where z is the longitudinal production point and z' is the point at which the state is absorbed. The nucleon absorption cross section, $\sigma_{\rm abs}$, typically depends on where the state is produced in the medium and how far it travels through nuclear matter. If absorption alone is active, *i.e.* no shadowing so that $S^j \equiv 1$, then an effective minimum bias A dependence is obtained after integrating $S^{\rm abs}$ over the spatial coordinates.

If $S^{\rm abs} = 1$ also, $\sigma_{{\rm d}A} = 2A\sigma_{pN}$. When $S^{\rm abs} \neq 1$, $\sigma_{{\rm d}A} = 2A^{\alpha}\sigma_{pN}$ where, if $\sigma_{\rm abs}$ is a constant, independent of the production mechanism for a nucleus of $\rho_A = \rho_0\theta(R_A - b)$, $\alpha = 1 - 9\sigma_{\rm abs}/(16\pi r_0^2)$, where $r_0 = 1.2$ fm. The contribution to the full A dependence of α from absorption alone is only constant if $\sigma_{\rm abs}$ is constant and independent of the production mechanism [90]. The observed J/ψ yield includes feed down from χ_{cJ} and ψ' decays, giving

$$S^{\rm abs}_{J/\psi}(b,z) = 0.58 S^{\rm abs}_{J/\psi,\,{\rm dir}}(b,z) + 0.3 S^{\rm abs}_{\chi_{cJ}}(b,z) + 0.12 S^{\rm abs}_{\psi'}(b,z) \ . \tag{4}$$

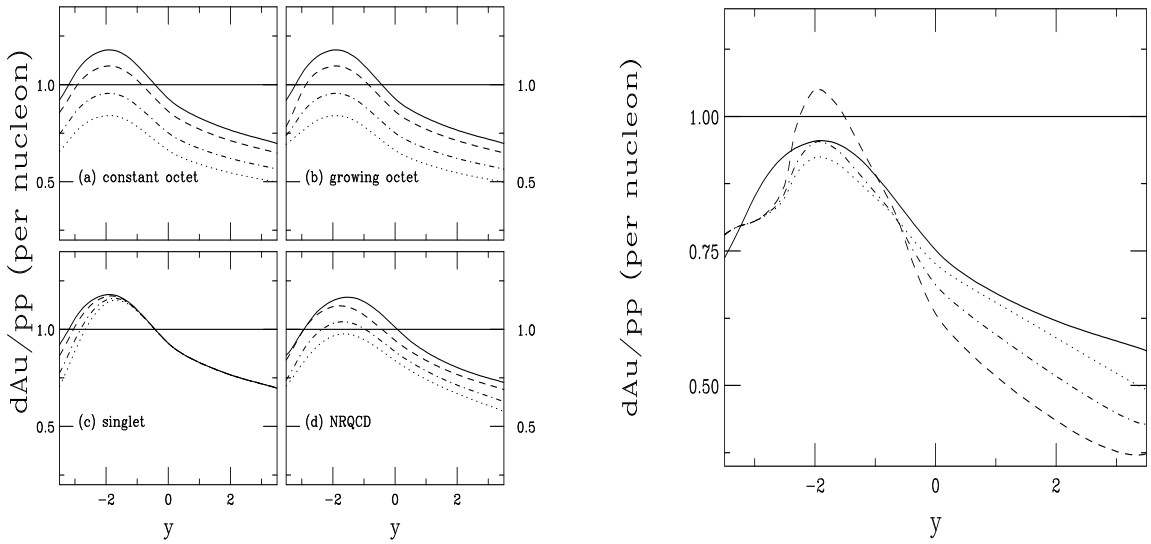

Fig. 14: Left-hand side: The J/ψ d+Au/p+p ratio with EKS98 at 200 GeV as a function of rapidity for (a) constant octet, (b) growing octet, (c) singlet, all calculated in the CEM and (d) NRQCD. For (a)-(c), the curves are no absorption (solid), $\sigma_{\rm abs} = 1$ (dashed), 3 (dot-dashed) and 5 mb (dotted). For (d), we show no absorption (solid), 1 mb octet/1 mb singlet (dashed), 3 mb octet/3 mb singlet (dot-dashed), and 5 mb octet/3 mb singlet (dotted). Right-hand side: The J/ψ d+Au/p+p ratio at 200 GeV for a growing octet with $\sigma_{\rm abs} = 3$ mb is compared for four shadowing parameterizations. We show the EKS98 (solid), FGSo (dashed), FGSh (dot-dashed) and FGSl (dotted) results as a function of rapidity.

The J/ψ may be produced as a color singlet, a color octet or in a combination of the two. In color singlet production, the final state absorption cross section depends on the size of the $c\overline{c}$ pair as it traverses the nucleus, allowing absorption to be effective only while the cross section is growing toward its asymptotic size inside the target. On the other hand, if the $c\overline{c}$ is only produced as a color octet, hadronization will occur only after the pair has traversed the target except at very backward rapidity. We have considered a constant octet cross section, as well as one that reverts to a color singlet at backward rapidities. For singlets, $S^{\rm abs}_{J/\psi,\,{\rm dir}} \neq S^{\rm abs}_{\chi_{cJ}} \neq S^{\rm abs}_{\psi'}$ but, with octets, one assumes that $S^{\rm abs}_{J/\psi,\,{\rm dir}} = S^{\rm abs}_{\chi_{cJ}} = S^{\rm abs}_{\psi'}$. As can be seen in Fig. 14, the difference between the constant and growing octet assumptions is quite small at large $\sqrt{s_{NN}}$ with only a small singlet effect at $y < -2$. Singlet absorption is also important only at similar rapidities and is otherwise not different from shadowing alone. Finally, we have also considered a combination of octet and singlet absorption in the context of the NRQCD approach, see Ref. [90] for more details. The combination of nonperturbative singlet and octet parameters changes the shape of the shadowing ratio slightly. Including the singlet contribution weakens the effective absorption. The results are shown integrated over impact parameter. The calculations use the EKS98 shadowing parameterization [91] since it gives good agreement with the trend of the PHENIX data. For results with other shadowing parameterizations, see Refs. [2, 3].

Several values of the asymptotic absorption cross section, $\sigma_{\rm abs} = 1$, 3 and 5 mb, corresponding to $\alpha = 0.98$, 0.95 and 0.92 respectively using Eqs. (3) and (4), are shown in Figs. 14 and 15 for d+Au

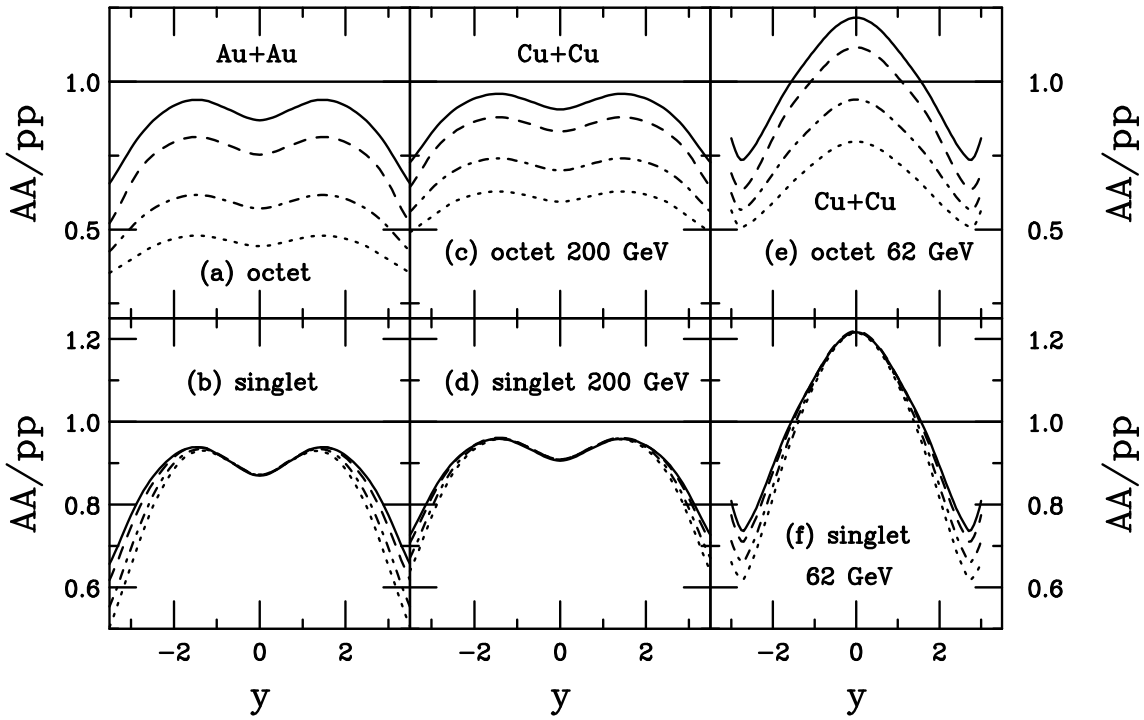

Fig. 15: The AA/pp ratio with the EKS98 parameterization as a function of y for octet (upper) and singlet (lower) absorption. In (a) and (b) we show the Au+Au results at 200 GeV while the Cu+Cu results are shown at 200 GeV (c) and (d) as well as at 62 GeV (e) and (f). The curves are $\sigma_{\rm abs} = 0$ (solid), 1 (dashed), 3 (dot-dashed) and 5 mb (dotted).

and A+A collisions respectively. These values of $\sigma_{\rm abs}$ are somewhat smaller than those obtained for the sharp sphere approximation. The diffuse surface of a real nucleus and the longer range of the density distribution result in a smaller value of $\sigma_{\rm abs}$ than a spherical nucleus. As will be seen later, there is good agreement with the trend of the preliminary PHENIX data [92] for $\sigma_{\rm abs} = 0 - 3$ mb. Work is in progress to quantify the shadowing parameterization and absorption cross section more precisely [93].

We note that the current RHIC data are not sufficiently precise to distinguish between J/ψ production and absorption in the CEM relative to that in the NRQCD approach. However, a measurement of the χ_c A dependence may be able to clarify the situation [90]. In the CEM, the J/ψ and χ_c distributions differ only in the value of F_C. In the NRQCD approach, the J/ψ is produced primarily in a color octet state while the χ_c is produced as a color singlet state. Thus while the production of both states would exhibit the same shadowing effect, a difference in the J/ψ and χ_c d+Au/p+p ratios due to octet relative to singlet absorption may be measurable.

We now turn to the centrality dependence of J/ψ production in d+Au and A+A collisions. In central collisions, inhomogeneous (spatially dependent) shadowing is stronger than the homogeneous (minimum bias) result. The stronger the homogeneous shadowing, the larger the inhomogeneity. In peripheral collisions, inhomogeneous effects are weaker than the homogeneous results but some shadowing is still present. Shadowing persists in peripheral collisions in part because the density in a heavy nucleus is large and approximately constant except close to the surface and because the deuteron wave function has a long tail. We also expect absorption to be stronger in central collisions.

To study the centrality dependence of shadowing and absorption, we present the d+Au/p+p and A+A/p+p ratios as a function of $N_{\rm coll}$,

$$N_{\rm coll}(b) = \sigma_{NN}^{\rm in} \int d^2s T_A(s) T_B(|\vec{b} - \vec{s}|) \ ,$$

where T_A and T_B are the nuclear thickness functions and the inelastic nucleon-nucleon cross section,

$\sigma_{NN}^{\rm in}$, is 42 mb at 200 GeV. In Figs. 16 and 17, we show the $N_{\rm coll}$ dependence for several representative rapidities, $y = -2$, 0 and 2 for RHIC. We have chosen an inhomogeneous shadowing parameterization proportional to the path length of the parton through the nucleus [89]. For more results, see Refs. [2, 3].

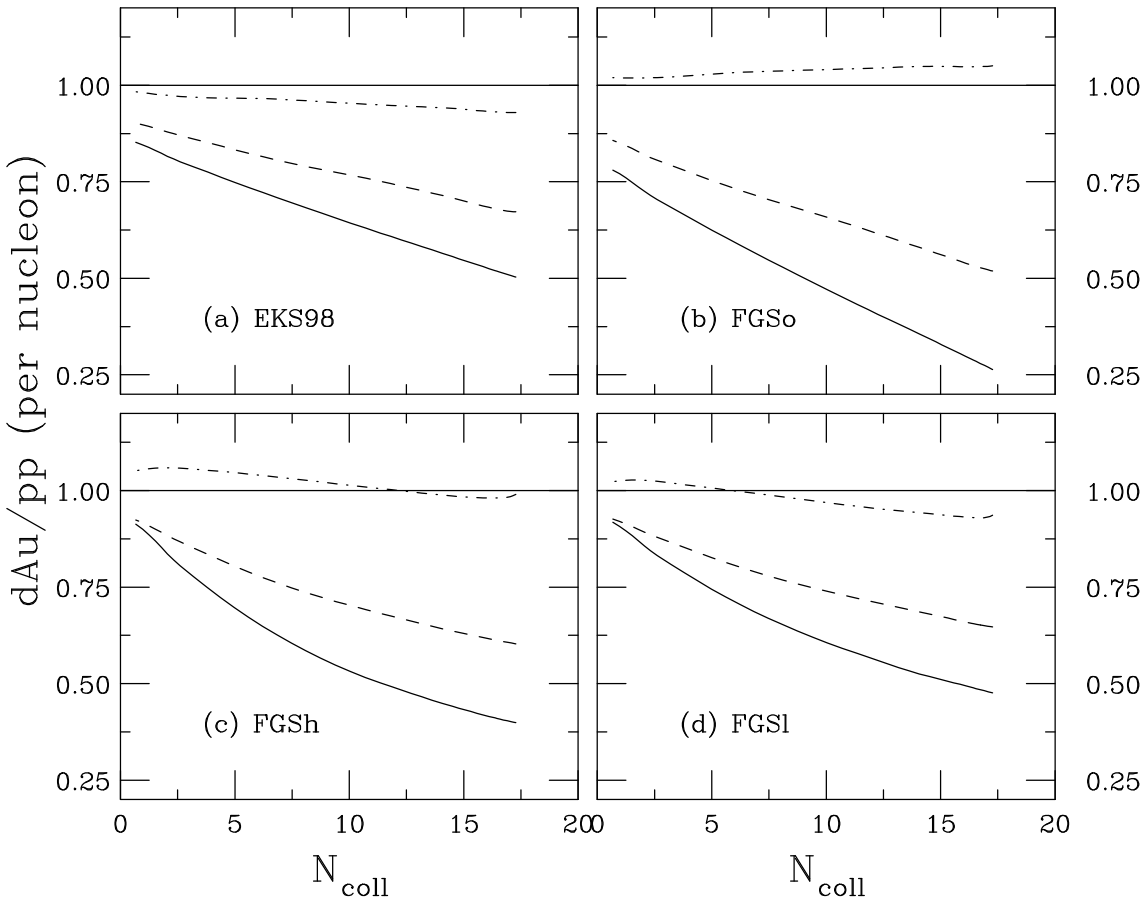

Fig. 16: The ratio d+Au/p+p as a function of $N_{\rm coll}$ for the EKS98 (a), FGSo (b), FGSh (c) and FGSl (d) shadowing parameterizations. The calculations with EKS98 and FGSo use the inhomogeneous path length parameterization while that obtained by FGS is used with FGSh and FGSl. Results are given for $y = -2$ (dot-dashed), $y = 0$ (dashed) and $y = 2$ (solid) at 200 GeV for a growing octet with $\sigma_{\rm abs} = 3$ mb.

The dependence of the RHIC ratios on $N_{\rm coll}$ is almost linear, as seen in Figs. 16 and 17. We do not show results for $N_{\rm coll} < 1$. The weakest $N_{\rm coll}$ dependence occurs in the antishadowing region, illustrated by the $y = -2$ result (dot-dashed curve). The overall dependence on $N_{\rm coll}$ is stronger than that obtained from shadowing alone, described in Ref. [89], where inhomogeneous shadowing effects depend strongly on the amount of homogeneous shadowing. Relatively large effects at low x are accompanied by the strongest impact parameter, b, dependence. In the transition region around mid rapidity at RHIC, the b dependence of the ratio d+Au/p+p due to shadowing is nearly negligible and almost all of the $N_{\rm coll}$ dependence at $y \sim 0$ can be attributed to absorption. The $y = -2$ results for color singlet production and absorption, in the antishadowing region, are fairly independent of $N_{\rm coll}$.

5.13 Models of quarkonium production in heavy ion collisions

In-medium properties of quarkonium from lattice QCD:

Properties of heavy quarks have been used to characterize "thermal properties of the QCD vacuum" ever since the first lattice calculations at non-zero temperature [94]. Modifications of the inter-

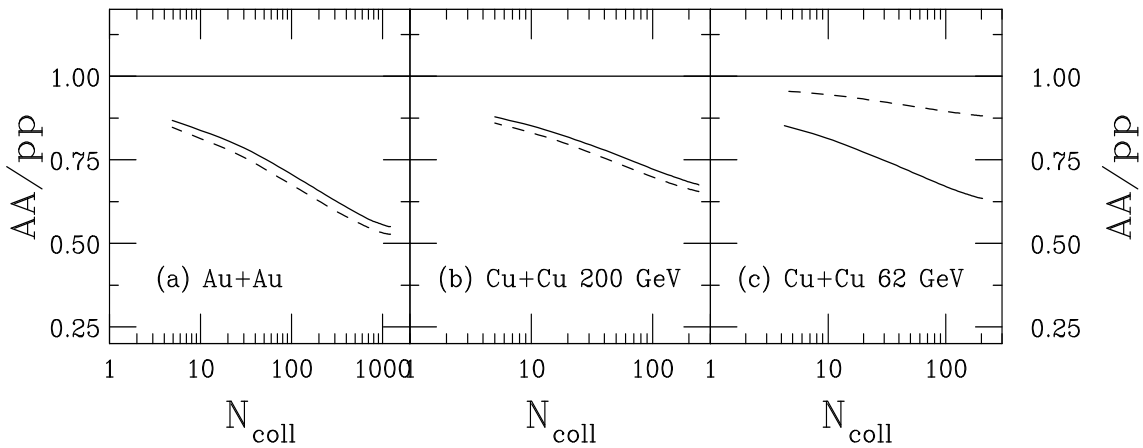

Fig. 17: The ratio AA/pp as a function of $N_{\rm coll}$ for a 3 mb octet absorption cross section and the EKS98 parameterization at $y = 0$ (dashed) and $y = 2$ (solid) for Au+Au at 200 GeV (a) and Cu+Cu at 200 GeV (b) and 62 GeV (c).

actions between heavy, static quarks in a thermal heat bath are clearly reflected by changes of the free energy which, in the zero temperature limit, reduces to the heavy quark potential [95]. To use this information to analyze thermal modifications of quarkonia requires an intermediate, phenomenological step: the construction of a temperature dependent effective potential which then can be used in a nonrelativistic Schrödinger equation [96, 97, 98] or a more refined coupled-channel analysis [99, 100]. Quite generically, the potential model analyses suggest a sequential suppression pattern where heavy quark bound states dissociate at temperatures at which their bound state radii become comparable to the Debye screening radius, illustrated in Fig. 18.

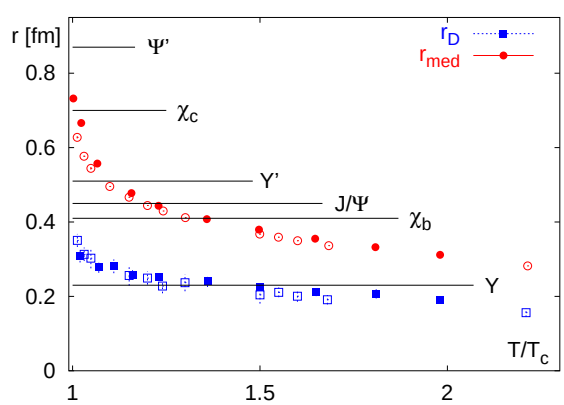

Fig. 18: Mean squared charge radii of some charmonium and bottomonium states compared to the Debye screening radius, $r_D \equiv 1/m_D$ and a related scale, $r_{\rm med}$, an estimate of the distance beyond which the force between a static $Q\overline{Q}$ pair is strongly modified by temperature effects [101]. Open (closed) symbols correspond to SU(3) (2-flavor QCD) calculations [95].

More recently, the calculation of thermal hadron correlation functions and their spectral analysis [102] eliminated some of the ambiguities inherent in the potential model approach. The spectral analysis, at least in principle, provides an ab-initio approach to the calculation of in-medium properties of heavy quark bound states. Its predictive power is reduced only by the application of statistical tools like the Maximum Entropy Method (MEM) which, however, can be steadily improved with further improvement of the available computing resources and numerical techniques. Predictions based on potential model calculations as well as the spectral analysis have been reviewed in recent studies that have been performed

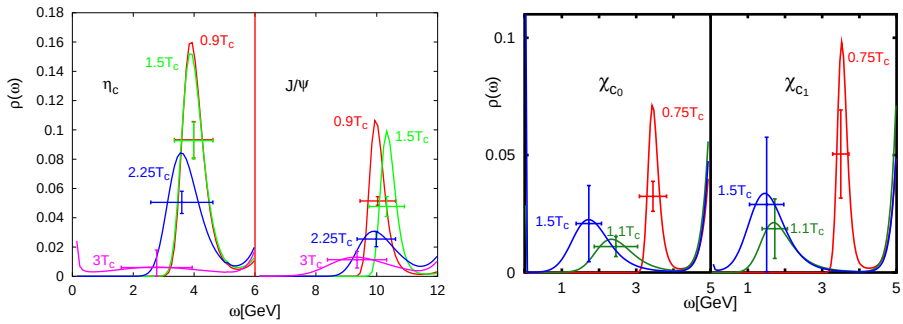

Fig. 19: Spectral functions of η_c, J/ψ (left) and and χ_{c0}, χ_{c1} (right) at various temperatures below and above T_c. The vertical bars give the error on the average value of the spectral function in the bin indicated by the horizontal bars.

to analyze prospects for quarkonium studies at the LHC [37, 103]. In the following, we concentrate on the analysis of thermal hadron correlation functions and the spectral functions extracted from them.

The finite temperature, Euclidean time correlation functions

$$G_H(\tau,\vec{r},T) = \langle J_H(\tau,\vec{r}) J_H^\dagger(0,\vec{0})\rangle \tag{5}$$

of hadronic currents, $J_H = \bar{q}(\tau,\vec{r})\Gamma_H q(\tau,\vec{r})$, where Γ_H denotes a suitable product of gamma matrices that projects onto the appropriate quantum numbers H, are directly related to spectral functions, $\sigma_H(\omega,T)$. These spectral functions encompass all the information about thermal modifications of the hadron spectrum in channel H, so that

$$G_H(\tau,\vec{r},T) = \int_0^\infty d\omega\, \frac{d^3\vec{p}}{(2\pi)^3}\sigma_H(\omega,\vec{p},T)\, \mathrm{e}^{i\vec{p}\cdot\vec{r}} \frac{\cosh(\omega(\tau-1/2T))}{\sinh(\omega/2T)} \quad , \tag{6}$$

are directly related to experimental observables. In particular, the spectral function in the vector channel, $\sigma_V(\omega,\vec{p},T)$, is directly related to the differential cross section for thermal dilepton production,

$$\frac{dW}{d\omega d^3p} = \frac{5\alpha^2}{27\pi^2}\frac{\sigma_V(\omega,\vec{p},T)}{\omega^2(\mathrm{e}^{\omega/T}-1)} \quad . \tag{7}$$

Note that the rates obtained using this method do not include any contributions arising from the feed down of other channels into the vector channel [104, 87].

Some generic aspects of the influence of a thermal medium on states with different quantum numbers can already be deduced from the temperature dependence of the thermal correlation functions themselves and does not require the additional step of applying the MEM analysis which, after all, is based on probabilistic assumptions. Such comparisons show that zero-momentum, thermal hadron correlation functions in the ground state channels, *i.e.* the vector (J/ψ, Υ) and pseudo-scalar (η_c, η_b) channels show only little modification in a thermal medium up to temperatures $T \gtrsim 1.5 T_c$. Correlation functions corresponding to radially excited charmonium states (χ_c), however, are modified strongly already close to or at T_c.

These generic features are reflected by the spectral functions. Although results from different groups currently still differ in details, the current understanding of charmonium at high temperature may be summarized as follows:

- The J/ψ and η_c remain unaffected by the thermal medium up to $T = 1.5\ T_c$. At higher temperatures it is unclear whether the J/ψ already disappears at $\simeq 1.9 T_c$ [105] or persists as a strongly modified resonance up to $2.25 T_c$ [23], see the left-hand side of Fig. 19.
- The $\chi_{c,0}$ and $\chi_{c,1}$ both disappear at $T \lesssim 1.1 T_c$, see the right-hand side of Fig. 19.

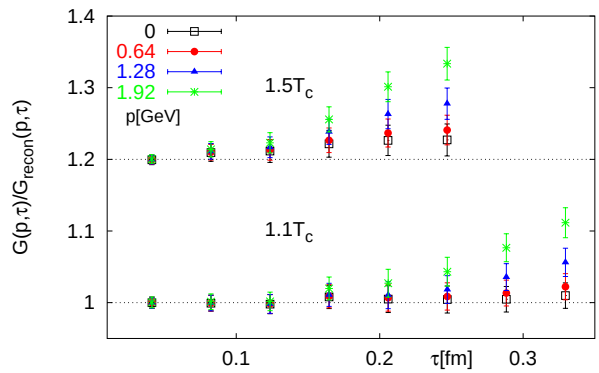

Fig. 20: The ratio of pseudoscalar correlation functions at non-zero momentum above and below T_c [106]. Deviations from the horizontal dotted lines for increasing p indicate stronger temperature-dependent modifi cations of the spectral functions.

- The J/ψ states with finite momentum show statistically significant but still small modifications for $T \lesssim 1.5T_c$ [106] due to collision broadening by higher momentum gluons seen by bound states moving relative to the heat bath, see Fig. 20 [107].

The strong binding of the J/ψ above T_c also is supported by the analysis of spatial correlations [108] and the observed insensitivity of the thermal vector and pseudo-scalar correlation functions to spatial boundary conditions [109].

Bottomonium studies are considerably more difficult since a larger lattice cut off is required to properly resolve these states, particularly for temperatures well above $2T_c$ where the Υ states are expected to be dissolved. First exploratory finite temperature results on bottomonium have been reported for temperatures up to $\sim 1.5T_c$. At this temperature, no thermal modifications of the Υ and η_b have been observed, as expected. The χ_b correlation functions are, however, modified at $T \sim 1.5T_c$, similar to the scalar charmonium case at $T \simeq 1.1T_c$. To firmly establish the onset of medium modifications in bottomonium states, however, requires further refined studies of the bottomonium system at lower and higher temperatures.

Dynamical Coalescence:

The production of multiple $c\overline{c}$ pairs in a single collision introduces a new charmonium formation mechanism [17]. In-medium charmonium formation utilizes a c and a $\overline{c}$ quark from independently produced $c\overline{c}$ pairs to form a J/ψ. Predictions of J/ψ production by this dynamical coalescence suffer from substantial uncertainties due to the dependence on the charm quark distributions in the medium. In fact, it possible to turn this uncertainty into an advantage and probe the medium properties using the observed J/ψ momentum distributions.

Two extremes can be considered [16]. If the charm quark distributions in the medium are identical to those of the initial production process, the interactions of charm quarks with the medium would be very weak. In this case, both the J/ψ rapidity and p_T distributions will be narrower than if no plasma is formed simply because the center of mass energy of secondary J/ψ production is lower than that of the initial nucleon-nucleon interactions. The lower energy results in a reduced $\langle p_T^2 \rangle$ and a narrower rapidity distribution. Thus, instead of the transverse momentum broadening expected from initial-state multiple scattering going from p+p to p+A to A+A, the average p_T^2 in A+A would no longer exhibit the monotonic increase seen in p+p and p+A interactions for increasing A. On the other hand, if the charm quarks are assumed to be in thermal equilibrium with the surrounding medium, the charm interaction

with the medium would be very strong. Any J/ψ's produced from thermalized charm quarks flowing with the medium would have a p_T distribution with a slope characteristic of the temperature of the system at the time they were formed, resulting in considerably narrower rapidity and p_T distributions. In either case, the effect would be largest in central collisions, reverting to "normal" broadening in peripheral collisions where on the order of one or fewer $c\overline{c}$ pairs will be produced since the number of $c\overline{c}$ pairs scales approximately with the number of collisions.

In order to extract the medium properties from secondary J/ψ production, systematic studies of J/ψ production in p+p, p+A and A+A interactions are necessary. The p+p data sets the intrinsic transverse momentum scale for a particular energy while the p+A results determine the level of broadening due to cold nuclear matter effects which would then apply to A+A interactions.

Models of regeneration, of course, also include J/ψ suppression. In addition to the screening effects discussed previously, the J/ψ can scatter with quarks and gluons in the plasma which may break it up more effectively than screening effects alone, especially if temperatures significantly above T_c are needed for screening to dissociate the directly produced J/ψ, as discussed in Ref. [28]. At low temperatures, relevant for SPS energies, $gJ/\psi \rightarrow c\overline{c}$ with a thermal gluon is effective for J/ψ breakup. However, at higher temperatures where the J/ψ should be more loosely bound, inelastic parton scattering, $g(q,\overline{q})J/\psi \rightarrow g(q,\overline{q})c\overline{c}$, calculated using the leading order matrix elements for gc and gq scattering, is more effective.

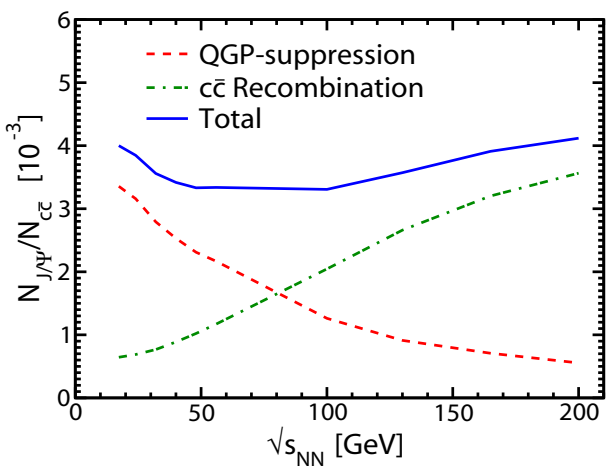

Fig. 21: Excitation function of the ratio of produced J/ψ to the number of $c\overline{c}$ pairs in central heavy-ion collisions for $N_{\rm part} = 360$ [110].

The relative importance of J/ψ suppression and regeneration will change as a function of energy, as shown in Fig. 21 for central collisions of heavy nuclei, $N_{\rm part} = 360$, from Ref. [110]. The J/ψ yield is dominated by primordial production at SPS energy, and dominated by regeneration at RHIC full energy.

5.2 Status of Quarkonium Physics at the CERN/SPS

The prospects of a "clean" QGP signature, destruction of the J/ψ by color screening, was discussed in the landmark paper by Matsui and Satz in 1986 [21]. This triggered an extensive experimental program at the CERN/SPS. HELIOS-III [111] and NA38 [112] (subsequently NA50 [113] and currently NA60 [114]) conducted detailed measurements of the dimuon invariant mass spectrum around mid rapidity. Despite early enthusiasm and enormous statistics (see Fig. 22) the picture that evolved is still rather ambiguous. The SPS measurements must also be understood in light of the many results on quarkonium production in p+A collisions from fixed target experiments. The status of the SPS program can be summarized as follows:

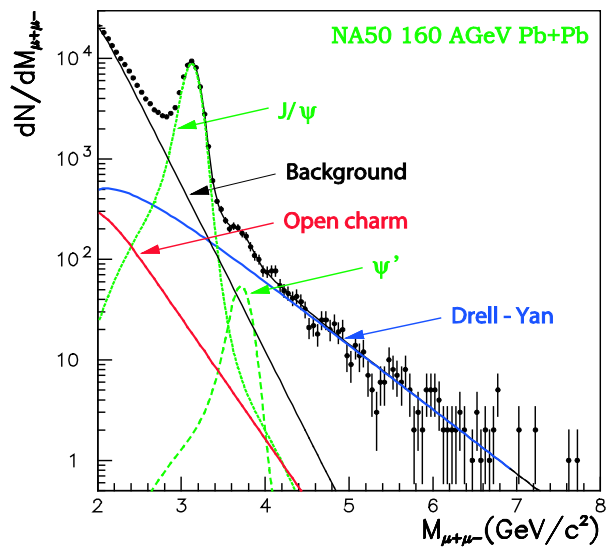

Fig. 22: Dimuon invariant mass spectrum from 158 AGeV Pb+Pb collisions at NA50 [113].

Feed down contributions (see Fig. 23) from higher charmonium states, $\chi_c \to J/\psi\gamma$ ($\sim$ 30%) and $\psi' \to J/\psi\pi\pi$ ($\sim$ 10%), are important [115, 116]. The χ_c has not yet been measured by the heavy ion detectors at the SPS, although it has been seen in other experiments there. These measurements are extremely difficult and the large scatter of available data depicted in Fig. 24 indicates that better measurements are desperately needed. The NA60 experiment is planning to conduct this analysis, although the feasibility with the present data sets still has to be verified.

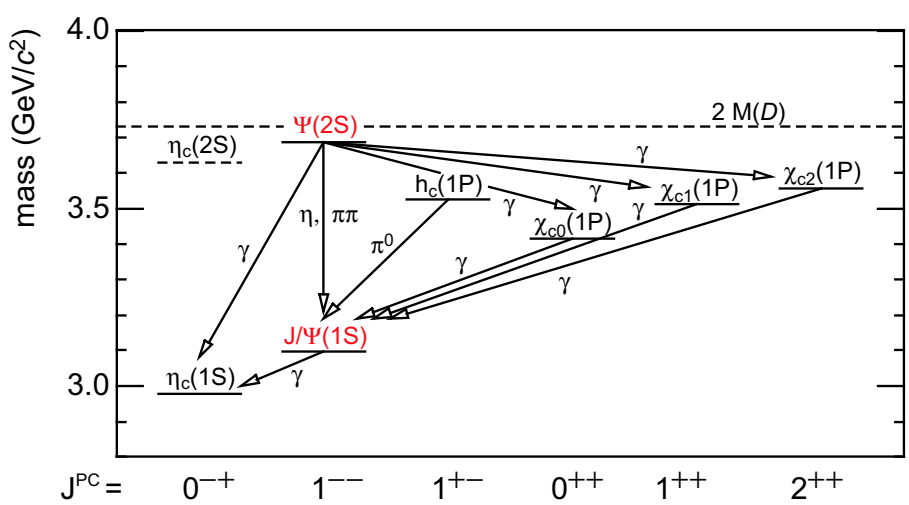

Fig. 23: Charmonium mass levels and spin states. The common feed down channels are indicated.

The J/ψ and ψ' have substantial absorption cross sections in normal nuclear matter, 4.2 and 9.6 mb respectively at midrapidity, determined from fits to p+A data (see Fig. 25) [117]. Studies of the A dependence were made at Fermilab over a large x_F range, but with higher $\sqrt{s}$ [118]. A very strong x_F dependence was observed for $x_F > 0.2$, as depicted in Fig. 26. Effects like shadowing, absorption and energy loss play varying roles at different x_F, resulting in the observed dependence [2].

The J/ψ is suppressed in semi-central and central Pb+Pb collisions [113] beyond absorption by nucleons alone, as shown in Fig. 27. Shadowing has not yet been included in the SPS analysis. The

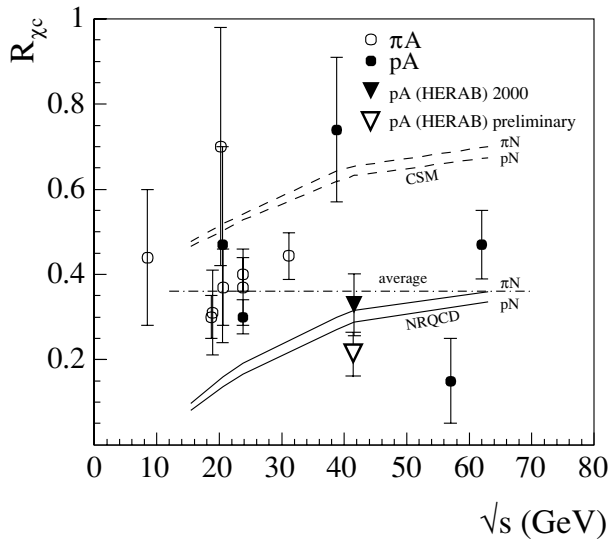

Fig. 24: The fraction, R_{χ_c}, of observed J/ψ's originating from radiative $\chi_{c1,2} \to J/\psi\gamma$ decays as a function of energy for proton and pion beams [115].

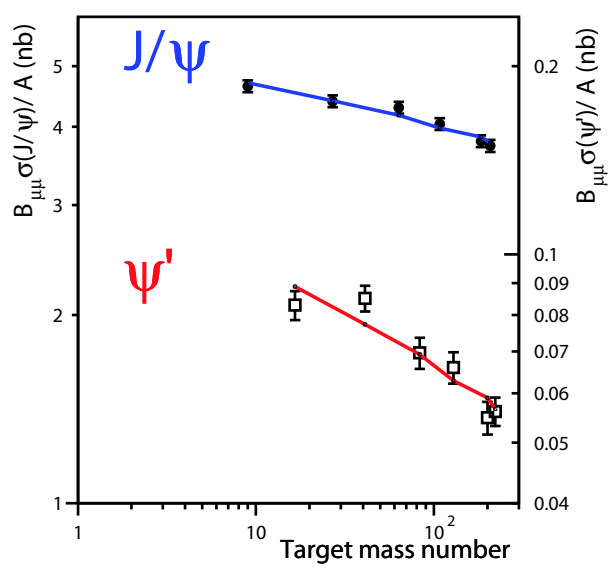

Fig. 25: NA50 measurements of J/ψ and ψ' absorption in 450 AGeV fi xed-target p+A collisions. The respective absorption cross sections are $\sigma_{\rm abs}(J/\psi) = 4.2 \pm 0.5$ mb and $\sigma_{abs}(\psi') = 9.6 \pm 1.6$ mb. These measurements are at midrapidity only ($x_F \sim 0$) [117].

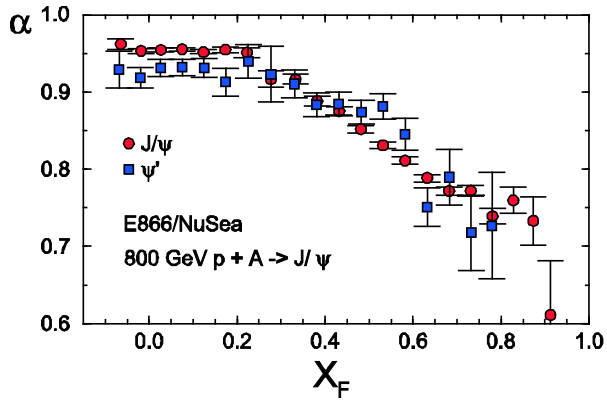

Fig. 26: Measurement of J/ψ and ψ' absorption in 800 GeV p+A collisions as a function of x_F at the Tevatron [118].

suppression observed in A+A interactions at the SPS can, for the most part, be accounted for by the assumption that the more loosely bound ψ' and χ_c states are both suppressed by plasma production, eliminating their contribution to the inclusive J/ψ measurement. The direct J/ψ contribution is assumed not to be suppressed at the SPS [119, 120].

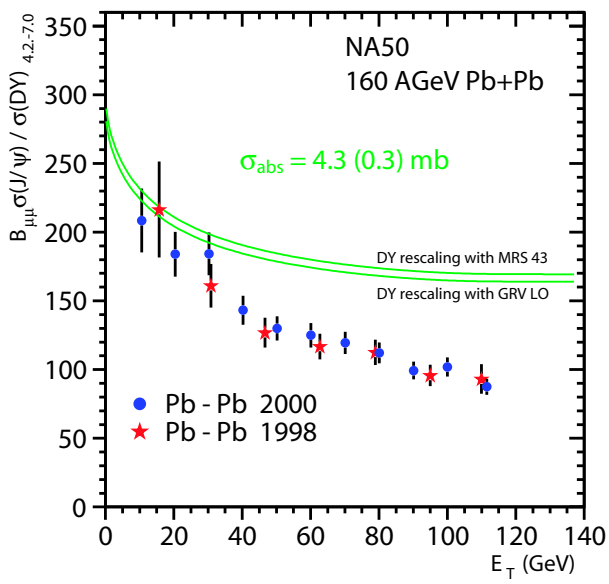

Fig. 27: J/ψ absorption in 158 AGeV Pb+Pb collisions from NA50 [117]. The curves depict predictions assuming nuclear absorption alone. Suppression beyond cold nuclear matter effects is seen starting from $E_T \sim 30$ GeV, increasing to a factor of ~ 1.7 at the highest E_T [117].

Alternatives to the QGP models are able to describe the observed J/ψ suppression by assuming breakup of the bound state by comoving matter [35]. Although these approaches make some unrealistic assumption about the hadron density, it is possible that a fraction of the observed suppression is due to comover absorption.

The ψ', lying 50 MeV below the $D\overline{D}$ threshold, is rather loosely bound and can more easily be broken up by interactions in the medium. The strong ψ' suppression measured by NA50 has been interpreted as both total suppression of the ψ' by color screening [120] and a larger interaction cross section for comovers [121].

Charm production is so far not measured in heavy ion experiments at the SPS. Fig. 22 shows that

the open charm contribution to the dilepton continuum in the J/ψ mass region is negligible at the SPS. Open charm measurements are, however, key to understanding the intermediate mass dilepton region. The NA60 experiment has used displaced vertices to separate charm decays from prompt dileptons. They are currently analyzing In+In data that will shed light on open charm contributions to the dilepton continuum.

The picture emerging from SPS studies is still somewhat inconclusive. The missing pieces of vital information have made quarkonium suppression seem to be an interesting but inconclusive study. Measurements from NA60 might provide some of the missing pieces although the future of the SPS program is currently rather uncertain. On the other hand, the vast experience gained at the SPS can and should be taken into account at RHIC. The main lesson learned is that a simple J/ψ measurement in A+A collisions as a function of centrality is insufficient to draw unique conclusions. Rather, a systematic and detailed study of all related aspects, *i.e.*, a systematic study of open charm, J/ψ, ψ', and χ_c production in p+p, p+A, and A+A collisions is required. Centrality, rapidity, and A dependence studies are mandatory.

5.3 Quarkonium measurements to date at RHIC

All of the published quarkonium results from RHIC to date are from PHENIX. However some proof-of-principle J/ψ results in a STAR poster at Quark Matter 2005 indicate that STAR will have J/ψ results from RHIC Runs 4 and 5, as well as from future RHIC runs. PHENIX measures quarkonium yields by reconstructing their invariant mass from decays to dileptons. Dielectrons are used in the central arms ($|\eta| < 0.35$) and dimuons are used in the muon arms ($1.2 < |\eta| < 2.2$). STAR uses dielectrons within the TPC acceptance ($|\eta| < 1$).

PHENIX has measured J/ψ yields at 200 GeV from p+p [122, 123], d+Au [123], Au+Au [34] and Cu+Cu collisions [34]. There are also J/ψ measurements from Cu+Cu collisions at 62 GeV [34].

At Quark Matter 2005, PHENIX also reported an observation of $\Upsilon \to \mu^+\mu^-$ in 200 GeV p+p collisions from RHIC Run 5 [124]. This was a very low statistics measurement (27 counts in both muon arms). Clearly it will be difficult to make definitive Υ yield measurements at RHIC I luminosities, but a crude Υ measurement from both PHENIX and STAR may be possible at RHIC I with about 10 times the existing integrated luminosity.

5.31 Baseline quarkonium measurements at RHIC

PHENIX has measured J/ψ cross sections in p+p and d+Au collisions at 200 GeV [123]. The rapidity dependence is summarized in Fig. 28. The left side shows the invariant J/ψ yields in p+p collisions while the right side shows the nuclear modification factor, R_{dAu}, for minimum bias d+Au collisions.

The curves on the right-hand side of Fig. 28 show the results of several calculations that include absorption and shadowing [2, 89, 125], discussed in the previous section. The data favor the relatively modest shadowing of the EKS98 parameterization with moderate nuclear absorption.

The nuclear modification factor is shown as a function of centrality in Fig. 29 for the forward, mid and backward rapidity regions covered by the three PHENIX arms. The curves are calculated with the XXX and YYY shadowing parameterization and a 3 mb absorption cross section. (Check with Mike L. about absorption).

5.32 Quarkonium measurements in heavy ion collisions at RHIC

At RHIC, preliminary measurements with the statistical precision needed to provide a strong test of models of J/ψ production in heavy ion collisions were released for the first time at Quark Matter 2005. These preliminary PHENIX measurements were from the Run 4 Au+Au and the Run 5 Cu+Cu data sets. The main features are summarized here.

Figure 30 shows the J/ψ nuclear modification factor, R_{AA}, measured in 200 GeV Au+Au and

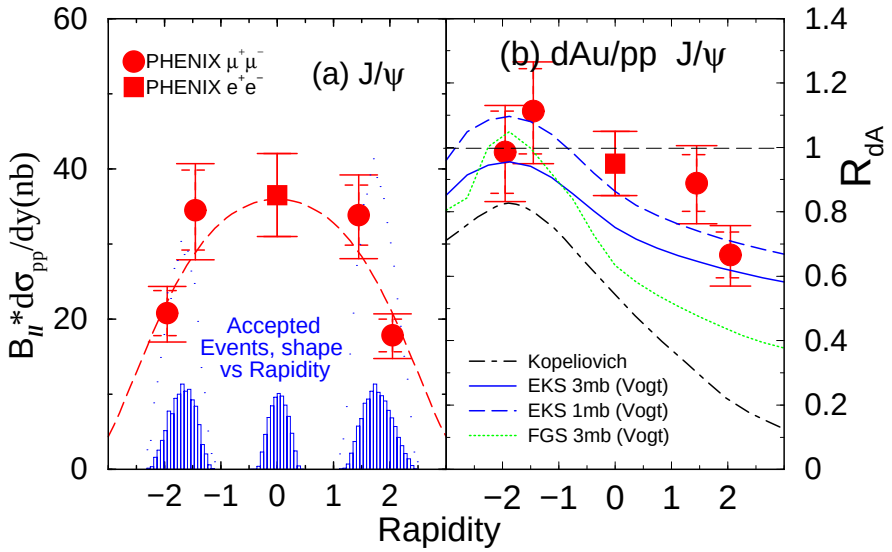

Fig. 28: The rapidity dependence of the p+p J/ψ invariant yield at 200 GeV (left side) and the nuclear modification factor for minimum bias d+Au collisions (right side) [123]. The curve on the left side is a fit used to extract the total cross section. The curves on the right side are theory calculations discussed in the text. The deuteron is defined to moving towards positive rapidity.

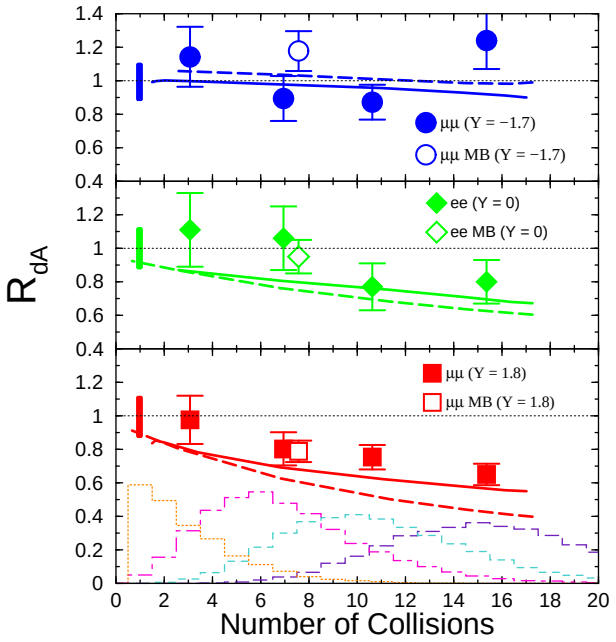

Fig. 29: Nuclear modification factor as a function of centrality for d+Au collisions measured at forward (bottom), mid (center) and backward (top) rapidity. The deuteron is moving toward forward rapidity. The theoretical curves, including both shadowing and nuclear absorption, are discussed in the text.

Cu+Cu collisions at both mid- and forward/backward-rapidities [34]. These preliminary data have substantial systematic and statistical errors which are expected to decrease somewhat in the final analysis. It is thus difficult to claim that there are any differences in behavior between the Au+Au and Cu+Cu systems or between the two different rapidity regions for each species combination. The data in Fig. 30 are compared with baseline calculations of cold nuclear matter effects [3] assuming EKS98 shadowing. These calculations are intended to show the expected result if the J/ψ did not interact with the hot, dense medium beyond the cold nuclear matter effects observed in d+Au collisions. The calculations also assume a 3 mb final state J/ψ absorption cross section. Note that these same parameters slightly overpredict the suppression in the d+Au data (see Fig. 29).

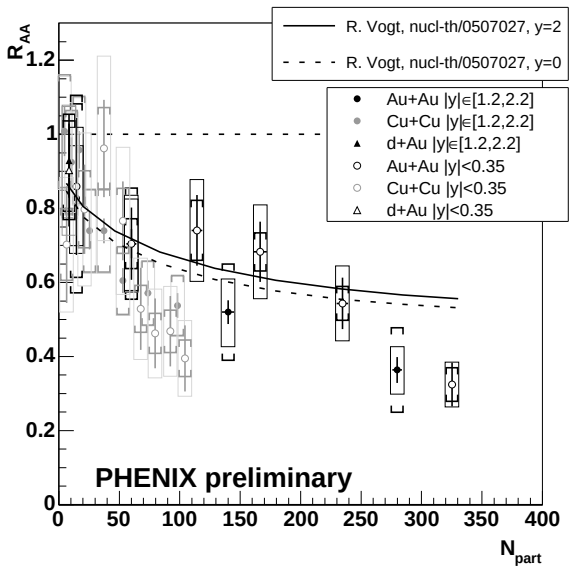

Fig. 30: *Replace this with a figure showing 1 mb absorption cross section* The nuclear modification factor as a function of centrality for 200 GeV Au+Au and Cu+Cu collisions measured at forward, mid and backward rapidity. The calculations [3] include cold nuclear effects from shadowing and "normal" final state absorption.

the same data are compared with models tuned to the CERN SPS results and extrapolated to RHIC on the left-hand side of Fig. 31, either by comover absorption [35] or by color screening in a plasma [36, 19]. No secondary production by $c\overline{c}$ coalescence is included. All the calculations strongly overestimate the suppression. The model calculations on the right-hand side of Fig. 31 include either coalescence [36, 19, 18] or transport in the medium [126, 127]. These latter calculations all differ in the way cold nuclear absorption is accounted for, in the J/ψ p+p production cross section used for normalization, and in the $c\overline{c}$ production cross section used in the coalescence calculation.

Figure 31 shows that the data are consistent with models that include both plasma screening and quark coalescence in the final state, although some of these models tend to underpredict the suppression in the most central collisions, as is the case for cold nuclear matter effects alone in Fig. 30. The calculations in Refs. [36, 19] tuned to the SPS data require either some coalescence not to overpredict the suppression [36] or a significant open charm enhancement [19] not supported by the NA60 data [128]. At RHIC, the coalescence component completely dominates J/ψ production in central Au+Au collisions and contributes about 50% of the J/ψ yield in central Cu+Cu collisions.

PHENIX also showed a measurement of the nuclear modification factor for 62 GeV Cu+Cu collisions at Quark Matter 2005. While this measurement has relatively low statistics, the 62 GeV data exhibit similar, and perhaps slightly stronger, suppression in the most central collisions than that seen in

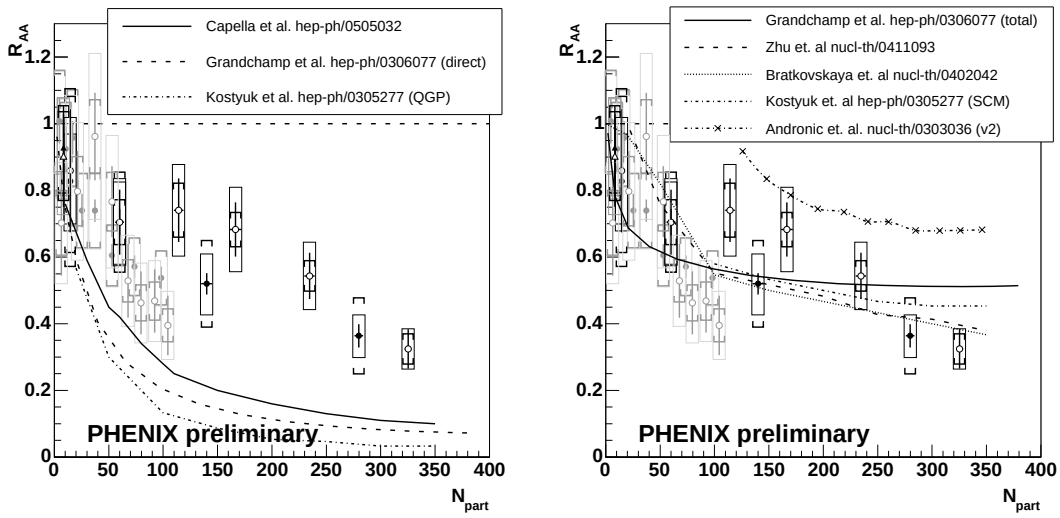

Fig. 31: Nuclear modification factor versus centrality for 200 GeV Au+Au and Cu+Cu collisions measured at forward, mid and backward rapidity, compared with various models that include final state effects from QGP suppression or comover absorption (left-hand side) and including coalescence or J/ψ transport (right-hand side). See the text for details. *The Bratkovskaya curve has problems - kink is a plotting error. Also which Andronic curve is used?*

200 GeV Cu+Cu collisions, consistent with predictions of a model with color screening and coalescence [110] (see Fig. 21).

Recently there have been theoretical efforts to predict the effects of $c\bar{c}$ coalescence on the rapidity and p_T dependence of the J/ψ yield. Figure 32 shows experimental results on the average p_T^2 of the J/ψ, $\langle p_T^2 \rangle$, as a function of $N_{\rm coll}$ for Au+Au and Cu+Cu collisions in the rapidity regions covered by PHENIX. The predictions are from a coalescence model by Thews and Mangano [16] (check!) and show the predicted $N_{\rm coll}$ dependence of $\langle p_T^2 \rangle$ for Au+Au and Cu+Cu collisions assuming either no coalescence or a realistic coalescence contribution. A large coalescence component results in a narrower J/ψ p_T distribution since a c and $\overline{c}$ with large relative p_T cannot coalesce. Thus $\langle p_T^2 \rangle$ is reduced by a factor of two in central collisions. The data have substantial statistical errors, particularly at midrapidity, but they do appear to favor the coalescence calculation.

No realistic model calculations of the rapidity distribution are available for heav y ion collisions so far, but it is predicted [16] that a strong charm coalescence component of J/ψ production will lead to a narrowing of the rapidity distribution, just as for the p_T distribution. Figure 33 shows the p+p, Cu+Cu, and Au+Au rapidity distributions measured by PHENIX. The Au+Au rapidity distributions are shown in three centrality bins while the Cu+Cu distributions are shown in four bins. The statistical and systematic uncertainties are fairly large, particularly for Au+Au. Within those uncertainties, the data show no evidence of a narrowing of the rapidity distribution in central heavy ion collisions.

A more quantitative statement can be made for Cu+Cu collisions, the highest statistics heavy ion data set so far. Figure 34 shows preliminary PHENIX results on the rapidity dependence of R_{AA} for minimum bias collisions (left) and for the muon data in several centrality bins (right). Also shown on these plots are baseline cold nuclear matter calculations [3], discussed in the previous section, and predictions of the effect on the R_{AA} rapidity dependence [16] if all of the J/ψ yield was due to coalescence. There are considerable uncertainties but the data show little evidence of strong rapidity narrowing in central Cu+Cu collisions. The uncertainties will be reduced significantly when the higher statistics Run 5 p+p reference data analysis is completed. The coalescence contribution is predicted to be approximately 50%

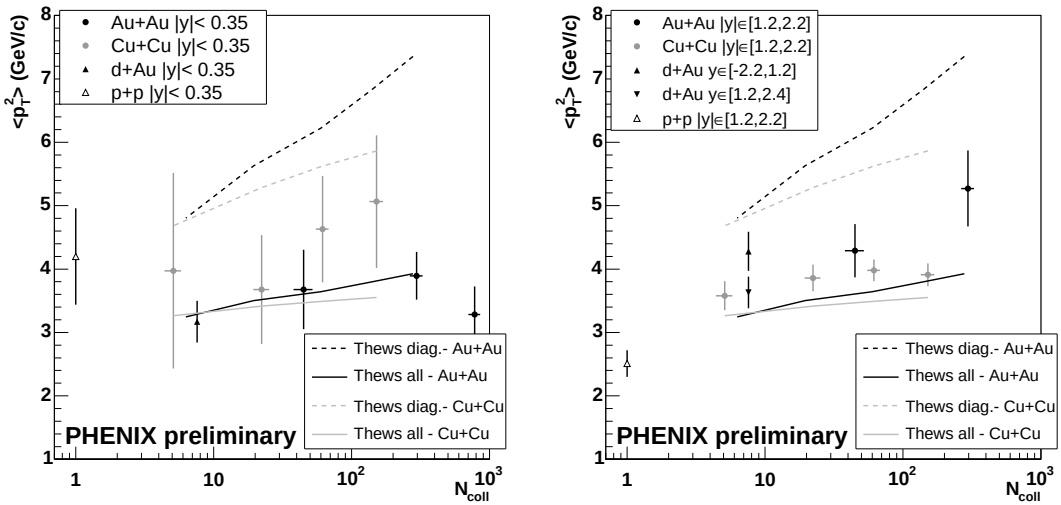

Fig. 32: The measured $\langle p_T^2 \rangle$ for 200 GeV Au+Au and Cu+Cu collisions measured at (left) and away from (right) midrapidity. The calculations [16] show the predicted behaviour with (upper curves, labeled 'all') and without (lower curves, labeled 'diag.' for correlated $c\overline{c}$ pairs only) the coalescence contribution.

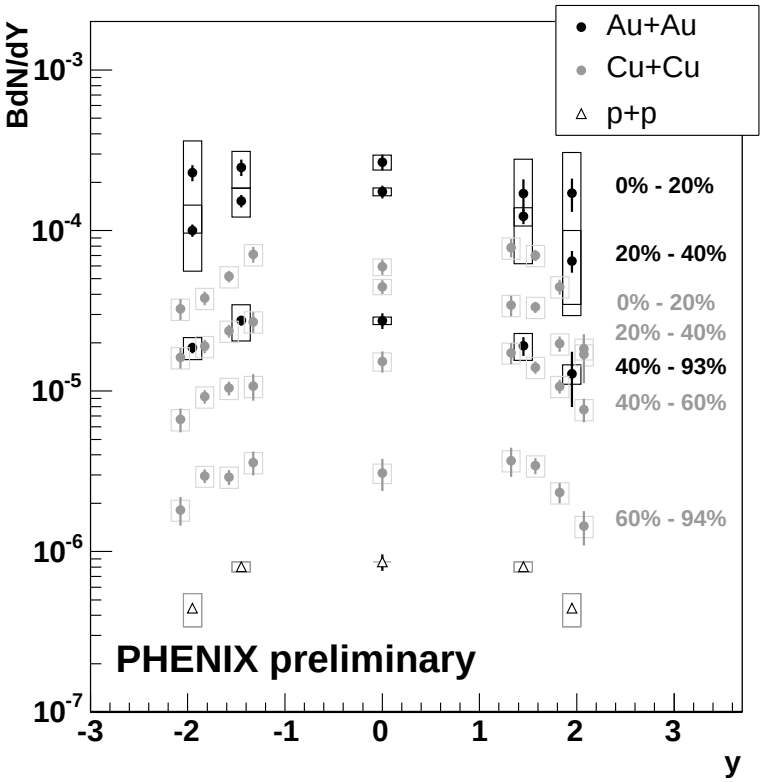

Fig. 33: The rapidity dependence of the J/ψ invariant yield for p+p, Cu+Cu and Au+Au collisions at various centralities (with 0% being most central). Within the large uncertainties, there is no signifi cant change in the shape of the rapidity distribution

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for the number of participants in central Cu+Cu collisions [36] so that some rapidity narrowing is expected in central Cu+Cu. But the results of Ref. [3], presented in Fig. 15, show significant modification of the J/ψ rapidity distribution (and thus presumably of the open charm distribution) due to shadowing, causing the midrapidity yield to be smaller relative to that at $|\eta| \sim 1.7$ in central collisions. Thus the observation that there is no dramatic rapidity narrowing may be consistent with coalescence models that include shadowing in the initial $c\bar{c}$ production.

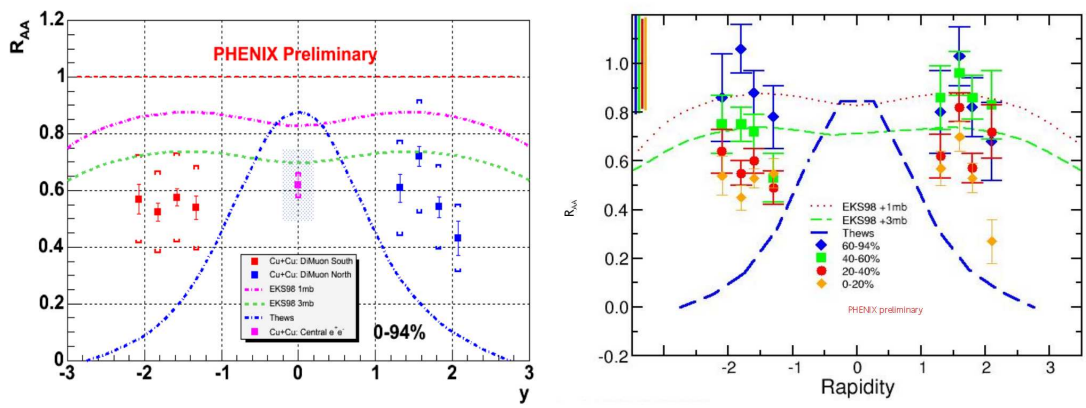

Fig. 34: The J/ψ R_{AA} as a function of rapidity for minimum bias Cu+Cu collisions (left) and binned in centrality in the muon arm (right). Within the large uncertainties, there is no signifi cant change of R_{AA} with rapidity. The curves are discussed in the text.

In summary, the evidence for a strong coalescence component in the J/ψ yield in central collisions is mixed. The suppression as a function of centrality shown in Fig. 31 is reasonably consistent with models that include coalescence, as is $\langle p_T^2 \rangle$ in Fig. 32. But the predicted narrowing of the rapidity distribution due to coalescence, based on the assumption of an underlying charm distribution peaked near midrapidity, does not seem to be supported by the data in Figs. 33 and 34. The open charm rapidity distribution has not yet been very well determined experimentally. As mentioned earlier, shadowing is predicted to modify the underlying charm rapidity distribution. The prediction of p_T narrowing, on the other hand, is based on a steeply falling charm p_T distribution, well established by the existing data.

A J/ψ result of great interest would be R_{AA} as a function of p_T. A useful result is expected from the PHENIX data once the RHIC Run 5 p+p analysis is completed, providing much better statistical precision for the high p_T p+p reference than presently available. There are also analyses underway in PHENIX to extract the J/ψ v_2 but the limited statistics of the present heavy ion data sets will likely preclude a meaningful result. Similarly, a statistically meaningful J/ψ polarization measurement is unlikely from the present data sets.

5.4 Proposed RHIC II quarkonia measurements

Unlike other probes, quarkonia measurements are guided by predictions from lattice QCD calculations. Color screening modifies the linear rise of the QCD potential at large distances. The quarkonia spectral functions quantify the temperature dependence of the potential. Since quarkonia suppression is determined by the plasma temperature and the binding energy (equivalently the quarkonium size and the Debye screening length), measuring the sequential disappearance of these states acts as a QCD thermometer.

Thus the importance of a comprehensive study of **all** experimentally accessible quarkonium states cannot be overstated. A systematic study of heavy quarkonium spectroscopy, with a complete determi-

nation of the suppression pattern of the quarkonium states, remains the most **direct probe of deconfinement**. It is also the signature that most closely acts as a thermometer of the hot initial state, which, with future improved lattice calculations, can be directly compared to QCD.

While J/ψ physics at RHIC is as compelling as it was in 1986 when first proposed by Matsui and Satz [21], the systematic study of all quarkonia states, and especially bottomonium, feasible at RHIC II, provides a more complete QGP probe than heretofore possible.

Table 10 relates the main physics topics to the relevant probes and subsequent detector requirements. The ability of a program at RHIC II to make these measurements can be judged from the yields given in Tables 2, 3, 4. The measurements that are possible at RHIC without the luminosity upgrade are the J/ψ rapidity and p_T distributions at full energy. The measurements that are newly possible at RHIC II are those for the excited charmonium states ($\psi\prime$ and χ_c) and the bottomonium states ($\Upsilon(1S)$, $\Upsilon(2S)$ and $\Upsilon(3S)$). For the J/ψ, measurements of v_2, polarization, and excitation functions of heavy flavor distributions will be possible only at RHIC II. It is evident that a comprehensive program to use quarkonium as a QCD thermometer to provide direct evidence of deconfinement is possible only with RHIC II luminosity.

The measurements needed to study the excited charmonium states, χ_c and $\psi\prime$, have quite different problems. The $\psi\prime$ measurement uses the same method as that for the J/ψ, reconstruction of dilepton decays, but requires $\sim$ 100 times as much integrated luminosity for the same yield. In addition, the $\psi\prime$ yield measurement is made more difficult by the existence of a significant background under the peak in the invariant mass spectrum. The signal to background is worse than for the J/ψ, and this raises the integrated luminosity needed for a measurement of a given precision. But measuring $\psi\prime$ yields at RHIC II is feasible. The χ_c measurement is made using the $\chi_c \rightarrow J/\psi\ \gamma$ channel, where the J/ψ is detected normally by reconstructing dilepton decays, and the photon is detected in an electromagnetic calorimeter. The yields are larger than for the $\psi\prime$, but the need to make an invariant mass peak by combining each J/ψ candidate with a large number of photons means that combinatorial backgrounds will be quite large. Thus the χ_c measurement will be difficult. There are ongoing simulation studies by PHENIX and STAR to determine how difficult this will be.

Like the J/ψ, the bottomonium states are studied using their dilepton decays. The bottomonium measurements require very large integrated liminosity and good invariant mass resolution. PHENIX expects to be able to resolve the $\Upsilon(1\mathrm{S})$, $\Upsilon(2\mathrm{S})$, and $\Upsilon(3\mathrm{S})$ states. Because of its larger acceptance, STAR will have $\sim$ 10 times larger Υ yields than PHENIX, but the states will not be cleanly resolved and fitting will be required to extract individual yields. Although the yields are small relative to the J/ψ, bottomonium measurements are quite clean. The states are massive ($\sim$ 10 GeV/c^2) so that their decay leptons have relatively large momenta and are thus easily distinguished from background leptons. The combinatorial background in this mass range is extremely small and multiple scattering is of less concern. While the interpretation of charmonium suppression is made more difficult by the rather large cross section for nucleon and comover absorption, the situation for bottonium is considerably better. Calculations of the absorption of directly produced bottomonium by hadronic comovers show that this effect is essentially negligible [129].

The dependence on $\sqrt{s}$ of the ratio of produced J/ψ relative to the number of $c\bar{c}$ pairs, depicted in Fig. 21 [110] is striking. A measurement of an excitation function of this ratio over a range of $30 < \sqrt{s_{NN}} < 200$ GeV could help to disentangle suppression from enhancement mechanisms such as recombination/coalescence. Such measurements, however, are extremely demanding in terms of statistics since both heavy quarks and quarkonia will need to be measured with good statistics over a wide range of beam energies.

A measurement of the quarkonium nuclear modification factor at high-p_T can provide a unique experimental probe for studying energy loss and color diffusion [130]. At relatively large transverse momentum, suppression due to color screening and coalescence, become negligible. Instead, the quarkonium state is a hard probe that interacts with the medium. In particular, any color octet can suffer energy

Physics Motivation	Probes	Measurements	Requirements
Baseline measurements	J/ψ, ψ', $\Upsilon(1\mathrm{S})$, $\Upsilon(2\mathrm{S})$, and $\Upsilon(3\mathrm{S})$ decays to dileptons	Rapidity and p_T spectra in p+A, and p+p as a function of $\sqrt{s}$	High luminosity and acceptance for sufficient statistics, especially for the Υ family. Good mass resolution to resolve ψ. and Υ states.
Deconfinement, Initial Temperature	J/ψ, ψ', $\Upsilon(1\mathrm{S})$, $\Upsilon(2\mathrm{S})$, and $\Upsilon(3\mathrm{S})$ decays to dileptons	Suppression patterns in A+A vs $\sqrt{s}$ and A	High luminosity, acceptance and mass resolution. Triggers that work for quarkonium in Au+Au collisions.
Thermalization and Transport	J/ψ	J/ψ v_2 as function of $\sqrt{s}$ and A.	High luminosity for good statistics in short runs for $\sqrt{s}$and A scans

Table 10: Main physics goals of the RHIC-II quarkonium program and corresponding probes, studies, and requirements.

loss. The relative abundance of charmonium resonances can provide an experimental handle on studying such phenomena as each resonance may have a different octet contribution. We must exercise caution, however, as competing charmonium production models exist. In parallel with any nucleus-nucleus studies, it is therefore important to investigate and compare production mechanisms in p+p and p+A interactions, all at low (central region) and high x_F (forward region) [131, 132, 133].

It is crucial for the interpretation of the quarkonia yields from A+A collisions to understand the feed down contributions from the χ_c states (see Fig. 23). By far the best measurement of this will be made using 500 GeV p+p collisions, because the increased luminosity and increased charmonium production cross sections lead to $\sim$ 10 times larger charmonium yields than in 200 GeV p+p collisions. It is not expected that the χ_c feeddown will change very much between 200 and 500 GeV, so the increased yield at 500 GeV is expected to yield a very definitive baseline measurement of the χ_c feeddown in p+p collisions.

Recently the measurement of quarkonium polarization was also suggested as a possible signature for QGP formation [134]. The quarkonium yields at RHIC II will be large enough to permit a J/ψ polarization measurement at low p_T by both PHENIX and STAR.

6. Relationship to the LHC program

The major differences for quarkonium studies at RHIC II and at the LHC are:

- The temperature in 5.5A TeV/c central Pb+Pb collisions at the LHC is expected to be $\sim 3.5\ T_C$ [136], while in 200A GeV/c central Au+Au collisions at RHIC the temperature is expected to be $\sim 2\ T_C$ [135].
- Heavy flavour production cross sections are much larger at the LHC. The cross sections for open charm and bottom production are $\sim\ 15$ and $\sim\ 100$ times higher, respectively [9] at the LHC. For charmonium and bottomonium they are ~ 13 and ~ 55 times higher, respectively [37].
- At LHC the number of $c\bar{c}$ and $b\bar{b}$ pairs created in a central Pb+Pb collision is estimated to be ~ 115 and ~ 5, respectively[137]. At RHIC, the number of $c\bar{c}$ and $b\bar{c}$ pairs created in a central Au+Au collision is estimated to be ~ 10 and ~ 0.05, respectively [138].
- The Au+Au luminosity at RHIC II is projected to be 14 times larger than the Pb+Pb luminosity at the LHC (7 $10^{27} cm^{-2} s^{-1}$ versus 5 $10^{26} cm^{-2} s-1$).
- The heavy ion running times per year at RHIC II are expected to be considerably longer than at the LHC. Taking into account that there will also be a polarized p+p program at RHIC II, the heavy ion program is expected to get $\sim\ 12$ weeks of physics running on average. The heavy ion program at the LHC will be allocated 1 month of physics running.

The larger heavy flavor cross sections at the LHC are approximately balanced by the increased luminosity and running times at RHIC II, making the heavy flavor yields per year similar. This means that the kinds of measurements that can be made at the two facilities will also be similar, and will be of similar quality (see Tables 2, 3 and 5). But there will be important differences in the physics environments prevailing at the two facilities, and these differences will make the two programs complementary to each other.

The higher initial energy density at the LHC means that the QGP will be created at a significantly higher temperature, and thus there is a strong potential for new physics effects to be observed at the LHC. In addition, the factor of ten increase in $c\bar{c}$ pairs and the factor of 100 increase in $b\bar{b}$ pairs per central collision at the LHC will have a major impact on the interpretation of heavy flavor measurements. We will discuss some of those differences here.

Lattice calculations suggest that the J/ψ may remain bound at the highest RHIC temperatures, while the excited charmonium states are predicted to be unbound. At the LHC all of the quarkonium states should be unbound at the highest temperatures. This implies that almost all charmonium production in central Pb+Pb collisions at the LHC will be due to coalescence of $c\bar{c}$ pairs as the QGP cools below the relevant charmonium dissociation temperature, and finally hadronizes. Thus the prompt charmonium yields at the LHC should reflect only the coalescence mechanism, with no contribution from the primordial J/ψ production (except for J/ψ produced in very peripheral collisions). The measurements at RHIC and the LHC will thus provide very different windows on charmonium suppression in the QGP that will help to resolve the ambiguities in interpreting data caused by the balance between destruction and coalescence formation of charmonium at RHIC.

Because of its higher binding energy, bottomonium should behave at the LHC similarly to charmonium at RHIC. The $\Upsilon(1S)$ state is expected to be bound at the highest temperature at the LHC, with the other bottomonium states all unbound. Given the estimated $b\bar{b}$ pair yield of ~ 5 in a central Pb+Pb collision (compare ~ 10 $c\bar{c}$ pairs at RHIC), the Υ yield at the LHC is predicted [139] to reflect a balance between dissociation and coalescence reminiscent of the RHIC J/ψ production models. At RHIC the situation will be significantly different. The $\Upsilon(1S)$ state is predicted to be firmly bound, the $\Upsilon(2S)$ state may be just bound, and the $\Upsilon(3S)$ state is likely unbound. Also, the $b\bar{b}$ pair yield at RHIC of $\sim\ 0.05$ per Au+Au central collision is so low that no significant coalescence yield of bottomonium is expected. Thus the bottomonium yields at RHIC II should reflect only the suppression in the QGP. The measurements at RHIC II and the LHC will thus provide very different windows on bottomonium suppression in the QGP

that will help to resolve the ambiguities in interpreting data caused by the balance between destruction and coalescence formation of bottomonium at the LHC.

The open heavy flavor programs at RHIC II and the LHC will consist of similar measurements with similar goals. They will study energy loss, thermalization and flow of heavy quarks in systems that will have different energy densities, interaction cross sections and lifetimes. There will be some differences in the problems that have to be dealt with. At $\sqrt{s} = 200$ GeV, bottom decays to leptons begin to dominate the single electron spectrum at $p_T \sim 4$ GeV/c. As the collision energy increases, the lepton spectra from B and D decays move closer together rather than further apart. Thus, the large increase in the $b\overline{b}$ cross section relative to $c\overline{c}$ does not make single leptons from B and D decays easier to separate. Preliminary calculations show that the $B \to e$ decay does become larger than that of $D \to e$ but at $p_T > 10$ GeV/c. The two lepton sources differ by less than a factor of two to $p_T \sim 50$ GeV/c in the range $|y| \leq 1$. Separation of single leptons from charm and bottom will require statistical separation using the difference in displaced vertex distributions at all p_T at the LHC. *Boost issues different at LHC?*. Thus interpretation of single lepton data from heavy flavor decays will be more difficult at the LHC.

For direct open charm measurements, ALICE can reconstruct D^0 decays from $p_T \sim 0$ to $p_T \sim 25$ GeV/c [38]. Like STAR, ALICE will not be able to trigger on D decays and will have to get these events from the minimum bias sample. Thus the longer running times at RHIC are an advantage, since they allow more minimum bias data to be taken (see Tables 3 and 5). While it is not yet clear what CMS and ATLAS will do to reconstruct charm, they should be able to do b jets well, similar to the Tevatron measurements. As at RHIC, B mesons can be measured cleanly at the LHC through their decays to J/ψ, although triggering on lower p_T J/ψ at LHC is difficult.

It has also been suggested that the $B\overline{B}$ contribution to the dimuon continuum, the dominant contribution above the Υ mass, can be used to measure energy loss [39]. That channel would be fairly clean at the LHC but more difficult at RHIC. *ADF: Yield for this? Sounds pretty tough to me!*

7. Summary

Add this section

Summarize what is the new physics we can expect to get from RHIC II.

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