

# Diffractive phenomena in pA and eA scattering.

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# Color fluctuations in the nucleon wave function & 3-dimensional mapping of the nucleon

$$|N\rangle = \textcircled{3q} + \textcircled{3qg} + \textcircled{3q + \pi} + \dots$$

Fluctuations of interaction strength are likely to be correlated with quark content of the hadron. Smallest configurations are most likely the minimal Fock state ones.

**Important feature of QCD:** QCD sum rules, form factors at large  $Q$ , chiral dynamics.

**Experimental evidence:** Pion diffraction into two jets, inelastic coherent diffraction off nuclei.

At RHIC energies the variance of the interaction strength is still large

$$\omega_\sigma \sim 0.25 \div 0.3 \quad \text{At LHC it is expected to drop to} \quad \omega_\sigma \sim 0.06$$

## Three methods of experimental study in coherent pA scattering

- Hard exclusive diffraction of proton into three jets
- Soft coherent diffraction:  $p + A \rightarrow X + A$
- Hard diffraction pA:  $p + A \rightarrow \text{"two jets"} + X + A$   
as a function of  $x$  of the parton in the proton.

Another option is study of the associated hadron production in hard scattering processes involving a leading quark or gluon in the proton like DY production at large  $x_p$

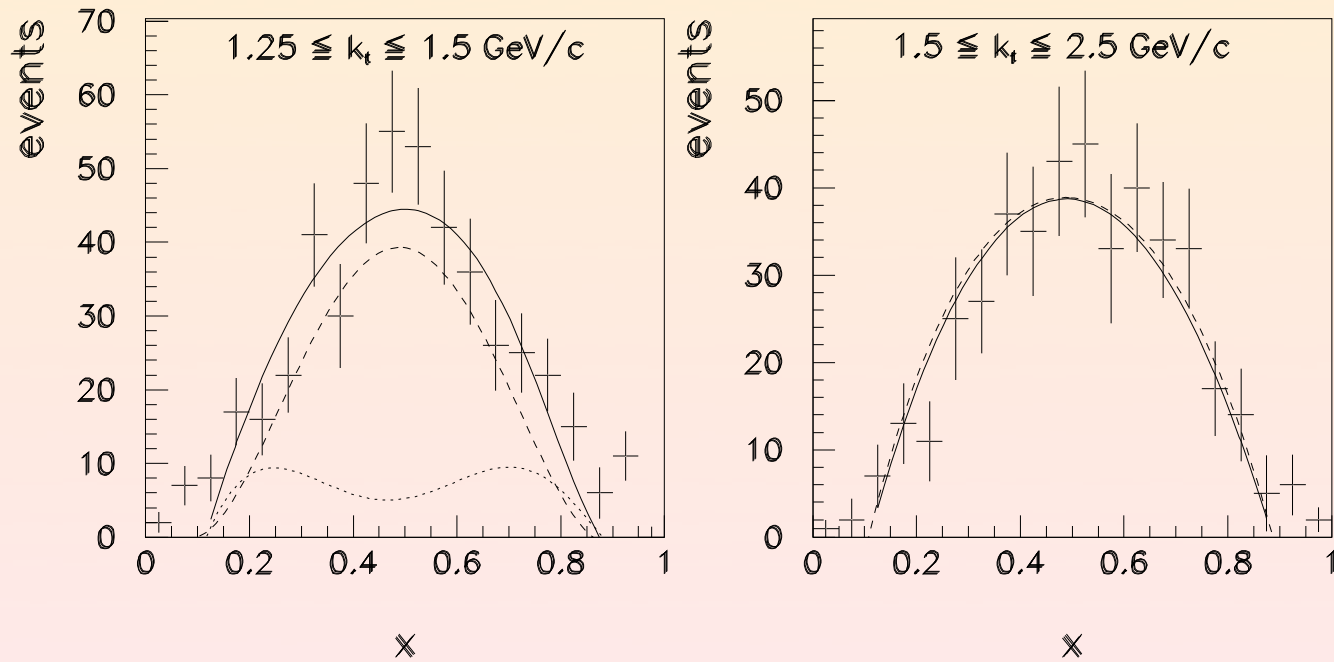
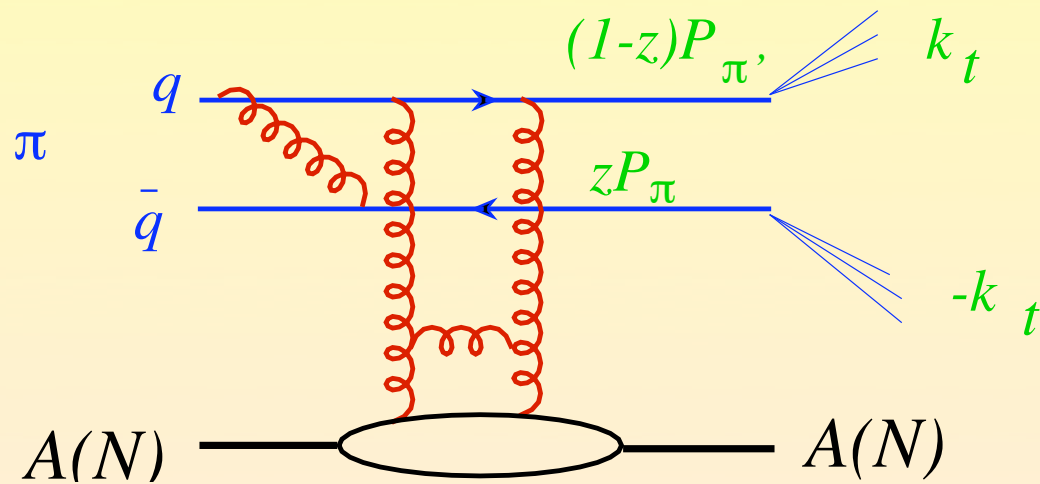
Very forward physics:  $p + A \rightarrow 3 \text{ jets} + A$

☀ Measurement of the proton wave function in  $|3q\rangle$  configuration with small transverse distance separation.

☀ Color transparency & Color opacity

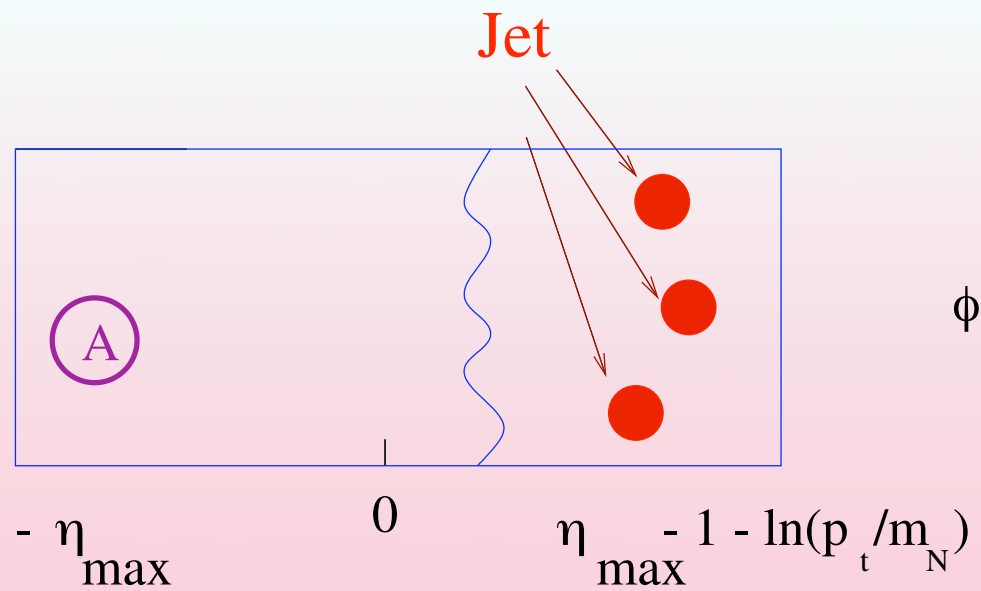
Extension of the study of  $\pi + N(A) \rightarrow "2 \text{ high } p_t \text{ jets}" + N(A)$

by FNAL 791 (2001), which confirmed Miller & F&S 93 prediction for the  $A$ -dependence,  $z$ -,  $p_t$ -dependence of the dijet coherent cross section.



Solid lines are the asymptotic (large  $p_t$ ) prediction:

$$\sigma(z) \propto \phi_\pi^2(z) \propto (1-z)^2 z^2$$



Lego plot for coherent 3 jet production  
in proton -nucleus scattering

$$\frac{d\sigma(pA \rightarrow (jet_1 + jet_2 + jet_3) + A)}{dt \prod_{i=1}^3 dx_i d^2 p_{ti}} \propto [\alpha_s x_A G_A(x_A, p_t^2)]^2.$$

$$\cdot \frac{\phi_N^2(x_1, x_2, x_3)}{\prod_{i=1}^3 p_i^4} \delta^2(\sum_{i=1}^3 \vec{p}_{ti} - q_t) \delta(\sum_{i=1}^3 x_i - 1) G_N^2(t) F_A^2(t),$$

where  $t = -q_t^2$ ,  $x_A = M_{3jet}^2/2s$ ,  $\phi_N$  nucleon 3 q wave function.

Advantages of the collider kinematics: it is easy to select coherent processes using zero angle neutron calorimeter. Energy dependence is very strong if there is no taming of gluon densities since

$$xG_A(x, Q^2 \sim 40 \text{ GeV}^2) \propto \frac{1}{\sqrt{x}} \quad \Longrightarrow \quad \sigma_{3 \text{ jet}} \propto s !!!$$

At RHIC one can hope to distinguish three jet configurations with

$$p_t \geq 3 \text{ GeV}/c \quad \text{corresponding to} \quad x = \frac{M_{3 \text{ jet}}^2}{s} = \frac{9p_t^2}{s} \sim 10^{-3}$$

For such transverse momenta shadowing is rather weak function of A

$$xG_A \propto A^{-0.1} \quad \Longrightarrow \quad \sigma_{3 \text{ jet}}(A) \propto A^{1.1} \text{ as compared to } \sigma_{\text{soft diffr.}} \propto A^{0.6}$$

The magnitude of the cross section is  $10^{-5}$  mb per nucleon.

## Comments:

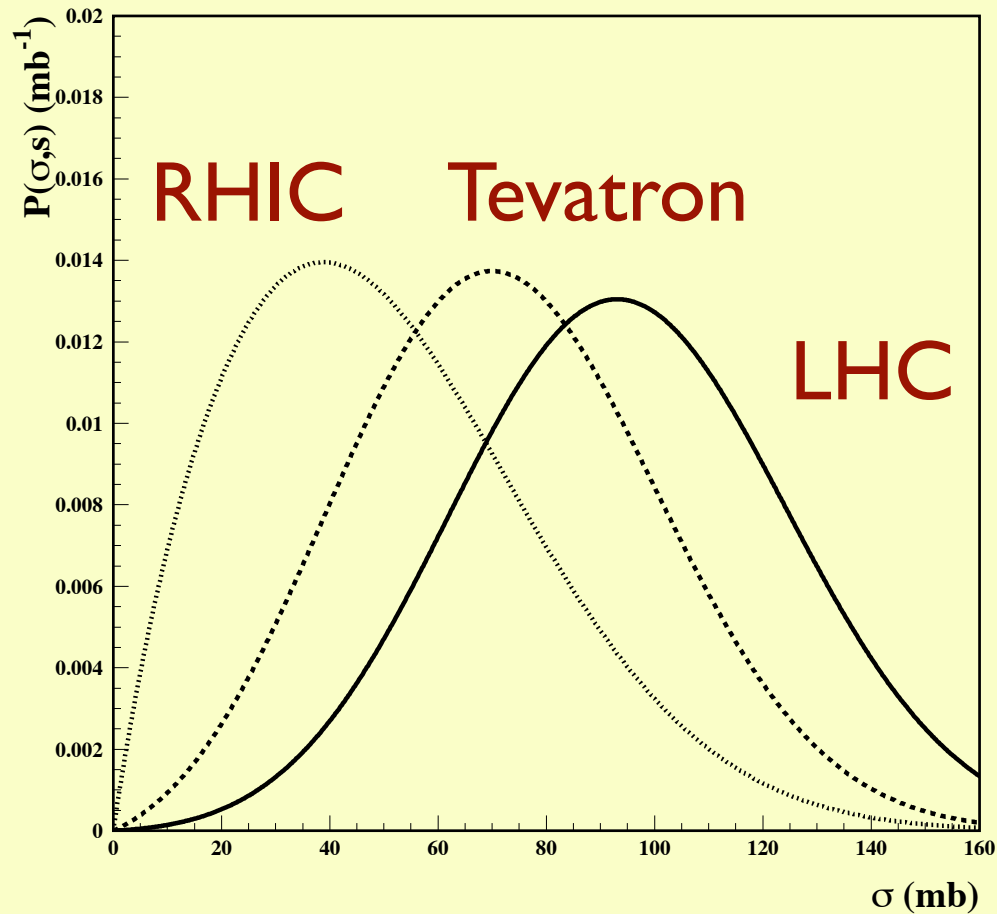
- If diquark-quark correlations are present in nucleons, one can look for 2 jet fragmentation.
- Interesting to search for exclusive channels at smaller  $p_t$  like  $p \rightarrow p\pi^+\pi^-, \Lambda K^+$  to find an effective trigger for interaction in point like configuration. If this would work can use double diffraction in pp scattering as a trigger for the scattering of two small dipoles.



The inelastic small  $t$  coherent diffraction off nuclei provides one of the most stringent tests of the presence of the fluctuations of the strength of the interaction in  $NN$  interactions. The answer is expressed through  $P(\sigma)$  - probability distribution for interaction with the strength  $\sigma$ . (Miller & FS 93)

$$\sigma_{diff}^{hA} = \int d^2b \left( \int d\sigma P_h(\sigma) |\langle h | F^2(\sigma, b) | h \rangle| - \left( \int d\sigma P(\sigma) |\langle h | F(\sigma, b) | h \rangle| \right)^2 \right).$$

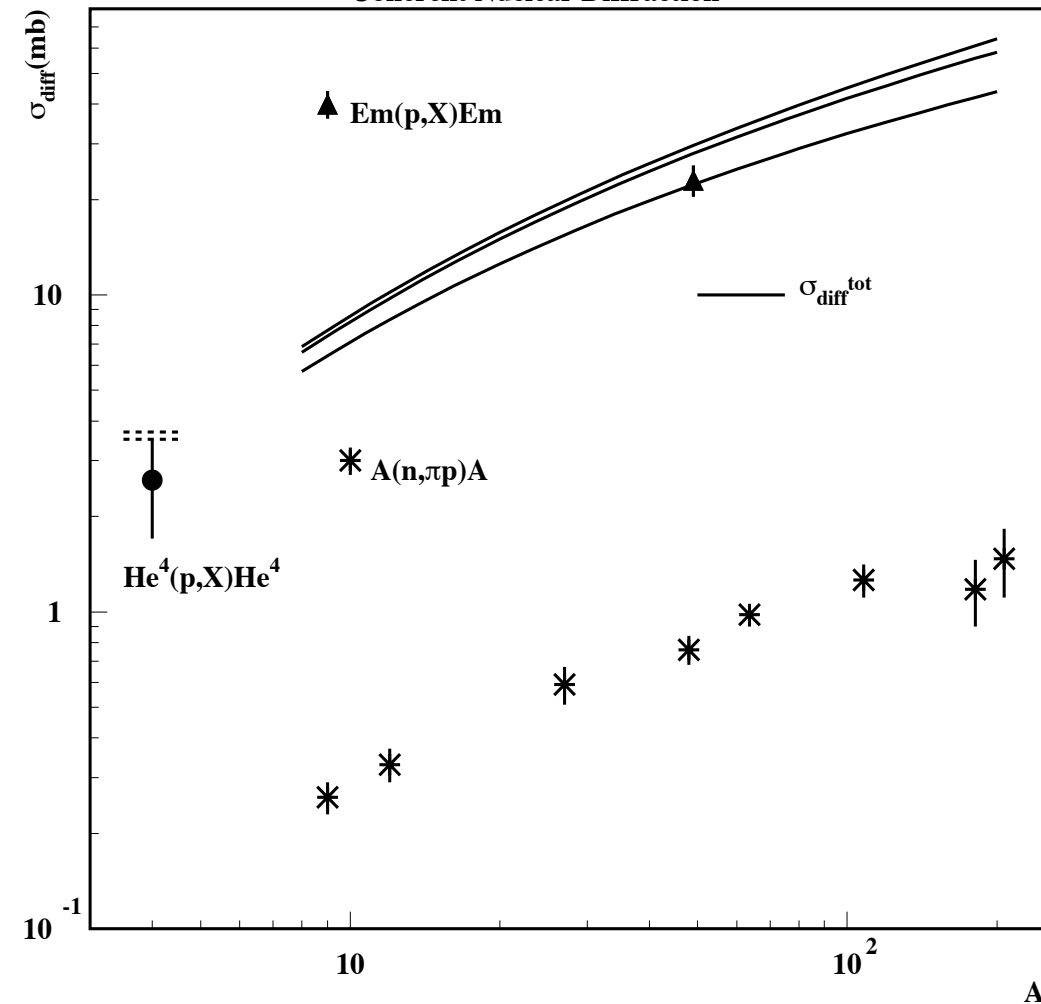
Here  $F(\sigma, b) = 1 - e^{-\sigma T(b)/2}$ ,  $T(b) = \int_{-\infty}^{\infty} \rho_A(b, z) dz$ , and  $\rho_A(b, z)$  is the nuclear density.



**Comment:** Fluctuations of strength of nucleon interaction lead to strong modifications of the distribution over the number of wounded nucleons (Baym, LF, MS 93)

Guzey and MS 05

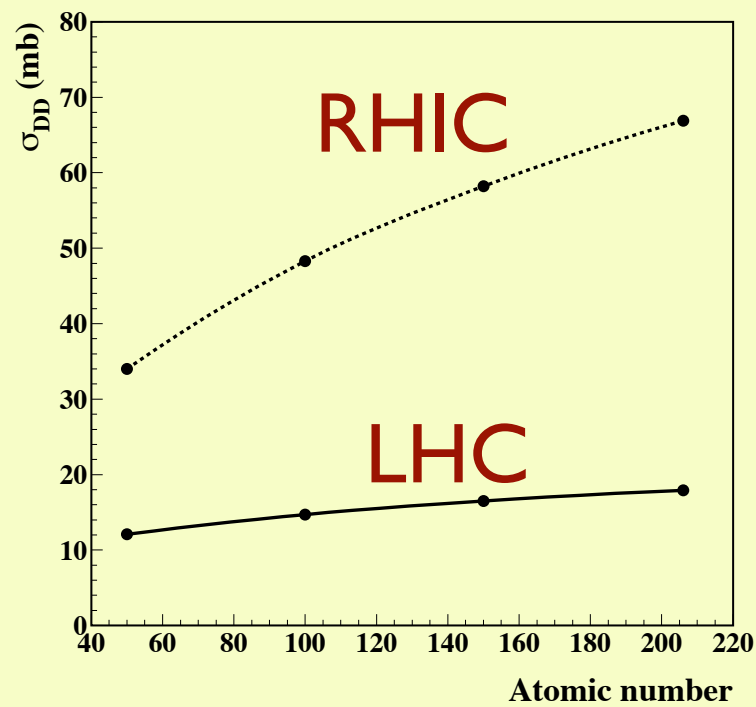
# Coherent Nuclear Diffraction



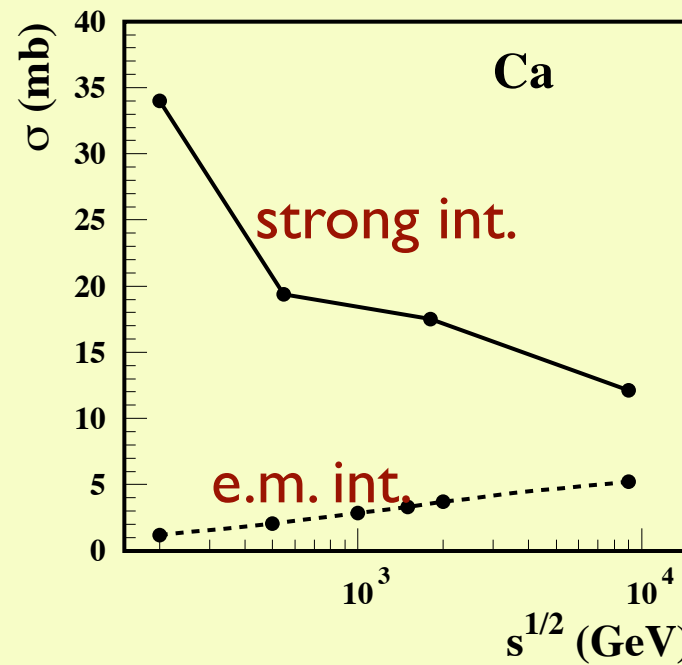
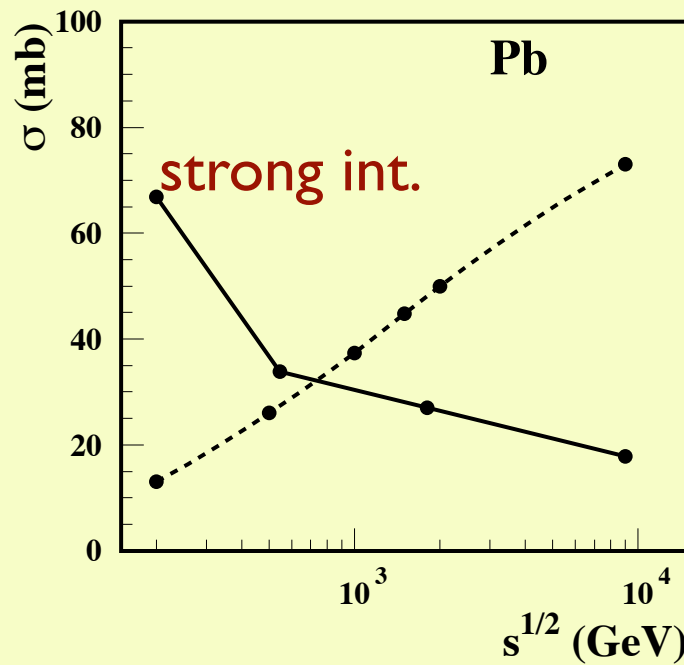
The cross section of coherent diffraction dissociation of protons and neutrons on nuclei as a function of the atomic number  $A$ . The solid lines are the theoretical prediction based on the above eqn. Total cross section data are from the FNAL emulsion and  $^4\text{He}$  jet target experiments. The FNAL data on the reaction  $n + A \rightarrow p\pi^- + A$  a small fraction of the total diffractive cross section is presented as stars have similar  $A$ -dependence for all masses and provide a good indication of the overall trend of the  $A$ -dependence. The theoretical prediction for coherent diffraction on  $^4\text{He}$  is given by the dashed lines.

⇒ Studies of inelastic diffraction in pA scattering at RHIC would allow to determine color fluctuations in protons at very high energies in the region of turnover to a much more black regime of TeV energies. The FMS prediction for RHIC  $\sigma^{inel.dif.} \propto A^{0.5-0.6}$  - significantly weaker A-dependence than for fixed target energies as a consequence of increase of  $\sigma_{tot}(pp)$ .

→ Note that a veto on production of forward neutrons is a very effective indicator of nucleus remaining intact.



Guzey and MS 05



e.m. int. = inelastic interaction of protons with the Coulomb field of the nucleus

E.m. mechanism is small at RHIC, however one will need to investigate whether there maybe a bigger relative contribution if one would investigate rare channels like those mentioned above.

For exclusive diffraction into three jets MS & LF found correction to be small even at LHC. So situation is likely to be the same for RHIC for wide range of channels. Still further studies are necessary. Clearly not a problem for lighter ions like Ca.

**Conclusions:** Diffraction is one more reason  
*for expanding the acceptance of the RHIC detectors (new  
detector?) in the very forward region.*

## *Diffraction in eA scattering*

Total cross section is closer to elastic case in pA than to inelastic simply because due to small e.m. coupling the process

$\gamma^* \Rightarrow h \Rightarrow \gamma^*$  is very small.

Virtual photon converts to configurations of very different size / strength of interaction. In inclusive DIS - bulk sum. In diffraction they enter separately.

Coherent diffraction: several limits probing different physics

In the limit of strong absorption - black body limit (which is not the same as saturation) - diffractive would reach 1/2 of the total cross section.



$$\frac{d\sigma_{diff}^{eA}}{dx dQ^2 dM^2 d\cos\theta} \propto \frac{\pi R_A^2}{12\pi^3} \frac{Q^4 \rho(M^2)}{(M^2 + Q^2)^2} \left( \frac{3}{8} \frac{M^2}{Q^2} (1 + \cos^2 \theta) + \frac{3}{4} \varepsilon \sin^2 \theta \right)$$

## Consequences:

Dominant final state - two jets with distribution over  $\cos \theta$   
like in  $e^+e^-$  annihilation (three-jets events are also like in  $e^+e^-$ )

Production of jets with large transverse momenta is enhanced relative to DGLAP

BBL for coherent exclusive vector meson (VM) electroproduction: *restoration of the vector dominance answer due to diagonality of transitions in the BBL !!!*

$$\frac{d\sigma^{\gamma_T^* + A \rightarrow V + A}}{dt} = \frac{M_V^2}{Q^2} \frac{d\sigma^{\gamma_L^* + A \rightarrow V + A}}{dt} = \frac{(2\pi R_A^2)^2}{16\pi} \frac{3\Gamma_V M_V^3}{\alpha(M_V^2 + Q^2)^2} \frac{4}{-t R_A^2} \left| J_1(\sqrt{-t} R_A) \right|^2$$

Q -dependence e.g. for  $\sigma_L$ :  $1/Q^2$  for BBL vs  $1/Q^6$  in the leading twist !!!

## Hard Exclusive Diffraction

**Motivation:** ♥ Establishing the  $Q^2$  range where QCD factorization is valid for exclusive processes. At what  $Q^2$  nucleon GPD's can be measured (C.Weiss talk)

♥ Comparison of nuclear GPDs and nuclear pdfs.

## Semiinclusive hard diffraction

**Motivation:** ♥ Much deeper understanding of small  $x$  nuclear dynamics via measurement of Pomeron- Pomeron interactions, nuclear diffractive pdfs.

☞ Exclusive diffraction - photon & vector meson production.  
 $x \geq 0.02$  Parton densities practically A-independent. Leading twist expectation for the amplitudes:

$$M(\gamma^* + A \rightarrow \gamma(V) + A) = A F_A(t) M(\gamma^* + A \rightarrow \gamma(V) + A)$$

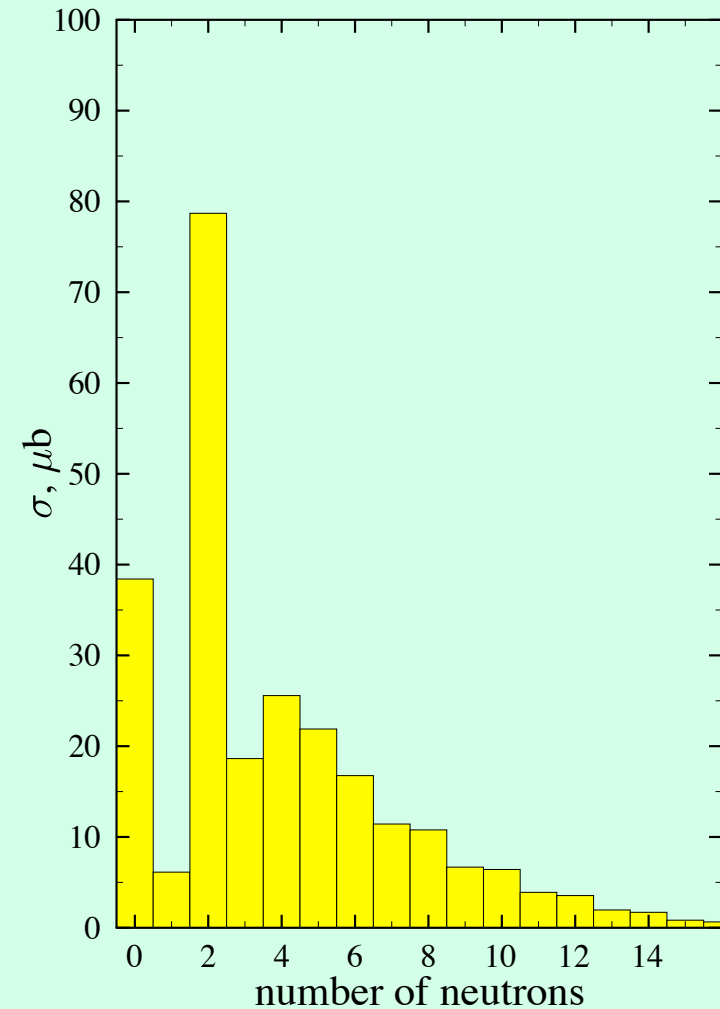
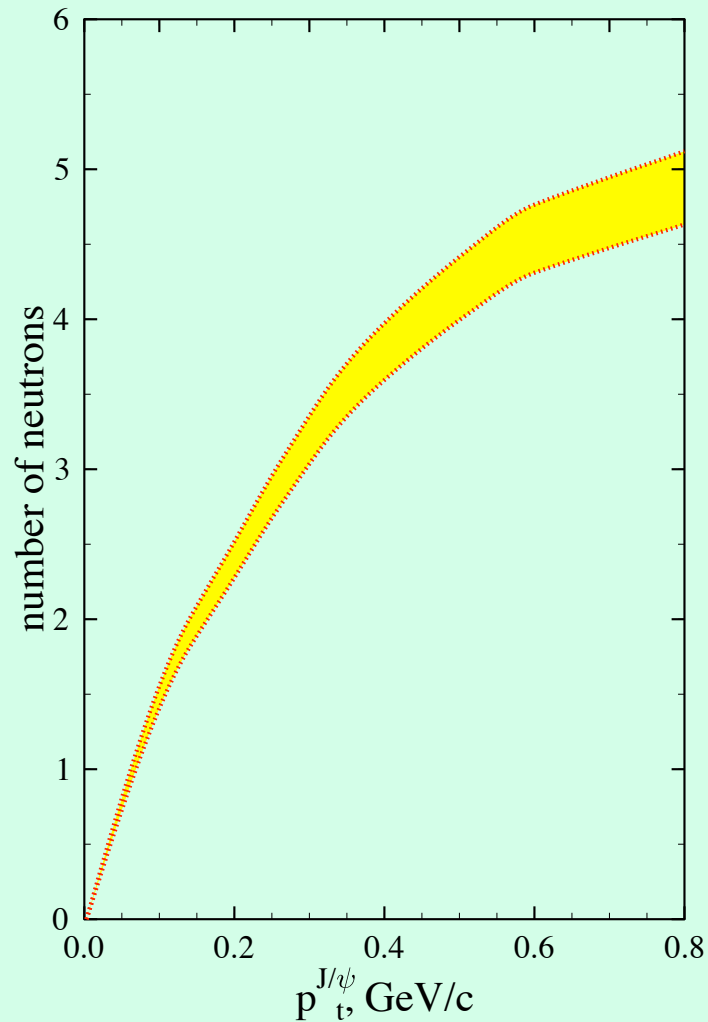
Observed for the  $J/\psi$  photoproduction at FNAL.

At what  $Q^2$  LT works for vector mesons? Leading twist dominates in the elementary reaction for  $Q^2 \geq 10 - 15 \text{ GeV}^2$  only. However, squeezing may work already for  $Q^2 \geq 5 \text{ GeV}^2$  (FKS95) leading to the color transparency phenomenon for much larger range of  $Q$  and hence to a possibility to measure the ratios of nucleon GPDs for a broad range of reactions.

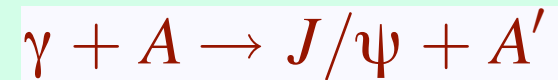
Alternative probe: incoherent exclusive processes for  $|t| \geq 0.1 \text{ GeV}^2$ :

$$\begin{aligned} d\sigma(\gamma^* + A \rightarrow \gamma(V, \pi, K, \dots) + A')/dt = \\ = Z d\sigma(\gamma^* + p \rightarrow \gamma(V, \pi, K, \dots) + N)/dt + (A-Z) d\sigma(\gamma^* + n \rightarrow \gamma(V, \pi, K, \dots) + N)/dt \end{aligned}$$

*Important advantage of collider kinematics - possibility to use ZDC to select the processes without nucleon excitation*



*Number of neutrons produced in the process*



Tverskoi, Zhalov, MS 05

☞ Exclusive diffraction -  $\gamma$  & VM production for  $x \leq 0.02$ .



DVCS on nuclei is doable - possible to measure both imaginary and real part.  
Large nuclear effects extend to  $x \sim 0.1$  for the real part. (Freund and MS 03)



Onset of perturbative color opacity:

$$M(\gamma^* + A \rightarrow V + A) = A F_A(t) M(\gamma^* + A \rightarrow V + A) G_A(x, Q^2) / G_N(x, Q^2)$$

First checks via ultraperipheral ion-ion collisions at LHC -

**$J/\psi$  photoproduction off nuclei - at  $x \geq 10^{-3}$**

EIC will be able to measure nuclear GPDs via DVCS and next to measure VM  
to see whether suppression is similar, and study it as a function of  $Q^2$  -  
interplay of HT and LT shadowing.

# Nuclear Diffractive parton densities

Collins factorization theorem: consider hard processes like

$$\gamma^* + T \rightarrow X + T(T'), \quad \gamma + T \rightarrow jet_1 + jet_2 + X + T$$

one can define conditional (fractional) parton distributions

$$f_j^D\left(\frac{x}{x_{\mathbb{P}}}, Q^2, x_{\mathbb{P}}, t\right) \quad \text{where} \quad x_{\mathbb{P}} \equiv 1 - x_{T_f}$$

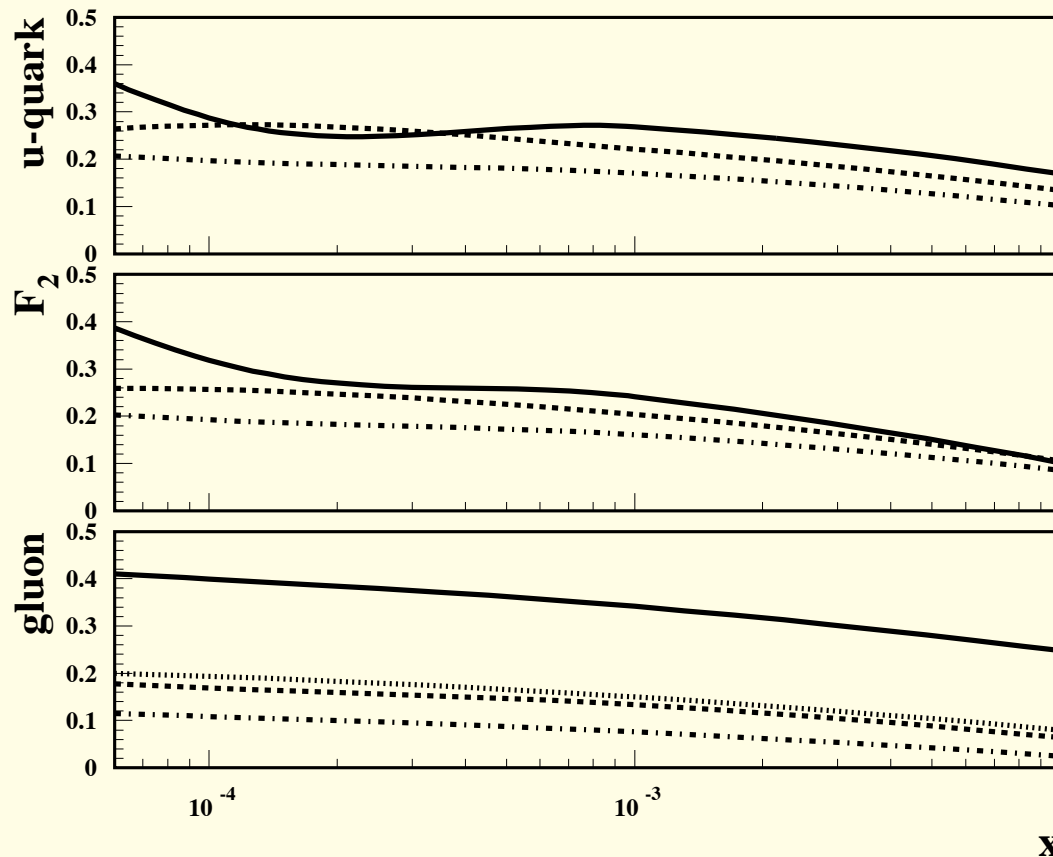
Theorem: for fixed  $x_{\mathbb{P}}, t$

the same Q evolution as for normal pdf's.

Nuclear diffractive pdfs were calculated by Guzey et al 03 in the same approximations as LT nuclear pdf's (no model necessary for double rescattering)



*Measurement of the leading twist inclusive diffraction off nuclei - expected to be a significant fraction of the cross section for  $x < 0.01$*



The ratios of diffractive and inclusive pdfs for the u-quarks and gluons and NLO  $F_{2A}$ 's for Pb. The solid curves

correspond to  $Q=2$  GeV; the dashed curves correspond to  $Q=10$  GeV; the dot-dashed curves correspond to  $Q=100$  GeV. In addition, for the gluons the dotted curve corresponds to  $Q=5$  GeV (Guzey, FS -03)

Note that maximal value of the ratio is  $1/2$  - black body limit. Both quark and gluon interactions with heavy nuclei are close to this limit for  $Q=2$  GeV !!!

EIC kinematics: To ensure coherence one needs

$$x/\beta < 0.01$$

Scaling - reasonably large produced mass  $> 2\text{GeV}$

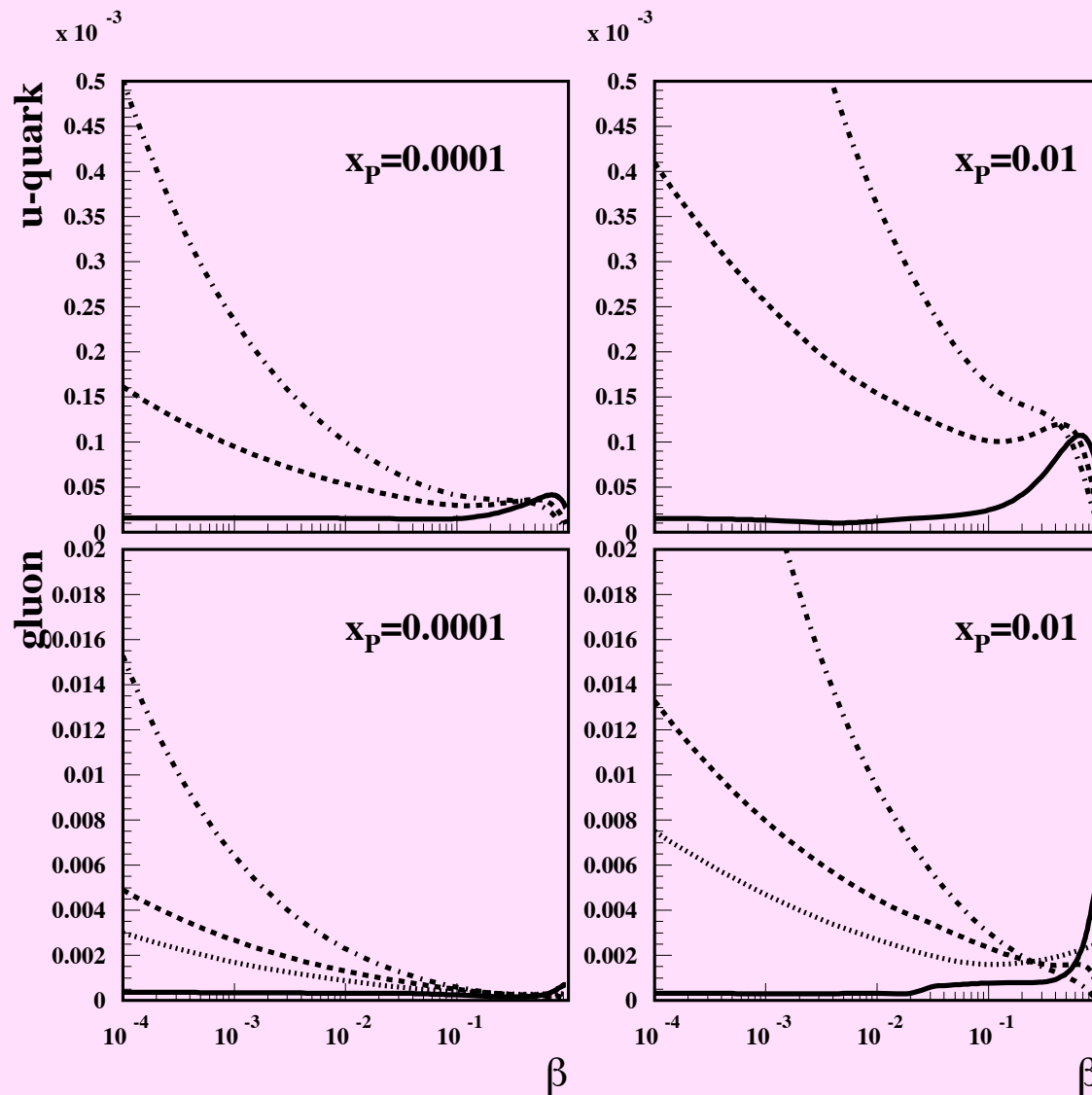
take  $x=0.005$ , which allows  $Q^2$  up to  $15\text{ GeV}^2$

$$\beta=0.5, \quad M^2 = Q^2$$

There is a phase space to study LT nuclear diffractive pdfs !!!  
Though it is rather limited would complement measurements  
at LHC using ultraperipheral collisions.

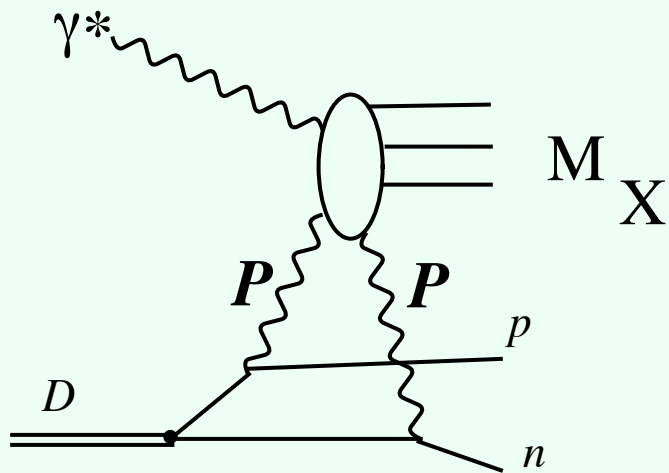


One also predicts a complicated dependence of the distribution on the diffractive mass as a function of *flavor*,  $Q$ , and  $\beta = Q^2 / (M^2 + Q^2)$ .



## 👉 SemiInclusive diffraction at $x \leq 0.02$ off deuteron

👉 *Double tagging for the deuteron in special kinematics - measure of “Pomeron-Pomeron” interactions.*



*Similar for the proton beams*

👉 *Coherent diffraction of the deuteron at large  $-t > 0.5 \text{ GeV}^2$*

# Conclusions

- ☺ EIC would allow to study diffraction in several kinematics relevant to the issues of color transparency, onset of color opacity ; leading twist diffraction in the kinematics which could not be probed at LHC; look for onset of black body limit regime.
- ☺ Various experimental tests of conjectured connection between diffraction in the scattering off nucleon/nuclei and nuclear shadowing
- ☺ Color transparency/opacity phenomena, first determinations of nuclear GPDs, establishing the  $Q^2$  range for studies of GPDs of nucleon in hard exclusive processes.
- ☺ pA diffractive processes: color fluctuations in nucleons should be studied - relevant for inelastic pA; looking for effective methods of triggering on small configurations in nucleons.



Many of the discussed measurements though doable from the angle of luminosity and kinematics are a challenge for detector design, systematics, etc.